THE OPERATIONAL TAU-ADOMIAN METHOD AND PADÉ APPROXIMANTS FOR SOLVING GENERALIZED NON-LINEAR VOLTERRA INTEGRO-DIFFERENTIAL EQUATIONS

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ABSTRACT. In this paper we use the operational Tau and Adomian methods with Padé approximants for solving the general Non-linear Volterra Integro-Differential Equations (NVIDE). We will present our method based on the matrix form of (NVIDE). The corresponding unknown coefficients of our method have been determined by using computational aspects of matrices. To this end, the Padé approximants have been used to accurately determine the numerical solution. Finally, accuracy of the method has been verified by presenting some numerical computations.

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1. INTRODUCTION

In 1981, Ortiz and Samara [22] proposed an operational technique for finding a numerical solution of non-linear ordinary differential equations with some supplementary conditions based on the Tau method [9]. During the last last years considerable works have been done both in development of the above mentioned technique, its theoretical analysis and numerical applications. Various techniques have been described in a series of papers [11,12,14,15,21] for the case of linear ordinary differential eigenvalue problems. In [13,16,18-20] numerical solution of partial differential equations and their related eigenvalue problems have discussed. The object of this paper is to present a simpler operational approach by using the Adomian decomposition method for the general form of non-linear Volterra integro-differential equations of the second kind with initial conditions. This method leads to an algorithm with remarkable simplicity, while retaining the accuracy of results. The obtained series solution is converted into Padé approximants to study the behavior of the solution.

2. Non-linear Volterra integro-differential equations

Consider the non-linear Volterra integro-differential equation

$$y^{(m)}(x) + G(x, y, y', \dots, y^{(m-1)}) - \int_0^x F(x, t, y, y', \dots, y^{(m)}) dt = f(x), \ x \in [0, a]$$
(1)

whit the initial conditions

$$y^{(j)}(0) = d_j, \quad j = 0, 1, \cdots, m-1.$$
 (2)

Here we assume that f(x) is polynomial, otherwise it can be approximated by polynomial to any degree of accuracy (by Taylor series or any other suitable method). Moreover, we suppose that $y_n(x)$ to be a polynomial approximation of degree n for y(x). Then one can write

$$f(x) = \sum_{j=0}^{n} f_j x^j = \underline{\mathbf{f}} \, \underline{\mathbf{X}}$$
$$y_n(x) = \sum_{j=0}^{n} a_j x^j = \underline{\mathbf{a}}_n \underline{\mathbf{X}}$$
(3)

where $\mathbf{f} = [f_0, f_1, \ldots, f_n, 0, 0, \ldots], \mathbf{a}_n = [a_0, a_1, \ldots, a_n, 0, 0, \ldots]$ and $\mathbf{X} = [1, x, x^2, \ldots]^T$ are the coefficients vectors of right-hand side of equation (1), unknown coefficients vector and the basis vector respectively. Without loss of generality we have taken all polynomials of degree n, because if f(x) and $y_n(x)$ are respectively of different degrees n_f and n_y then we can set $n = \max\{n_f, n_y\}$.

3. Converting NVE to a system of algebraic equations

The effect of differentiation on the coefficients $\underline{\mathbf{a}}_n = [a_0, \cdots, a_n, 0, \cdots]$ of a polynomial $y_n(x) = \underline{\mathbf{a}}_n \underline{\mathbf{X}}$ is the same as that of post-multiplication of $\underline{\mathbf{a}}_n$ by the matrix η by

$$\eta = \begin{bmatrix} 0 & & & \\ 1 & 0 & & & \\ 0 & 2 & 0 & \vdots \\ 0 & 0 & 3 & 0 & \\ & & \dots & \ddots \end{bmatrix}.$$

Lemma 1. If $y_n(x)$ be a polynomial of the form

$$y_n(x) = \sum_{i=0}^n a_i x^i = \sum_{i=0}^\infty a_i x^i.$$

Then

$$\frac{d^r}{dx^r}y_n(x) = \underline{\mathbf{a}}_n \eta^r \underline{\underline{\mathbf{X}}} , \quad r = 1, 2, 3, \cdots$$

where $\underline{\mathbf{a}}_n = [a_0, a_1, \dots, a_n, 0, 0, \dots].$

The proof follows immediately by induction. \Box

By using the Adomian decomposition method (See [13]), one can simplify the non-linear terms of (1) as follows.

