# EXISTENCE OF SOLUTIONS TO NONLINEAR PARABOLIC PROBLEMS WITH DELAY 

DELIANG HSU


#### Abstract

We prove the existence and uniqueness of global solutions to semilinear parabolic equations with a nonlinear delay term. We study these problems in the whole space $\mathbb{R}^{n}$, obtain classic solutions, and give a mass decay result of the solution.


## 1. Introduction

This paper is devoted to the study of existence and uniqueness of solutions to parabolic problems with delay in unbounded domains. Let $\mathbb{R}^{n}$ be Euclidean space, $n \geq 1$ and consider the problem

$$
\begin{gather*}
\frac{\partial u}{\partial t}=\Delta u+\mu|\nabla u|^{p}+f\left(t, x, u_{t}\right) \quad x \in \mathbb{R}^{n}  \tag{1.1}\\
u(x, s)=\phi(x, s) \quad-r \leq s \leq 0, x \in \mathbb{R}^{n}
\end{gather*}
$$

where $f: \mathbb{R}^{+} \times \mathbb{R}^{n} \times C \rightarrow \mathbb{R}$ is a locally Lipschitz functions with respect to $u_{t}, C$ denotes the phase space will be defined in section $2, r>0, \mu \neq 0$, and $u_{t}=u(x, t+\theta)$, for $-r \leq \theta \leq 0$, denotes the delay term, $\phi(x, s) \in C^{2}\left(\mathbb{R}^{n} \times[-r, 0]\right)$.

In the recent years many authors have been concerned with the nonlinear partial differential equations involving delay, for the examples we refer to the book [8] and references cited therein. Generally speaking, this type of problems can be described as an abstract nonlinear partial functional differential equation

$$
\begin{gathered}
\frac{d u}{d t}=A_{T} u(t)+F\left(t, u, u_{t}\right) \\
u(\theta)=\phi(\theta), \quad \text { for }-r \leq \theta \leq 0
\end{gathered}
$$

and then it can be written as an integral equation

$$
u(t)=T(t) \phi(0)+\int_{0}^{t} T(t-s) F\left(s, u, u_{s}\right) d s
$$

where $T(t)$ is a strongly continuous semigroup of bounded linear operators with $A_{T}$ its infinitesimal generator. Then many methods similar to those adopted in ordinary equations can be used to study this type of problems in abstract space, for the details we refer to [8]. However, the problem (1.1) we investigate here

[^0]cannot be carried on along this line for at least in two reasons: first, the operator $\Delta$ on $\mathbb{R}^{n}$ is not compact, so the related semigroup bear some more complexity than the compact cases; second, the general nonlinear term $F\left(t, u, u_{t}\right)$ depend on $\nabla u$, which may cause some difficulty for the study of our problem. Many mechanic and physical problems reduced to the problem (1.1), for the case of $f\left(t, x, u_{t}\right)=0$, (1.1) has emerged in recent years in a number of interesting, and quite different models, for example in the one-dimensional case and $1<p<2$, it has been known as the KPZ-equation, serving as a physical model of growth of surfaces, see Krug and Spohn [6, [7]. In [1] the authors studied this problem in high dimension and proved the existence results, and get some useful mass decay results. There also has been a strong interest in the related equation
\[

$$
\begin{equation*}
u_{t}=\Delta u+\mu|\nabla u|^{p}+u^{q} \tag{1.2}
\end{equation*}
$$

\]

where $p, q>1, \mu<0$, both in bounded and unbounded domains, see [4, 8] and the references cited therein. In this paper, we focus on studying problem 1.1) by using some estimates in parabolic equation involving the nonlinearity $|\nabla u|^{p}$, which is called the damping nonlinear gradient term. To our knowledge there are no results about the problem we present here, even if in the case of $\mu=0$ the solution of (1.1) may blow up in finite time. This paper is motivated by the recent results in 1] and [4. Our main results are the following theorem.

