

Spin Models Constructed from Hadamard matrices

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A new spin model M is constructed from an arbitrary Hadamard matrix H through a distance-regular graph which is called a Hadamard graph. F. Jaeger gives a formula for the link invariant of the model M , and V. F. R. Jones gives two links which have the same V-polynomial but different polynomials of M .

1 Definition of a Spin Model

The following definition is essentially due to V. F. R. Jones [8].

Definition 1 Let n be a positive integer, D be one of the square roots of n . A *spin model* with loop variable D is a pair (X, w) of a finite non-empty set X of size n , and a complex-valued symmetric function w on $X \times X$ which satisfy the following equations for all $\alpha, \beta, \gamma \in X$:

$$\frac{1}{n} \sum_{x \in X} \frac{w(\alpha, x)}{w(\beta, x)} = \delta_{\alpha, \beta} \quad (1)$$

$$\frac{1}{D} \sum_{x \in X} \frac{w(\alpha, x)w(\beta, x)}{w(\gamma, x)} = \frac{w(\alpha, \beta)}{w(\alpha, \gamma)w(\beta, \gamma)} \quad (2)$$

Each element of X is called a *spin*, and the function w is called *Boltzmann weight*. The $(n \times n)$ -matrix $W = (w(\alpha, \beta))$, is called the *weight matrix* of the spin model. The equation (2) is called *star-triangle relation*.

Example Let X be a finite set of size $n = D^2 > 1$ and let a, b be complex numbers such that

$$b^2 + \frac{1}{b^2} + D = 0, \quad a = -\frac{1}{b^3}.$$

Define a function w by

$$w(\alpha, \beta) = \begin{cases} a & \text{if } \alpha = \beta \\ b & \text{if } \alpha \neq \beta \end{cases}$$

As easily shown, (X, w) becomes a spin model with the weight matrix

$$M = (a - b)I + bJ.$$

This spin model is called *Potts model*.

Remark 1 If (X, w) is a spin model with $D = \sqrt{n}$, then $(X, \sqrt{-1}w)$ becomes a spin model with $D = -\sqrt{n}$.

Remark 2 Under (1), the star-triangle relation (2) is equivalent to:

$$\frac{1}{D} \sum_{x \in X} \frac{w(\alpha, x)}{w(\beta, x)w(\gamma, x)} = \frac{w(\alpha, \beta)w(\alpha, \gamma)}{w(\beta, \gamma)}. \quad (3)$$

Remark 3 By putting $\beta = \gamma$ in 2, we get

$$\frac{1}{D} \sum_{x \in X} w(\alpha, x) = \frac{1}{w(\beta, \beta)}.$$

This shows $w(\beta, \beta)$ is independent on the choice of $\beta \in X$:

$$w(\beta, \beta) = a$$

is a constant called *modulus* of the model. Thus we have

$$\frac{1}{D} \sum_{x \in X} w(\alpha, x) = \frac{1}{a}.$$

From 3, we have

$$\frac{1}{D} \sum_{x \in X} \frac{1}{w(\alpha, x)} = a.$$

Remark 4 The equation (1) is equivalent to

$$\sum_{x \in X} \frac{w(\alpha, x)}{w(\beta, x)} = 0 \quad \text{if } \alpha \neq \beta.$$

2 Spin Models on Distance-Regular Graphs

A connected graph Γ is said to be *distance-regular* if there are integers b_i, c_i ($i \geq 0$) such that for any two vertices u, x at distance $i = \partial(u, x)$, there are precisely c_i neighbours of x in $\Gamma_{i-1}(u)$ and b_i neighbours of x in $\Gamma_{i+1}(u)$. In particular, Γ is regular of valency $k = b_0$. The sequence

$$\iota(\Gamma) = \{b_0, b_1, \dots, b_{d-1}; c_1, c_2, \dots, c_d\},$$

where d is the diameter of Γ , is called the intersection array of G . For two vertices u, v , the size

$$p_{ij}^\alpha = |\Gamma_i(u) \cap \Gamma_j(v)|$$

depends only on the distance $\alpha = \partial(u, v)$, rather than the individual vertices u, v with $\partial(u, v) = \alpha$ (see [4] 4.1). In particular $k_i = |\Gamma_i(u)|$, which is called the *i-th valency*, does not depend on the choice of a vertex u . For three vertices u, v, w , put

$$P_{ij\ell}(u, v, w) = |\Gamma_i(u) \cap \Gamma_j(v) \cap \Gamma_\ell(w)|.$$

More precise descriptions about distance-regular graphs will be found in [3], [4].

