

# A representation of unital completely positive maps

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## Abstract

Let  $M_n(\mathbb{C})$  be the algebra of  $n \times n$  complex matrices, and let  $\Phi$  be a unital completely positive map of  $M_n(\mathbb{C})$  to  $M_k(\mathbb{C})$ . With the notion of the von Neumann entropy for a state in mind, we give a model of  $r$ -tuple  $\{v_j\}_{j=1}^r$  so that  $\Phi(x) = v_1^* x v_1 + \cdots + v_r^* x v_r$ , ( $x \in M_n(\mathbb{C})$ ). The  $r$  is uniquely determined for  $\Phi$  and the  $r$ -tuple is also unique up to a  $r \times r$  unitary matrix.

## 1 Introduction

In the framework of the theory of operator algebras, the notion of entropy for automorphisms was introduced by Connes-Størmer in [8], Connes-Narnhofer-Thirring in [9] and Voiculescu in [14] (which is extended by Brown [4]). The Connes-Størmer entropy  $H(\theta)$  is defined for a  $*$ -automorphism  $\theta$  of finite von Neumann algebra  $M$  with  $\tau = \tau \circ \theta$ , where  $\tau$  is a fixed given finite trace of  $M$ . After then, the Connes-Narnhofer-Thirring entropy  $h_\phi(\theta)$  is given as an extended version of  $H(\theta)$  for a  $*$ -automorphism  $\theta$  of a  $C^*$ -algebra  $A$  by replacing the trace  $\tau$  to a state  $\phi$  of  $A$ , and if  $A$  is a finite von Neumann algebra then  $h_\tau(\theta) = H(\theta)$ . Voiculescu's topological entropy  $ht(\theta)$  is defined as an independent version of any state of  $A$ .

We studied these entropies in [6] and [7] for not only  $*$ -automorphisms but also  $*$ -endomorphisms like so called canonical shifts. As one of interesting such  $*$ -endomorphisms, we picked up the Cuntz canonical endomorphism  $\Phi_n$  on the Cuntz algebra  $O_n$  which has a strong connection to Longo's canonical shift (cf. [6]). The  $O_n$  is the  $C^*$ -algebra generated by isometries  $\{S_1, \dots, S_n\}$  such that  $S_1 S_1^* + \cdots + S_n S_n^* = 1$ , and the  $\Phi_n$  is defined as

$$\Phi_n(x) = S_1 x S_1^* + \cdots + S_n x S_n^*, \quad (x \in O_n). \quad (1.1)$$

Such maps given by the form as the right hand side of (1.1) are unital completely positive maps, and the above notions  $H(\cdot)$ ,  $h_\phi(\cdot)$  and  $ht(\cdot)$  are available for unital completely positive maps too.

Conditional expectations are the most typical examples of unital completely positive maps, and states of the matrix algebras  $M_n(\mathbb{C})$  are considered as the most elementary example of conditional expectations. However, for a conditional expectation  $E$ , by their definitions it holds always that  $H(E) = h_\phi(E) = ht(E) = 0$ . On the other hand, in the case of the von Neumann entropy  $S(\phi)$  for a state  $\phi$  of  $M_n(\mathbb{C})$ , it is possible that  $S(\phi) \neq 0$ .

In order to define "entropy", we need the notion of "finite partition of unity" (see for example, [11]). The most generalized one of "finite partition of unity" was introduced by Lindblad ([10]), and it is called the "finite operational partition of unity".

With these facts in mind, here we give a method to induce the finite operational partition of unity for a given unital completely positive map. That is, let  $A$  and  $B$  be unital  $C^*$ -algebras and let  $\Phi$  be a unital completely positive map of  $A$  to  $B$ . We give a method to get a model of  $r$ -tuple  $v(\Phi) = \{v_1, v_2, \dots, v_r\}$  such that

$$\Phi(x) = v_1^* x v_1 + \dots + v_r^* x v_r, \quad (x \in A). \quad (1.2)$$

When  $A$  and  $B$  are matrix algebras, such a representation is called Kraus representation (cf. Appendix in [13]), or obtained as a straightforward application of Stinespring's theorem (see for example, [1, 3]). We note that this representation is not unique

Our main purpose in this note is to show, for a given completely positive map  $\Phi$ , a unique  $r$ -tuple  $v(\Phi) = \{v_1, \dots, v_r\}$  which is suitable to extend the notion of von Neumann entropy  $S(\phi)$  for a state  $\phi$  of matrix algebras to the entropy  $S(\Phi)$  for a unital completely positive map  $\Phi$ .

