

Eigenvalue problems associated with 1-dimensional scalar field equations

Tohru Wakasa
 Kyushu Institute of Technology

1 Introduction

Let us consider the boundary value problem

$$\begin{cases} \varepsilon^2 u_{xx} + f(u) = 0, & \text{in } (0, 1), \\ u_x(0) = u_x(1) = 0, \\ u > 0, & \text{in } (0, 1). \end{cases} \quad (\text{NP})$$

Here ε is a positive parameter and

$$f(u) = -u + u^p$$

with $p > 1$. The problem (NP) is called the scalar field equation, which is motivated by the parabolic PDEs modelling pattern formation and chemotaxis, and also by PDEs of Klein-Gordon/nonlinear Schrödinger type ([2] and [4]).

In the 1-D problem (NP) a classical shooting argument solves the global bifurcation structure of all solutions ([5]). In particular, one can observe the families of periodic solutions: for any n , there exist $\varepsilon_n > 0$ such that there exist $1/n$ -periodic solutions $u_{n,\varepsilon}^\pm(x)$ for $\varepsilon \in (0, \varepsilon_n)$;

$$u_{n,\varepsilon}^+(0) = \max_{x \in [0,1]} u_{n,\varepsilon}^+(x), \quad u_{n,\varepsilon}^-(0) = \min_{x \in [0,1]} u_{n,\varepsilon}^-(x). \quad (1.1)$$

In view of pattern formation, we say $u_{n,\varepsilon}^+$ is *interior spike* solution, and $u_{n,\varepsilon}^-$ is *boundary spike* solution.

We fix $n \in \mathbf{N}$ arbitrarily and assume $\varepsilon \in (0, \varepsilon_n)$. Consider the linearized eigenvalue problems associated with $u_{n,\varepsilon}^+$ and $u_{n,\varepsilon}^-$

$$\begin{cases} \varepsilon^2 \varphi_{xx}(x) + f_u(u_{n,\varepsilon}^\pm(x))\varphi(x) + \lambda\varphi(x) = 0 & \text{in } (0, 1), \\ \varphi_x(0) = \varphi_x(1) = 0. \end{cases} \quad (\text{LP}_\pm)$$

For $j \in \mathbf{N} \cup \{0\}$ we denote by $\lambda_j^\pm = \lambda_{j,\varepsilon}^\pm$ and $\varphi_j^\pm(x) = \varphi_{j,\varepsilon}^\pm(x)$, the $(j + 1)$ -th eigenvalue and the corresponding eigenfunction, respectively. By the Sturm-Liouville theory shows that

$$\begin{aligned} \lambda_{0,\varepsilon}^+ &< \lambda_{1,\varepsilon}^+ < \cdots < \lambda_{j,\varepsilon}^+ < \cdots < +\infty, \\ \lambda_{0,\varepsilon}^- &< \lambda_{1,\varepsilon}^- < \cdots < \lambda_{j,\varepsilon}^- < \cdots < +\infty, \end{aligned}$$

and the both $\varphi_{j,\varepsilon}^\pm$ have exactly j zeros in $(0, 1)$.

We are interested in the ε -dependence of the both eigenpairs $\lambda_{j,\varepsilon}^\pm$ and $\varphi_{j,\varepsilon}^\pm$ for (LP_\pm) , and the qualitative difference between $\{\varphi_{j,\varepsilon}^+\}$ and $\{\varphi_{j,\varepsilon}^-\}$. In the above problems we focus on the case $p = 3$:

$$f(u) = -u + u^3.$$

In this case we can give expressions of $u_{n,\varepsilon}^\pm$ in terms of the Jacobi elliptic function. We note that the similar expressions could be observed for the case $p = 2$. Moreover, *the method of representation equation*, which is proposed by the author and S. Yostsutani [7], [10] and [12] for analyzing the linearized eigenvalue problems to (NP) with bistable nonlinearities, can be applied to show the precise information on eigenvalues and eigenfunctions of (LP_\pm) :

- representation formulas on eigenfunctions,
- characterization of eigenvalues,
- asymptotic formulas of eigenpairs as $\varepsilon \rightarrow 0$.

More precisely, we obtain the three special eigenfunctions for both (LP_+) and (LP_-) when $p = 3$: $\varphi_{0,\varepsilon}^\pm$, $\varphi_{0,\varepsilon}^\pm$ and $\varphi_{2n,\varepsilon}^\pm$. They are given explicitly in the algebraic expression of elliptic function, and also possess fine symmetric and periodic properties. Moreover, the other eigenfunctions can be given by the Liouvillian form ([3]). The corresponding eigenvalues $\lambda_{j,\varepsilon}^\pm$ of (LP_\pm) is determined by *the characteristic equation*,

$$\mathcal{A}(k, \mu) = \frac{j\pi}{2n},$$

where \mathcal{A} is the characteristic function, which can be expressed by the complete elliptic integral (of the third kind).

