GLOBAL CONTINUATION BEYOND SINGULARITY ON THE BOUNDARY

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1. INTRODUCTION

We consider problems of the form

$u_t = u_{xx},$	0 < x < 1,	0 < t < T,
$u_x(0,t)=0,$	0 < t < T,	
$u_x(1,t) = f(u(1,t)),$	0 < t < T,	
$u(x,0) = u_0(x) > 0,$	$0 \le x \le 1,$	

where $f(u) = -u^{-p}$, p > 0, or $f(u) = u^{p}$, p > 1. We shall call them Problem (Q) and Problem (B), respectively. We discuss them separately.

1.1. Problem (Q) $(f(u) = -u^{-p})$. This problem was studied before by Fila & Levine(1993) where it was shown that that every solution quenches in a finite time $T = T(u_0)$ in the sense that u > 0 in $[0, 1] \times [0, T)$ and $u(1, t) \to 0$ as $t \to T$. The behavior of u near (1, T) for $t \leq T$ was also studied.

The question whether it is possible to continue the solution beyond t = T (in some suitable sense) was raised by Levine(1993). Since $u(\cdot, T) \in C([0, 1])$ and u(1, T) = 0, an obvious possibility of continuing the solution is to extend it for t > T by \tilde{u} which solves

$$\begin{split} \tilde{u}_t &= \tilde{u}_{xx}, & 0 < x < 1, \quad t > T, \\ \tilde{u}_x(0,t) &= 0, & t > T, \\ \tilde{u}(1,t) &= 0, & t > T, \\ \tilde{u}(x,T) &= u(x,T), & 0 \le x \le 1. \end{split}$$

We show that this continuation is natural since it can be obtained as a limit of a sequence of solutions of regularized problems. More precisely, if $\varepsilon > 0$ and $f_{\varepsilon} \in C^1([0,\infty))$ is such that $f_{\varepsilon}(0) = 0$ and

$$\begin{aligned} f_{\varepsilon}(s) &= -s^{-p} & \text{for } s \ge \varepsilon, \\ f(s) &\leq f_{\varepsilon_1}(s) \le f_{\varepsilon_2}(s) & \text{for } s > 0 & \text{and } \varepsilon_1 < \varepsilon_2, \end{aligned}$$

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then the solutions of (Q_{ε}) :

$$\begin{cases} u_t^{\varepsilon} = u_{xx}^{\varepsilon}, & 0 < x < 1, \quad 0 < t < \infty, \\ u_x^{\varepsilon}(0,t) = 0, & 0 < t < \infty, \\ u_x^{\varepsilon}(1,t) = f_{\varepsilon}(u^{\varepsilon}(1,t)), & 0 < t < \infty, \\ u^{\varepsilon}(x,0) = u_0(x), & 0 \le x \le 1, \end{cases}$$

converge to the extension of u by \tilde{u} .

The fact that solutions of Problem (Q) can be continued beyond t = T for all p > 0 is in contrast with the situation when quenching occurs in the interior. Namely, for the problem

$$u_t = u_{xx} - u^{-p}, \qquad 0 < x < 1, \quad 0 < t < T,$$

$$u_x(0,t) = 0, \qquad 0 < t < T,$$

$$u(1,t) = 1, \qquad 0 < t < T,$$

$$u(x,0) = u_0(x), \qquad 0 < x < 1,$$

solutions can be continued beyond quenching if and only if 0 (cf. Phillips(1987),Galaktionov & Vazquez(1995)).

Let us also mention here that a similar phenomenon when the continuation beyond gradient blow-up does not satisfy the original boundary condition was observed by Fila & Lieberman(1994).

1.2. Problem (B) $(f(u) = u^p)$. The study of blow-up of solutions of the heat equation with a nonlinear boundary condition was initiated by Levine & Payne(1974) and it has attracted considerable attention (see a survey paper of Fila & Filo(1996)). It was shown by Fila(1989) that every solution of Problem (B) blows up in a finite time $T = T(u_0)$ and it is also known (cf. López Gómez, Márquez, & Wolanski(1991)) that the only blow-up point is x = 1.

(By a blow-up point we mean a point $a \in [0,1]$ such that there are $\{x_n\} \subset [0,1]$ and $t_n \to T$ such that $x_n \to a$ and $u(x_n, t_n) \to \infty$ as $n \to \infty$.)

We show that for Problem (B) blow-up is always complete in the following sense. If

 $f^n_{}(s) = \min\{s^p, n^p\}, \quad s \ge 0, \quad n \in \mathbb{N},$ (1.1)

and u^n is the solution of (B^n) :

 $\begin{array}{l} & & & & \\ & & & & \\ & & & & \\$

then $u^n(x,t) \to \infty$ for $(x,t) \in [0,1] \times (T,\infty)$.

For results on complete blow-up for the problem when the nonlinearity occurs in the equation we refer to the papers of Baras & Cohen(1987), Lacev & Tzanetis(1988), Galaktionov & Vazquez(1995, 1997), Martel(1998), etc.

Our method is different and it is restricted to one space dimension since we are using an intersection-comparison (or zero number(cf. [14])) argument.