First, for a given completely positive map  $\Phi$  from  $B(H)$  to  $B(K)$  of finite dimensional Hilbert spaces  $H, K$ , we construct the Hilbert spaces  $H \otimes_\Phi K$ . Let  $r = \dim(H \otimes_\Phi K)$ . Next, we give a  $r$ -tuple  $v(\Phi) = \{v_1, v_2, \dots, v_r\}$  which satisfy that  $\Phi(x) = v_1^* x v_1 + \dots + v_r^* x v_r$ . The  $r$ -tuple is unique up to unitaries and induces  $S(\Phi)$ . After then, we apply these to the non-commutative Bernoulli shift  $\beta$  and we define the entropy  $S_\varphi(\Phi)$  with respect to a state  $\varphi$  with  $\varphi = \varphi \cdot \beta$ . These results in this note are in [5].

## 2 Preliminaries

Here, we denote some notations and terminologies which we use later.

We denote by  $M_n(\mathbb{C})$  the algebra of  $n \times n$  complex matrices, and by  $\text{Tr}_n$  the standard trace, that is, the sum of all diagonal components. A matrix  $D \in M_n(\mathbb{C})$  is called a *density matrix* if  $D$  is a positive operator with  $\text{Tr}_n(D) = 1$  (cf. [11] [12]).

The notation  $\eta$  is called the entropy function in usual, and it is the function defined by

$$\eta(t) = \begin{cases} -t \log t, & (0 < t \leq 1) \\ 0, & t = 0 \end{cases}$$

### 2.1 Finite partitions

The notion of "a finite partition of unity" is the starting point of our study.

#### 2.1.1 Finite partitions of 1

The first one is discussed in the real numbers  $\mathbb{R}$ . Let

$$\lambda = \{\lambda_1, \dots, \lambda_n\}$$

be the set of real numbers  $\lambda_i \geq 0$  with  $\sum_i \lambda_i = 1$ . We say that the  $n$ -tuple  $\lambda \subset \mathbb{R}$  is a *finite partition of 1*.

#### 2.1.2 Finite operational partition of unity

The terminology, a *finite operational partition of unity*, was first given by Lindblad ([10]) and after then it is used by Alicki-Fannes([2]).

Let  $A$  be a unital  $C^*$ -algebra. Let  $x = \{x_1, \dots, x_k\} \subset A$ . Then  $x$  is said to be a *finite operational partition of unity* of size  $k$  if

$$\sum_i^k x_i^* x_i = 1_A. \quad (2.1)$$

Such a finite operational partition of unity  $x = \{x_1, \dots, x_k\}$  in  $A$  induces an  $A$ -coefficient in  $M_k(A)$ , whose  $(i, j)$  coefficient  $x(j, i)$  is given by the following:

$$x(j, i) = x_i^* x_j, \quad (1 \leq i, j \leq k). \quad (2.2)$$

We denote this matrix by  $[x]$ . Then  $[x]$  is an  $A$ -coefficient density matrix in  $M_k(A)$ , that is,  $[x]$  is a positive operator with  $\text{Tr}([x]) = \sum_{i=1}^k x(i, i) = 1_A$ .

## 2.2 Entropy for finite partitions of unity

### 2.2.1 Entropy for finite partitions of 1

Let a  $n$ -tuple  $\lambda \subset \mathbb{R}$  be a a *finite partition of 1*. Let

$$H(\lambda) = \eta(\lambda_1) + \cdots + \eta(\lambda_n). \quad (2.3)$$

Then  $H(\lambda)$  is called the *entropy* for the finite partiton  $\lambda$  of 1.

### 2.2.2 Entropy for Finite operational partition of unity

Let  $x = \{x_1, \dots, x_k\}$  be an operational partition of unity in a unital  $C^*$ -algebra  $A$ , and let  $\varphi$  be a state of  $A$ . The  $\rho_\varphi[x]$  is the  $k \times k$  matrix whose  $\{i, j\}$ -component is defined by

$$\rho_\varphi[x](i, j) = \varphi(x_j^* x_i), \quad (i, j = 1, \dots, k). \quad (2.4)$$

Then  $\rho_\varphi[x]$  is a density matrix. We call  $\rho_\varphi[x]$  the *density matrix associate with  $x$  and  $\varphi$* . If  $\varphi$  is a unique tracial state, we denote  $\rho_\varphi[x]$  by  $\rho[x]$  simply.