The asymptotic formulas of every eigenvalue are obtained by asymptotic formulas on elliptic integrals, and the asymptotic profiles of every eigenfunction are obtained by the Floquet/Bloch type argument with the asymptotic formulas of eigenvalues ([11] and [13]).

It should be noted that in the both problem (LP_\pm) for $f(u) = -u + u^3$ the three special eigenpairs play an important role for the *limit classification*; every eigenpair is classified into the three classes, and roughly speaking, each class is provided by the corresponding special eigenpair. The limit classification above is described by the asymptotic formulas of eigenvalues and the asymptotic profiles of eigenfunctions when ε is sufficiently small. We will see that the limit classification on (LP_+) is essentially the same type as the limit classifications for the cases of nonlinear bistable f 's ([11] and [13]), which justify and generalize a conjecture by E. Yanagida on the asymptotic profiles of the first n -eigenfunctions. The limit classification for (LP_-) will show us that Yanagida's conjecture also hold for the first n -eigenfunctions. However, it appears a different asymptotic characteraton on the eigenfunctions from the case (LP_+) . which is due to the difference of the special eigenvalues between in (LP_+) and (LP_-) . See Figs. 1 and 2.

In this article we obtain the representation formulas of eigenfunctions and derive the characteristic functions for (LP_\pm) . Also, we discuss the asymptotic results on the eigenvalues and eigenfunctions. The organization of the article is as follows. In Section 2 we prepare the elliptic integrals and functions. In Section 3 we give the main results. In Section 4 we introduce some key lemmas in the analysis and give a sketch of proofs.

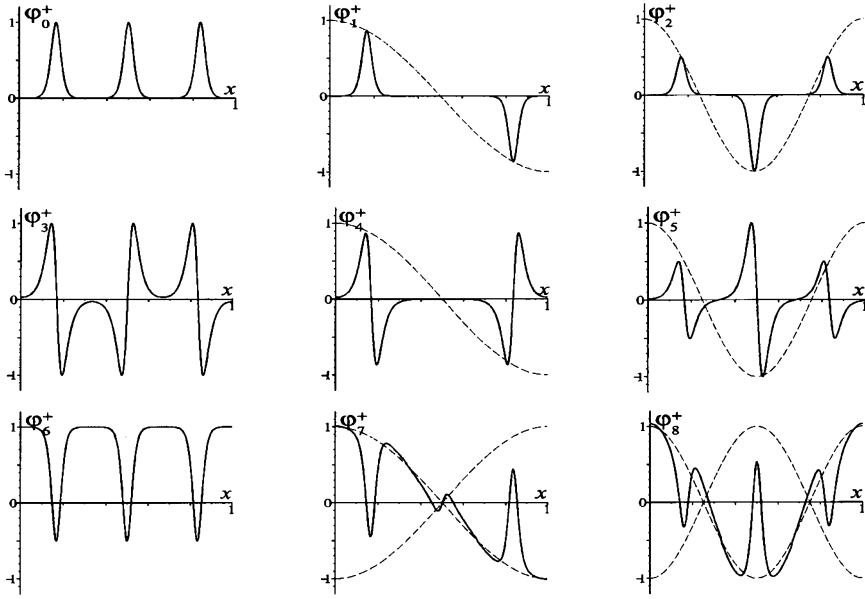


Figure 1: Profiles of $\varphi_{j,\varepsilon}^+$ ($j = 0, \dots, 8$) for $f(u) = -u + u^3$ with $n = 3$.

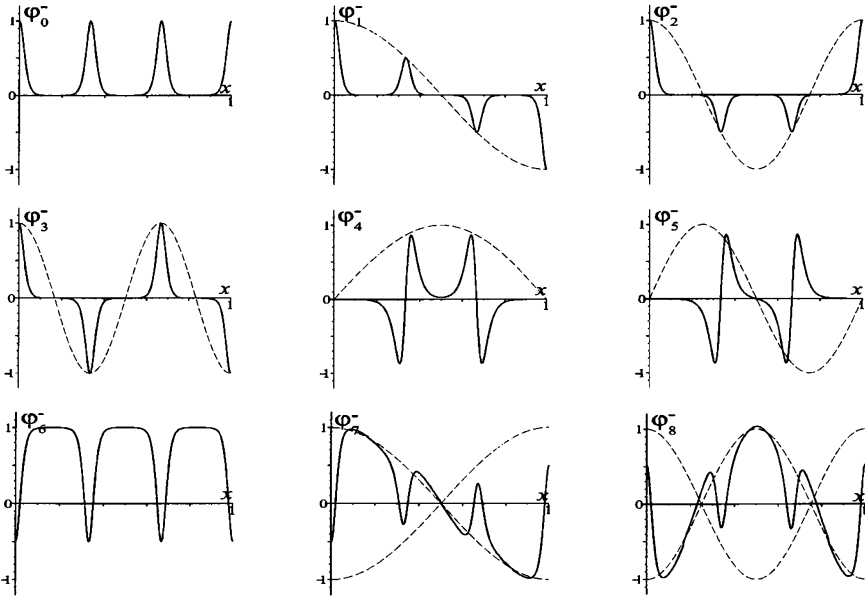


Figure 2: Profiles of $\varphi_{j,\varepsilon}^-$ ($j = 0, \dots, 8$) for $f(u) = -u + u^3$ with $n = 3$.

