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REMARKS ON REGULARITY CRITERIA FOR THE 3D NAVIER-STOKES EQUATIONS

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ABSTRACT. In this article, we study the regularity criteria for the 3D Navier-Stokes equations involving derivatives of the partial components of the velocity. It is proved that if $\nabla_h \tilde{u}$ belongs to Triebel-Lizorkin space, ∇u_3 or u_3 belongs to Morrey-Campanato space, then the solution remains smooth on [0, T].

1. INTRODUCTION

This article is devoted to the Cauchy problem for the following incompressible 3D Navier-Stokes equation:

$$u_t + (u \cdot \nabla)u + \nabla p = \Delta u, \quad x \in \mathbb{R}^3, \ t > 0$$

div $u = 0, \quad x \in \mathbb{R}^3, \ t > 0$ (1.1)

with initial data

$$u(x,0) = u_0, \quad x \in \mathbb{R}^3, \tag{1.2}$$

where $u = (u_1(x,t), u_2(x,t), u_3(x,t))$ and p = p(x,t) denote the unknown velocity vector and the unknown scalar pressure, respectively. In the last century, Leray [11] and Hopf [8] proved the global existence of a weak solution $u(x,t) \in$ $L^{\infty}(0, \infty; L^2(\mathbb{R}^3)) \cap L^2(0, \infty; H^1(\mathbb{R}^3))$ to (1.1)-(1.3) for any given initial datum $u_0(x) \in L^2(\mathbb{R}^3)$. However, whether or not such a weak solution is regular and unique is still a challenging open problem. From that time on, different criteria for regularity of the weak solutions has been proposed.

The classical Prodi-Serrin conditions (see [16, 18, 19]) say that if

$$u \in L^{t}(0,T;L^{s}(\mathbb{R}^{3})), \quad \frac{2}{t} + \frac{3}{s} = 1, \quad 3 \le s \le \infty,$$

then the solution is smooth. Similar results is showed by Beirão da Veiga [1] involving the velocity gradient growth condition:

$$\nabla u \in L^t(0, T; L^s(\mathbb{R}^3)), \quad \frac{2}{t} + \frac{3}{s} = 2, \quad \frac{3}{2} < s \le \infty.$$

Actually, whether the weak solution is smooth when a part of the velocity components is involved. As for this direction, later on, criteria just for one velocity

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component appeared. The first result in this direction is due to Neustupa et al [15] (see also Zhou [21]), where the authors showed that if

$$u_3 \in L^t(0,T; L^s(\mathbb{R}^3)), \quad \frac{2}{t} + \frac{3}{s} = \frac{1}{2}, \quad s \in (6,\infty],$$

then the solution is smooth. A similar result, for the gradient of one velocity component, is independently due to Zhou [22] and Pokorný [17]. In [22], Zhou proved that if

$$\nabla u_3 \in L^t(0,T; L^s(\mathbb{R}^3)), \quad \frac{2}{t} + \frac{3}{s} = \frac{3}{2}, \quad 3 \le s < \infty.$$

then the solution is smooth on [0, T]. This result is extended by Zhou and Pokorný [26]; that is,

$$\nabla u_3 \in L^t(0,T; L^s(\mathbb{R}^3)), \quad \frac{2}{t} + \frac{3}{s} = \frac{23}{12}, \quad 2 \le s \le 3.$$

Further criteria, including several components of the velocity gradient, pressure or other quantities, can be found, here we just list some. Zhou and Pokorný [25] proved the regularity condition

$$u_3 \in L^t(0,T;L^s(\mathbb{R}^3)), \quad \frac{2}{t} + \frac{3}{s} = \frac{3}{4} + \frac{1}{2s}, \quad s > \frac{10}{3}.$$

And in [10], Jia and Zhou proved that if a weak solution u satisfies one of the following two conditions:

$$u_3 \in L^{\infty}(0,T; L^{\frac{10}{3}}(\mathbb{R}^3)); \quad \nabla u_3 \in L^{\infty}(0,T; L^{30/19}(\mathbb{R}^3)),$$

then u is regular on [0, T]. Dong and Zhang [5] proved that if the horizontal derivatives of the two velocity components

$$\int_0^T \|\nabla_h \tilde{u}(s)\|_{\dot{B}^0_{\infty,\infty}} ds < \infty,$$

then the solution keeps smoothness up to time T, where $\tilde{u} = (u_1, u_2, 0)$, and $\nabla_h \tilde{u} = (\partial_1 \tilde{u}, \partial_2 \tilde{u}, 0)$. For other kinds of regularity criteria, see [2, 6, 7, 9, 23, 24, 28, 29, 30] and the references cited therein.