Let  $\lambda(\rho_\varphi[x]) = \{\lambda_1, \lambda_2, \dots, \lambda_k\}$  be the eigenvalues of the matrix  $\rho_\varphi[x]$ . Then  $\lambda(\rho_\varphi[x])$  is a finite partition of 1 because  $\rho_\varphi[x]$  is a density matrix. Hence we have the entropy  $H(\lambda(\rho_\varphi[x]))$ .

Let  $S(\rho_\varphi[x])$  be the von Neumann entropy (cf. [11, 12]) for the density matrix  $\rho_\varphi[x]$ . Then  $S(\rho_\varphi[x])$  is nothing else but  $H(\lambda(\rho_\varphi[x]))$ , that is,

$$S(\rho_\varphi[x]) = \text{Tr}_k(\eta(\rho[x])) = H(\lambda(\rho_\varphi[x])) = \sum_i \eta(\lambda_i). \quad (2.5)$$

## 3 Representation of completely positive maps

Let  $\Phi$  be a completely positive map of  $M_n(\mathbb{C})$  to  $M_k(\mathbb{C})$ . Put  $A = M_n(\mathbb{C})$ . We give a method to get a "finite" family  $v(\Phi) = \{v_1, v_2, \dots, v_r\}$  for  $\Phi$  which satisfies that  $\Phi(x) = v_1^* x v_1 + \cdots + v_r^* x v_r$  for all  $x \in A$ . We remark that if  $\Phi$  is unital, then  $v(\Phi)$  is a finite operational partition:

$$\sum_{j=1}^r v_j^* v_j = \Phi(1_A) = 1_{B(K)} \quad (3.1)$$

### 3.1 Hilbert space $H \otimes_{\Phi} K$

Let  $H$  be an  $n$ -dimensional Hilbert space, and let  $\Phi : A \rightarrow B(K)$  be a completely positive linear map. Let  $\{e_1, \dots, e_m\}$  be the set of mutually orthogonal minimal projections in  $B(H)$  with  $\Phi(e_i) \neq 0$  for all  $i$ . Let  $\xi_i \in e_i(H)$  be a vector with  $\|\xi_i\| = 1$  for  $i = 1, \dots, m$  and we extend  $\{\xi_1, \dots, \xi_m\}$  to an orthonormal basis of  $H$  as  $\{\xi_1, \dots, \xi_n\}$ . Let  $\{e_{ij}; i, j = 1, \dots, n\}$  be a matrix units of  $A$  with  $e_{ij}\xi_j = \xi_i$  so that  $e_{ii} = e_i$  for  $i = 1, \dots, m$ . Then each  $\zeta \in H \odot K$  (the algebraic tensor product  $H \odot K$  of  $H$  and  $K$ ) is written by

$$\zeta = \sum_{i=1}^n \xi_i \otimes \mu_i, \quad \text{for some } \mu_i \in K. \quad (3.2)$$

**Definition 3.1.1.** We define a sesquilinear form  $\langle \cdot, \cdot \rangle_{\Phi}$  on the space  $H \odot K$  by

$$\left\langle \sum_{i=1}^n \xi_i \otimes \mu_i, \sum_{j=1}^n \xi_j \otimes \nu_j \right\rangle_{\Phi} = \sum_{i,j} \langle \Phi(e_{ji})\mu_i, \nu_j \rangle_K, \quad (3.3)$$

where  $\langle \cdot, \cdot \rangle_K$  means the inner product of the Hilbert space  $K$ .