By setting $\hat{G}(x) = G(x, y(x), y'(x), \dots, y^{(m-1)}(x))$, we have

$$\hat{G}(x) = G(x, \sum_{j=0}^{\infty} a_j x^j, \sum_{j=0}^{\infty} (j+1)a_{j+1}x^j, \dots, \sum_{j=0}^{\infty} \frac{(j+m-1)!}{j!}a_{j+m-1}x^j)$$

$$= \sum_{i=0}^{\infty} A_i^G x^i$$

$$= \underline{\mathbf{A}}^G \underline{\mathbf{X}}$$
(4)

where $\underline{\mathbf{A}}^{G} = [A_{0}^{G}, A_{1}^{G}, \ldots]$ with

$$A_i^G = \frac{1}{i!} \{ \frac{d^i}{dx^i} G(x, \sum_{j=0}^\infty a_j x^j, \dots, \sum_{j=0}^\infty \frac{(j+m-1)!}{j!} a_{j+m-1} x^j) \}_{x=0} = \frac{\hat{G}^{(i)}(0)}{i!}$$

which is depend on $a_0, a_1, \ldots, a_{i+m-1}$ for $i = 0, 1, \ldots$

By using this method for the non-linear term under integral sign, we have

$$\hat{F}(x,t) = F(x,t,y,y',...,y^{(m)}) = \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} A_{ij}^{F} x^{i} t^{j}$$
(5)

where

$$A_{ij}^{F} = \frac{1}{i!j!} \{ \frac{\partial^{i+j}}{\partial x^{i} \partial t^{j}} F(x,t,y,y',\dots,y^{(m)}) \}_{(x,y)=(0,0)}$$
$$= \frac{\frac{\partial^{i+j}\hat{F}}{\partial x^{i} \partial t^{j}}(0,0)}{i!j!}$$

which is depend on $a_0, a_1, \ldots, a_{r+m}$ for $i, j = 0, 1, 2, \ldots$, and $r = \max\{i, j\}$. Now if we replace y(x) by $y_n(x) = \sum_{i=0}^n a_i x^i$ in (5), then we have

$$\int_{0}^{x} F(x,t,y,y',\ldots,y^{(m)})dt = \int_{0}^{x} \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} A_{ij}^{F} x^{i} t^{j} dt$$
$$= \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} A_{ij}^{F} \frac{x^{i+j+1}}{j+1}$$
$$= \sum_{i=1}^{\infty} \hat{A}_{i}^{F} x^{i}$$
$$= \underline{\hat{A}}^{F} \underline{\underline{X}},$$
(6)

where $\hat{\mathbf{A}}^{F} = [\hat{A}_{0}^{F}, \hat{A}_{1}^{F}, \ldots]$ with $\hat{A}_{0}^{F} = 0$ and $\hat{A}_{i}^{F} = \sum_{j=1}^{i} \frac{A_{i-j,j-1}^{F}}{j}$ for $i = 1, 2, \cdots$.

Therefore, the matrix form of (1) can be written as:

$$\underline{\mathbf{a}}_n \eta^m \underline{\underline{\mathbf{X}}} + \underline{\mathbf{A}}^G \underline{\underline{\mathbf{X}}} - \underline{\hat{\mathbf{A}}}^F \underline{\underline{\mathbf{X}}} = \underline{\mathbf{f}} \underline{\underline{\mathbf{X}}}$$

which yields

$$\underline{\mathbf{a}}_{n}\eta^{m} + \underline{\mathbf{A}}^{G} - \underline{\hat{\mathbf{A}}}^{F} = \underline{\mathbf{f}},\tag{7}$$

since $\underline{\mathbf{X}}$ is a base vector. Now, the unknown coefficients can be determined by (2) and $(\overline{7})$. Note that we use (2) to write

$$a_j = \frac{d_j}{j!}$$
 $j = 0, 1, \cdots, m - 1,$

which determines the unknowns for the index $j = 0, \ldots, m - 1$.