2 Preliminaries

We start with giving standard notations on the elliptic integrals and elliptic functions. See for instance, the Handbook by Byrd and Friedman[1]. Let $k \in (0, 1)$ and let $\nu \in \mathbf{C} \setminus (-\infty, -1]$. We denote by $K(k)$, $E(k)$ and $\Pi(\nu, k)$, the complete elliptic integrals of the first, the second and the third kind, respectively. The Jacobi elliptic function $\operatorname{sn}(x, k)$ is a periodic analytic function with the period $4K(k)$ as the function for the real domain, and is defined locally by

$$x = \int_0^{\operatorname{sn}(x, k)} \frac{1}{\sqrt{(1-s^2)(1-k^2s^2)}} ds$$

for $x \in [0, K(k)]$. In a standard manner the elliptic functions $\operatorname{cn}(x, k)$ and $\operatorname{dn}(x, k)$ is defined by

$$\operatorname{sn}(x, k)^2 + \operatorname{cn}(x, k)^2 = 1, \quad \operatorname{sn}(x, k)^2 + k^2 \operatorname{dn}(x, k)^2 = 1$$

for $x \in \mathbf{R}$ and $k \in (0, 1)$.

Here we introduce the modified form of Π (see [8] and [9])

$$\begin{aligned} \mathcal{M}(\nu, k) &:= \sqrt{\frac{(1+\nu)(k^2+\nu)}{\nu}} \int_0^1 \frac{1}{(1+\nu s^2)\sqrt{(1-s^2)(1-k^2s^2)}} ds \\ &\left(= \sqrt{\frac{(1+\nu)(k^2+\nu)}{\nu}} \Pi(\nu, k) \right) \end{aligned} \quad (2.1)$$

for $(\nu, k) \in \mathcal{D}$, where

$$\sqrt{z} := \exp\left(\frac{1}{2}\operatorname{Log} z\right), \quad z \in \mathbf{C} \setminus (-\infty, 0],$$

$$\mathcal{D} := \{(\nu, k) \in \mathbf{C} \times \mathbf{R} \mid \nu \notin (-\infty, -1] \cup [-k^2, 0]\}.$$

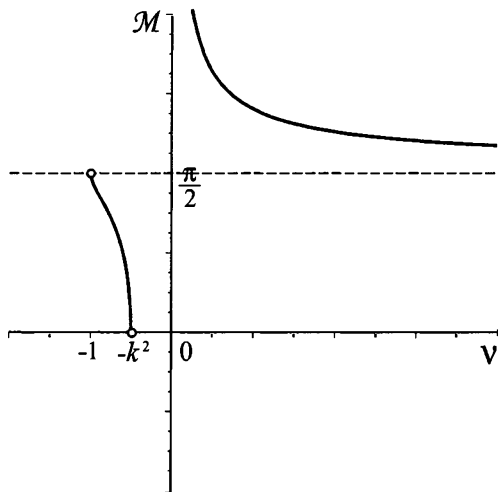


Figure 3: A Graph of $\mathcal{M}(\nu, k)$ for real ν ($k = 1/\sqrt{2}$).

3 Main results

Let us show the expressions of $u_{n,\varepsilon}^\pm$ of (NP) for $f(u) = -u + u^3$. A standard shooting argument can be applied to (NP) (see Section 4 of [12]), and in a situation that ε is small enough, one can obtain the monotone increasing and decreasing solutions of (NP), which are symmetric with the line $x = 1/2$. In addition, one can describe the global bifurcation structure of (NP) by using the two monotone solutions. In particular, we obtain the $1/n$ -periodic solutions $u_{n,\varepsilon}^\pm$ for each $n \in \mathbf{N}$.

Proposition 3.1. *Fix $n \in \mathbf{N}$ arbitrarily and assume $\varepsilon \in (0, \sqrt{2}/(n\pi))$. Let $k_\varepsilon = k_{n,\varepsilon} \in (0, 1)$ be a solution of*

$$\sqrt{2 - k^2}K(k) = \frac{1}{n\varepsilon} \quad (3.1)$$

(note that it is uniquely determined). Then,

$$u_{n,\varepsilon}^+(x) = \sqrt{\frac{2}{2 - k_\varepsilon^2}} \operatorname{dn}(K(k_\varepsilon)(1 + 2nx), k_\varepsilon),$$

and

$$u_{n,\varepsilon}^-(x) = \sqrt{\frac{2}{2 - k_\varepsilon^2}} \operatorname{dn}(2nK(k_\varepsilon)x, k_\varepsilon).$$

is the $1/n$ -periodic solutions of (NP) satisfying (1.1).