Since  $\Phi$  is completely positive, this form  $\langle \cdot, \cdot \rangle_{\Phi}$  turns out positive semidefinite. The value  $\langle \cdot, \cdot \rangle_{\Phi}$  depends on the choice of the orthonormal basis of  $H$ . However the kernel of this sesquilinear form  $\langle \cdot, \cdot \rangle_{\Phi}$  is unique up to unitaries on  $H \otimes K$  as follows:

**Proposition 3.1.2.** Let  $\{\xi_1, \dots, \xi_n\}$  (resp.,  $\{\xi'_1, \dots, \xi'_n\}$ ) be an orthonormal basis of  $H$ , and let  $u$  be a unitary on  $H \otimes K$  with  $\xi'_i = u\xi_i$  for all  $i = 1, \dots, n$ . Let  $Ker(\Phi)$  (resp.,  $Ker'(\Phi)$ ) be the kernel of this form via  $\{\xi_i\}_i$  (resp.,  $\{\xi'_i\}_i$ ) Then

$$Ker'(\Phi) = (\bar{u}u^* \otimes 1)Ker(\Phi)$$

where  $\bar{u}$  is the unitary matrix on  $H$  whose  $(i, j)$ -entry is the conjugate complex number of  $u(i, j)$  for all  $i, j = 1, \dots, n$ .

**Definition 3.1.3.** Now taking the quotient by the space  $Ker(\Phi)$ , we have a preHilbert space and complete to get the Hilbert space  $H \otimes_{\Phi} K$ .

We denote by  $(\sum_{i=1}^n \xi_i \otimes \mu_i)_{\Phi}$  the element in  $H \otimes_{\Phi} K$  corresponding to  $\sum_{i=1}^n \xi_i \otimes \mu_i \in H \odot K$ .

If  $K$  is finite dimensional, of course the  $H \otimes_{\Phi} K$  is finite dimensional. We can extend this method to get the Hilbert space  $H \otimes_{\Phi} K$  to (for an example) non-commutative Bernoulli shifts, and it induces finite dimensional  $H \otimes_{\Phi} K$  even if  $H$  and  $K$  are infinite dimensional. By Proposition 3.1.2, we have the following:

**Proposition 3.1.4.** *The dimension of  $H \otimes_{\Phi} K$  does not depend on the choice of orthonormal basis of  $H$ .*

**Example 3.1.5.**

1. If  $\phi$  is a state of  $M_n(\mathbb{C})$ , then

$$\dim(\mathbb{C}^n \otimes_{\phi} \mathbb{C}) = \text{rank of } \phi, \text{ (i.e., the rank of the density matrix of } \phi).$$

2. If  $E$  is the conditional expectation of  $M_n(\mathbb{C})$  to a maximal abelian subalgebra  $A$  of  $M_n(\mathbb{C})$ , then

$$\dim(\mathbb{C}^n \otimes_E \mathbb{C}^n) = n.$$

Here, we remark that  $A$  is isomorphic to the diagonal subalgebra  $D_n(\mathbb{C})$ .

3. Let  $B$  be a subfactor of  $M_n(\mathbb{C})$ . If  $E$  is the conditional expectation  $M_n(\mathbb{C})$  to  $B$ , then

$$\dim(\mathbb{C}^n \otimes_E \mathbb{C}^m) = \frac{n}{m}$$

Here, we remark that  $B$  is isomorphic to  $M_m(\mathbb{C})$  for some  $m$  by which  $n$  can be divided.

4. If  $\alpha$  is an automorphism of  $M_n(\mathbb{C})$ , then

$$\dim(\mathbb{C}^n \otimes_{\alpha} \mathbb{C}^n) = 1.$$

The following shows that  $\dim(H \otimes_{\Phi} K)$  can be finite, even if  $H$  and  $K$  are infinite dimensional.

**Example 3.1.6.** Let  $\beta$  be the non-commutative Bernoulli shift of  $A = \otimes_{i=1}^{\infty} M_n(\mathbb{C})$ . That is, let  $A_i = M_n(\mathbb{C})$  for all  $i = 1, 2, \dots$ , and for each  $m \in \mathbb{N}$ , let

$$A(m) = A_1 \otimes \cdots \otimes A_m \otimes 1 \otimes \cdots \subset A, \quad (3.4)$$

where  $1$  is the unit of  $M_n(\mathbb{C})$ .