Theorem 1.1. Let $\phi \in C$. Assume that $p \geq 1, f$ satisfies a Lipschitz-type condition,

$$
\begin{equation*}
\left|\nabla_{y} f(t, x, y)\right| \leq L|y| \tag{1.3}
\end{equation*}
$$

for some constant $L>0$, then there exists a unique classical solution of (1.1) in $\Omega_{T_{0}}$, where $T_{0}=\min \left\{\left[2^{p+1} C_{e}\right]^{-2}, \frac{1}{2}\right\}>0$, and $C_{e}$ will be defined later.
Theorem 1.2. Let $f\left(t, x, u_{t}\right)=g(t, u(t-r)), \phi \in C$, here $g: \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{R}$ satisfies the strong Lipschitz-condition

$$
\begin{equation*}
\left|g_{u}(t, u)\right| \leq L \tag{1.4}
\end{equation*}
$$

for any $(t, u) \in \mathbb{R} \times \mathbb{R}$. Then (1.1) has a unique classical solution $u(x, t)$ in $\mathbb{R}^{n} \times$ $[0, T]$, where $T>0$ is any real number, and for every $t \in[0, T], u(\cdot, t) \in C_{b}^{2}\left(\mathbb{R}^{n}\right)$.
Theorem 1.3. Assume that $0 \leq \phi(0, x) \in C_{b}^{2}\left(\mathbb{R}^{n}\right) \cap L^{1}\left(\mathbb{R}^{n}\right), \mu<0, p<(n+$ $2) /(n+1)$ and $f(t, x, u) \leq 0$ for any $(t, u) \in \mathbb{R}^{+} \times C$. If $u(x, t)$ is a $C^{2}$ solution of (1.1) in $\mathbb{R}^{n} \times[0, \infty)$, then it decays as $t \rightarrow+\infty$ in the sense that

$$
\lim _{t \rightarrow+\infty} \int_{\mathbb{R}^{n}} u(x, t) d x=0
$$

## 2. Proof of Main Theorems

In what follows we set $C=\left\{u(\cdot, t) \in C_{b}^{2}\left(\mathbb{R}^{n}\right)\right.$, for any $\left.t \in[-r, 0]\right\}$ as the phase space with the norm

$$
\|u(x, t)\|_{C}=\max _{t}\left\{\|u(x, t)\|_{L^{\infty}}+\|\nabla u(x, t)\|_{L^{\infty}}\right\}
$$

where

$$
\begin{gathered}
C_{b}^{2}\left(\mathbb{R}^{n}\right)=\left\{u: u \in C^{2}\left(\mathbb{R}^{n}\right), u, \nabla u, \nabla^{2} u \in L^{\infty}\left(\mathbb{R}^{n}\right)\right\} \\
\Omega_{t}=\mathbb{R}^{n} \times[-r, t], \quad t \geq 0
\end{gathered}
$$

To prove our theorems, we begin by establishing the existence of the solution for a short time. The idea of the proof is inspired by the papers of Amour and Ben-Artzi [5] and [1].
Proof of Theorem 1.1. Define

$$
u^{0}(x, t)= \begin{cases}\phi(x, t) & \text { for } t \in[-r, 0] \\ \int_{\mathbb{R}^{n}} G(x-y, t) \phi(y, 0) d y & \text { for } t \in\left(0, T_{0}\right)\end{cases}
$$

where the heat kernel is,

$$
G(x, t)=(4 \pi t)^{-\frac{n}{2}} e^{-\frac{|x|^{2}}{4 t}}, \quad t>0
$$

which satisfies,

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} G(x, t) d x=1, \quad \int_{\mathbb{R}^{n}}|\nabla G(x, t)| d x=\beta t^{-1 / 2} \tag{2.1}
\end{equation*}
$$

with $\beta=\int_{\mathbb{R}^{n}}|\nabla G(x, 1)| d x$. Then we may solve by iterations the following linear heat equation, with delay,

$$
\begin{gather*}
\frac{\partial u^{k}}{\partial t}-\Delta u^{k}=\mu\left|\nabla u^{k-1}\right|^{p}+f\left(t, x, u_{t}^{k-1}\right)  \tag{2.2}\\
u^{k}(x, s)=\phi(x, s) \quad \text { for }-r \leq s \leq 0
\end{gather*}
$$

$k=1,2,3, \ldots$ By Duhamel's principle from 2.2, we get

$$
\begin{align*}
u^{k}(x, t)= & \int_{\mathbb{R}^{n}} G(x-y, t) \phi(y, 0) d y+\mu \int_{\mathbb{R}^{n}} \int_{0}^{t} G(x-y, t-s)\left|\nabla u^{k-1}(y, s)\right|^{p} d y d s \\
& +\int_{\mathbb{R}^{n}} \int_{0}^{t} G(x-y, t-s) f\left(s, x, u_{s}^{k-1}\right) d y d s \\
u^{k}(x, t)= & \phi(x, t) \quad \text { for } t \in[-r, 0] \tag{2.3}
\end{align*}
$$