For a convenience, we introduce a notation of the two modified sn-functions

$$\operatorname{SN}^+(x; k) = \operatorname{sn}(K(k)(1 + 2nx), k), \quad \operatorname{SN}^-(x; k) = \operatorname{sn}(2nK(k)x, k).$$

Both functions have the same period $1/n$ for any $k \in (0, 1)$. The functions CN^\pm and DN^\pm are defined in a similar manner;

$$u_{n,\varepsilon}^\pm(x) = \sqrt{\frac{2}{2 - k_\varepsilon^2}} \operatorname{DN}^\pm(x; k_\varepsilon).$$

The representation formulas of eigenvalues and eigenfunctions are given by the following two theorems.

Theorem 1. *The Linearized problem (LP₊) has the following pairs of eigenvalues and eigenfunctions:*

- (i) $\lambda_{0,\varepsilon}^+ = -1 - \frac{2\sqrt{1 - k_\varepsilon^2 + k_\varepsilon^4}}{2 - k_\varepsilon^2}$,
 $\varphi_{0,\varepsilon}^+(x) = 1 - (1 + k_\varepsilon^2 - \sqrt{1 - k_\varepsilon^2 + k_\varepsilon^4})\operatorname{SN}^+(x; k_\varepsilon)^2$,
- (ii) $\lambda_{n,\varepsilon}^+ = -\frac{3(1 - k_\varepsilon^2)}{2 - k_\varepsilon^2}$, $\varphi_{n,\varepsilon}^+(x) = \operatorname{SN}^+(x, k_\varepsilon)\operatorname{DN}^+(x, k_\varepsilon)$,

$$(iii) \quad \lambda_{2n,\varepsilon}^+ = -1 + \frac{2\sqrt{1-k_\varepsilon^2+k_\varepsilon^4}}{2-k_\varepsilon^2},$$

$$\varphi_{2n,\varepsilon}^+(x) = -1 + (1+k_\varepsilon^2 + \sqrt{1-k_\varepsilon^2+k_\varepsilon^4})\text{SN}^+(x; k_\varepsilon)^2,$$

where k_ε is the unique solution of (3.1).

Theorem 2. *The Linearized problem (LP₋) has the following pairs of eigenvalues and eigenfunctions:*

$$(i) \quad \lambda_{0,\varepsilon}^- = -1 - \frac{2\sqrt{1-k_\varepsilon^2+k_\varepsilon^4}}{2-k_\varepsilon^2},$$

$$\varphi_{0,\varepsilon}^-(x) = 1 - (1+k_\varepsilon^2 - \sqrt{1-k_\varepsilon^2+k_\varepsilon^4})\text{SN}^-(x; k_\varepsilon)^2,$$

$$(ii) \quad \lambda_{n,\varepsilon}^- = -\frac{3}{2-k_\varepsilon^2}, \quad \varphi_{n,\varepsilon}^-(x) = \text{CN}(x, k_\varepsilon)\text{DN}(x, k_\varepsilon),$$

$$(iii) \quad \lambda_{2n,\varepsilon}^- = -1 + \frac{2\sqrt{1-k_\varepsilon^2+k_\varepsilon^4}}{2-k_\varepsilon^2},$$

$$\varphi_{2n,\varepsilon}^-(x) = -1 + (1+k_\varepsilon^2 + \sqrt{1-k_\varepsilon^2+k_\varepsilon^4})\text{SN}^-(x; k_\varepsilon)^2,$$

where k_ε is the unique solution of (3.1).

Here we introduce the following notations

$$\left\{ \begin{array}{l} \Lambda_0(k) := -1 - \frac{2\sqrt{1-k^2+k^4}}{2-k^2}, \\ \Lambda_1^-(k) := -\frac{3}{2-k^2}, \\ \Lambda_1^+(k) := -\frac{3(1-k^2)}{2-k^2}, \\ \Lambda_2(k) := -1 + \frac{2\sqrt{1-k^2+k^4}}{2-k^2}. \end{array} \right. \quad (3.2)$$

Note that $\Lambda_0(k) < \Lambda_1^-(k) < \Lambda_1^+(k) < 0 < \Lambda_2(k)$ for $k \in (0, 1)$. In addition, the functions Λ_0 , Λ_1^- and the functions Λ_1^+ , Λ_2 are monotone decreasing and increasing in $k \in (0, 1)$, respectively. Also, $\Lambda_1(0) = -2$, $\Lambda^\pm(0) = -3/2$, $\Lambda_2(0) = 0$ and $\Lambda_0(1) = \Lambda_1^-(1) = -3$, $\Lambda_1^+(1) = 0$, $\Lambda_2(1) = 1$. We see from Theorems 1 and 2 that

$$\lambda_{0,\varepsilon}^\pm = \Lambda_0(k_\varepsilon), \quad \lambda_{2n,\varepsilon}^\pm = \Lambda_2(k_\varepsilon),$$

$$\lambda_{n,\varepsilon}^- = \Lambda_1^-(k_\varepsilon), \quad \lambda_{n,\varepsilon}^+ = \Lambda_1^+(k_\varepsilon),$$

and in particular,

$$\lambda_{0,\varepsilon}^- = \lambda_{0,\varepsilon}^+, \quad \lambda_{+, \varepsilon}^- < \lambda_{n,\varepsilon}^+, \quad \lambda_{2n,\varepsilon}^- = \lambda_{2n,\varepsilon}^+.$$

