

A NOTE ON A 3-DIMENSIONAL STATIONARY SCHRÖDINGER–POISSON SYSTEM

KHALID BENMLIH

ABSTRACT. In a previous paper we have proved existence of a ground state for a stationary Schrödinger–Poisson system in the whole space \mathbb{R}^3 under appropriate assumptions on the data, namely the dopant-density n^* and the effective potential \tilde{V} . In this note we show that the same result remains true under less restrictive hypotheses.

1. INTRODUCTION

We are concerned with existence of standing waves (i.e. solutions of the form $u(t, x) = e^{i\omega t}u(x)$ with a real constant ω) for a time-dependent Schrödinger equation where the electric potential V satisfies a linear Poisson equation. This leads to solving the stationary Schrödinger–Poisson system

$$-\frac{1}{2}\Delta u + (V + \tilde{V})u + \omega u = 0 \quad \text{in } \mathbb{R}^3 \quad (1.1)$$

$$-\Delta V = |u|^2 - n^* \quad \text{in } \mathbb{R}^3 \quad (1.2)$$

where the dopant-density n^* and the effective potential \tilde{V} are given real functions. An existence result of a solution for (1.1)–(1.2) has been established by Lions [3] in the particular case where $\tilde{V}(x) = -2/|x|$ and $n^* \equiv 0$, by Nier [4] under some assumptions on the data essentially when $\|\tilde{V}\|_{L^2}$ and $\|n^*\|_{L^2}$ are small enough and also recently by the author [1] under appropriate assumptions on \tilde{V} and n^* .

In this note, we show existence of a ground state of (1.1)–(1.2) as in [1] but under less restrictive assumptions. More precisely, an adequate modification on the proof of the main result in [1, theorem 1.3] allows us to avoid the condition where $n^* \in L^1(\mathbb{R}^3)$.

Let us recall firstly the principal theorem and the several steps of its proof given in [1]: after solving explicitly the Poisson equation for any fixed $u \in H^1(\mathbb{R}^3)$, we substitute the unique solution then obtained $V = V(u)$ in the Schrödinger equation (1.1) and show existence of a ground state of

$$-\frac{1}{2}\Delta u + (V(u) + \tilde{V})u + \omega u = 0 \quad \text{in } \mathbb{R}^3. \quad (1.3)$$

2000 *Mathematics Subject Classification.* 35J50, 35Q40.

Key words and phrases. Schrödinger equation, Poisson equation, standing wave, variational method.

©2004 Texas State University - San Marcos.

Submitted July 29, 2003. Published February 24, 2004.

To this end, we show that the energy functional corresponding to (1.3) is exactly the expression

$$E(\varphi) := \frac{1}{4} \int_{\mathbb{R}^3} |\nabla \varphi|^2 dx + \frac{1}{4} \int_{\mathbb{R}^3} |\nabla V(\varphi)|^2 dx + \frac{1}{2} \int_{\mathbb{R}^3} \tilde{V} \varphi^2 dx + \frac{\omega}{2} \int_{\mathbb{R}^3} \varphi^2 dx \quad (1.4)$$

and a solution of (1.3) is obtained as a minimizer of E on $H^1(\mathbb{R}^3)$.

Before giving the assumptions imposed on \tilde{V} and n^* to solve the system (1.1)-(1.2), we recall the following concepts.

Definition 1.1. We say that g satisfies the decomposition (1.5) if:

- (i) $g \in L^1_{\text{loc}}(\mathbb{R}^3)$,
- (ii) $g \geq 0$, and
- (iii) There exists $q_0 \in [3/2, \infty]$ such that for all $\lambda > 0$ there exists $g_{1\lambda} \in L^{q_0}(\mathbb{R}^3)$, $g_{2\lambda} \in L^{q_\lambda}(\mathbb{R}^3)$ such that

$$g = g_{1\lambda} + g_{2\lambda} \quad \text{and} \quad \lim_{\lambda \rightarrow 0} \|g_{1\lambda}\|_{L^{q_0}} = 0. \quad (1.5)$$

As interesting examples of this definition we may consider $g(x) = 1/|x|^\alpha$ for some $0 < \alpha < 2$ or $g \in L^r(\mathbb{R}^3)$ for some $r > 3/2$ (taking $|g|$ if g is negative).

In what follows we will denote by $\|\cdot\|$ the norm $\|\cdot\|_{L^2}$ on $L^2(\mathbb{R}^3)$ and by $[E \leq c]$ the set $\{\varphi; E(\varphi) \leq c\}$.