The following Proposition is obtained directly from the definition and remarks in the previous section.

Proposition 1 Let Γ be a distance-regular graph of diameter d with the vertex set X . Put $|X| = n$ and let D be one of the square roots of n . Let t_0, t_1, \dots, t_d be non-zero complex numbers and let w be the complex valued function on $X \times X$ defined by $w(u, v) = t_i$ where $i = \partial(u, v)$. Then (X, w) becomes a spin model if and only if the following conditions hold:

$$(C1) \sum_{i=0}^d k_i t_i = D t_0^{-1},$$

$$(C2) \sum_{i=0}^d k_i t_i^{-1} = D t_0,$$

$$(C3) \sum_{i=0}^d \sum_{j=0}^d p_{ij}^\alpha t_i t_j^{-1} = 0 \quad (\alpha = 1, 2, \dots, d),$$

(C4) For all vertices u, v, w in X ,

$$\sum_{\ell=0}^d \sum_{i=0}^d \sum_{j=0}^d P_{ij\ell}(u, v, w) t_i t_j t_\ell^{-1} = D t_\alpha t_\beta^{-1} t_\gamma^{-1},$$

where $\alpha = \partial(u, v)$, $\beta = \partial(u, w)$, $\gamma = \partial(v, w)$.

Remark 5 Though conditions (C1) and (C2) can be removed in the above, these are useful to find solutions of the equations.

3 Result

A distance-regular graph having the intersection array

$$\{4m, 4m - 1, 2m, 1; 1, 2m, 4m - 1, 4m\}$$

is called a *Hadamard graph* of order $4m$. There is a natural one-to-one correspondence between Hadamard graphs of order $4m$ and Hadamard matrices of order $4m$ (see [4] 1.8). Now our main result follows:

Theorem 2 Let Γ be a Hadamard graph of order $4m$. Let s, t_0, t_1 be complex numbers such that

$$s^2 + 2(2m - 1)s + 1 = 0, \quad t_0^2 = \frac{2\sqrt{m}}{(4m - 1)s + 1}, \quad t_1^4 = 1.$$

Put $t_2 = s t_0$, $t_3 = -t_1$ and $t_4 = t_0$. Then t_0, \dots, t_4 satisfy the conditions in Proposition 1 with $D = 4\sqrt{m}$.

Theorem 2 can be described without using distance-regular graphs as follows:

Theorem 3 Let H be a Hadamard matrix of order n , $n \equiv 0 \pmod{4}$, and let M be the weight matrix of the Potts model of size n . Let ω be one of the 4-th roots of 1, $\omega^4 = 1$. Define a $4n \times 4n$ -matrix W as:

$$W = \begin{pmatrix} M & M & \omega H & -\omega H \\ M & M & -\omega H & \omega H \\ \omega H^t & -\omega H^t & M & M \\ -\omega H^t & \omega H^t & M & M \end{pmatrix}$$

Then W becomes the weight matrix of a spin model having $4n$ spins.

4 Proof of Theorem 2

Let H be a Hadamard graph of order $4m$ and let s, t_0, \dots, t_4 be complex numbers such that

$$s^2 + 2(2m - 1)s + 1 = 0, \quad t_0^2 = \frac{2\sqrt{m}}{(4m - 1)s + 1},$$

$$t_1^4 = 1, \quad t_2 = st_0, \quad t_3 = -t_1, \quad t_4 = t_0.$$

By $k_{i-1}b_{i-1} = k_i c_i$, we get

$$k_0 = 1, \quad k_1 = 4m, \quad k_2 = 8m - 2, \quad k_3 = 4m, \quad k_4 = 1.$$

So (C1) becomes

$$t_0 + 4mt_1 + (8m - 2)t_2 + 4mt_3 + t_4 = 4\sqrt{m}t_0^{-1}.$$

By $t_3 = -t_1, t_0 = t_4$ and $t_2 = st_0$, this becomes

$$2t_0 + (8m - 2)st_0 = 4\sqrt{m}t_0^{-1}.$$

Clearly this holds by the assumption $t_0^2 = 2\sqrt{m}((4m - 1)s + 1)^{-1}$.