For $j \neq 0, n, 2n$, let us show the representation formula of $\varphi_{j,\varepsilon}^\pm$. Set

$$\begin{aligned} \Sigma_0 &:= \{(k, \mu) \mid k \in (0, 1), \mu/(2-k^2) \in (\Lambda_0(k), \Lambda_1^-(k))\}, \\ \Sigma_1 &:= \{(k, \mu) \mid k \in (0, 1), \mu/(2-k^2) \in (\Lambda_1^+(k), 0)\}, \\ \Sigma_2 &:= \{(k, \mu) \mid k \in (0, 1), \mu/(2-k^2) \in (\Lambda_2(k), +\infty)\}. \end{aligned} \quad (3.3)$$

and set

$$\Sigma := \Sigma_0 \cup \Sigma_1 \cup \Sigma_2.$$

For $(k, \mu) \in \Sigma$, the characteristic function is defined by

$$\mathcal{A}(\mu, k) := |\mathcal{M}(\nu_-(k, \mu), k) - \mathcal{M}(\nu_+(k, \mu), k)|, \quad (3.4)$$

where \mathcal{M} is defined by (2.1),

$$\nu_{\pm}(k, \mu) := \frac{3k^2}{2} \cdot \frac{\mu - 3k^2 \pm \sqrt{-3\mu^2 + 6(k^2 - 2)\mu + 9k^4}}{\mu(\mu + 3 - 3k^2)}. \quad (3.5)$$

Theorem 3. *Suppose $j \neq 0, n, 2n$. Let k_{ε} be the solution of (3.1) and let $\mu_j(k)$ be the unique solution of*

$$\mathcal{A}(k, \mu) = \frac{j\pi}{n}, \quad (3.6)$$

where \mathcal{A} is given by (3.4). Then,

$$\lambda_{j,\varepsilon}^{\pm} = \frac{\mu_j(k_{\varepsilon})}{2 - k_{\varepsilon}^2}.$$

Moreover,

$$\varphi_{j,\varepsilon}^{\pm}(x) = \sqrt{|h(u_{n,\varepsilon}^{\pm}(x), \lambda_{j,\varepsilon}^{\pm}(k_{\varepsilon}); k_{\varepsilon})|} \cos \left(\frac{1}{\varepsilon} \int_0^x \frac{\sqrt{\rho(\lambda_{j,\varepsilon}^{\pm}, k_{\varepsilon})}}{|h(u_{n,\varepsilon}^{\pm}(\xi), \lambda_{j,\varepsilon}^{\pm}(k_{\varepsilon}); k_{\varepsilon})|} d\xi \right),$$

where

$$h(u, \lambda, k) := - \left[\frac{k^2}{(1+k^2)^2} - \frac{1}{2}u^2 + \frac{1}{4}u^4 \right] + \frac{\lambda}{6}(u^2 - 2) - \frac{\lambda^2}{9} \quad (3.7)$$

and

$$\rho(\lambda, k) = \frac{1}{81}(\lambda - \Lambda_0(k))(\lambda - \Lambda_1^-(k))(\lambda - \Lambda_1^+(k))\lambda(\lambda - \Lambda_2(k)). \quad (3.8)$$

Combining Theorems 1-3 we have all eigenvalues and eigenfunctions to the both (LP_+) and (LP_-) .

4 Key lemmas on characteristic function

Proofs of Theorems 1-3 are done by an algorithm of the representation equation for the linearized equation of (LP_{\pm}) (see [12]). We would like to describe whole procedure in the forth-coming papers.

Here we only focus on the two key lemmas for justifying existence and uniqueness of solution to (3.6).

4.1 Fundamental properties on the modified elliptic integral

We prepare a fundamental properties on the modified elliptic integral \mathcal{M} ([8] and [9]).

Proposition 4.1. *Let $(\nu, k) \in \mathcal{D}$ and let \mathcal{M} be defined by (2.1) . Then*

$$(i) \quad \frac{\partial \mathcal{M}}{\partial k}(\nu, k) = \sqrt{\frac{(1+\nu)(k^2+\nu)}{\nu}} \frac{kE(k)}{(k^2+\nu)(1-k^2)},$$

$$(ii) \quad \frac{\partial \mathcal{M}}{\partial \nu}(\nu, k) = \sqrt{\frac{(1+\nu)(k^2+\nu)}{\nu}} \left[-\frac{K(k)}{2\nu(1+\nu)} + \frac{E(k)}{2(1+\nu)(k^2+\nu)} \right].$$

Remark 4.1. It is well known that $(1-k^2)K(k) < E(k) < K(k)$ for $k \in (0, 1)$. Se we see from (ii) of Proposition 4.1 that for each $k \in (0, 1)$ the real function $\mathcal{M}(\cdot, k)$ is decreasing for $\nu \in (-1, k^2) \cup (0, \infty)$ (see Fig. 3).