Other coefficients are determined by solving (7), as:

$$a_{m+j} = \frac{f_j + \hat{A}_j^F - A_j^G}{\frac{(m+j)!}{j!}}, \quad j = 0, 1, \cdots, n-m.$$

The obtained numerical solution by above mentioned procedure can be improved by the use of the Padé approximants.

4. PADÉ APPROXIMANTS

A Padé approximant is the ratio of two polynomials constructed from the coefficients of the Taylor series expansion of a function u(x). The [L/M] Padé approximants to a function y(x) are given by [16].

$$\left[\frac{L}{M}\right] = \frac{P_L(x)}{Q_M(x)},\tag{8}$$

where $P_L(x)$ is polynomial of degree at most L and $Q_M(x)$ is a polynomial of degree at most M. The formal power series

$$y(x) = \sum_{i=1}^{\infty} a_i x^i,\tag{9}$$

$$y(x) - \frac{P_L(x)}{Q_M(x)} = O(x^{L+M+1}), \tag{10}$$

determine the coefficients of $P_L(x)$ and $Q_M(x)$ by the equation.

Since we can clearly multiply the numerator and denominator by a constant and leave [L/M] unchanged, we imposed the normalization condition

$$Q_M(x) = 1.0.$$
 (11)

Finally, we require that $P_L(x)$ and $Q_M(x)$ have non common factors. If we write the coefficient of $P_L(x)$ and $Q_M(x)$ as

$$P_{L}(x) = p_{0} + p_{1}x + \dots + p_{L}x^{L},$$

$$Q_{M}(x) = q_{0} + q_{1}x + \dots + q_{M}x^{M},$$
(12)

Then by (11) and (12), we may multiply (8) by $Q_M(x)$, which linearizes the coefficient equations. We can write out (10) in more details as

$$\begin{cases} a_{L+1} + a_L q_1 + \dots + a_{L-M} q_M = 0, \\ a_{L+2} + a_{L+1} q_1 + \dots + a_{L-M+2} q_M = 0, \\ \vdots \\ a_{L+M} + a_{L+M-1} q_1 + \dots + a_L q_M = 0, \end{cases}$$
(13)
$$\begin{cases} a_0 = p_0, \\ a_1 + a_0 q_1 = p_1, \\ \vdots \\ a_L + a_{L-1} q_1 + \dots + a_0 q_L = p_L, \end{cases}$$
(14)

To solve these equations, we start with Eq.(13), which is a set of linear equations for all the unknown q's. Once the q's are known, then Eq.(14) gives an explicit formula for the unknown p's, which complete the solution. If Eqs.(13) and (14) are nonsingular, then we can solve them directly and obtain Eq.(14), where Eq.(14)holds, and if the lower index on a sum exceeds the upper, the sum is replaced by zero:

$$\left[\frac{L}{M}\right] = \frac{det \begin{bmatrix} a_{L-M+1} & a_{L-M+2} & \cdots & a_{L+1} \\ \vdots & \vdots & \ddots & \vdots \\ a_{L} & a_{L+1} & \cdots & a_{L+M} \\ \sum_{j=M}^{L} a_{j-M} x^{j} & \sum_{j=M-1}^{L} a_{j-M+1} x^{j} & \cdots & \sum_{j=0}^{L} a_{j} x^{j} \end{bmatrix}}{det \begin{bmatrix} a_{L-M+1} & a_{L-M+2} & \cdots & a_{L+1} \\ \vdots & \vdots & \ddots & \vdots \\ a_{L} & a_{L+1} & \cdots & a_{L+M} \\ x^{M} & x^{M-1} & \cdots & 1 \end{bmatrix}}$$
(15)