Consider now the following hypotheses:

$$\tilde{V}^+ \in L^1_{\text{loc}}(\mathbb{R}^3) \quad \text{and} \quad \tilde{V}^- \text{ satisfies the decomposition (1.5)} \quad (1.6)$$

$$n^* \in L^1 \cap L^{6/5}(\mathbb{R}^3) \quad (1.7)$$

$$\inf \left\{ \int_{\mathbb{R}^3} (|\nabla \varphi|^2 + \varrho(x)\varphi^2) dx, \int_{\mathbb{R}^3} |\varphi|^2 = 1 \right\} < 0 \quad (1.8)$$

where $\varrho(x) := 2\tilde{V}(x) - \frac{1}{2\pi} \int_{\mathbb{R}^3} \frac{n^*(y)}{|x-y|} dy$.

The main result in [1] is as follows.

Theorem 1.2. *Assuming (1.6), (1.7) and (1.8) there exists $\omega_* > 0$ such that for all $0 < \omega < \omega_*$ the equation (1.3) has a nonnegative solution $u \not\equiv 0$ which minimizes the functional E given by (1.4):*

$$E(u) = \min_{\varphi \in H^1(\mathbb{R}^3)} E(\varphi).$$

The proof of this theorem is divided into the four following Lemmas.

Lemma 1.3. *Let $\omega \geq 0$ and $c \in \mathbb{R}$. If the set $[E \leq c]$ is bounded in $L^2(\mathbb{R}^3)$ then it is also bounded in $H^1(\mathbb{R}^3)$.*

Lemma 1.4. *For all $\omega > 0$ and $c \in \mathbb{R}$ the set $[E \leq c]$ is bounded in $L^2(\mathbb{R}^3)$.*

Lemma 1.5. *For any $\omega > 0$ the functional E is weakly lower semicontinuous on $H^1(\mathbb{R}^3)$ and attains its minimum on $H^1(\mathbb{R}^3)$ at $u \geq 0$.*

Lemma 1.6. *There exists $\omega_* > 0$ such that if $0 < \omega < \omega_*$ then $E(u) < E(0)$ and thus $u \not\equiv 0$.*

After analyzing the proofs of the four Lemmas above given in [1], we remark that theorem 1.2 remains true even if we replace the condition (1.7) by

$$n^* \in L^{6/5}(\mathbb{R}^3). \quad (1.9)$$

In the sequel we shall minimize the energy functional E on the space

$$H := \left\{ u \in H^1(\mathbb{R}^3) : \int_{\mathbb{R}^3} \tilde{V}^+ u^2 dx < \infty \right\}$$

which is a Hilbert space, continuously embedded in $H^1(\mathbb{R}^3)$, when endowed it with its natural scalar product and norm

$$(\varphi|\psi) := \int_{\mathbb{R}^3} \left(\nabla\varphi \cdot \nabla\psi + \varphi\psi + \tilde{V}^+\varphi\psi \right) dx, \quad \|\varphi\|_H := (\varphi|\varphi)^{1/2}.$$

Consequently Theorem 1.2 becomes

Theorem 1.7. *Assuming (1.6), (1.8) and (1.9) there exists $\omega_* > 0$ such that for all $0 < \omega < \omega_*$ the equation (1.3) has a nonnegative solution $u \not\equiv 0$ which minimizes on the space H the functional E :*

$$E(u) = \min_{\varphi \in H} E(\varphi).$$

2. PRELIMINARIES

Here we recall the three following Lemmas which will be useful in the sequel.

Lemma 2.1. *Let $n^* \in L^{6/5}(\mathbb{R}^3)$. For all $\varphi \in H^1(\mathbb{R}^3)$ the Poisson equation*

$$-\Delta V = |\varphi|^2 - n^* \quad \text{in } \mathbb{R}^3 \tag{2.1}$$

has a unique solution $V := V(\varphi) \in \mathcal{D}^{1,2}(\mathbb{R}^3)$ given by

$$V(\varphi)(x) = \frac{1}{4\pi} \int_{\mathbb{R}^3} \frac{(|\varphi|^2 - n^*)(y)}{|x - y|} dy. \tag{2.2}$$

Moreover if we denote by

$$I(\varphi) := \frac{1}{4} \int_{\mathbb{R}^3} |\nabla V(\varphi)|^2 dx,$$

then I is C^1 on $H^1(\mathbb{R}^3)$ and its derivative satisfies

$$\langle I'(\varphi), \psi \rangle = \int_{\mathbb{R}^3} V(\varphi) \varphi \psi dx \quad \forall \psi \in H^1(\mathbb{R}^3).$$

For the proof of this lemma see [1, Lemma 2.1, Lemma 2.2].