Condition (C2) becomes

$$t_0^{-1} + 4mt_1^{-1} + (8m - 2)t_2^{-1} + 4mt_3^{-1} + t_4^{-1} = 4\sqrt{m}t_0,$$

and it becomes

$$2t_0^{-1} + (8m - 2)t_2^{-1} = 4\sqrt{m}t_0,$$

$$1 + (4m - 1)s^{-1} = 2\sqrt{m}t_0^2.$$

By the assumption $t_0^2 = 2\sqrt{m}((4m - 1)s + 1)^{-1}$, it is equivalent to

$$1 + (4m - 1)s^{-1} = 2\sqrt{m} \cdot 2\sqrt{m}((4m - 1)s + 1)^{-1}.$$

This is implied by the assumption $s^2 + 2(2m - 1)s + 1 = 0$.

Next consider condition (C3). The values of p_{ij}^α are easily computed by the following formula ([4] 4.1.7).

$$p_{j+1,\ell}^\alpha = \frac{1}{c_{j+1}}(p_{j,\ell-1}^\alpha b_{\ell-1} + p_{j,\ell+1}^\alpha c_{\ell+1} - p_{j-1,\ell}^\alpha b_{j-1}).$$

Case $\alpha = 1$:

(i, j)	p_{ij}^1
$(0, 1), (1, 0), (3, 4), (4, 3)$	1
$(1, 2), (2, 1), (2, 3), (3, 2)$	$4m - 1$

Condition (C3) becomes

$$t_0 t_1^{-1} + t_1 t_0^{-1} + t_3 t_4^{-1} + t_4 t_3^{-1} + (4m - 1)(t_1 t_2^{-1} + t_2 t_1^{-1} + t_2 t_3^{-1} + t_3 t_2^{-1}) = 0.$$

This holds by $t_3 = -t_1$ and $t_0 = t_4$.

Case $\alpha = 2$:

(i, j)	p_{ij}^2
$(0, 2), (2, 0), (2, 4), (4, 2)$	1
$(1, 1), (1, 3), (3, 1), (3, 3)$	$2m$
$(2, 2)$	$8m - 4$

(C3) becomes

$$t_0 t_2^{-1} + t_2 t_0^{-1} + t_2 t_4^{-1} + t_4 t_2^{-1} + 2m(t_1 t_1^{-1} + t_1 t_3^{-1} + t_3 t_1^{-1} + t_3 t_3^{-1}) + (8m - 4) = 0.$$

This is implied by $t_3 = -t_1$, $t_0 = t_4$, $t_2 = st_0$ and $s^2 + 2(2m - 1)s + 1 = 0$.

Case $\alpha = 3$:

(i, j)	p_{ij}^3
$(0, 3), (3, 0), (1, 4), (4, 1)$	1
$(1, 2), (2, 1), (2, 3), (3, 2)$	$4m - 1$

$$t_0 t_3^{-1} + t_3 t_0^{-1} + t_1 t_4^{-1} + t_4 t_1^{-1} + (4m - 1)(t_1 t_2^{-1} + t_2 t_1^{-1} + t_2 t_3^{-1} + t_3 t_2^{-1}) = 0.$$

This holds by $t_3 = -t_1$ and $t_0 = t_4$.

Case $\alpha = 4$:

(i, j)	p_{ij}^4
$(0, 4), (4, 0)$	1
$(1, 3), (3, 1)$	$4m$
$(2, 2)$	$8m - 2$

$$t_0 t_4^{-1} + t_4 t_0^{-1} + 4m(t_1 t_3^{-1} + t_3 t_1^{-1}) + (8m - 2)t_2 t_2^{-1} = 0.$$

Clearly this holds.