The  $\beta$  is given as the shift as the followings:

$$\beta(x) = 1 \otimes x \otimes 1 \otimes \cdots, \quad \text{for all } m, x \in A(m). \quad (3.5)$$

Let  $H_i = \mathbb{C}^n$  for all  $i = 1, 2, \dots$ , and for each  $m \in \mathbb{N}$ . Fix an vector  $\Omega \in \mathbb{C}$  with  $\|\Omega\| = 1$ , and let

$$H(m) = H_1 \otimes \cdots \otimes H_m \otimes \Omega \otimes \cdots \subset H = \otimes_{i=1}^{\infty} H_i, \quad (3.6)$$

The  $\beta$  is a unital completely positive map from  $A \subset B(H)$  to  $B(H)$ , and the restriction  $\beta|_{A(m)}$  of  $\beta$  to  $A(m)$  is a unital completely positive map from  $A(m) \subset B(H(m))$  to  $B(H(m+1))$ . Apply the above method to  $\beta|_{A(m)}$ , we have always that

$$\dim(H(m) \underset{\beta|_{A(m)}}{\otimes} H(m+1)) = n, \quad \text{for all } m$$

As a result, we have that

$$\dim(H \underset{\beta}{\otimes} H) = n = \dim(H \underset{E}{\otimes} H),$$

where  $E : A \rightarrow \beta(A)$  is the conditional expectation.

Now, we call the dimension of  $H \otimes_{\Phi} K$  the *rank* of  $\Phi$ .

**A phenomenon** As an example, we show a phenomenon of the above discussion in the case of a state of  $M_2(\mathbb{C})$  which indicates how the dimension of  $H \otimes_{\phi} K$  coincides with the rank of a state  $\phi$  of the usual sense.

**Example 3.1.7.** Let  $\{\xi_1, \xi_2\}$  be an orthonormal basis of  $\mathbb{C}^2$  and let  $\{e_{ij}; i, j = 1, 2\}$  be a matrix units of  $\mathbb{C}^2$  with  $e_{ij}\xi_j = \xi_i$ . We give a vector representation for each  $\xi \in \mathbb{C}^2$  relative to this  $\{\xi_i; i = 1, 2\}$  and a matrix representation for each  $x \in M_2(\mathbb{C})$  relative to this  $\{e_{ij}; i, j = 1, 2\}$ :

$$\xi_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \xi_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad (3.7)$$

and

$$e_{11} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}, \quad e_{12} = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}, \quad e_{21} = \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}, \quad e_{22} = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}. \quad (3.8)$$

Assume that  $\phi$  is a state given by

$$\phi(x) = \phi\left(\begin{bmatrix} x_{11} & x_{12} \\ x_{21} & x_{22} \end{bmatrix}\right) = \frac{1}{2} \sum_{i,j=1}^2 x_{ij}. \quad (3.9)$$

Then  $\phi(e_i) = 1/2$  for  $i = 1, 2$  and the discussion with respect to  $\{\xi_1, \xi_2\}$  and  $\{e_{ij}; i, j = 1, 2\}$  is as follows:

Let  $\Omega$  be a fixed unit vector in  $\mathbb{C}$  and let

$$\zeta_i = \sqrt{2}\xi_i \otimes \Omega \in \mathbb{C}^2 \otimes_{\phi} \mathbb{C}. \quad (3.10)$$

Then

$$\langle \zeta_i, \zeta_j \rangle_{\phi} = 2\phi(e_{ij}) = 1, \quad \text{for all } i, j = 1, 2. \quad (3.11)$$

This implies that  $\zeta_i \in \mathbb{C}^2 \otimes_{\phi} \mathbb{C}$  has norm 1 for  $i = 1, 2$  and

$$\zeta_1 = \sqrt{2}\xi_1 \otimes \Omega = \sqrt{2}\xi_2 \otimes \Omega = \zeta_2. \quad (3.12)$$

Next we choose another family of minimal projections with  $\phi(p) \neq 0$  and orthonormal basis of  $\mathbb{C}^2$ . Let

$$e'_{11} = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}, \quad (3.13)$$

then  $\phi(e'_{11}) = 1$  so that the set containing  $e'_{11}$  of minimal projections with  $\phi(\cdot) \neq 0$  is the one point set  $\{e'_{11}\}$ . The corresponding orthonormal basis of  $\mathbb{C}^2$  and the corresponding matrix units are as follows:

$$\xi'_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}, \quad \xi'_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} \quad (3.14)$$

and

$$e'_{11}, e'_{12} = \frac{1}{2} \begin{bmatrix} 1 & -1 \\ 1 & -1 \end{bmatrix}, \quad e'_{21} = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ -1 & -1 \end{bmatrix}, \quad e'_{22} = \frac{1}{2} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \quad (3.15)$$

