

# RANDOM CHORDS AND POINT DISTANCES IN REGULAR POLYGONS 

U. BÄSEL


#### Abstract

In this paper we obtain the chord length distribution function for any regular polygon. From this function we conclude the density function and the distribution function of the distance between two uniformly and independently distributed random points in the regular polygon. The method calculating the chord length distribution function is quite different from those of Harutyunyan and Ohanyan. It uses only elementary methods and provides the result with only a few natural case distinctions.


## 1. Introduction

A random line $g$ intersecting a convex set $\mathcal{K}$ in the plane produces a chord of $\mathcal{K}$. The length $s$ of this chord is a random variable. If the motion invariant line measure (see below) is used for the definition of the line, the expectation of the chord length is equal to $\pi A / u$ where $A$ is the area of $\mathcal{K}$ and $u$ is the length of its perimeter [19, p. 30]. The chord length distribution function of a

Go back

Full Screen

Close

[^0]

Go back

Full Screen

Close
The distance $t$ between two points chosen independently and uniformly at random from $\mathcal{K}$ is also a random variable. Borel [4] considered this distance in elementary geometric figures such as triangles, squares and so on (see [17, p. 163]). The expectations for the distance between two random points for an equilateral triangle and a rectangle are to be found in [19, p. 49]. Ghosh [11] derived the distance distribution for a rectangle. Bailey, Borwein \& Crandall [3] studied the expected distance between two random points in the unit $n$-cube giving closed form expressions for the cases $n=1, \ldots, 5$ (see also [5], especially Example 14). There are a lot of results concerning the distance $t$ within a convex set or in two convex sets (see Chapter 2 in [16], and [7]).

The moments of $s$ and $t$, resp., are closely connected by a simple formula[19, pp. 46-47]. The second moments of the chord length for regular polygons were obtained by Heinrich [14].

For practical applications of chord lenghts and point distances of convex sets in physics, material sciences, operations research and other fields see [12] and [15].

The first aim of the present paper is to derive the chord length distribution function for any regular polygon in a simple form with only a few natural case distinctions using a method that requires only elementary geometric considerations and elementary integrations (especially not using Dirac's $\delta$-function in Pleijel's identity as done in [13]). Our method is also suitable for irregular and even (with slight modifications) non-convex polygons as shown in [2]. The second aim is to conclude the density function and the distribution function of the distance between two random points in every regular polygon. This result is new to the author's knowledge.

By $\mathcal{P}_{n, r}$ we denote the regular polygon with $n$ sides and circumscribed circle with radius $r$ and centre point in the origin $O$ (see Fig. 1). A straight line $g$ in the plane is determined by the angle $\phi, 0 \leq \phi<2 \pi$, between the direction perpendicular to $g$ and a fixed direction (e.g., the $x$-axis), and by its distance $p, 0 \leq p<\infty$, from the origin $O$

$$
g=g(p, \phi)=\left\{(x, y) \in \mathbb{R}^{2}: x \cos \phi+y \sin \phi=p\right\}
$$



Go back

The measure $\mu$ of a set of lines $g(p, \phi)$ is defined by the integral, over the set, of the differential form $\mathrm{d} g=\mathrm{d} p \mathrm{~d} \phi$. Up to a constant factor, this measure is the only one that is invariant under motions in the Euclidean plan [19, p. 28].

* 4 4 $\mid$ • $\mid$

Go back

Full Screen
The chord length distribution function of $\mathcal{P}_{n, r}$ is usually defined as

$$
F(s)=\frac{1}{u} \mu\left(\left\{g: g \cap \mathcal{P}_{n, r} \neq \emptyset,|\chi(g)| \leq s\right\}\right),
$$

where $\chi(g)=g \cap \mathcal{P}_{n, r}$ is the chord of $\mathcal{P}_{n, r}$, produced by the line $g,|\chi(g)|$ is the length of $\chi(g)$, and $u$ is the length of the perimeter of $\mathcal{P}_{n, r}$. (The measure of all lines $g$ that intersect a convex set is equal to its perimeter [19, p. 30].) We use the distribution function in the form

$$
\begin{equation*}
F(s)=1-\frac{1}{u} \mu\left(\left\{g: g \cap \mathcal{P}_{n, r} \neq \emptyset,|\chi(g)|>s\right\}\right) \tag{1}
\end{equation*}
$$

(cf. [1, p. 161]). So it remains to calculate the measure of all lines that produce a chord of length $|\chi(g)|>s$. Using the abbreviation

$$
\mathcal{S}(g: s):=\left\{g: g \cap \mathcal{P}_{n, r} \neq \emptyset,|\chi(g)|>s\right\},
$$

we have

$$
\mu(\mathcal{S}(g: s))=\int_{\mathcal{S}(g: s)} \mathrm{d} g=\int_{\mathcal{S}(g: s)} \mathrm{d} p \mathrm{~d} \phi
$$

We consider all lines $g$, having a direction perpendicular to a fixed angle $\phi \in[0, \pi)$ with $g \cap \mathcal{P}_{n, r} \neq \emptyset$. In almost all cases among these lines there are two lines $g_{1}$ and $g_{2}$ with chords of equal length $s$ (see Fig. 1). All parallel lines $g$ lying in the strip between $g_{1}$ and $g_{2}$ have a chord with length $|\chi(g)|>s$. The breadth of this strip is equal to $d(s, \phi)+d(s, \phi+\pi)$, where $d(s, \phi)$ and $d(s, \phi+\pi)$ are the distances between $O$ and $g_{1}$ and $O$ and $g_{2}$, repectively. So we have

$$
\begin{equation*}
\mu(\mathcal{S}(g: s))=\int_{0}^{\pi}[d(s, \phi)+d(s, \phi+\pi)] \mathrm{d} \phi \tag{2}
\end{equation*}
$$

## 2. The distance function

In the following we determine the distance function in formula (2)

$$
d:[0, \max (s)] \times[0, \infty) \rightarrow[0, r], \quad(s, \phi) \mapsto d(s, \phi),
$$

where $\max (s)$ is the maximum chord length $s$ in $\mathcal{P}_{n, r}$. If no chord of length $s$ in the direction perpendicular to $\phi$ exists, we put $d(s, \phi)=0$. Of course for fixed value of $s, d(s, \cdot)$ is a $2 \pi / n$ periodic function.