Proposition 4.2. *Fix $k \in (0, 1)$ and consider the real function $\mathcal{M}(\cdot, k) : (-1, -k^2) \cup (0, \infty) \rightarrow \mathbf{R}$. Then the following formulas hold:*

$$\lim_{\nu \rightarrow -1} \mathcal{M}(\nu, k) = \frac{\pi}{2}, \quad (4.1)$$

$$\lim_{\nu \rightarrow -k^2} \mathcal{M}(\nu, k) = 0, \quad (4.2)$$

$$\lim_{\nu \rightarrow 0} \mathcal{M}(\nu, k) = \infty, \quad (4.3)$$

$$\lim_{\nu \rightarrow \infty} \mathcal{M}(\nu, k) = \frac{\pi}{2}. \quad (4.4)$$

4.2 Fundamental properties of ν_{\pm}

Recall the characteristic function

$$\mathcal{A}(k, \mu) := |\mathcal{M}(\nu_+(k, \mu), k) - \mathcal{M}(\nu_-(k, \mu), k)|$$

for $(k, \mu) \in \Sigma$. We first remark that $\nu_{\pm}(k, \mu)$ of (3.5) is characterized by

$$\nu_+(k, \mu) + \nu_-(k, \mu) = \frac{3k^2(\mu - 3k^2)}{\mu(\mu - 3 + 3k^2)}$$

and

$$\nu_+(k, \mu)\nu_-(k, \mu) = \frac{9k^4}{\mu(\mu - 3 + 3k^2)}.$$

Then an elementary algebra gives

$$(\nu_+(k, \mu) - \nu_-(k, \mu))^2 = \frac{-27k^4(\mu - (2 - k^2)\Lambda_0(k))(\mu - (2 - k^2)\Lambda_2(k))}{\mu^2(\mu + 3 - 3k^2)^2}.$$

Moreover, we see

$$\begin{cases} \frac{1}{\nu_+(k, \mu)} + \frac{1}{\nu_-(k, \mu)} = \frac{\mu - 3k^2}{3k^2}, \\ \frac{1}{\nu_+(k, \mu)} \cdot \frac{1}{\nu_-(k, \mu)} = \frac{\mu(\mu + 3 - 3k^2)}{9k^4}, \end{cases} \quad (4.5)$$

$$\begin{cases} \left(1 + \frac{1}{\nu_+(k, \mu)}\right) + \left(1 + \frac{1}{\nu_-(k, \mu)}\right) = \frac{\mu + 3k^2}{3k^2}, \\ \left(1 + \frac{1}{\nu_+(k, \mu)}\right) \left(1 + \frac{1}{\nu_-(k, \mu)}\right) = \frac{\mu(\mu + 3)}{9k^4}, \end{cases} \quad (4.6)$$

and

$$\begin{cases} \left(1 + \frac{k^2}{\nu_+(k, \mu)}\right) + \left(1 + \frac{k^2}{\nu_-(k, \mu)}\right) = \frac{\mu + 6 - 3k^2}{3k^2}, \\ \left(1 + \frac{k^2}{\nu_+(k, \mu)}\right) \left(1 + \frac{k^2}{\nu_-(k, \mu)}\right) = \frac{(\mu + 3)(\mu + 3 - 3k^2)}{9}, \end{cases} \quad (4.7)$$

respectively.

The following proposition helps some calculations on the characteristic functions together with \mathcal{M} .

Proposition 4.3. For $\nu_{\pm} = \nu_{\pm}(k, \mu)$ of (3.5)

$$\frac{1}{\nu_{\pm}} \left(1 + \frac{1}{\nu_{\pm}}\right) \left(1 + \frac{k^2}{\nu_{\pm}}\right) = \frac{\mu(\mu - 3)(\mu + 3 - 3k^2)}{27k^4}.$$

4.3 Several limit for characteristic function

Lemma 4.1. Let $k \in (0, 1)$ be fixed and $(k, \mu) \in \Sigma$. Then the following (i) and (ii) hold true:

$$(i) \quad \lim_{\mu \rightarrow (2-k^2)\Lambda_0(k)} \mathcal{A}(k, \mu) = 0, \quad \lim_{\mu \rightarrow -3} \mathcal{A}(k, \mu) = \frac{\pi}{2}.$$

$$(ii) \quad \lim_{\mu \rightarrow -3(1-k^2)} \mathcal{A}(k, \mu) = \frac{\pi}{2}, \quad \lim_{\mu \rightarrow 0} \mathcal{A}(k, \mu) = \pi.$$