5. Estimation of error function

In this section, an error function is obtained for the approximate solution of Eqs.(1) and (2). Let $e_n(x) = y(x) - y_n(x)$ be called the error function of Tau approximation $y_n(x)$ to y(x) where y(x) is the exact solution. Hence $y_n(x)$ satisfies the following problem:

$$y_n^{(m)}(x) + G(x, y_n, y'_n, \dots, y_n^{(m-1)}) - \int_0^x F(x, t, y_n, y'_n, \dots, y_n^{(m)}) dt$$

= $f(x) + H_n(x), \quad x \in [0, a].$ (16)

with

$$y_n^{(j)}(0) = d_j, \quad j = 0, 1, \cdots, m-1.$$
 (17)

The function $H_n(x)$ is the perturbation term associated with $y_n(x)$. Hence

$$H_n(x) = y_n^m(x) + G(x, y_n, y'_n, \dots, y_n^{(m-1)}) - \int_0^x F(x, t, y_n, y'_n, \dots, y_n^{(m)}) dt - f(x)$$

we proceed to find an approximation $e_{n,N}(x)$ to the error function $e_n(x)$ in the same way as we did before for the solution of problems in Eqs.(1) and (2). Subtracting (16) and (17) from (1) and (2) respectively and taking a term of expansions G'(x) and F'(x,t) around $y_n(t)$, the error function $e_n(x)$ satisfies the problem

$$e_{n}^{(m)}(x) + \sum_{i=0}^{m-1} e_{n}^{(i)} \frac{\partial G}{\partial y^{(i)}}(x, y_{n}, y'_{n}, \dots, y_{n}^{(m-1)}) - \int_{0}^{x} \sum_{i=0}^{m} \frac{\partial F}{\partial y^{(i)}} F(x, t, y_{n}, y'_{n}, \dots, y_{n}^{(m)}) dt = -H_{n}(x), \quad x \in [0, a]$$

with

$$e_n^{(j)}(0) = 0, \quad j = 0, 1, \cdots, m - 1.$$

6. NUMERICAL EXAMPLES

6.1. Population problem

The study of Volterra integral equations originated with the work of Volterra on population dynamics. The equation actually used by Volterra, has the form

$$\frac{dN(t)}{dt} = N(t)\{\alpha - \beta N(t) - \int_0^t k(t-s)N(s)d(s)\},\$$

where the term $-\beta N^2(t)$ introduced to account for the competition between individuals in the population and tends to inhibit the growth of the population [10], and k(t-s) is the survival function [1].

6.2. Problem of polymer rheology The equation

$$\mu u'(t) = u^{3}(t)g(t) + \int_{0}^{t} k(t-s) \{\frac{u^{3}(t)}{u^{2}(s)} - u(s)\} ds$$

models the elongation of filament of a certain polyethylene which is stretched on the time interval $-\infty < t \leq 0$, then released and allowed to undergo elastic recovery for t > 0[10].

Both problems can be solved by the method of this paper directly.

7. Numerical results

The following examples are given to clarify accuracy of the presented method and also shows the importance of the after-treatment method used to improve the accuracy of the approximate solution in Tau-Padé method. Example[3,4] are selected

from different references, so their numerical results obtained here can be compared with the other numerical methods. The computations associated with the examples were performed using Maple 13 on a Personal computer.

Example 1.

$$\begin{cases} y^{''}(x) - Ln(y(x)y'(x)) + \int_0^x \left\{ \frac{2y'(t)}{y(x)} \left(\sin(t)e^{-t}y(t) + \cos(t) \right) - y(t) \right\} dt \\ = 2sin(x) - 2x + 1, \\ y(0) = 1, \quad y'(0) = 1, \qquad 0 \le x \le 1. \end{cases}$$

The exact solution is given by $y(x) = e^x$. For the numerical results with n = 10, 15 see Table 1.