We put

$$
K=\left\lfloor\frac{n-2}{2}\right\rfloor
$$

where $\lfloor\cdot\rfloor$ is the integer part of $\cdot$, and define the function $m: \mathbb{N} \times \mathbb{N} \rightarrow \mathbb{N}$ by

$$
m(k, n)= \begin{cases}k \bmod n & \text { if } \quad k \bmod n \neq 0 \\ n & \text { if } \quad k \bmod n=0 .\end{cases}
$$

The angle $\delta_{k}$ (see Fig. 2) between the lines $i$ and $m(i+k, n)$ containing the sides $i, i=1, \ldots, n$, and $m(i+k, n)$ of $\mathcal{P}_{n, r}$ is given by

$$
\delta_{k}=\left(1-\frac{2 k}{n}\right) \pi, \quad k=1, \ldots, K^{*},
$$

where

$$
K^{*}= \begin{cases}K+1 & \text { if } n \text { is odd } \\ K & \text { if } n \text { is even }\end{cases}
$$

The distance $\ell_{k}$ between the vertices $i$ and $m(i+k, n)$ is for $k=0, \ldots, K+1$ given by

$$
\ell_{k}=2 r \sin \frac{k \pi}{n} .
$$



Go back

Full Screen

Close
The maximum chord length in $\mathcal{P}_{n, r}$ is equal to $\ell_{K+1}$. For the distance $x$ between one point of side $i$ and one point of side $m(i+k, n), k=1, \ldots, K^{*}$, we have $\ell_{k-1} \leq x \leq \ell_{k+1}$, and $\ell_{k} \leq x \leq \ell_{k+2}$ for the analogous distance of the sides $i$ and $m(i+k+1, n)$. Therefore, a chord of length $s$, $\ell_{k} \leq s \leq \ell_{k+1}$, is a chord between two sides $i$ and $m(i+k, n)$ or two sides $i$ and $m(i+k+1, n)$.
$44|4| \gg$
Go back

Full Screen

$$
\ell_{k} \leq s<\ell_{k+1} \quad \text { with } \quad\left\{\begin{array}{l}
k \in\{1, \ldots, K-1\} \text { if } n \text { is even, } \\
k \in\{1, \ldots, K\} \text { if } n \text { is odd and } s \leq 2 r \cos ^{2} \frac{\pi}{2 n}
\end{array}\right.
$$

For $0 \leq \psi \leq \alpha_{k}(s)$ (e.g., for position $\chi_{1}$ of $\chi$ in Fig. 2) the distance $d_{k}^{*}(s, \psi)$ between $O$ and $\chi$ is equal to

$$
\begin{equation*}
q_{k}(s, \psi):=r \cos \frac{\pi}{n} \sec \frac{k \pi}{n} \cos \psi-\frac{s}{2}\left(\tan \frac{k \pi}{n} \cos ^{2} \psi-\cot \frac{k \pi}{n} \sin ^{2} \psi\right) . \tag{5}
\end{equation*}
$$

The angle $\alpha_{k}$ is determined by the position $\chi_{2}$ of $\chi$ with the upper end-point in the vertex $m(i+$ $k+1, n$ )

$$
\begin{equation*}
\alpha_{k}(s)=\arcsin \left(\frac{2 r}{s} \sin \frac{k \pi}{n} \sin \frac{(k+1) \pi}{n}\right)-\frac{k \pi}{n} . \tag{6}
\end{equation*}
$$

For $\alpha_{k}(s) \leq \psi \leq \pi / n, \chi$ is a chord between the sides $i$ and $m(i+k+1, n)$. So we find

Case 1a

$$
d_{k}^{*}(s, \psi)=\left\{\begin{array}{lll}
q_{k}(s, \psi) & \text { if } \quad 0 \leq \psi \leq \alpha_{k}(s)  \tag{7}\\
q_{k+1}(s, \psi-\pi / n) & \text { if } \quad \alpha_{k}(s)<\psi \leq \pi / n
\end{array}\right.
$$

$$
\begin{equation*}
0=\ell_{0} \leq s<\ell_{1} \quad \text { and } \quad s<2 r \cos ^{2} \frac{\pi}{2 n} \tag{8}
\end{equation*}
$$

We have $\alpha_{0}(s)=0$ if $s \neq 0$, and the limit of $\alpha_{0}(s)$ at $s=0$ is 0 . Therefore, we get

$$
\begin{equation*}
d_{0}^{*}(s, \psi)=q_{1}(s, \psi-\pi / n) \quad \text { for } \quad 0 \leq \psi \leq \pi / n \tag{9}
\end{equation*}
$$

as special case of case 1 with $\alpha_{0}(s)=0$ in formula (7).
Case 2
(10)

$$
n \text { is even and } \ell_{K} \leq s \leq \ell_{K+1}
$$

A chord $\chi$ in the direction perpendicular to $\psi$ does not exist if $\alpha_{K}(s)<\psi \leq \pi / n$, therefore

$$
d_{K}^{*}(s, \psi)=\left\{\begin{array}{lll}
q_{K}(s, \psi) & \text { if } & 0 \leq \psi \leq \alpha_{K}(s),  \tag{11}\\
0 & \text { if } & \alpha_{K}(s)<\psi \leq \pi / n
\end{array}\right.
$$

Case 3

$$
\begin{equation*}
n \text { is odd and } \ell_{K} \leq s \leq \ell_{K+1} \text { and } s \geq 2 r \cos ^{2} \frac{\pi}{2 n} \tag{12}
\end{equation*}
$$

A chord $\chi$ in the direction perpendicular to $\psi$ does not exist if

$$
\beta(s)<\psi<\frac{\pi}{n}-\beta(s),
$$

with

$$
\begin{equation*}
\beta(s)=\frac{\pi}{2 n}-\arccos \left(\frac{2 r}{s} \cos ^{2} \frac{\pi}{2 n}\right) \tag{13}
\end{equation*}
$$

therefore,

$$
d_{K}^{*}(s, \psi)= \begin{cases}q_{K}(s, \psi) & \text { if } 0 \leq \psi<\alpha_{K}(s)  \tag{14}\\ q_{K+1}(s, \psi-\pi / n) & \text { if } \alpha_{K}(s) \leq \psi \leq \beta(s) \\ 0 & \text { if } \beta(s)<\psi<\pi / n-\beta(s) \\ q_{K+1}(s, \psi-\pi / n) & \text { if } \pi / n-\beta(s) \leq \psi \leq \pi / n\end{cases}
$$

Due to the symmetry of the graph of $d_{k}^{*}(s, \psi)$ with respect to the line $\psi=\pi / n$, the values in the interval $\pi / n<\psi \leq 2 \pi / n$ can be easily calculated from (7), (9), (11) and (14) with

$$
\begin{equation*}
d_{k}^{*}(s, \psi)=d_{k}^{*}(s, 2 \pi / n-\psi) \tag{15}
\end{equation*}
$$

Since $d_{k}^{*}(s, \psi)$ is a $2 \pi / n$-periodic function, we get the values for $2 \pi / n<\psi<\infty$ with the translation

$$
\psi \mapsto \psi-\delta(\psi) \quad \text { with } \quad \delta(\psi)=\left\lfloor\frac{n \psi}{2 \pi}\right\rfloor \frac{2 \pi}{n}
$$

In the case of even $n$, the substitution $\psi=\phi+\pi / n$ yields the distances for angle $\phi$ starting from a vertex as shown in Fig. 1. In the case of odd $n$ we have $\psi=\phi$. So we get the following lemma.