This Lemma shows in particular that the energy functional corresponding to (1.3) is exactly the expression given in (1.4), namely

$$E(\varphi) := \frac{1}{4} \int_{\mathbb{R}^3} |\nabla\varphi|^2 dx + I(\varphi) + \frac{1}{2} \int_{\mathbb{R}^3} \tilde{V}\varphi^2 dx + \frac{\omega}{2} \int_{\mathbb{R}^3} \varphi^2 dx.$$

Lemma 2.2. *Let $\theta \in L^r(\mathbb{R}^3)$ for some $r \geq 3/2$ then for all $\delta > 0$ there exists $C_\delta > 0$ such that*

$$\int_{\mathbb{R}^3} \theta(x) |\varphi(x)|^2 dx \leq \delta \|\nabla\varphi\|^2 + C_\delta \|\varphi\|^2 \quad \forall \varphi \in H^1(\mathbb{R}^3).$$

For the proof of this lemma see [1] or [2].

Remark that since \tilde{V}^- satisfies the decomposition (1.5) then for any fixed $\lambda > 0$ we have $\tilde{V}^- = \tilde{V}_{1\lambda}^- + \tilde{V}_{2\lambda}^-$ where for $i = 1, 2$, $\tilde{V}_{i\lambda}^- \in L^s(\mathbb{R}^3)$ for some $s \in [3/2, \infty]$

($s = q_0$ or $s = q_\lambda$). Hence taking $\theta := \tilde{V}_{i\lambda}^-$ the inequality of Lemma 2.2 holds for $i = 1, 2$ and consequently for all $\delta > 0$ there exists $C_\delta > 0$ so that

$$\int_{\mathbb{R}^3} \tilde{V}^-(x) |\varphi(x)|^2 dx \leq \delta \|\nabla \varphi\|^2 + C_\delta \|\varphi\|^2 \quad \forall \varphi \in H^1(\mathbb{R}^3). \quad (2.3)$$

Lemma 2.3. *Let $\psi \in L^r(\mathbb{R}^3)$ for some $r > 3/2$. If $v_n \rightharpoonup 0$ weakly in $H^1(\mathbb{R}^3)$ then*

$$\int_{\mathbb{R}^3} \psi(x) v_n^2(x) dx \rightarrow 0 \quad \text{as } n \rightarrow +\infty$$

For the proof of this lemma see [1, Lemma 2.5].

3. PROOF OF THEOREM 1.7

We will use once again the same steps as in [1]. Remark at first that the proofs given in [1] for Lemma 1.5 and Lemma 1.6 do not require the hypothesis $n^* \in L^1(\mathbb{R}^3)$ and consequently remain valid assuming (1.9) instead of (1.7).

Proof of Lemma 1.3. We show here that if the set

$$[E \leq c] := \{\varphi \in H; E(\varphi) \leq c\}$$

is bounded in $L^2(\mathbb{R}^3)$ then it is also bounded in H . Indeed since $I(\varphi)$ and ω are both nonnegative, the inequality $E(\varphi) \leq c$ gives in particular

$$\frac{1}{4} \|\nabla \varphi\|^2 + \frac{1}{2} \int \tilde{V}^+ \varphi^2 dx - \frac{1}{2} \int \tilde{V}^- \varphi^2 dx \leq c.$$

Now using the estimate (2.3) with $\delta = 1/4$ we get

$$\frac{1}{8} \|\nabla \varphi\|^2 + \frac{1}{2} \int \tilde{V}^+ \varphi^2 dx \leq K_0 \|\varphi\|^2 + c.$$

for some constant $K_0 > 0$. □

Let us recall that in [1] we have decomposed the expression of $E(\varphi)$ as

$$E(\varphi) = E_1(\varphi) - E_2(\varphi) + E_3(\varphi) + E(0)$$

where

$$E_1(\varphi) := \frac{1}{4} \int |\nabla \varphi|^2 dx + \frac{1}{2} \int \tilde{V}^+ \varphi^2 dx + \frac{\omega}{2} \int \varphi^2 dx$$

$$E_2(\varphi) := \frac{1}{2} \int \tilde{V}^- \varphi^2 dx + \frac{1}{8\pi} \iint \frac{n^*(y)}{|x-y|} \varphi^2(x) dx dy$$

$$E_3(\varphi) := \frac{1}{16\pi} \iint \frac{\varphi^2(x) \varphi^2(y)}{|x-y|} dx dy$$

$$E(0) := \frac{1}{16\pi} \iint \frac{n^*(x) n^*(y)}{|x-y|} dx dy.$$

Indeed, for the term $I(\varphi)$ it suffices to multiply the equation (2.1) by $V(\varphi)$, integrate by parts and use the formula (2.2).