Proof of Lemma 4.1. (i) Suppose that $(k, \mu) \in \Sigma_0$. Then, it follows from (4.5)-(4.7) that $-1 < \nu_-(k, \mu) < \nu_+(k, \mu) < -k^2$ and

- $\lim_{\mu \rightarrow (2-k^2)\Lambda_0(k)} \nu_{\pm}(k, \mu) = \nu_*(k) \in (-1, k^2),$
- $\lim_{\mu \rightarrow -3} \nu_-(k, \mu) = -1, \quad \lim_{\mu \rightarrow -3} \nu_+(k, \mu) = -k^2.$

Therefore, by using monotonicity of \mathcal{M}

$$\begin{aligned} \lim_{\mu \rightarrow (2-k^2)\Lambda_0(k)} \mathcal{A}(k, \mu) &= \lim_{\mu \rightarrow (2-k^2)\Lambda_0(k)} (\mathcal{M}(\nu_-(k, \mu), k) - \mathcal{M}(\nu_+(k, \mu), k)) \\ &= \mathcal{M}(\nu_*(k), k) - \mathcal{M}(\nu_*(k), k) \\ &= 0. \end{aligned}$$

In the similar manner we see from (4.1) and (4.2) of Proposition 4.2 that

$$\begin{aligned}\lim_{\mu \rightarrow -3} \mathcal{A}(k, \mu) &= \lim_{\mu \rightarrow -3} (\mathcal{M}(\nu_-(k, \mu), k) - \mathcal{M}(\nu_+(k, \mu), k)) \\ &= (\mathcal{M}(-1, k) - \mathcal{M}(-k^2, k)) \\ &= \frac{\pi}{2}.\end{aligned}$$

Thus it complete a proof.

(ii) of the lemma is similarly done. So we omit it. \square

Previous results for the cases of bistable nonlinearities suggest that we will also have

$$(iii) \quad \lim_{\mu \rightarrow (2-k^2)\Lambda_2(k)} \mathcal{A}(k, \mu) = \pi.$$

However, a different type of asymptotic formulas should be applied to show the above claim. So we omit it.

4.4 Monotonicity of characteristic function

Lemma 4.2. *Let $(k, \mu) \in \Sigma$. For each $i = 0, 1, 2$, $\frac{\partial \mathcal{A}}{\partial \mu}(k, \mu) > 0$ in Σ_i .*

Proof. Denote $\nu_{\pm} := \nu_{\pm}(k, \mu)$ and set

$$\mathcal{R}(k, \mu) := \frac{\mu(\mu - 3)(\mu + 3 - 3k^2)}{27k^4}.$$

For simplicity we only consider the case $(k, \mu) \in \Sigma_0$. By using (ii) of Proposition 4.1 and Proposition 4.3, we are led to

$$\begin{aligned}\frac{\partial \mathcal{A}}{\partial \mu} &= \frac{\partial \mathcal{M}}{\partial \nu}(\nu_-, k) \cdot \frac{\partial \nu_-}{\partial \mu} - \frac{\partial \mathcal{M}}{\partial \nu}(\nu_+, k) \cdot \frac{\partial \nu_+}{\partial \mu} \\ &= \sqrt{\frac{(1 + \nu_-)(k^2 + \nu_-)}{\nu_-}} \left[-\frac{K(k)}{2\nu_-(1 + \nu_-)} + \frac{E(k)}{2(1 + \nu_-)(k^2 + \nu_-)} \right] \cdot \frac{\partial \nu_-}{\partial \mu} \\ &\quad - \sqrt{\frac{(1 + \nu_+)(k^2 + \nu_+)}{\nu_+}} \left[-\frac{K(k)}{2\nu_+(1 + \nu_+)} + \frac{E(k)}{2(1 + \nu_+)(k^2 + \nu_+)} \right] \cdot \frac{\partial \nu_+}{\partial \mu} \\ &= \sqrt{\mathcal{R}(k, \mu)} \left(-\nu_- \frac{\partial \nu_-}{\partial \mu} \right) \left[-\frac{K(k)}{2\nu_-(1 + \nu_-)} + \frac{E(k)}{2(1 + \nu_-)(k^2 + \nu_-)} \right] \\ &\quad - \sqrt{\mathcal{R}(k, \mu)} \left(-\nu_+ \frac{\partial \nu_+}{\partial \mu} \right) \left[-\frac{K(k)}{2\nu_+(1 + \nu_+)} + \frac{E(k)}{2(1 + \nu_+)(k^2 + \nu_+)} \right] \\ &= \frac{\sqrt{\mathcal{R}(k, \mu)}}{2} \left[\left(\frac{1}{1 + \nu_-} \frac{\partial \nu_-}{\partial \mu} - \frac{1}{1 + \nu_+} \frac{\partial \nu_+}{\partial \mu} \right) K(k) \right. \\ &\quad \left. - \left(\frac{\nu_-}{(1 + \nu_-)(k^2 + \nu_-)} \frac{\partial \nu_-}{\partial \mu} - \frac{\nu_+}{(1 + \nu_+)(k^2 + \nu_+)} \frac{\partial \nu_+}{\partial \mu} \right) E(k) \right]\end{aligned}$$