Go back

Full Screen

Close
Lemma 1. The restrictions $d_{k}(s, \phi)=\left.d(s, \phi)\right|_{\ell_{k} \leq s<\ell_{k+1}}$ of the distance function $d$ are given by

$$
d_{k}(s, \phi)= \begin{cases}d_{k}^{*}(s, \phi-\delta(\phi)) & \text { if } n \text { is odd } \\ d_{k}^{*}\left(s, \phi+\frac{\pi}{n}-\delta\left(\phi+\frac{\pi}{n}\right)\right) & \text { if } n \text { is even }\end{cases}
$$

for $k=0, \ldots, K$, where

$$
\delta(\cdot)=\left\lfloor\frac{n \cdot}{2 \pi}\right\rfloor \frac{2 \pi}{n}
$$



Go back

Full Screen

Close
and $d_{k}^{*}$ according to the formulas (7), (11), (14) and (15) with $\alpha_{k}$ and $\beta$ according to (6) and (13), respectively.

## 3. Chord length distribution function

So we can write the chord length distribution function (1) in the form

$$
F(s)= \begin{cases}0 & \text { if } \quad-\infty<s<\ell_{0}=0 \\ H_{k}(s) & \text { if } \quad \ell_{k} \leq s<\ell_{k+1} \text { for } k=0, \ldots, K, \\ 1 & \text { if } \quad \ell_{K+1} \leq s<\infty,\end{cases}
$$

where

$$
H_{k}(s)=1-\frac{\mu_{k}(s)}{2 n r \sin (\pi / n)} \text { with } \mu_{k}(s):=\int_{0}^{\pi}\left[d_{k}(s, \phi)+d_{k}(s, \phi+\pi)\right] \mathrm{d} \phi .
$$

With ( $\star$ ) the $2 \pi / n$-periodicity of $d_{k}(s, \phi)$ and $(\diamond)$ the symmetry of $d_{k}(s, \phi)$ with respect to the line $\phi=\pi / n$, for odd and even $n$, we find

$$
\begin{aligned}
\mu_{k}(s) & =\int_{0}^{\pi}\left[d_{k}(s, \phi)+d_{k}(s, \phi+\pi)\right] \mathrm{d} \phi \\
& \stackrel{\star}{=} \frac{n}{2} \int_{0}^{2 \pi / n}\left[d_{k}(s, \phi)+d_{k}(s, \phi+\pi)\right] \mathrm{d} \phi \\
& =\frac{n}{2} \int_{0}^{2 \pi / n}\left[d_{k}(s, \phi)+d_{k}\left(s, \phi+n \frac{2 \pi}{n}\right)\right] \mathrm{d} \phi \\
& \stackrel{\star}{=} n \int_{0}^{2 \pi / n} d_{k}(s, \phi) \mathrm{d} \phi \stackrel{ }{=} 2 n \int_{0}^{\pi / n} d_{k}(s, \phi) \mathrm{d} \phi .
\end{aligned}
$$

Note that this integral formula (together with the piecewise definition of the distance function) allows us to calculate the distribution function of every regular polygon in a rather simple way.

The indefinite integral of the function $q_{k}$ (see (5)) is given by

$$
\begin{align*}
J_{k}(s, \phi)= & \int q_{k}(s, \phi) \mathrm{d} \phi=r \cos \frac{\pi}{n} \sec \frac{k \pi}{n} \sin \phi \\
& -\frac{s}{8}\left[(2 \phi+\sin 2 \phi) \tan \frac{k \pi}{n}-(2 \phi-\sin 2 \phi) \cot \frac{k \pi}{n}\right] . \tag{16}
\end{align*}
$$

In case 1 (see (4)) we get with $\alpha_{k}=\alpha_{k}(s)$

$$
\begin{aligned}
\frac{\mu_{k}(s)}{2 n} & =\int_{0}^{\alpha_{k}} d_{k}(s, \phi) \mathrm{d} \phi+\int_{\alpha_{k}}^{\pi / n} d_{k}(s, \phi) \mathrm{d} \phi \\
& =\int_{0}^{\alpha_{k}} q_{k}(s, \phi) \mathrm{d} \phi+\int_{\alpha_{k}}^{\pi / n} q_{k+1}\left(s, \phi-\frac{\pi}{n}\right) \mathrm{d} \phi
\end{aligned}
$$

and with the substitution $\phi^{*}=\phi-\pi / n$

$$
\begin{aligned}
\frac{\mu_{k}(s)}{2 n} & =\int_{0}^{\alpha_{k}} q_{k}(s, \phi) \mathrm{d} \phi+\int_{\alpha_{k}-\pi / n}^{0} q_{k+1}\left(s, \phi^{*}\right) \mathrm{d} \phi^{*} \\
& =J_{k}\left(s, \alpha_{k}\right)-J_{k}(s, 0)+J_{k+1}(s, 0)-J_{k+1}\left(s, \alpha_{k}-\frac{\pi}{n}\right) \\
& =J_{k}\left(s, \alpha_{k}\right)-J_{k+1}\left(s, \alpha_{k}-\frac{\pi}{n}\right),
\end{aligned}
$$

since $J_{k}(s, 0)=0$. In case 1a (see (8)), one finds with $\alpha_{0}(s)=0$ and $J_{k}(s,-\phi)=-J_{k}(s, \phi)$

$$
\frac{\mu_{0}(s)}{2 n}=-J_{1}\left(s,-\frac{\pi}{n}\right)=J_{1}\left(s, \frac{\pi}{n}\right) .
$$

Putting $J_{0}(s, 0)=0$, this formula can be considered as a special case of the formula for case 1 . In case 2 (see (10)), one easily finds

$$
\frac{\mu_{K}(s)}{2 n}=J_{K}\left(s, \alpha_{K}\right)
$$

and in case 3 (see (12)) with $\beta=\beta(s)$, we have

$$
\begin{aligned}
\frac{\mu_{K}(s)}{2 n} & =\left[\int_{0}^{\alpha_{K}}+\int_{\alpha_{K}}^{\beta}+\int_{\beta}^{\pi / n-\beta}+\int_{\pi / n-\beta}^{\pi / n}\right] d_{K}(s, \phi) \mathrm{d} \phi \\
& =\left[\int_{0}^{\alpha_{K}}+\int_{\alpha_{K}}^{\beta}+\int_{\pi / n-\beta}^{\pi / n}\right] d_{K}(s, \phi) \mathrm{d} \phi \\
& =\int_{0}^{\alpha_{K}} q_{K}(s, \phi) \mathrm{d} \phi+\left[\int_{\alpha_{K}}^{\beta}+\int_{\pi / n-\beta}^{\pi / n}\right] q_{K+1}(s, \phi-\pi / n) \mathrm{d} \phi \\
& =\int_{0}^{\alpha_{K}} q_{K}(s, \phi) \mathrm{d} \phi+\left[\int_{\alpha_{K}-\pi / n}^{\beta-\pi / n}+\int_{-\beta}^{0}\right] q_{K+1}\left(s, \phi^{*}\right) \mathrm{d} \phi^{*} \\
& =J_{K}\left(s, \alpha_{K}\right)-J_{K+1}\left(s, \alpha_{K}-\frac{\pi}{n}\right)+J_{K+1}\left(s, \beta-\frac{\pi}{n}\right)+J_{K+1}(s, \beta) .
\end{aligned}
$$

Full Screen
The function $J_{k}$ (see (16)) can be written as

$$
J_{k}(s, \phi)=r \cos \frac{\pi}{n} \sec \frac{k \pi}{n} \sin \phi+\frac{s}{4}\left(2 \phi \cot \frac{2 k \pi}{n}-\sin (2 \phi) \csc \frac{2 k \pi}{n}\right)
$$