Now we consider condition (C4). Since (C4) is symmetric in u, v , we may assume $\partial(u, w) \leq \partial(v, w)$. Fix three vertices u, v, w . Put $\partial(u, v) = \alpha$, $\partial(u, w) = \beta$, $\partial(v, w) = \gamma$ and $P_{ij\ell} = P_{ij\ell}(u, v, w)$. If $\beta = 0$, we have $u = w$, $\alpha = \gamma$, and $P_{ij\ell} = 0$ for $i \neq \ell$. Therefore

$$\sum_{i,j,\ell} P_{ij\ell} t_i t_j t_\ell^{-1} = \sum_j \sum_i P_{ij i} t_j = \sum_j k_j t_j,$$

and (C4) is equivalent to (C1) in the case $\beta = 0$. So we must verify (C4) in each of the following cases of (α, β, γ) :

$$\begin{array}{ccccccc} (0, 1, 1) & (0, 2, 2) & (0, 3, 3) & (0, 4, 4) & & & \\ (1, 1, 2) & (1, 2, 3) & (1, 3, 4) & & & & \\ (2, 1, 1) & (2, 1, 3) & (2, 2, 2) & (2, 2, 4) & (2, 3, 3) & & \\ (3, 1, 2) & (3, 1, 4) & (3, 2, 3) & & & & \\ (4, 1, 3) & (4, 2, 2) & & & & & \end{array}$$

In the case $(\alpha, \beta, \gamma) \neq (2, 2, 2)$, the values of $P_{ij\ell}$ are easily computed. We need the following Lemma for the case $(\alpha, \beta, \gamma) = (2, 2, 2)$.

Lemma 4 *If $\partial(u, v) = \partial(u, w) = \partial(v, w) = 2$, then w has precisely m neighbours in $\Gamma_1(u) \cap \Gamma_1(v)$.*

Proof. Put $D_j^i = \Gamma_i(u) \cap \Gamma_j(v)$. We have $w \in D_2^2$. Put $e(w, D_1^1) = r$, $e(w, D_3^1) = s$, $e(w, D_1^3) = s'$, $e(w, D_3^3) = r'$. Notice that every vertex $x \in X$ has the unique *opposite* vertex x' such that $\partial(x, x') = 4$, since we have $k_4 = 1$. Since the opposite vertex x' of $x \in D_1^1 \cap \Gamma_1(w)$ is in D_3^3 , we get $r' \leq |D_3^3| - r = 2m - r$. Similarly we get $s' \leq 2m - s$. On the other hand, we have $r + s = 2m$ since w has precisely $2m$ neighbours in $\Gamma_1(u)$. We have also $s + r' = 2m$ since w has $2m$ neighbours in $\Gamma_3(v)$. These imply $r = r'$. By the same reason, we get $s = s'$. Therefore we must have $r = s = r' = s' = m$.

Case $(\alpha, \beta, \gamma) = (0, 1, 1)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 0, 1), (1, 1, 0), (3, 3, 4), (4, 4, 3)$	1
$(1, 1, 2), (2, 2, 1), (2, 2, 3), (3, 3, 2)$	$4m - 1$

So, condition (C4) becomes

$$t_0^2 t_1^{-1} + t_1^2 t_0^{-1} + t_3^2 t_4^{-1} + t_4^2 t_3^{-1} + (4m - 1)(t_1^2 t_2^{-1} + t_2^2 t_1^{-1} + t_2^2 t_3^{-1} + t_3^2 t_2^{-1}) = Dt_0 t_1^{-2},$$

$$2t_1^2 t_0^{-1} + (8m - 2)t_1^2 t_2^{-1} = Dt_0 t_1^{-2}.$$

By $t_1^4 = 1$, this is equivalent to (C2).

Case $(\alpha, \beta, \gamma) = (0, 2, 2)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 0, 2), (2, 2, 0), (2, 2, 4), (4, 4, 2)$	1
$(1, 1, 1), (1, 1, 3), (3, 3, 1), (3, 3, 3)$	$2m$
$(2, 2, 2)$	$8m - 4$

Then condition (C4) becomes

$$2(t_0^2 t_2^{-1} + t_2^2 t_0^{-1}) + (8m - 4)t_2 = Dt_0 t_2^{-2},$$

$$s^{-1} + s^2 + (4m - 2)s = 2\sqrt{m} s^{-2} t_0^{-2}.$$

By the assumption $t_0^2 = 2\sqrt{m}((4m - 1)s + 1)^{-1}$, this becomes

$$s^{-1} + s^2 + (4m - 2)s = (4m - 1)s^{-1} + s^{-2}.$$

This is implied by the assumption $s^2 + 2(2m - 1)s + 1 = 0$.