and furthermore,

$$\begin{aligned}
 \frac{\partial \mathcal{A}}{\partial \mu}(k, \mu) &= \frac{1}{2\sqrt{\mathcal{R}(k, \mu)}} \left[\left(\frac{k^2 + \nu_-}{\nu_-^2} \frac{\partial \nu_-}{\partial \mu} - \frac{k^2 + \nu_+}{\nu_+^2} \frac{\partial \nu_+}{\partial \mu} \right) K(k) \right. \\
 &\quad \left. - \left(\frac{1}{\nu_-^2} \frac{\partial \nu_-}{\partial \mu} - \frac{1}{\nu_+^2} \frac{\partial \nu_+}{\partial \mu} \right) E(k) \right] \\
 &= \frac{1}{2\sqrt{\mathcal{R}(k, \mu)}} \left[\frac{\partial}{\partial \mu} \left[\left(\frac{1}{\nu_+} - \frac{1}{\nu_-} \right) \left(1 + \frac{k^2}{2} \left(\frac{1}{\nu_+} + \frac{1}{\nu_-} \right) \right) \right] K(k) \right. \\
 &\quad \left. - \frac{\partial}{\partial \mu} \left(\frac{1}{\nu_+} - \frac{1}{\nu_-} \right) E(k) \right] \\
 &= \frac{\mathcal{S}_1(k, \mu)K(k) + \mathcal{S}_2(k, \mu)E(k)}{6k^2\sqrt{\mathcal{R}(k, \mu)}\sqrt{-3\mu^2 + 6(k^2 - 2)\mu + 9k^4}},
 \end{aligned}$$

where

$$\mathcal{S}_1 := \mu^2 + 3(2 - k^2)\mu + 6(1 - k^2), \quad \mathcal{S}_2 := -3(\mu - k^2 + 2).$$

Finally, the claim of the lemma is proved by repeating argument in [12, Proof of Lemma 6.2]. \square

References

- [1] P. F. Byrd and M. D. Friedman, “Handbook of Elliptic Integrals for Engineers and Scientists”, Springer-Verlag, 1981.
- [2] H. Berestycki and P.-L. Lions, *Nonlinear scalar field equations I: existence of a ground state*, Arch. Rational Mech. Anal., **82** (1983), 313-345.
- [3] J. Kovacic, *An algorithm for solving second order linear homogeneous differential equations*, J. Symbolic Comput., **2** (1986), 3–43.
- [4] C.-S. Lin, W. Ni and I. Takagi, *Large amplitude stationary solutions to a Chemotaxis system*, J. Differential equations, **72** (1988), 1-27.
- [5] Y. Miyamoto and K. Yagasaki, *Monotonicity of the first eigenvalue and the global bifurcation diagram for the branch of interior peak solutions*, J. Differential equations, **254** (2013), 342-367.
- [6] K. Takemura, *The Heun equation and the Calogero- Moser-Sutherland system I: the Bethe Ansatz method*, Commun. Math. Phys. **235** (2003), 467–494.
- [7] T. Wakasa, *Exact eigenvalue and eigenfunction associated with linearization for Chafee-Infante problem*, Funkcialaj Ekvacioj, **49** (2006), 321–336.
- [8] T. Wakasa, *Note on parameter dependence of eigenvalues for a linearized eigenvalue problem*, Bull. Kyushu Inst. Tech., **49** (2016), 1–12.

- [9] T. Wakasa, *Analysis on the linearized eigenvalue problems and its application the global bifurcation problems*, Lecture note on “Study meeting on Applied Mathematics 2016” (Japanese).
- [10] T. Wakasa and S. Yotsutani, *Representation formulas for some 1-dimensional linearized eigenvalue problems*, Commun. Pure Appl. Anal. **7** (2008), 745–763.
- [11] T. Wakasa and S. Yotsutani, *Asymptotic profiles of eigenfunctions for some 1-dimensional linearized eigenvalue problems*, Commun. Pure Appl. Anal. **9** (2010).
- [12] T. Wakasa and S. Yotsutani, *Limiting classification on linearized eigenvalue problem for 1-dimensional Allen-Cahn Equation I -Asymptotic formulas of eigenvalues* J. Differential equations, **258** (2015), 3960-4006.
- [13] T. Wakasa and S. Yotsutani, *Limiting classification on linearized eigenvalue problem for 1-dimensional Allen-Cahn Equation II -Asymptotic profiles of eigenfunctions* J. Differential equations, **261** (2016), 5465-5498.
- [14] E. T. Whittaker and G. N. Watson, “A Course of Modern Analysis”, Fourth Edition, Cambridge University Press, New York, (1962).