Example 2.

$$\begin{cases} y^{(4)}(x) + (1+x)e^{Ln(y''(x))} + \int_0^x 2y'(t)Ln(y''(t))dt = e^{-x} - x^2 + 1, & 0 \le x \le 1\\ y(0) = 1, & y'(0) = 0, & y''(0) = 1, & y'''(0) = -1. \end{cases}$$

The exact solution is given by $y(x) = x + e^{-x}$. For the numerical results with n = 10, 15 see Table 2.

Example 3. [3,4,17]

$$\begin{cases} y'(x) + \int_0^x \cos(x-t)y^2(t)dt = \sin(2x), \\ y(0) = 1, \quad 0 \le x \le 1. \end{cases}$$

The exact solution is given by $y(x) = \cos(x)$. Table 3 shows the numerical results and comparison with the other numerical solutions.

Example 4. [2,6,8]

$$\begin{cases} y'(x) = 1 + \int_0^x y(t)y'(t)dt, \\ y(0) = 0, \quad 0 \le x \le 1. \end{cases}$$

The exact solution is given by $y(x) = \sqrt{2} \tan(\frac{x}{\sqrt{2}})$. For the numerical results with n = 8 see Table 4.

Remark 1. Note that in the following tables, the notations Exact, App, Err. Tau, Est. Err. and Err. Tau-Pade[m, n], have been used for exact solution, approximate solution obtained by our method and absolute error and estimate error of approximate solution and absolute error of approximate solution improved by use of the Padé approximants respectively.

x	Exact	App	Err.Tau	Est.Err.	Err.Tau - Pade[5, 5]	
n = 10						
0.0	1.000000	1.000000	0	0	0	
0.2	1.221403	1.221403	5.21752e - 16	5.13067e - 16	2.48910e - 18	
0.4	1.491825	1.491825	1.08869e - 12	1.05076e - 12	6.24388e - 15	
0.6	1.822119	1.822119	9.56518e - 11	9.08333e - 11	6.62761e - 13	
0.8	2.225541	2.225541	2.30479e - 09	2.15196e - 09	1.92928e - 11	
1.0	2.718282	2.718282	2.73127e - 08	2.50521e - 08	2.76650e - 10	
x	Exact	App	Err.Tau	Est. Err.	Err.Tau - Pade[8,7]	
n = 15						
0.0	1.000000	1.000000	0	0	0	
0.2	1.221403	1.221403	3.16954e - 25	3.13228e - 25	5.87483e - 29	
0.4	1.491825	1.491825	2.10217e - 20	2.05277e - 20	4.65275e - 24	
0.6	1.822119	1.822119	1.39757e - 17	1.34834e - 17	3.69789e - 21	
0.8	2.225541	2.225541	1.41155e - 15	1.34530e - 15	4.47002e - 19	
1.0	2.718282	2.718282	5.07711e - 14	4.77948e - 14	1.92645e - 17	
Table 1.	Table 1. Numerical results of Example 1.					

x	Exact	App	Err.Tau	Est.Err.	Err.Tau - Pade[5, 5]
n = 10					
0.0	1.000000	1.000000	0	0	0
0.2	1.018731	1.018731	5.04646e - 16	5.13067e - 16	3.03697e - 18
0.4	1.070320	1.070320	1.01678e - 12	1.05076e - 12	5.31185e - 15
0.6	1.148812	1.148812	8.65450e - 11	9.08883e - 11	3.93157e - 13
0.8	1.249329	1.249329	2.01685e - 09	2.15196e - 09	7.98071e - 12
1.0	1.367879	1.367879	2.31143e - 08	2.50521e - 08	7.98044e - 11
x	Exact	App	Err.Tau	Est.Err.	Err.Tau - Pade[8,7]
n = 15					
0.0	1.000000	1.000000	0	0	0
0.2	1.018731	1.018731	3.09583e - 25	3.13228e - 25	4.03810e - 29
0.4	1.070320	1.070320	2.00552e - 20	2.05277e - 20	2.19830e - 24
0.6	1.148812	1.148812	1.30229e - 17	1.34834e - 17	1.20086e - 21
0.8	1.249329	1.249329	1.28469e - 15	1.34530e - 15	9.97743e - 20
1.0	1.367879	1.367879	4.51317e - 14	4.77948e - 14	2.95549e - 18