44 4 • $\mid$ •
Go back

Full Screen
Furthermore, we write both functions $\alpha_{k}$ (see (6)) and $\beta$ (see (13)) in the form

$$
\arcsin \frac{a}{s}-b
$$

with

$$
\begin{equation*}
a=A_{1}(k)=2 r \sin \frac{k \pi}{n} \sin \frac{(k+1) \pi}{n}, \quad b=B_{1}(k)=\frac{k \pi}{n} \tag{17}
\end{equation*}
$$

for $\alpha_{k}$, and

$$
\begin{equation*}
a=A_{2}=2 r \cos ^{2} \frac{\pi}{2 n}, \quad b=B_{2}=\frac{\pi}{2}\left(1-\frac{1}{n}\right) \tag{18}
\end{equation*}
$$

for $\beta$. Using some easy algebraic manipulations, one finds

$$
\begin{aligned}
\frac{1}{r} J_{k}\left(s, \arcsin \frac{a}{s}-b\right)= & \Theta_{1}(k, a, b) s+\Theta_{2}(k, a, b) \frac{1}{s}+\Theta_{3}(k, a, b) \frac{\sqrt{s^{2}-a^{2}}}{s} \\
& +\Theta_{4}(k, a, b) s \arcsin \frac{a}{s}=: h_{k}(s, a, b),
\end{aligned}
$$

where

$$
\left\{\begin{array}{l}
\Theta_{1}(k, a, b)=\frac{1}{4 r} \csc \frac{\pi}{n}\left(\sin (2 b) \csc \frac{2 k \pi}{n}-2 b \cot \frac{2 k \pi}{n}\right), \\
\Theta_{2}(k, a, b)=a\left(\cos b \cot \frac{\pi}{n} \sec \frac{k \pi}{n}-\frac{a}{2 r} \sin (2 b) \csc \frac{\pi}{n} \csc \frac{2 k \pi}{n}\right),  \tag{19}\\
\Theta_{3}(k, a, b)=-\left(\sin b \cot \frac{\pi}{n} \sec \frac{k \pi}{n}+\frac{a}{2 r} \cos (2 b) \csc \frac{\pi}{n} \csc \frac{2 k \pi}{n}\right), \\
\Theta_{4}(k, a, b)=\frac{1}{2 r} \csc \frac{\pi}{n} \cot \frac{2 k \pi}{n} .
\end{array}\right.
$$

In summary, we have proved the following theorem.
Theorem 1. The chord length distribution function $F$ of the regular polygon $\mathcal{P}_{n, r}$ is given by

$$
F(s)= \begin{cases}0 & \text { if }-\infty<s<\ell_{0}=0 \\ H_{k}(s) & \text { if } \ell_{k} \leq s<\ell_{k+1} \text { for } k=0, \ldots, K, \\ 1 & \text { if } \ell_{K+1} \leq s<\infty,\end{cases}
$$

where

$$
\ell_{k}=2 r \sin \frac{k \pi}{n}, \quad K=\left\lfloor\frac{n-2}{2}\right\rfloor
$$

and

$$
H_{k}(s)=\left\{\begin{array}{l}
1-h_{k}\left(s, A_{1}(k), B_{1}(k)\right)+h_{k+1}\left(s, A_{1}(k), B_{1}(k)+\pi / n\right) \\
\quad \text { if }(n \text { is even } \wedge k \in\{0, \ldots, K-1\}) \vee(n \text { is odd } \wedge s<\lambda), \\
1-h_{K}\left(s, A_{1}(K), B_{1}(K)\right)+h_{K+1}\left(s, A_{1}(K), B_{1}(K)+\pi / n\right) \\
-h_{K+1}\left(s, A_{2}, B_{2}+\pi / n\right)-h_{K+1}\left(s, A_{2}, B_{2}\right) \\
\quad \text { if } n \text { is odd } \wedge s \geq \lambda, \\
1-h_{K}\left(s, A_{1}(K), B_{1}(K)\right) \text { if } n \text { is even } \wedge k=K
\end{array}\right.
$$

with

$$
\lambda=2 r \cos ^{2} \frac{\pi}{2 n}
$$

$A_{1}(k)$ and $B_{1}(k)$ according to (17), $A_{2}$ and $B_{2}$ according to (18), and

$$
h_{k}(s, a, b)= \begin{cases}0 & \text { if } k=0 \\ \sum_{i=1}^{4} \Theta_{i}(k, a, b) L_{i}(s, a) & \text { if } k=1,2, \ldots\end{cases}
$$

with $\Theta_{i}(k, a, b)$ according to (19), and

$$
L_{1}(s, a)=s, \quad L_{2}(s, a)=\frac{1}{s}, \quad L_{3}(s, a)=\frac{\sqrt{s^{2}-a^{2}}}{s}, \quad L_{4}(s, a)=s \arcsin \frac{a}{s} .
$$

$F$ can be written in the form

$$
F(s)=H_{0}(s)=\left[\left(1-\frac{\pi}{n} \cot \frac{\pi}{n}\right) \csc \frac{\pi}{n}+\frac{\pi}{n} \sec \frac{\pi}{n}\right] \frac{s}{4 r}
$$

for $0 \leq s \leq \lambda$ if $n=3$, and $0 \leq s \leq \ell_{1}$ if $n=4,5, \ldots$. Note that this is a linear equation of $s$ (cf. [10, pp. 866-867]).

From [20, p. 55, Satz 2], it follows that the chord length distribution function of a regular polygon is a continuous function.

## 4. Point distances

In the following, we consider the distance between two uniformly and independently distributed random points within the polygon $\mathcal{P}_{n, r}$ with perimeter $u$ and area $A$

$$
u=2 n r \sin \frac{\pi}{n}, \quad A=\frac{1}{2} n r^{2} \sin \frac{2 \pi}{n} .
$$

Theorem 2. The density function $g$ of the distance $t$ between two random points in $\mathcal{P}_{n, r}$ is given by

$$
g(t)= \begin{cases}\frac{2 t}{A}\left[\pi+\frac{u}{A}\left(\phi^{*}(t)-t\right)\right] & \text { if } t \in\left[0, \ell_{K+1}\right), \\ 0 & \text { if } t \in \mathbb{R} \backslash\left[0, \ell_{K+1}\right),\end{cases}
$$

where

$$
\phi^{*}(t)=\sum_{\nu=0}^{k-1} J_{\nu}^{*}\left(\ell_{\nu}, \ell_{\nu+1}\right)+J_{k}^{*}\left(\ell_{k}, t\right) \quad \text { if } \quad \ell_{k} \leq t<\ell_{k+1}, k=0, \ldots, K
$$