In the proof of the similar lemma [1, Lemma 3.1] we have estimated $E_2(\varphi)$ instead of $\int \tilde{V}^- \varphi^2 dx$. More precisely we have estimated the second term of $E_2(\varphi)$ by using a certain inequality of type *Hardy* and the fact that $n^* \in L^1(\mathbb{R}^3)$.

We point out finally that the decomposition of $E(\varphi)$ as above remains useful for the rest of proofs.

Proof of Lemma 1.4. Assume by contradiction that there exists a sequence $(u_j)_j \subset H$ such that $E(u_j) \leq c$ and $\|u_j\| \rightarrow +\infty$ as $j \rightarrow +\infty$. Let $v_j := u_j/\|u_j\|$ then $\|v_j\| = 1$ and from $E(u_j) \leq c$ we get

$$\frac{1}{4} \int |\nabla v_j|^2 dx - E_2(v_j) + E_3(v_j)\|u_j\|^2 + \frac{\omega}{2} \leq \frac{c_0}{\|u_j\|^2} \quad (3.1)$$

where $c_0 = c - E(0)$. To estimate $E_2(v_j)$ it suffices to use (2.3) for the first term $\int \tilde{V}^- v_j^2 dx$. As to the second, unlike the proof in [1] we do not require here the assumption $n^* \in L^1(\mathbb{R}^3)$. Indeed, setting

$$V^*(x) := \iint_{\mathbb{R}^3 \times \mathbb{R}^3} \frac{n^*(y)}{|x-y|} dy = -V(0)(x) \quad (3.2)$$

as denoted in Lemma 2.1 we may write

$$\iint_{\mathbb{R}^3 \times \mathbb{R}^3} \frac{n^*(y)}{|x-y|} v_j^2(x) dx dy = \int_{\mathbb{R}^3} V^*(x) v_j^2(x) dx.$$

Knowing that $V(0) \in L^6(\mathbb{R}^3)$ we can use once more Lemma 2.2 with $\theta := V^*$.

On the whole, we obtain in particular

$$E_2(v_j) \leq \frac{1}{8} \|\nabla v_j\|^2 + K_0$$

for some positive constant K_0 and consequently we infer from the inequality (3.1) that

$$\frac{1}{8} \|\nabla v_j\|^2 + E_3(v_j)\|u_j\|^2 + \frac{\omega}{2} \leq \frac{c_0}{\|u_j\|^2} + K_0.$$

For the remainder of the proof, we conclude exactly as in of [1, Lemma 3.2]. Precisely we show first that, up to a subsequence, $v_j \rightharpoonup 0$ weakly in $H^1(\mathbb{R}^3)$. Next, from (3.1) it follows in particular that

$$\frac{\omega}{2} - E_2(v_j) \leq \frac{c_0}{\|u_j\|^2}. \quad (3.3)$$

Using the decomposition $\tilde{V}^- = \tilde{V}_{1\lambda}^- + \tilde{V}_{2\lambda}^-$ and (3.2), we show according to Lemma 2.3 that $E_2(v_j) \rightarrow 0$ as $j \rightarrow \infty$. Finally, letting j go to infinity in (3.3) we obtain a contradiction since ω is positive. Consequently, all $(u_j)_j \subset H$ such that $E(u_j) \leq c$ is bounded in $L^2(\mathbb{R}^3)$. \square

REFERENCES

- [1] Kh. Benmlih, Stationary Solutions for a Schrödinger–Poisson System in \mathbb{R}^3 ; Electron. J. Diff. Eqns., Conf. 09 (2002), pp. 65-76. (<http://ejde.math.swt.edu>)
- [2] H. Brezis & T. Kato: Remarks on the Schrödinger operator with singular complex potentials; J. Math. Pures & Appl. N 2, 58 (1979), 137-151.
- [3] P. L. Lions: Some remarks on Hartree equation; Nonlinear Anal., Theory Methods Appl. 5 (1981), 1245-1256.
- [4] F. Nier: Schrödinger–Poisson systems in dimension $d \leq 3$, the whole space case; Proceedings of the Royal Society of Edinburgh, 123A (1993), 1179-1201.

KHALID BENMLIH

DEPARTMENT OF ECONOMIC SCIENCES, UNIVERSITY OF FEZ

P.O. BOX 42A, FEZ, MOROCCO

E-mail address: kbenmlih@hotmail.com