$$
\nabla u^{k}(x, t)=\int_{\mathbb{R}^{n}} G(x-y, t) \nabla \phi(y, 0) d y+\mu \int_{\mathbb{R}^{n}} \int_{0}^{t} \nabla_{x} G(x-y, t-s)\left|\nabla u^{k-1}(y, s)\right|^{p} d y d s
$$

$$
\begin{equation*}
+\int_{\mathbb{R}^{n}} \int_{0}^{t} \nabla_{x} G(x-y, t-s) f\left(s, x, u_{s}^{k-1}\right) d y d s \tag{2.4}
\end{equation*}
$$

Setting $M_{k}(t)=\sup _{\Omega_{t}}\left|\nabla u^{k}(x, t)\right|$, and $U_{k}(t)=\sup _{\Omega_{t}}\left|u^{k}(x, t)\right|$, in view of 2.1), (2.3), and (2.4), we have

$$
\begin{gather*}
M_{k}(t) \leq M_{0}(t)+\beta|\mu| \int_{0}^{t}(t-s)^{-1 / 2} M_{k-1}^{p}(s) d s+L \int_{0}^{t}(t-s)^{-1 / 2} U_{k-1}(s) d s  \tag{2.5}\\
U_{k}(t) \leq U_{0}(t)+\beta|\mu| \int_{0}^{t} M_{k-1}^{p}(s) d s+L \int_{0}^{t} U_{k-1}(s) d s \tag{2.6}
\end{gather*}
$$

Since $M_{0}(t) \leq\|\nabla \phi(x, t)\|_{C}$ and $U_{k}(t) \leq\|\phi(x, t)\|_{C}$, it follows inductively from 2.5 that for $t \leq T_{0}$,

$$
\begin{gather*}
\left\|\nabla u^{k}(x, t)\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \leq 2\|\phi(x, t)\|_{C}  \tag{2.7}\\
\left\|u^{k}(x, t)\right\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \leq 2\|\phi(x, t)\|_{C}, \quad k=0,1,2, \ldots
\end{gather*}
$$

To prove the convergence of the iterations, we need the following inequality

$$
\begin{equation*}
\left|\left|\nabla u^{k}(x, t)\right|^{p}-\left|\nabla u^{k-1}(x, t)\right|^{p}\right| \leq C_{p}\|\phi(x, t)\|_{C}^{p-1}\left|\nabla u^{k}(x, t)-\nabla u^{k-1}(x, t)\right| \tag{2.8}
\end{equation*}
$$

which can be derived by using (2.7) and the classical inequality: $a^{p}-b^{p} \leq C_{p}(a-$ $b)\left(a^{p-1}-b^{p-1}\right), a>0, b>0$. Now setting

$$
N_{k}(t)=\sup _{\Omega_{t}}\left|\nabla u^{k}(x, t)-\nabla u^{k-1}(x, t)\right|, \quad V_{k}(t)=\sup _{\Omega_{t}}\left|u^{k}(x, t)-u^{k-1}(x, t)\right|,
$$

then from (2.3), 2.4 and (2.8), we have

$$
\begin{align*}
& N_{k}(t) \\
& \leq\left.|\mu| \sup \int_{0}^{t} \int_{\mathbb{R}^{n}}| | \nabla u^{k-1}(y, s)\right|^{p}-\left|\nabla u^{k-2}(y, s)\right|^{p} \| \nabla_{x} G(x-y, t-s) \mid d y d s \\
& \quad+L \sup \int_{0}^{t} \int_{\mathbb{R}^{n}}\left|\nabla_{x} G(x-y, t-s) \| u^{k}(y, s)-u^{k-1}(y, s)\right| d y d s \\
& \leq C_{p}|\mu|\|\phi(x, t)\|_{C}^{p-1} \beta \int_{0}^{t}(t-s)^{-1 / 2} N_{k-1}(s) d s+L \beta \int_{0}^{t}(t-s)^{-1 / 2} V_{k-1}(s) d s  \tag{2.9}\\
& V_{k}(t) \leq C_{p}|\mu|\|\phi(x, t)\|_{C}^{p-1} \int_{0}^{t} N_{k-1}(s) d s+L \int_{0}^{t} V_{k-1}(s) d s \\
& \quad \leq C_{p}|\mu|\|\phi(x, t)\|_{C}^{p-1} \int_{0}^{t}(t-s)^{-1 / 2} N_{k-1}(s) d s+L \int_{0}^{t}(t-s)^{-1 / 2} V_{k-1}(s) d s \tag{2.10}
\end{align*}
$$