Throughout this paper C will denote a generic positive constant which can vary from line to line. For simplicity, we shall use $\int f(x) dx$ to denote $\int_{R^3} f(x) dx$, use $\|\cdot\|_p$ to denote $\|\cdot\|_{L^p}$.

The purpose of this article is to improve and extend above known regularity criterion of weak solution for the equations (1.1), (1.2) to the Triebel-Lizorkin space and Morrey-Campanato spaces. The main results of this paper read:

Theorem 1.1. Assume that $u_0 \in H^1(\mathbb{R}^3)$ with div $u_0 = 0$. u(x,t) is the corresponding weak solution to (1.1) and (1.2) on [0,T). If additionally

$$\int_{0}^{T} \|\nabla_{h} \widetilde{u}(\cdot, t)\|_{\dot{F}^{0}_{q, \frac{2}{3}q}}^{p} dt < \infty, \quad \text{with } \frac{2}{p} + \frac{3}{q} = 2, \ \frac{3}{2} < q \le \infty,$$
(1.3)

then the solution remains smooth on [0, T].

Theorem 1.2. Assume that $u_0 \in H^1(\mathbb{R}^3)$ with div $u_0 = 0$. u(x,t) is the corresponding weak solution to (1.1) and (1.2) on [0,T). If additionally

$$\int_{0}^{T} \|\nabla u_{3}(\cdot, t)\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{8}{3(2-r)}} dt < \infty, \quad with \quad 0 < r \le 1, \ 2 \le p \le \frac{3}{r}, \tag{1.4}$$

then the solution remains smooth on [0, T].

Theorem 1.3. Assume that $u_0 \in H^1(\mathbb{R}^3)$ with div $u_0 = 0$. u(x,t) is the corresponding weak solution to (1.1) and (1.2) on [0,T). If additionally

$$\int_{0}^{T} \left\| u_{3}(\cdot, t) \right\|_{\dot{M}^{p, \frac{3}{r}}}^{\frac{8}{3-4r}} dt < \infty, \quad with \ 0 < r < \frac{3}{4}, \ 2 \le p \le \frac{3}{r}, \tag{1.5}$$

then the solution remains smooth on [0, T].

Remark 1.4. Noticing that the classical Riesz transformation is bounded in $\dot{B}^0_{\infty,\infty}$, if we take $q = \infty$ in Theorem 1.1, then the classical Beal-Kato-Majda criterion for the Navier-Stokes equations is obtained; that is, if

$$\int_0^T \|\nabla_h \widetilde{u}(\cdot, t)\|_{\dot{B}^0_{\infty,\infty}} dt < \infty$$

then the solution remains smooth on [0, T].

Remark 1.5. Since (it is proved in [12, 13])

$$L^{q}(\mathbb{R}^{3}) = \dot{M}^{q,q}(\mathbb{R}^{3}) \subset \dot{M}^{p,q}(\mathbb{R}^{3}), \quad 1
$$L^{\frac{3}{r}}(\mathbb{R}^{3}) \subset \dot{M}^{p,\frac{3}{r}}(\mathbb{R}^{3}) \subset \dot{X}_{r}(\mathbb{R}^{3}) \subset \dot{M}^{2,\frac{3}{r}}(\mathbb{R}^{3}), \quad 2$$$$

the result of Theorem 1.2 is an improved version of [27, Theorem 2]. Also we obtain, if

$$\int_0^1 \|u_3(\cdot, t)\|_{\dot{X}_r}^{\frac{8}{3-4r}} dt < \infty, \quad \text{with} \quad 0 < r < \frac{3}{4},$$

then the solution remains smooth on [0, T].