Case $(\alpha, \beta, \gamma) = (0, 3, 3)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 0, 3), (1, 1, 4), (3, 3, 0), (4, 4, 1)$	1
$(1, 1, 2), (2, 2, 1), (2, 2, 3), (3, 3, 2)$	$4m - 1$

(C4) becomes

$$t_0^2 t_3^{-1} + t_1^2 t_4^{-1} + t_3^2 t_0^{-1} + t_4^2 t_1^{-1} + (4m - 1)(t_1^2 t_2^{-1} + t_2^2 t_1^{-1} + t_2^2 t_3^{-1} + t_3^2 t_2^{-1}) = Dt_0 t_3^{-2}.$$

This is equivalent to Case $(\alpha, \beta, \gamma) = (0, 1, 1)$.

Case $(\alpha, \beta, \gamma) = (0, 4, 4)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 0, 4), (4, 4, 0)$	1
$(1, 1, 3), (3, 3, 1)$	$4m$
$(2, 2, 2)$	$8m - 2$

$$t_0^2 t_4^{-1} + t_4^2 t_0^{-1} + 4m(t_1^2 t_3^{-1} + t_3^2 t_1^{-1}) + (8m - 2)t_2^2 t_2^{-1} = Dt_0 t_4^{-2}.$$

Case $(\alpha, \beta, \gamma) = (1, 1, 2)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 1, 1), (1, 0, 2), (1, 2, 0), (3, 2, 4), (3, 4, 2), (4, 3, 3)$	1
$(2, 1, 1), (2, 3, 3)$	$2m - 1$
$(2, 1, 3), (2, 3, 1)$	$2m$
$(1, 2, 2), (3, 2, 2)$	$4m - 2$

$$t_0 + t_4 + t_0 t_1 t_2^{-1} + t_1 t_2 t_0^{-1} + t_2 t_3 t_4^{-1} + t_3 t_4 t_2^{-1} + 2m(t_1 t_2 t_3^{-1} + t_2 t_3 t_1^{-1}) \\ + (4m - 2)(t_1 + t_2 + t_3) = Dt_2^{-1}.$$

Case $(\alpha, \beta, \gamma) = (1, 2, 3)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 1, 2), (1, 0, 3), (2, 1, 4), (2, 3, 0), (3, 4, 1), (4, 3, 2)$	1
$(1, 2, 3), (3, 2, 1)$	$2m - 1$
$(1, 2, 1), (3, 2, 3)$	$2m$
$(2, 1, 2), (2, 3, 2)$	$4m - 2$

$$t_0 t_1 t_2^{-1} + t_0 t_1 t_3^{-1} + t_1 t_2 t_4^{-1} + t_2 t_3 t_0^{-1} + t_3 t_4 t_1^{-1} + t_3 t_4 t_2^{-1} \\ + (2m - 1)(t_1 t_2 t_3^{-1} + t_2 t_3 t_1^{-1}) + (4m - 2)(t_1 + t_3) + 4m t_2 = Dt_1 t_2^{-1} t_3^{-1}.$$

Case $(\alpha, \beta, \gamma) = (1, 3, 4)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 1, 3), (1, 0, 4), (3, 4, 0), (4, 3, 1)$	1
$(1, 2, 2), (2, 1, 3), (2, 3, 1), (3, 2, 2)$	$4m - 1$

$$t_0 t_1 t_3^{-1} + t_0 t_1 t_4^{-1} + t_3 t_4 t_0^{-1} + t_3 t_4 t_1^{-1} + (4m - 1)(t_1 + t_3 + t_1 t_2 t_3^{-1} + t_2 t_3 t_1^{-1}) \\ = Dt_1 t_3^{-1} t_4^{-1}.$$

Case $(\alpha, \beta, \gamma) = (2, 1, 1)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 2, 1), (2, 0, 1), (2, 4, 3), (4, 2, 3), (1, 1, 0), (3, 3, 4)$	1
$(1, 1, 2), (3, 3, 2)$	$2m - 1$
$(1, 3, 2), (3, 1, 2)$	$2m$
$(2, 2, 1), (2, 2, 3)$	$4m - 2$

$$t_1^2 t_0^{-1} + t_3^2 t_4^{-1} + 2(t_0 t_2 t_1^{-1} + t_2 t_4 t_3^{-1}) + (2m - 1)(t_1^2 t_2^{-1} + t_3^2 t_2^{-1}) \\ + (4m - 2)(t_2^2 t_1^{-1} + t_2^2 t_3^{-1}) + 4m t_1 t_3 t_2^{-1} = Dt_2 t_1^{-2}.$$