Table 2. Numerical results of Example 2.

x	Exact	Tau(n=15)	Tau - Pade[8,7]	BPF	Adomian
0.0	1.000000	1.000000	1.000000	1.000000	1.000000
0.1	0.995004	0.995004	0.995004	0.995141	0.994951
0.2	0.980067	0.980067	0.980067	0.975784	0.980303
0.3	0.955336	0.955336	0.955336	0.960386	0.955685
0.4	0.921061	0.921061	0.921061	0.918443	0.921165
0.5	0.877583	0.877583	0.877583	0.862193	0.877048
0.6	0.825336	0.825336	0.825336	0.828963	0.822596
0.7	0.764842	0.764842	0.764842	0.752929	0.755333
0.8	0.696707	0.696707	0.696707	0.710418	0.667739
0.9	0.621610	0.621610	0.621610	0.617232	0.547241
1.0	0.540302	0.540302	0.540302	0.566917	0.364798

Table 3. Numerical results of Example 3.

x	Exact	Tau(n=8)	Err.Tau	Err.T P.[4, 4]	Err.WGM	Err.HPM
0	0	0	0	0	0	0
0.0625	0.062541	0.062541	1.99e - 14	9.18e - 18	5.93e - 05	0
0.1250	0.125327	0.125327	1.02e - 11	4.73e - 15	2.65e - 05	0
0.1875	0.188606	0.188606	3.94e - 10	1.83e - 13	6.41e - 06	0
0.2500	0.252637	0.252637	5.28e - 09	2.47e - 12	6.29e - 05	0
0.3125	0.317688	0.317688	3.96e - 08	1.87e - 11	1.24e - 05	0
0.3750	0.384043	0.384043	2.06e - 07	9.84e - 11	5.65e - 05	1.00e - 05
0.4375	0.452013	0.452012	8.35e - 07	4.03e - 10	1.25e - 05	2.00e - 05
0.5000	0.521931	0.521928	2.81e - 06	1.38e - 09	6.59e - 05	5.00e - 05
0.5625	0.594169	0.594160	8.23e - 06	4.10e - 09	3.14e - 05	1.20e - 04
0.6250	0.669142	0.669120	2.16e - 05	1.10e - 08	5.81e - 05	2.40e - 04
0.6875	0.747320	0.747268	5.19e - 05	2.69e - 08	1.98e - 05	4.70e - 04
0.7500	0.829239	0.829123	1.16e - 04	6.15e - 08	6.10e - 05	8.60e - 04
0.8125	0.915520	0.915276	2.43e - 04	1.33e - 07	8.03e - 05	1.51e - 03
0.8750	1.006886	1.006400	4.87e - 04	2.73e - 07	1.38e - 05	2.56e - 03
0.9375	1.104193	1.103263	9.30e - 04	5.39e - 07	6.82e - 06	4.17e - 03
1.0000	1.208460	1.206746	1.71e - 03	1.03e - 06	3.98e - 05	6.61e - 03

Table 4. Numerical results of Example 4.

8. Conclusions

In this paper, we have solved a special class of NVIDEs which is important in practical problems, see the given practical problems in section 6. For solving this type of problems, we have designed remarkably simple method which has high accuracy in comparison with other existing methods and clarified the accuracy through solving numerical examples (see Tables[3,4]). Some of the advantages of this method are as follows: 1. It solves NVIDEs without linearizing the nonlinear terms;

2. It gives an error estimator as a polynomial and gives more accurate solution by increasing n.

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