with

$$
J_{k}^{*}(s, t)=H_{k}^{*}(t)-H_{k}^{*}(s),
$$

where

$$
H_{k}^{*}(t)=\left\{\begin{array}{l}
t-h_{k}^{*}\left(t, A_{1}(k), B_{1}(k)\right)+h_{k+1}^{*}\left(t, A_{1}(k), B_{1}(k)+\pi / n\right) \\
\quad \text { if }(n \text { is even } \wedge k \in\{0, \ldots, K-1\}) \vee(n \text { is odd } \wedge t<\lambda), \\
t-h_{K}^{*}\left(t, A_{1}(K), B_{1}(K)\right)+h_{K+1}^{*}\left(t, A_{1}(K), B_{1}(K)+\pi / n\right) \\
-h_{K+1}^{*}\left(t, A_{2}, B_{2}+\pi / n\right)-h_{K+1}^{*}\left(t, A_{2}, B_{2}\right) \\
+h_{K+1}^{*}\left(\lambda, A_{2}, B_{2}+\pi / n\right)+h_{K+1}^{*}\left(\lambda, A_{2}, B_{2}\right) \\
\quad \text { if } n \text { is odd } \wedge t \geq \lambda, \\
t-h_{K}^{*}\left(t, A_{1}(K), B_{1}(K)\right) \text { if } n \text { is even } \wedge k=K
\end{array}\right.
$$

with

$$
\begin{aligned}
& h_{k}^{*}(t, a, b)=\left\{\begin{array}{l}
0 \quad \text { if } \quad k=0 \vee(k \neq 0 \wedge t=0), \\
\sum_{i=1}^{4} \Theta_{i}(k, a, b) L_{i}^{*}(t, a) \quad \text { if } k \neq 0 \wedge t>0,
\end{array}\right. \\
& L_{1}^{*}(t, a)=\frac{t^{2}}{2}, \quad L_{2}^{*}(t, a)=\ln t, \quad L_{3}^{*}(t, a)=\sqrt{t^{2}-a^{2}}+a \arcsin \frac{a}{t}, \\
& L_{4}^{*}(t, a)=\frac{1}{2}\left(a \sqrt{t^{2}-a^{2}}+t^{2} \arcsin \frac{a}{t}\right) .
\end{aligned}
$$

Proof. According to Piefke [18, p. 130], the density function of the distance is given by

$$
g(t)=\frac{2 u t}{A^{2}} \int_{t}^{\ell_{K+1}}(s-t) f(s) \mathrm{d} s,
$$

where $f$ is the density function of the chord length. From integral geometry it is well-known that

$$
\int_{0}^{\ell_{K+1}} s f(s) \mathrm{d} s=\frac{\pi A}{u}
$$

(see [19, p. 47], [16, p. 94]), hence

$$
\int_{t}^{\ell_{K+1}} s f(s) \mathrm{d} s=\frac{\pi A}{u}-\int_{0}^{t} s f(s) \mathrm{d} s
$$

Using integration by parts, we have

$$
\int_{0}^{t} s f(s) \mathrm{d} s=\left.s F(s)\right|_{0} ^{t}-\int_{0}^{t} F(s) \mathrm{d} s=t F(t)-\int_{0}^{t} F(s) \mathrm{d} s
$$

Therefore, we obtain

$$
g(t)=\frac{2 t}{A}\left[\pi-\frac{u}{A}\left(t-\int_{0}^{t} F(s) \mathrm{d} s\right)\right]=\frac{2 t}{A}\left[\pi+\frac{u}{A}\left(\phi^{*}(t)-t\right)\right]
$$

with

$$
\phi^{*}(t):=\int_{0}^{t} F(s) \mathrm{d} s
$$

For $\ell_{k} \leq t<\ell_{k+1}, k=0, \ldots, K$, this yields

$$
\phi^{*}(t)=\sum_{\nu=0}^{k-1} \int_{\ell_{\nu}}^{\ell_{\nu+1}} H_{\nu}(s) \mathrm{d} s+\int_{\ell_{k}}^{t} H_{k}(s) \mathrm{d} s
$$

(in case $k=0$, the sum is empty). With

$$
H_{k}^{*}(t):=\int H_{k}(t) \mathrm{d} t \quad \text { und } \quad J_{k}^{*}(s, t)=H_{k}^{*}(t)-H_{k}^{*}(s)
$$

it follows that

$$
\begin{aligned}
\phi^{*}(t) & =\sum_{\nu=0}^{k-1}\left[H_{\nu}^{*}\left(\ell_{\nu+1}\right)-H_{\nu}^{*}\left(\ell_{\nu}\right)\right]+H_{k}^{*}(t)-H_{k}^{*}\left(\ell_{k}\right) \\
& =\sum_{\nu=0}^{k-1} J_{\nu}^{*}\left(\ell_{\nu}, \ell_{\nu+1}\right)+J_{k}^{*}\left(\ell_{k}, t\right)
\end{aligned}
$$

Furthermore, if $k \neq 0$ and $t>0$,

$$
\begin{aligned}
h_{k}^{*}(t, a, b) & :=\int h_{k}(t, a, b) \mathrm{d} t=\sum_{i=1}^{4} \Theta_{i}(k, a, b) \int L_{i}(t, a) \mathrm{d} t \\
& =\sum_{i=1}^{4} \Theta_{i}(k, a, b) L_{i}^{*}(t, a)
\end{aligned}
$$

with the indefinite integrals

$$
\begin{aligned}
& L_{1}^{*}(t, a)=\int t \mathrm{~d} t=\frac{t^{2}}{2}, \quad L_{2}^{*}(t, a)=\int \frac{1}{t} \mathrm{~d} t=\ln t \\
& L_{3}^{*}(t, a)=\int \frac{\sqrt{t^{2}-a^{2}}}{t} \mathrm{~d} t=\sqrt{t^{2}-a^{2}}+a \arcsin \frac{a}{t}
\end{aligned}
$$

(see [6, p. 48, Eq. 217]) and using integration by parts,

$$
\begin{aligned}
L_{4}^{*}(t, a) & =\int L_{4}(t, a) \mathrm{d} t=\int t \arcsin \frac{a}{t} \mathrm{~d} t \\
& =\frac{1}{2}\left(t^{2} \arcsin \frac{a}{t}+a \int \frac{t}{\sqrt{t^{2}-a^{2}}} \mathrm{~d} t\right) \\
& =\frac{1}{2}\left(t^{2} \arcsin \frac{a}{t}+a \sqrt{t^{2}-a^{2}}\right) .
\end{aligned}
$$

For odd $n$, the function

$$
H_{K}^{*}(t)=\left\{\begin{array}{lll}
H_{K, 1}^{*}(t) & \text { if } \quad t<\lambda, \\
H_{K, 2}^{*}(t) & \text { if } \quad t \geq \lambda
\end{array}\right.
$$

with

$$
\begin{aligned}
H_{K, 1}^{*}(t):= & t-h_{K}^{*}\left(t, A_{1}(k), B_{1}(k)\right)+h_{K+1}^{*}\left(t, A_{1}(K), B_{1}(K)+\pi / n\right), \\
H_{K, 2}^{*}(t):= & t-h_{K}^{*}\left(t, A_{1}(K), B_{1}(K)\right)+h_{K+1}^{*}\left(t, A_{1}(K), B_{1}(K)+\pi / n\right) \\
& -h_{K+1}^{*}\left(t, A_{2}, B_{2}+\pi / n\right)-h_{K+1}^{*}\left(t, A_{2}, B_{2}\right)
\end{aligned}
$$

is not continuous in $t=\lambda$. This causes a false result when calculating the integral

$$
J_{K}^{*}\left(\ell_{K}, \lambda\right)=H_{K}^{*}(\lambda)-H_{K}^{*}\left(\ell_{K}\right) .
$$