Let

$$\zeta'_i = \frac{1}{\sqrt{2}}\xi'_i \otimes \Omega \in \mathbb{C}^2 \otimes_{\phi} \mathbb{C}, \quad (i = 1, 2). \quad (3.16)$$

Then

$$\langle \zeta'_1, \zeta'_1 \rangle_{\phi} = \frac{1}{2}\phi(e'_{11}) = 1, \quad (3.17)$$



and

$$\langle \zeta'_2, \zeta'_2 \rangle_\phi = \frac{1}{2} \phi(e'_{22}) = 0 \quad (3.18)$$

This means that

$$\mathbb{C}^2 \otimes_\phi \mathbb{C} = \mathbb{C} \zeta'_i = \xi'_i \otimes \mathbb{C}.$$

We remark that  $\{\xi_1, \xi_2\}$  and  $\{\xi'_1, \xi'_2\}$  are combined as  $u\xi_i = \xi'_i$ , ( $i = 1, 2$ ) by the unitary  $u$ :

$$u = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}. \quad (3.19)$$

**Some relation to the Choi matrix.** For a completely positive map  $\Phi$  of  $M_n(\mathbb{C})$ , it is given the so called the Choi matrix  $C_\Phi$ .

In the case of  $\Phi$  is a state  $\phi$  of  $M_n(\mathbb{C})$ , we have the following relation between the sesquilinear form  $\langle \cdot, \cdot \rangle_\phi$  and the coefficient of  $C_\phi$  :

$$\langle \xi_i \otimes \Omega, \xi_j \otimes \Omega \rangle_\phi = C_\phi(j, i) \quad (3.20)$$

where  $\{\xi_1, \dots, \xi_n\}$  is a orthonormal basis of  $\mathbb{C}^n$  and  $\Omega$  is 1 considered as the vector in  $\mathbb{C}$ .

### 3.2 Operators $\{v_j; j = 1, \dots, r\}$

As the above section, let  $\Phi : B(H) \rightarrow B(K)$  be a completely positive linear map. Here, we assume that  $H$  and  $K$  be finite dimensional, and let  $r = \dim(H \otimes_\Phi K)$ .

**Definition 3.2.1.** Let  $\{\xi_i; i = 1, \dots, n\}$  be an orthonormal basis of  $H$ , and let  $\{\zeta_j; j = 1, \dots, r\}$  be an orthonormal basis of  $H \otimes_\Phi K$ . Define  $v_j : K \rightarrow H$  by

$$v_j(\mu) = \sum_{i=1}^n \langle \xi_i \otimes \mu, \zeta_j \rangle_\Phi \xi_i, \quad (\mu \in K) \quad (3.21)$$

**Proposition 3.2.2.** Let  $\{v_1, v_2, \dots, v_r\}$  be the tuple obtained by (3.17). Then

(i) They satisfies the desired following property:

$$\Phi(x) = v_1^* x v_1 + \dots + v_r^* x v_r, \quad (x \in B(H)) \quad (3.22)$$

(ii)  $\|v_i\| \leq 1$ , for all  $i = 1, \dots, r$ ;

(iii)  $\{v_1, v_2, \dots, v_r\}$  are linearly independent.

**Proposition 3.2.3.** *The tuple  $\{v_1, v_2, \dots, v_r\}$  for unital completely positive map  $\Phi$  satisfy the following convenient properties to compute the von Neumann type entropy  $S(\Phi)$ .*

1. *If  $E$  is a conditional expectation of  $M_n(\mathbb{C})$  onto  $D_n(\mathbb{C})$  then  $\{v_j : 1 \leq j \leq r\}$  are mutually orthogonal minimal projections.*

2. *If  $\Phi$  is a  $*$ -homomorphism, then  $\{v_j : 1 \leq j \leq r\}$  are isometries with*

$$v_i v_j^* = \delta_{ij} 1. \quad (3.23)$$

We remark that the following results are well known (see for example [1, 13]) so that our tuple  $\{v_1, v_2, \dots, v_r\}$  for unital completely positive map  $\Phi$  is unique up to a unitary matrix:

**Proposition 3.2.4.** *Let  $H$  and  $K$  be finite dimensional Hilbert spaces. Assume that  $v = \{v_1, v_2, \dots, v_r\}$  and  $w = \{w_1, \dots, w_r\}$  are two families of operators in  $B(K, H)$ . Then*

$$v_1^* x v_1 + \dots + v_r^* x v_r = w_1^* x w_1 + \dots + w_r^* x w_r, \quad \text{for all } x \in B(H)$$

*if and only if there is a unitary matrix  $[u(i, j)] \in M_r(\mathbb{C})$  such that*

$$v_i = \sum_{j=1}^r u(i, j) w_j, \quad i = 1, \dots, r.$$

**Example 3.2.5.** Let  $\beta$  be the non-commutative Bernoulli shift on  $A = \otimes_{i=1}^{\infty} M_n(\mathbb{C})$ . The  $n$ -tuple  $\{v_j\}_j$  of  $\beta$  are as follows. We use the same notation as Example 3.1.6. Let

$$W_m = \{\alpha = (\alpha_1, \alpha_2, \dots, \alpha_m), \quad \alpha_j \in \{1, 2, \dots, n\}\}. \quad (3.24)$$

and let

$$\xi_\alpha = \xi_{\alpha_1} \otimes \xi_{\alpha_2} \otimes \dots \otimes \xi_{\alpha_m} \otimes \Omega \otimes \dots \in H(m) \quad (3.25)$$

Then  $\{\xi_\alpha; \alpha \in W_m\}$  is an orthonormal basis of  $H(m)$ , and

$$\{\xi_\alpha \otimes_\beta (\xi_i \otimes \xi_\alpha); i = 1, 2, \dots, n\} \quad (3.26)$$

is an orthonormal basis of  $H(m) \otimes_\beta H(m+1)$ . Then our tuple for  $\beta$  is given by

$$v_j(\xi_i \otimes \xi_\beta) = \delta_{ij} \xi_\beta, \quad \text{for all } m, \beta \in W_m. \quad (3.27)$$

We remark that  $\{v_j^*\}_j$  are isometries satisfying the Cuntz relation.

### 3.3 Relation to Stinespring's theorem

Let  $\Phi : A \subset B(H) \rightarrow B(K)$  be a completely positive map, where  $H$  and  $K$  are finite dimensional. Let  $\{v_j : K \rightarrow H\}_{j=1}^r$  be the tuple by (3.17). We denote by  $L$  the Hilbert space of  $r$ -direct sum of  $H$ , i.e.,  $H \oplus \cdots \oplus H$ . Let

$$V(\xi) = (v_1\xi, \cdots, v_r\xi) \in L, \quad \xi \in K \quad (3.28)$$

and let

$$\pi(x) = \begin{bmatrix} x & 0 & \cdots & 0 \\ \vdots & x & \vdots & 0 \\ \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & \cdots & x \end{bmatrix}, \quad x \in A. \quad (3.29)$$

Then we have the followings:

1. The  $\pi$  is a representation of the  $C^*$ -algebra  $A$  on the Hilbert space  $L$ .
2. The property that  $\sum_{i=1}^r v_i^* v_i = 1_K$  means that the operator  $V$  defined by (3.20) is an isometry from  $K$  to  $L$ .
3. The property  $\Phi(x) = \sum_{j=1}^r v_j^* x v_j$  is written as

$$\Phi(x) = (v_1^*, \cdots, v_r^*) \begin{bmatrix} x & 0 & \cdots & 0 \\ \vdots & x & \vdots & 0 \\ \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & \cdots & x \end{bmatrix} \begin{pmatrix} v_1 \\ \vdots \\ \vdots \\ v_r \end{pmatrix} = V^* \pi(x) V \quad (3.30)$$

4. These imply that  $(\pi, V)$  can be considered as the pair obtained by the Stinespring's representation.
5. The property that  $\{v_j\}_j$  are linearly independent satisfies that  $(\pi, V)$  is the minimal pair in the sense of Arveson [1].