Choosing $C_{e}=\max \left\{C_{p}|\mu|\|\phi(x, t)\|_{C}^{p-1} \beta, L \beta, C_{p}|\mu|\|\phi(x, t)\|_{C}^{p-1}, L\right\}>0$, from 2.9, (2.10), we have

$$
\begin{equation*}
N_{k}(t)+V_{k}(t) \leq C_{e} \int_{0}^{t}(t-s)^{-1 / 2}\left[N_{k-1}(s)+V_{k-1}(s)\right] d s \tag{2.11}
\end{equation*}
$$

then it follows inductively that (see, [10, Chapter 3]),

$$
\begin{equation*}
N_{k}(t)+V_{k}(t) \leq C_{e}^{k} t^{\frac{k}{2}} \Gamma\left(\frac{k+2}{2}\right)^{-1} \tag{2.12}
\end{equation*}
$$

In particular, $\sum_{k}\left(N_{k}\left(T_{0}\right)+V_{k}\left(T_{0}\right)\right)<\infty$, we conclude that $\left\{\nabla u^{k}\right\}_{k=1}^{\infty}$ and $\left\{u^{k}\right\}_{k=1}^{\infty}$ converge uniformly in $\Omega_{T_{0}}$. Then we set

$$
\begin{equation*}
u(x, t)=\lim _{k \rightarrow \infty} u^{k}(x, t) \tag{2.13}
\end{equation*}
$$

Next, we prove the uniform boundedness of $\left\{\nabla^{2} u^{k}\right\}_{k=1}^{\infty}$ in $\Omega_{T_{0}}, f \in C^{2}$ one has $|\nabla| \nabla f\left|\left|\leq C_{n}\right| \nabla^{2} f\right|$, hence also $\left.\left.|\nabla| \nabla f\right|^{p}\left|\leq C_{p, n}\right| \nabla f\right|^{p-1}\left|\nabla^{2} f\right|$, here $C_{n}$ and $C_{p, n}$ depend only on $n$ and $p, n$ respectively. Denoting

$$
L_{k}(t)=\sup _{\Omega_{t}}\left|\nabla^{2} u^{k}(x, t)\right|
$$

it follows from Duhamel's principle and previous estimates that

$$
\begin{align*}
\nabla^{2} u^{k}(x, t)= & \int_{\mathbb{R}^{n}} G(x-y, t) \nabla^{2} \phi(y, 0) d y \\
& +\mu \int_{\mathbb{R}^{n}} \int_{0}^{t} \nabla_{x} G(x-y, t-s) \nabla\left|\nabla u^{k-1}(y, s)\right|^{p} d y d s  \tag{2.14}\\
& +\int_{\mathbb{R}^{n}} \int_{0}^{t} \nabla_{x} G(x-y, t-s) \nabla f\left(s, x, u_{s}^{k-1}\right) d y d s
\end{align*}
$$

then

$$
\begin{equation*}
L_{k}(t) \leq L_{0}(t)+C_{p, n}\|\nabla \phi\|_{C}^{p-1} \beta|\mu| \int_{0}^{t}(t-s)^{-1 / 2} L_{k-1}(s) d s+2 L \beta\|\phi\|_{C} t^{1 / 2} \tag{2.15}
\end{equation*}
$$

If $\Lambda>0$ is large so that

$$
C_{p, n}\|\phi\|_{C}^{p-1} \beta|\mu| \int_{0}^{T_{0}} s^{-1 / 2} e^{-\Lambda s} d s<\frac{1}{2}
$$

then it follows inductively, using $L_{0}(t) \leq\left\|\nabla^{2} \phi\right\|_{C}$ that

$$
\begin{equation*}
L_{k}(t) \leq 2\left\|\nabla^{2} \phi\right\|_{C} \tag{2.16}
\end{equation*}
$$

for $t \in\left[-r, T_{0}\right], k=0,1,2, \ldots$.
Then the same argument as in the proof of [2, Proposition 2.4] shows that $\left\{\nabla^{2} u^{k}\right\}_{k=1}^{\infty}$ is equicontinuous in $\Omega_{T_{0}}$. Using the Arzela-Ascoli theorem and the convergence of 2.13 we conclude from the standard regularity results of parabolic equation that the solution of 2.13 satisfies the 1.1 in classic sense in $\Omega_{T_{0}}$.