In order to avoid this problem (and unnecessary case distinctions), we define

$$
\begin{aligned}
\widetilde{H}_{K, 2}^{*}(t) & =H_{K, 2}^{*}(t)-H_{K, 2}^{*}(\lambda)+H_{K, 1}^{*}(\lambda) \\
& =H_{K, 2}^{*}(t)+h_{K+1}^{*}\left(\lambda, A_{2}, B_{2}+\pi / n\right)+h_{K+1}^{*}\left(\lambda, A_{2}, B_{2}\right)
\end{aligned}
$$

and put

$$
H_{K}^{*}(t):=\left\{\begin{array}{lll}
H_{K, 1}^{*}(t) & \text { if } \quad t<\lambda, \\
\widetilde{H}_{K, 2}^{*}(t) & \text { if } \quad t \geq \lambda,
\end{array}\right.
$$

so that $H_{K}^{*}$ is now a continuous function. This completes the proof.
Corollary 1. The distribution function $G$ of the distance $t$ between two random points in $\mathcal{P}_{n, r}$ is given by

$$
G(t)= \begin{cases}0 & \text { if }-\infty<s<0, \\ \frac{1}{A}\left[t^{2}\left(\pi-\frac{2 u}{3 A} t\right)+\frac{2 u}{A} \phi^{\natural}(t)\right] & \text { if } 0 \leq t<\ell_{K+1}, \\ 1 & \text { if } t \geq \ell_{K+1}\end{cases}
$$

with

$$
\phi^{\natural}(t)=\sum_{\nu=0}^{k-1} K_{\nu}\left(\ell_{\nu+1}\right)+K_{k}(t) \quad \text { if } \quad \ell_{k} \leq t<\ell_{k+1}, k=0, \ldots, K,
$$

where

$$
K_{k}(t)=\frac{1}{2}\left(t^{2}-\ell_{k}^{2}\right)\left(\sum_{\nu=0}^{k-1} J_{\nu}^{*}\left(\ell_{\nu}, \ell_{\nu+1}\right)-H_{k}^{*}\left(\ell_{k}\right)\right)+J_{k}^{\natural}\left(\ell_{k}, t\right)
$$

with $J_{k}^{*}$ and $H_{k}^{*}$ according to Theorem 2 and

$$
J_{k}^{\natural}(s, t)=H_{k}^{\natural}(t)-H_{k}^{\natural}(s),
$$

where

$$
H_{k}^{\natural}(t)=\left\{\begin{array}{l}
\frac{t^{3}}{3}-h_{k}^{\natural}\left(t, A_{1}(k), B_{1}(k)\right)+h_{k+1}^{\natural}\left(t, A_{1}(k), B_{1}(k)+\pi / n\right) \\
\quad \text { if } \quad(n \text { is even } \wedge k \in\{0, \ldots, K-1\}) \vee(n \text { is odd } \wedge t<\lambda), \\
\frac{t^{3}}{3}-h_{K}^{\natural}\left(t, A_{1}(K), B_{1}(K)\right)+h_{K+1}^{\natural}\left(t, A_{1}(K), B_{1}(K)+\pi / n\right) \\
-h_{K+1}^{\natural}\left(t, A_{2}, B_{2}+\pi / n\right)-h_{K+1}^{\natural}\left(t, A_{2}, B_{2}\right) \\
+\frac{t^{2}}{2}\left[h_{K+1}^{*}\left(\lambda, A_{2}, B_{2}+\pi / n\right)+h_{K+1}^{*}\left(\lambda, A_{2}, B_{2}\right)\right] \\
\text { if } n \text { is odd } \wedge t \geq \lambda, \\
\frac{t^{3}}{3}-h_{K}^{\natural}\left(t, A_{1}(K), B_{1}(K)\right) \quad \text { if } n \text { is even } \wedge k=K
\end{array}\right.
$$

with $h_{K+1}^{*}$ from Theorem 2, and

$$
h_{k}^{\natural}(t, a, b)=\left\{\begin{array}{l}
0 \quad \text { if } \quad k=0 \vee(k \neq 0 \wedge t=0), \\
\sum_{i=1}^{4} \Theta_{i}(k, a, b) L_{i}^{\natural}(t, a) \quad \text { if } \quad k \neq 0 \wedge t>0,
\end{array}\right.
$$

$$
\begin{aligned}
& L_{1}^{\natural}(t, a)=\frac{t^{4}}{8}, \quad L_{2}^{\natural}(t, a)=\frac{t^{2}}{4}(2 \ln t-1), \\
& L_{3}^{\natural}(t, a)=\frac{1}{3}\left(t^{2}-a^{2}\right)^{3 / 2}+\frac{a}{2}\left(a \sqrt{t^{2}-a^{2}}+t^{2} \arcsin \frac{a}{t}\right), \\
& L_{4}^{\natural}(t, a)=\frac{1}{8}\left[\frac{5 a}{3}\left(t^{2}-a^{2}\right)^{3 / 2}+a^{3} \sqrt{t^{2}-a^{2}}+t^{4} \arcsin \frac{a}{t}\right] .
\end{aligned}
$$

Proof. For $0 \leq t<\ell_{K+1}$, one gets

$$
\begin{aligned}
G(t) & =\int_{0}^{t} g(\tau) \mathrm{d} \tau=\int_{0}^{t}\left(\frac{2 \pi \tau}{A}-\frac{2 u \tau^{2}}{A^{2}}+\frac{2 u \tau}{A^{2}} \int_{0}^{\tau} F(s) \mathrm{d} s\right) \mathrm{d} \tau \\
& =\frac{\pi t^{2}}{A}-\frac{2 u t^{3}}{3 A^{2}}+\frac{2 u}{A^{2}} \int_{0}^{t} \tau\left(\int_{0}^{\tau} F(s) \mathrm{d} s\right) \mathrm{d} \tau \\
& =\frac{\pi t^{2}}{A}-\frac{2 u t^{3}}{3 A^{2}}+\frac{2 u}{A^{2}} \int_{0}^{t} \tau \phi^{*}(\tau) \mathrm{d} \tau=\frac{1}{A}\left[t^{2}\left(\pi-\frac{2 u}{3 A} t\right)+\frac{2 u}{A} \phi^{\natural}(t)\right]
\end{aligned}
$$

with

$$
\phi^{\natural}(t):=\int_{0}^{t} s \phi^{*}(s) \mathrm{d} s
$$

It remains to calculate $\phi^{\natural}(t)$. For $\ell_{k} \leq t<\ell_{k+1}, k=0, \ldots, K$, we have

$$
\begin{aligned}
\phi^{\natural}(t) & =\int_{0}^{\ell_{k}} s \phi^{*}(s) \mathrm{d} s+\int_{\ell_{k}}^{t} s \phi^{*}(s) \mathrm{d} s \\
& =\sum_{\nu=0}^{k-1} \int_{\ell_{\nu}}^{\ell_{\nu+1}} s \phi^{*}(s) \mathrm{d} s+\int_{\ell_{k}}^{t} s \phi^{*}(s) \mathrm{d} s
\end{aligned}
$$