## 4 Entropy for unital completely positive maps

In this section, we denote an application to the notion of entropy. Let  $\Phi$  be a unital completely positive map of a  $C^*$ -algebra  $A \subset B(H)$  to  $B \subset B(K)$  and let  $v(\Phi) = \{v_1, v_2, \cdots, v_r\}$  be the tuple for  $\Phi$  obtained by (3.21). Then the tuple  $v(\Phi)$  is a finite operational partition of unity in  $B(K)$ . Hence we can apply the discussion in the section 2.2.2 to unital completely positive map.

## 4.1 Case of a state $\phi$ of $M_n(\mathbb{C})$

First, we consider the case of a state  $\phi$  of  $M_n(\mathbb{C})$ . Let  $\phi$  be a state of  $M_n(\mathbb{C})$ , and let  $v(\phi) = \{v_j : 1 \leq j \leq r\}$  be the tuple associated with  $\phi$  defined by (3.20).

Let  $\tau$  be the unique tracial state of  $M_n(\mathbb{C})$ , that is  $\tau(x) = \text{Tr}(x)/n$  for all  $x \in M_n(\mathbb{C})$ . The density matrix  $\tau[v(\phi)]$  is given by (2.4). Let  $\{e_i\}_{i=1}^m$  be the mutually orthogonal minimal projections in  $M_n(\mathbb{C})$  such that  $\phi(e_i) \neq 0$ . Then we see that

$$\tau[v(\phi)](i, j) = \delta_{ij} \sqrt{\phi(e_i)\phi(e_j)}. \quad (4.1)$$

It is clear that  $\tau[v(\phi)]$  is a diagonal matrix, and the entropy  $S(\tau[v(\phi)])$  in the section 2.2.2 is nothing else but the von Neumann entropy  $S(\phi)$  of  $\phi$ :

$$S(\tau[v(\phi)]) = \sum_{j=1}^r \eta(\phi(e_j)) = - \sum_{j=1}^r \phi(e_j) \log \phi(e_j) = S(\phi). \quad (4.2)$$

## 4.2 Entropy for unital completely positive maps

On the basis of the fact in the above section 4.1, we denote the  $S(\rho[v(\Phi)])$  for a unital completely positive map  $\Phi : A \rightarrow B$  by  $S(\rho(\Phi))$ , and in the case of the tracial state  $\rho$  we use the same notation  $S(\Phi)$  simply.

Here, we show the case of  $A$  which has a unique tracial state  $\tau$  and we number the values of typical examples of von Neumann type entropy  $S(\Phi)$  for unital completely positive maps  $\Phi$ .

1. If  $\phi$  is a state of  $M_n(\mathbb{C})$ , then

$$S(\phi) = \sum_{j=1}^r \eta(\lambda_j) \quad (4.3)$$

where  $\{\lambda_j\}$  are eigenvalues of  $\phi$ .

2. If  $E$  is the conditional expectation of  $M_n(\mathbb{C})$  to a maximal abelian subalgebra  $B$  then

$$S(E) = \log n. \quad (4.4)$$

Compare this fact to that  $H(E) = ht(E) = 0$ .

3. If  $E$  is the conditional expectation of  $M_n(\mathbb{C})$  to a subfactor  $B$  then

$$S(E) = \log \frac{n}{k}. \quad (4.5)$$

Here we remark that a subfactor  $B$  of  $M_n(\mathbb{C})$  is isomorphic to  $M_k(\mathbb{C})$  some  $k$ , and that  $n$  is divisible by  $k$ . Compare this fact to that  $H(E) = ht(E) = 0$ .

4. If  $\alpha$  is an automorphism of  $M_n(\mathbb{C})$ , then

$$S(\alpha) = 0 \quad (4.6)$$

and this coincides with the fact that  $H(\alpha) = ht(\alpha) = 0$ .

5. If  $\beta$  is the non-commutative Bernoulli shift on  $\otimes_{i=1}^{\infty} M_n(\mathbb{C})$ , then

$$S(\beta) = \log n \quad (4.7)$$

and this coincides with the fact that  $H(\beta) = ht(\beta) = \log n$ .

6. If  $\Phi_n$  is the Cuntz's canonical shift on  $O_n$ , then

$$S_{\Psi}(\Phi_n) = \log n \quad (4.8)$$

and this coincides with the fact that  $h_{\Psi}(\Phi_n) = ht(\Phi_n) = \log n$ . Here  $\Psi$  is the state of  $O_n$  which is given by the left inverse of  $\Phi$ .

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