If $v(\cdot, t) \in C^{2}$ is another classical solution in $\Omega_{T_{0}}, v(x, t)=\phi(x, t)$ for $-r \leq t \leq 0$, then setting $N_{1}(t)=\sup _{\Omega_{t}}|\nabla u-\nabla v|, N_{2}(t)=\sup _{\Omega_{t}}|u-v|$, we obtain as in 2.11) that

$$
N_{1}(t)+N_{2}(t) \leq C \int_{0}^{t}(t-s)^{-1 / 2}\left(N_{1}(s)+N_{2}(s)\right) d s
$$

for a sufficiently large constant $C$; this implies $N_{1}(t)+N_{2}(t) \equiv 0$ and $u \equiv v$. This completes the proof of the Lemma.

Proof of Theorem 1.2. By Theorem 1.1. there exists a local solution to (1.1). Hence we just need to get an estimate of $\nabla u$. Taking differential on both side we have

$$
\frac{\partial u_{j}}{\partial t}-\Delta u_{j}-g_{u}(t, u(t-r)) u_{j}(t-r)=\sum_{i=1}^{n} \psi_{i}(x, t) \frac{\partial u_{j}}{\partial x_{j}}
$$

where $u_{j}=\frac{\partial u}{\partial x_{j}}$, and $\psi_{i}(x, t)=\mu p|\nabla u|^{p-2} \frac{\partial u}{\partial x_{j}} \in L^{\infty}\left(\Omega_{T}\right)$. By a maximum principle of linear parabolic equation which was proved in the appendix of [1] we can obtain, for $0 \leq t \leq r$,

$$
\left|u_{j}(t, x)\right| \leq C\left(\left|u_{j}(0, x)\right|+\max _{t \in[-r, 0]}\left|\phi_{j}(t, x)\right|\right)
$$

where $C$ depends only on $L, p, n$. By reiterating the procedure above on $[r, 2 r]$, $[2 r, 3 r], \ldots$, we conclude that, for $t \in[0, T]$

$$
\left|u_{j}(t, x)\right| \leq C_{3} e^{r T}\left(\left|u_{j}(0, x)\right|+\max _{t \in[-r, 0]}\left|\phi_{j}(t, x)\right|\right)
$$

that is

$$
|\nabla u(t, x)| \leq C_{3} e^{r T}\left(|\nabla u(0, x)|+\max _{t \in[-r, 0]}|\nabla \phi(t, x)|\right) .
$$

So from Theorem 1, we get the existence of a global solution, and this completes the proof.

Proof of Theorem 1.3. Let $Q(t)=\int_{\mathbb{R}^{n}} u(x, t) d x$. Integrating 1.1 and using the Gauss formula, we have

$$
\begin{equation*}
\frac{d}{d t} Q(t) \leq 0 \tag{2.17}
\end{equation*}
$$

Let $\widetilde{u}(x, t)$ be the solution of the heat equation $\frac{\partial \widetilde{u}}{\partial t}-\Delta \widetilde{u}=0$, having the same initial data $\widetilde{u}(x, 0)=\phi(x, 0)$. Since $\mu<0$, and $f\left(t, x, u_{t}\right) \leq 0$, it follows by standard comparison principle that

$$
\begin{equation*}
0 \leq u(x, t) \leq \widetilde{u}(x, t) \tag{2.18}
\end{equation*}
$$

Integrating (1.1) with respect to $x$ and $t$ yield

$$
Q(T)=Q(0)+\mu \int_{0}^{T} \int_{\mathbb{R}^{n}}|\nabla u(y, s)|^{p} d y d s+\int_{0}^{T} \int_{\mathbb{R}^{n}} f\left(s, x, u_{s}\right) d y d s
$$

by 2.17) and $f(t, x, u) \leq 0$ we can easily conclude that

$$
\begin{equation*}
\int_{0}^{\infty} \int_{\mathbb{R}^{n}}|\nabla u(x, t)|^{p} d x d t<\infty \tag{2.19}
\end{equation*}
$$

Fix $\varepsilon>0$. It follows from (2.19) that there exists a sequence $1<t_{1}<t_{2}<\cdots<$ $t_{m} \rightarrow+\infty$ such that

$$
\int_{\mathbb{R}^{n}}\left|\nabla u\left(x, t_{j}\right)\right|^{p} d x \leq \varepsilon t_{j}^{-1}
$$