Case $(\alpha, \beta, \gamma) = (2, 1, 3)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 2, 1), (2, 0, 3), (2, 4, 1), (4, 2, 3), (1, 3, 0), (3, 1, 4)$	1
$(1, 3, 2), (3, 1, 2)$	$2m - 1$
$(1, 1, 2), (3, 3, 2)$	$2m$
$(2, 2, 1), (2, 2, 3)$	$4m - 2$

$$t_0 t_2 t_1^{-1} + t_0 t_2 t_3^{-1} + t_2 t_4 t_1^{-1} + t_2 t_4 t_3^{-1} + t_1 t_3 t_0^{-1} + t_1 t_3 t_4^{-1} + 2m(t_1^2 t_2^{-1} + t_3^2 t_2^{-1}) \\ + (4m - 2)(t_1 t_3 t_2^{-1} + t_2^2 t_1^{-1} + t_2^2 t_3^{-1}) = Dt_2 t_1^{-1} t_3^{-1}.$$

Case $(\alpha, \beta, \gamma) = (2, 2, 2)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 2, 2), (2, 0, 2), (2, 2, 0), (2, 2, 4), (2, 4, 2), (4, 2, 2)$	1
$(1, 1, 1), (1, 1, 3), (1, 3, 1), (1, 3, 3)$	m
$(3, 1, 1), (3, 1, 3), (3, 3, 1), (3, 3, 3)$	m
$(2, 2, 2)$	$8m - 6$

$$t_2^2 t_0^{-1} + t_2^2 t_4^{-1} + 2(t_0 + t_4) + m(t_1^2 t_3^{-1} + t_3^2 t_1^{-1}) + 3m(t_1 + t_3) \\ + (8m - 6)t_2 = Dt_2^{-1}.$$

Case $(\alpha, \beta, \gamma) = (2, 2, 4)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 2, 2), (2, 0, 4), (2, 4, 0), (4, 2, 2)$	1
$(1, 1, 3), (1, 3, 1), (3, 1, 3), (3, 3, 1)$	$2m$
$(2, 2, 2)$	$8m - 4$

$$t_0 + t_4 + t_0 t_2 t_4^{-1} + t_2 t_4 t_0^{-1} + 2m(t_1 + t_3 + t_1^2 t_3^{-1} + t_3^2 t_1^{-1}) + (8m - 4)t_2 = Dt_4^{-1}.$$

Case $(\alpha, \beta, \gamma) = (2, 3, 3)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 2, 3), (1, 1, 4), (2, 0, 3), (2, 4, 1), (3, 3, 0), (4, 2, 1)$	1
$(1, 1, 2), (3, 3, 2)$	$2m - 1$
$(1, 3, 2), (3, 1, 2)$	$2m$
$(2, 2, 1), (2, 2, 3)$	$4m - 2$

$$t_1^2 t_4^{-1} + t_3^2 t_0 + 2(t_0 t_2 t_3^{-1} + t_2 t_4 t_1^{-1}) + (2m - 1)(t_1^2 t_2^{-1} + t_3^2 t_2^{-1}) \\ + (4m - 2)(t_2^2 t_1^{-1} + t_2^2 t_3^{-1}) + 4m t_1 t_3 t_2^{-1} = D t_2 t_3^{-2}.$$

Case $(\alpha, \beta, \gamma) = (3, 1, 2)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 3, 1), (1, 2, 0), (3, 0, 2), (1, 4, 2), (3, 2, 4), (4, 1, 3)$	1
$(2, 1, 3), (2, 3, 1)$	$2m - 1$
$(2, 1, 1), (2, 3, 3)$	$2m$
$(1, 2, 2), (3, 2, 2)$	$4m - 2$

$$t_0 t_3 t_1^{-1} + t_0 t_3 t_2^{-1} + t_1 t_2 t_0^{-1} + t_1 t_4 t_2^{-1} + t_1 t_4 t_3^{-1} + t_2 t_3 t_4^{-1} \\ + (2m - 1)(t_1 t_2 t_3^{-1} + t_2 t_3 t_1^{-1}) + (4m - 2)(t_1 + t_3) + 4m t_2 = D t_3 t_1^{-1} t_2^{-1}.$$