with

$$
\begin{aligned}
\int_{\ell_{k}}^{t} s \phi^{*}(s) \mathrm{d} s & =\int_{\ell_{k}}^{t} s\left(\sum_{\nu=0}^{k-1} J_{\nu}^{*}\left(\ell_{\nu}, \ell_{\nu+1}\right)+J_{k}^{*}\left(\ell_{k}, s\right)\right) \mathrm{d} s \\
& =\sum_{\nu=0}^{k-1} J_{\nu}^{*}\left(\ell_{\nu}, \ell_{\nu+1}\right) \int_{\ell_{k}}^{t} s \mathrm{~d} s+\int_{\ell_{k}}^{t} s J_{k}^{*}\left(\ell_{k}, s\right) \mathrm{d} s \\
& =\sum_{\nu=0}^{k-1} J_{\nu}^{*}\left(\ell_{\nu}, \ell_{\nu+1}\right) \int_{\ell_{k}}^{t} s \mathrm{~d} s+\int_{\ell_{k}}^{t} s\left[H_{k}^{*}(s)-H_{k}^{*}\left(\ell_{k}\right)\right] \mathrm{d} s \\
& =\left(\sum_{\nu=0}^{k-1} J_{\nu}^{*}\left(\ell_{\nu}, \ell_{\nu+1}\right)-H_{k}^{*}\left(\ell_{k}\right)\right) \int_{\ell_{k}}^{t} s \mathrm{~d} s+\int_{\ell_{k}}^{t} s H_{k}^{*}(s) \mathrm{d} s
\end{aligned}
$$

Putting

$$
H_{k}^{\natural}(t):=\int_{\ell_{k}}^{t} s H_{k}^{*}(s) \mathrm{d} s \quad \text { and } \quad J_{k}^{\natural}(s, t):=H_{k}^{\natural}(t)-H_{k}^{\natural}(s),
$$

it follows that

$$
\begin{aligned}
\int_{\ell_{k}}^{t} s \phi^{*}(s) \mathrm{d} s & =\frac{1}{2}\left(t^{2}-\ell_{k}^{2}\right)\left(\sum_{\nu=0}^{k-1} J_{\nu}^{*}\left(\ell_{\nu}, \ell_{\nu+1}\right)-H_{k}^{*}\left(\ell_{k}\right)\right)+H_{k}^{\natural}(t)-H_{k}^{\natural}\left(\ell_{k}\right) \\
& =\frac{1}{2}\left(t^{2}-\ell_{k}^{2}\right)\left(\sum_{\nu=0}^{k-1} J_{\nu}^{*}\left(\ell_{\nu}, \ell_{\nu+1}\right)-H_{k}^{*}\left(\ell_{k}\right)\right)+J_{k}^{\natural}\left(\ell_{k}, t\right)=: K_{k}(t)
\end{aligned}
$$

and hence

$$
\phi^{\natural}(t)=\sum_{\nu=0}^{k-1} K_{\nu}\left(\ell_{\nu+1}\right)+K_{k}(t), \quad \ell_{k} \leq t<\ell_{k+1}, \quad k=0, \ldots, K .
$$

If $k \neq 0$ and $t>0$, one finds

$$
\begin{aligned}
h_{k}^{\natural}(t, a, b) & :=\int t h_{k}^{*}(t, a, b) \mathrm{d} t=\sum_{i=1}^{4} \Theta_{i}(k, a, b) \int t L_{i}^{*}(t, a) \mathrm{d} t \\
& =\sum_{i=1}^{4} \Theta_{i}(k, a, b) L_{i}^{\natural}(t, a)
\end{aligned}
$$

with the indefinite integrals

$$
L_{1}^{\natural}(t, a)=\int \frac{t^{3}}{2} \mathrm{~d} t=\frac{t^{4}}{8}, \quad L_{2}^{\natural}(t, a)=\int t \ln t \mathrm{~d} t=\frac{t^{2}}{4}(2 \ln t-1)
$$

and

$$
\begin{aligned}
L_{4}^{\natural}(t, a) & =\int t L_{4}^{*}(t, a) \mathrm{d} t=\frac{1}{2} \int t\left(a \sqrt{t^{2}-a^{2}}+t^{2} \arcsin \frac{a}{t}\right) \mathrm{d} t \\
& =\frac{1}{2}\left(a \int t \sqrt{t^{2}-a^{2}} \mathrm{~d} t+\int t^{3} \arcsin \frac{a}{t} \mathrm{~d} t\right)
\end{aligned}
$$

with

$$
\int t \sqrt{t^{2}-a^{2}} \mathrm{~d} t=\frac{1}{3}\left(t^{2}-a^{2}\right)^{3 / 2} \quad[6, \text { p. } 47, \text { Eq. 214] }
$$

Using integration by parts, we find

$$
\int t^{3} \arcsin \frac{a}{t} \mathrm{~d} t=\frac{1}{4}\left(t^{4} \arcsin \frac{a}{t}+a \int \frac{t^{2}}{\sqrt{1-(a / t)^{2}}} \mathrm{~d} t\right)
$$

Since $t \geq a>0$ in the present cases,

$$
\int \frac{t^{2}}{\sqrt{1-(a / t)^{2}}} \mathrm{~d} t=\int \frac{t^{3}}{\sqrt{t^{2}-a^{2}}} \mathrm{~d} t
$$

and

$$
\int \frac{t^{3}}{\sqrt{t^{2}-a^{2}}} \mathrm{~d} t=\frac{1}{3}\left(t^{2}-a^{2}\right)^{3 / 2}+a^{2} \sqrt{t^{2}-a^{2}} \quad[6, \text { p. 48, Eq. 223]. }
$$

This yields

$$
L_{4}^{\natural}(t, a)=\frac{1}{8}\left[\frac{5 a}{3}\left(t^{2}-a^{2}\right)^{3 / 2}+a^{3} \sqrt{t^{2}-a^{2}}+t^{4} \arcsin \frac{a}{t}\right] .
$$

Furthermore,

$$
\begin{aligned}
L_{3}^{\natural}(t, a) & =\int t \sqrt{t^{2}-a^{2}} \mathrm{~d} t+a \int t \arcsin \frac{a}{t} \mathrm{~d} t \\
& =\frac{1}{3}\left(t^{2}-a^{2}\right)^{3 / 2}+\frac{a}{2}\left(a \sqrt{t^{2}-a^{2}}+t^{2} \arcsin \frac{a}{t}\right)
\end{aligned}
$$

(see the calculations of $L_{4}^{\natural}(t, a)$ and $L_{4}^{*}(t, a)$ ).