$j=1,2, \ldots$ Using the Sobolev inequality we have

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} u\left(x, t_{j}\right)^{p^{\star}} d x \leq C_{p}\left(\varepsilon t_{j}^{-1}\right)^{p^{\star} / p} \tag{2.20}
\end{equation*}
$$

where $p^{\star}=n p /(n-p)$. Now using the condition $p<(n+2) /(n+1)$ it is easy to see that one can choose $\delta>0$ such that

$$
\begin{equation*}
-\frac{1}{p}+\left(\frac{1}{2}+\delta\right) n\left(1-\frac{1}{p^{\star}}\right)<0 \tag{2.21}
\end{equation*}
$$

then the Hölder inequality and 2.20 imply

$$
\int_{|x| \leq t_{j}^{\frac{1}{2}+\delta}} u\left(x, t_{j}\right) d x \leq C_{p, n} \varepsilon^{\frac{1}{p}} t_{j}^{-\frac{1}{p}} t_{j}^{\left(\frac{1}{2}+\delta\right) n\left(1-\frac{1}{p^{\star}}\right)}
$$

$j=1,2, \ldots$, so that by 2.21 we have

$$
\begin{equation*}
\int_{|x| \leq t_{j}^{\frac{1}{2}+\delta}} u\left(x, t_{j}\right) d x \leq C_{p, n} \varepsilon^{1 / p} \tag{2.22}
\end{equation*}
$$

By standard linear parabolic theory, it is easy to see that

$$
\begin{equation*}
\int_{|x| \leq t_{j}^{\frac{1}{2}+\delta}} \widetilde{u}\left(x, t_{j}\right) d x \rightarrow 0 \tag{2.23}
\end{equation*}
$$

as $j \rightarrow \infty$, which in conjunction with (2.17), 2.18) and (2.22) gives

$$
\lim _{t \rightarrow \infty} \int_{\mathbb{R}^{n}} u(x, t) d x=0
$$

This completes the proof.

Remark 2.1. In the paper [4, the authors give the decay results in the critical case $p=(n+2) /(n+1)$ when $f(t, u) \equiv 0$. We think the same result will be true in the critical case, we will study this problem in a later work.

Acknowledgments. The author thanks Professor Shunian Zhang for his useful discussion and suggestions about this paper. Professor Zhang died two years ago, for which the author feels much sorrow.

## References

[1] L.Amour and M.Ben-Artzi, Global existence and decay for viscous Hamilton-Jacobi equation. Nonl. Anal. TMA, 1998, 31:621-628.
[2] M.Ben-Artzi, Global existence and decay for a nonlinear parabolic equation. Nonlinar Analysis, 1992, 19 763-768.
[3] M.Ben-Artzi, J. Goodman and A. Levy, Remark on nonlinear parabolic equation. Trans. $A M S$., to appear.
[4] M.Ben-Artzi and H.Koch, Decay of mass for a semilinear parabolic equation. Commun. in Partial Differential Equations, 1999, 24, 869-881.
[5] L. Alfonsi and F.B. Weissler, Blow-up in $\mathbb{R}^{n}$ for a parabolic equation with a damping nonlinear gradient term. Progress in nonlinear Differential Equations, N. G. Lloyd et al (Ed.). Birkhauser 92.
[6] J. Krug and H. Spohn, Universality classes for deterministic surface growth. Phys. Rev. A, 1998, 38, 4271-4283.
[7] J. Krug and H. Spohn, Kinetic roughening of growing surfaces. Solids far from equilibrium, Ed. C.Godreche, Cambridge Univ. Press, 1991, 479-582.
[8] Jianhong Wu, Theorey and Applications of Partial Functional Differential Equations. Springer-Verlag, 1996.
[9] P. Souplet and F.B.Weissler, Poincaré's inequality and global solutions of a nonlinear parabolic equation. Annales Inst. H. Poinaré Anal. Nonlin. to appear.
[10] Brezis, H., Analyse Fonctionnelle. Masson, Paris, 1983.
Deliang Hsu
Department of Applied Mathematics, Shanghai Jiaotong University, 200240, Shanghai, China

E-mail address: hsudl@online.sh.cn


[^0]:    2000 Mathematics Subject Classification. 35R10, 58F40.
    Key words and phrases. Nonlinear parabolic equation, delay, global existence, mass decay. © 2004 Texas State University - San Marcos.
    Submitted November 18, 2003. Published May 11, 2004.