2. Incomplete Quenching

In this section we show that if u(x,t) is the solution of the problem

$$\begin{cases} u_t = u_{xx}, & 0 < x < 1, \quad 0 < t < T, \\ u_x(0,t) = 0, & 0 < t < T, \\ u_x(1,t) = -u^{-p}(1,t), & 0 < t < T, \\ u(x,0) = u_0(x) > 0, & 0 \le x \le 1, \end{cases}$$
(Q)

where p > 0 and T is the quenching time for u then there is a natural continuation of u beyond T. We shall assume that $u_0 \in C^1([0, 1])$ and the compatibility conditions

$$u_0'(0) = 0, \qquad u_0'(1) = -u_0^{-p}(1)$$

are satisfied.

Assume that $0 < \varepsilon < u_0(1)$. Then there exists a unique global (in time) solution u^{ε} of (Q_{ε}) such that $u^{\varepsilon} \in C^{2,1}([0,1] \times [0,\tau])$ for any $\tau > 0$ and

- (i) $u^{\varepsilon} > 0$ for $(x, t) \in [0, 1] \times [0, \infty)$,
- (ii) $u^{\varepsilon_1} \leq u^{\varepsilon_2}$ for $0 < \varepsilon_1 < \varepsilon_2$ and $(x,t) \in [0,1] \times [0,\infty)$,
- (iii) $u^{\varepsilon} \ge u$ for $(x,t) \in [0,1] \times [0,T)$.

Also, by the maximum principle, it is clear that

$$u^{\varepsilon} \leq K \equiv \max_{0 \leq x \leq 1} u_0(x)$$

for all $\varepsilon > 0$.

Now, let

$$v(x,t) = \lim_{\varepsilon \to 0} u^{\varepsilon}(x,t), \qquad (x,t) \in [0,1] \times [0,\infty).$$

$$(2.1)$$

Then v is well-defined and $0 \le v \le K$ in $[0, 1] \times [0, \infty)$. It follows from the regularity theory for parabolic equations that v satisfies the heat equation in $(0, 1) \times (0, \infty)$. By the maximum principle, v > 0 in $(0, 1) \times (0, \infty)$. Also, it is clear that $v_x(0, t) = 0$ for t > 0. Furthermore, if $t \in (0, T)$, then

$$v_x(1,t) = -v^{-p}(1,t).$$

It follows that v is a solution of (Q). By uniqueness, v = u in $[0, 1] \times [0, T)$. For the boundary condition for v on $\{x = 1, t > T\}$, it can be shown that v(1, t) = 0 for $t \ge T$.

We summarize the above results as follows:

Theorem 2.1[15]. The function v defined by (2.1) satisfies

 $egin{aligned} & v_t = v_{xx}, & 0 < x < 1, \quad t > 0, \ & v_x(0,t) = 0, & t > 0, \ & v_x(1,t) = -v^{-p}(1,t), & 0 < t < T, \ & v(1,t) = 0, & t \ge T, \ & v(x,0) = u_0(x), & 0 \le x \le 1. \end{aligned}$

It coincides with the solution u of Problem (Q) for $t \leq T$.

3. Complete Blow-up

Consider the problem

$$\begin{cases} u_t = u_{xx}, & 0 < x < 1, \quad 0 < t < T, \\ u_x(0,t) = 0, & 0 < t < T, \\ u_x(1,t) = u^p(1,t), & 0 < t < T, \\ u(x,0) = u_0(x) > 0, & 0 \le x \le 1, \end{cases}$$
(B)

where p > 1, and T is the blow-up time for u. We assume further that $u'_0(0) = 0$ and $u'_0(1) = u^p_0(1)$.

Let $K = \max_{0 \le x \le 1} u_0(x)$. For any n > K, $n \in \mathbb{N}$, we define f^n as in (1.1). Note that f^n is Lipschitz and $u'_0(1) = f^n(u_0(1))$ if n > K. Hence, the solution of (\mathbb{B}^n) is C^1 up to the boundary. We show that there exists a unique global (in time) solution u^n of (\mathbb{B}^n) such that

- (i) $u^n > 0$ for $(x, t) \in [0, 1] \times [0, \infty)$,
- (ii) $u^n \le u^{n+1}$ for $(x,t) \in [0,1] \times [0,\infty)$,
- (iii) $u^n \le u$ for $(x, t) \in [0, 1] \times [0, T)$.

Define

$$v(x,t) = \lim_{n \to \infty} u^n(x,t), \qquad 0 \le x \le 1, \quad t \ge 0.$$
(3.1)

Similarly, one can show that $v_x(1,t) = v^p(1,t)$ for $t \in (0,T)$. Then it is clear that v(x,t) = u(x,t) for 0 < t < T. Note that $v(1,T) = \infty$. Furthermore, there holds $v(1,t) = \infty$ for $t \ge T$.

This proves the following:

Theorem 3.1[15]. The function v defined in (3.1) coincides with the solution u of Problem (B) for $t \leq T$ and $v(x,t) = \infty$ for $(x,t) \in [0,1] \times (T,\infty)$.

Acknowledgment. This is a joint work with Marek Fila.

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