## 5. Examples

Fig. 3 shows examples of chord length distribution functions $F$.
As special case of Theorem 2, the distance density function for an equilateral triangle $\mathcal{P}_{3, r}$ with circumscribed circle of radius $r$ is given by

$$
g(t)=\left\{\begin{array}{lll}
\frac{2 t}{A}\left[\pi+\frac{u}{A}(\phi(t)-t)\right] & \text { if } & t \in[0, \sqrt{3} r), \\
0 & \text { if } & t \in \mathbb{R} \backslash[0, \sqrt{3} r)
\end{array}\right.
$$

with $u=3 \sqrt{3} r, A=\frac{3}{4} \sqrt{3} r^{2}$ and

$$
\phi(t)= \begin{cases}\frac{(3 \sqrt{3}+2 \pi) t^{2}}{36 r} & \text { if } 0 \leq t<\frac{3 r}{2}, \\ \frac{3}{2}\left[t \sqrt{1-\left(\frac{3 r}{2 t}\right)^{2}}-\frac{\pi r}{2}\right]+\left(\frac{1}{4 \sqrt{3}}-\frac{\pi}{9}\right) \frac{t^{2}}{r} & \\ +\left(\frac{3 r}{2}+\frac{t^{2}}{3 r}\right) \arcsin \frac{3 r}{2 t} & \text { if } \frac{3 r}{2} \leq t<\sqrt{3} r .\end{cases}
$$

Fig. 4 shows the function $r \times g(t)$ for $\mathcal{P}_{3, r}$ and some other examples.
For the expectation of the distance for $\mathcal{P}_{3, r}$ one finds

$$
\begin{aligned}
\mathrm{E}[t]= & \int_{0}^{\sqrt{3} r} t g(t) \mathrm{d} t=\left(\int_{0}^{3 r / 2}+\int_{3 r / 2}^{\sqrt{3} r}\right) t g(t) \mathrm{d} t \\
= & \frac{r}{20}(27-90 \sqrt{3}+26 \sqrt{3} \pi)+\frac{r}{20}(-27+94 \sqrt{3}-26 \sqrt{3} \pi \\
& +\sqrt{3} \ln 27)=\frac{\sqrt{3} r}{20}(4+3 \ln 3) .
\end{aligned}
$$

Since the side length $a$ of $\mathcal{P}_{3, r}$ is equal to $\sqrt{3} r$, we get

$$
\mathrm{E}[t]=\frac{a}{20}(4+3 \ln 3)=\frac{3 a}{5}\left(\frac{1}{3}+\frac{1}{4} \ln 3\right) .
$$

This is the result from [19, p. 49].



Go back
Full Screen

Close
For the square $\mathcal{P}_{4, r}$, the side length is equal to $\sqrt{2} r$. For $a=1$, a calculation using Mathematica yields the expected distance $\mathrm{E}[t]=0.5214054331 \ldots$. This result can be found in $[9, \mathrm{p} .479]$. It

Figure 3. $F$ for $\mathcal{P}_{7,1}$ (thick), $\mathcal{P}_{8,1}$ (dashed) and circle of radius $r=1$ (thin). also follows from [3, p. 13] and [7].


Figure 4. $\quad r \times g(t)$ for $\mathcal{P}_{3, r}$ (thick and dashed), $\mathcal{P}_{4, r}$ (thin), $\mathcal{P}_{10, r}$ (thick) and circle of radius $r$ (thin and dashed).

Acknowledgement. I wish to thank Lothar Heinrich (University of Augsburg) for bringing Piefke's paper to my attention and the anonymous referee for careful reading and valuable advice.

Go back

Full Screen

Close

1. Ambartzumjan R. V., Mecke J. and Stoyan D., Geometrische Wahrscheinlichkeiten und stochastische Geometrie, Akademie Verlag, Berlin, 1993.
2. Bäsel U. and Duma A., Verteilungsfunktionen der Sehnenlänge eines nichtkonvexen Polygons, Fernuniversität Hagen: Seminarberichte aus der Fakultät für Mathematik und Informatik, 84 (2011), 141-151.
3. Bailey D. H., Borwein J. M. and Crandall R. E., Box Integrals, Lawrence Berkeley National Laboratory (2006) Permalink: http://escholarship.org/uc/item/8xr051pf
4. Borel E., Principes et Formules Classiques du Calcul des Probabilités, Gauthier-Villars, Paris, 1925.
5. Borwein J. M. and Crandall R. E., Closed Forms: What They Are and Why We Care, Notices of the AMS, Vol. 60, No. 1, January 2013.
6. Bronstein I. N. and Semendjajew K. A., Taschenbuch der Mathematik, Verlag Nauka, Moskau und BSB B. G. Teubner Verlagsgesellschaft, Leipzig, 1989.
7. Burgstaller B. and Pillichshammer F., The average distance between two points, Bull. Austral. Math. Soc., 80 (2009), 353-359
8. Czuber E., Wahrscheinlichkeitsrechnung und ihre Anwendung auf Fehlerausgleichung, Statistik und Lebensversicherung, 1. Band, B. G. Teubner, Leipzig und Berlin, 1914.
9. Finch S., Mathematical Constants, Cambridge University Press, Cambridge, 2003.
10. Gates J., Some properties of chord length distributions, J. Appl. Prob., 24 (1987), 863-874.
11. Ghosh B., Random distances within a rectangle and between two rectangles, Bull. Calcutta Math. Soc., 43 (1951), 17-24.
12. Gille W., The chord length distribution of parallelepipeds with their limiting cases, Exp. Techn. Phys., 36 (1988), 197-208.
13. Harutyunyan H. S. and Ohanyan V. K., The chord length distribution function for regular polygons, Adv. Appl. Prob. (SGSA), 41 (2009), 358-366.
14. Heinrich L., Lower Bounds of Second-Order Chord Power Integrals of Convex Discs, Rend. Circ. Mat. Palermo, Serie II, Suppl., 81 (2009), 213-222.
15. Marsaglia G., Narasimhan B. G. and Zaman A., The distance between random points in rectangles, Commun. Statist. - Theory Meth., 19(11) (1990), 4199-4212.
16. Mathai A. M., An Introduction to Geometrical Probability, Gordon and Breach, Australia, 1999.


Go back

Full Screen

Close

Quit
17. Mathai A. M., Moschopoulos P. and Pederzoli G., Random points associated with rectangles, Rend. Circ. Mat. Palermo, Serie II, Tomo XLVIII (1999), 163-190.
18. Piefke F., Beziehungen zwischen der Sehnenlängenverteilung und der Verteilung des Abstandes zweier zufälliger Punkte im Eikörper, Z. Wahrscheinlichkeitstheorie verw. Gebiete, 43 (1978), 129-134.
19. Santaló L. A., Integral Geometry and Geometric Probability, Addison-Wesley, London, 1976.
20. Sulanke R., Die Verteilung der Sehnenlängen an ebenen und räumlichen Figuren, Math. Nachr., 23 (1961), 51-74.
U. Bäsel, HTWK Leipzig, Faculty of Mechanical and Energy Engineering, PF 3011 66, 04251 Leipzig, Germany, $e$-mail: uwe.baesel@htwk-leipzig.de


[^0]:    Received July 9, 2012; revised September 27, 2013.
    2010 Mathematics Subject Classification. Primary 60D05, 52A22.
    Key words and phrases. Geometric probability; random sets; integral geometry; chord length distribution function; random distances; distance distribution function; regular polygons; Piefke formula.