Case $(\alpha, \beta, \gamma) = (3, 1, 4)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 3, 1), (3, 0, 4), (1, 4, 0), (4, 1, 3)$	1
$(1, 2, 2), (2, 1, 3), (2, 3, 1), (3, 2, 2)$	$4m - 1$

$$t_0 t_3 t_1^{-1} + t_0 t_3 t_4^{-1} + t_1 t_4 t_0^{-1} + t_1 t_4 t_3^{-1} + (4m - 1)(t_1 + t_3) \\ + (4m - 1)(t_1 t_2 t_3^{-1} + t_2 t_3 t_1^{-1}) = D t_3 t_1^{-1} t_4^{-1}.$$

Case $(\alpha, \beta, \gamma) = (3, 2, 3)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 3, 2), (2, 1, 4), (3, 0, 3), (1, 4, 1), (2, 3, 0), (4, 1, 2)$	1
$(1, 2, 1), (3, 2, 3)$	$2m - 1$
$(1, 2, 3), (3, 2, 1)$	$2m$
$(2, 1, 2), (2, 3, 2)$	$4m - 2$

$$t_0 + t_4 + t_0 t_3 t_2^{-1} + t_2 t_3 t_0^{-1} + t_1 t_2 t_4^{-1} + t_1 t_4 t_2^{-1} + 2m(t_1 t_2 t_3^{-1} + t_2 t_3 t_1^{-1}) \\ + (4m - 2)(t_1 + t_2 + t_3) = D t_2^{-1}.$$

Case $(\alpha, \beta, \gamma) = (4, 1, 3)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 4, 1), (1, 3, 0), (3, 1, 4), (4, 0, 3)$	1
$(1, 3, 2), (2, 2, 1), (2, 2, 3), (3, 1, 2)$	$4m - 1$

$$t_0 t_4 t_1^{-1} + t_0 t_4 t_3^{-1} + t_1 t_3 t_0^{-1} + t_1 t_3 t_4^{-1} + (4m - 1)(t_2^2 t_1^{-1} + t_2^2 t_3^{-1}) \\ + (8m - 2)t_1 t_3 t_2^{-1} = D t_4 t_1^{-1} t_3^{-1}.$$

Case $(\alpha, \beta, \gamma) = (4, 2, 2)$:

(i, j, ℓ)	$P_{ij\ell}$
$(0, 4, 2), (2, 2, 0), (2, 2, 4), (4, 0, 2)$	1
$(1, 3, 1), (1, 3, 3), (3, 1, 1), (3, 1, 3)$	$2m$
$(2, 2, 2)$	$8m - 4$

$$t_2^2 t_0^{-1} + t_2^2 t_4^{-1} + 2t_0 t_4 t_2^{-1} + 4m(t_1 + t_3) + (8m - 4)t_2 = D t_4 t_2^{-2}.$$

参考文献

- [1] S. S. AGAIAN, "Hadamard matrices and their applications," Lecture Notes in Math. 1168, Springer, Berlin, 1985.
- [2] E. BANNAI, E. BANNAI, Spin models on finite cyclic groups, *preprint*.
- [3] E. BANNAI AND T. ITO, "Algebraic Combinatorics I," Benjamin, London, 1984.
- [4] A. E. BROUWER, A. M. COHEN AND A. NEUMAIER, Distance-regular graphs, Springer-Verlag, Berlin, Heidelberg, 1989.
- [5] P. DE LA HARPE, Spin models for link polynomials, strongly regular graphs and Jaeger's Higman-Sims model, *preprint*.
- [6] P. DE LA HARPE AND V. F. R. JONES, Graph invariants related to statistical mechanical models : examples and problems, *J. Combinatorial Theory (B)*, to appear
- [7] F. JAEGER, Strongly regular graphs and spin models for the Kauffman polynomial, *Geom. Dedicata*, to appear.
- [8] V. F. R. JONES, On knot invariants related to some statistical mechanical models, *Pac. J. Math.* **137** (1989), 311-336.