2. Preliminaries

In this section, we shall introduce the Littlewood-Paley decomposition theory, and then give some definitions of the homogeneous Besov space, homogeneous Triebel-Lizorkin space, Morrey-Campanato space and multiplier space as well as some relate spaces used throughout this paper. Before this, let us first recall the weak solutions of (1.1)-(1.3):

Let $u_0 \in L^2(\mathbb{R}^3)$ with $\nabla \cdot u_0 = 0$, a measurable \mathbb{R}^3 -valued vector u is said to be a weak solution of (1.1)-(1.3) if the following conditions hold:

- (1) $u(x,t) \in L^{\infty}(0,\infty; L^2(\mathbb{R}^3)) \cap L^2(0,\infty; H^1(\mathbb{R}^3));$
- (2) u solves (1.1)-(1.2) in the sense of distributions;
- (3) the energy inequality holds; i.e,

$$||u||_{2}^{2} + 2\int_{0}^{t} ||\Delta u(\cdot, \tau)||_{2}^{2} d\tau \le ||u_{0}||_{2}^{2}, \quad 0 \le t \le T.$$

Let us choose a nonnegative radial function $\varphi \in C^{\infty}(\mathbb{R}^3)$ be supported in the annulus $\{\xi \in \mathbb{R}^3 : \frac{3}{4} \leq |\xi| \leq \frac{8}{3}\}$, such that $\sum_{l=-\infty}^{\infty} \varphi(2^{-l}\xi) = 1, \forall \xi \neq 0$. For $f \in S'(\mathbb{R}^3)$, the frequency projection operators Δ_l is defined as

$$\Delta_l f = \mathscr{F}^{-1}(\varphi(2^{-j}\cdot)) * f,$$

where $\mathscr{F}^{-1}(g)$ is the inverse Fourier transform of g. The formal decomposition

$$f = \sum_{l=-\infty}^{\infty} \triangle_l f.$$
(2.1)

is called the homogeneous Littlewood-Paley decomposition. Noticing

$$\triangle_l f = \sum_{j=l-1}^{l+1} \triangle_j(\triangle_l f)$$

and using the Young inequality, we have the following class Bernstein inequality:

Lemma 2.1 ([3]). Let $\alpha \in N$, then for all $1 \leq p \leq q \leq \infty$, $\sup_{|\alpha|=k} \|\partial^{\alpha} \triangle_l f\|_q \leq C2^{lk+3l(\frac{1}{p}-\frac{1}{q})} \|\triangle_l f\|_p$. and C is a constant independent of f, l.

For $s\in\mathbb{R}$ and $(p,q)\in[1,\infty]\times[1,\infty],$ the homogeneous Besov space $\dot{B}^s_{p,q}$ is defined by

$$\hat{S}_{p,q}^{s} = \{ f \in Z'(R^{3}) : \|f\|_{\dot{S}_{p,q}^{s}} < \infty \},\$$

where

$$||f||_{\dot{B}^{s}_{p,q}} = \begin{cases} \left(\sum_{j \in z} 2^{jsq} ||\Delta_{j}f(\cdot)||_{p}^{q} \right)^{1/q}, & 1 \le q < \infty, \\ \sup_{j \in z} 2^{js} ||\Delta_{j}f(\cdot)||_{p}, & q = \infty. \end{cases}$$

and $Z'(\mathbb{R}^3)$ denote the dual space of

$$Z'(\mathbb{R}^3) = \{ f \in S(\mathbb{R}^3) : D^{\alpha} \hat{f}(0) = 0. \ \forall a \in N^3 \}.$$

On the other hand, for $s \in \mathbb{R}$, $(p,q) \in [1,\infty) \times [1,\infty]$, and for $s \in \mathbb{R}$, $p = q = \infty$, the homogenous Triebel-Lizorkin space is defined as

$$F_{p,q}^{s} = \{ f \in Z'(\mathbb{R}^{3}) : \|f\|_{\dot{F}_{p,q}^{s}} < \infty \},$$
$$\|f\|_{\dot{F}_{p,q}^{s}} = \begin{cases} \|(\sum_{j \in z} 2^{jsq} |\Delta_{j}f(\cdot)|^{q})^{1/q}\|_{p}, & 1 \le q < \infty, \\ \|\sup_{j \in z} (2^{js} |\Delta_{j}f(\cdot)|)\|_{p}, & q = \infty. \end{cases}$$

Notice that by Minkowski inequality, we have the following two imbedding relations:

$$\dot{B}_{p,q}^s \subset \dot{F}_{p,q}^s, \quad q \le p; \\ \dot{F}_{p,q}^s \subset \dot{B}_{p,q}^s, \quad p \le q.$$

and the following two inclusions:

$$\dot{H}^{s} = \dot{B}^{s}_{2,2} = \dot{F}^{s}_{2,2}, \quad L^{\infty} \subset \dot{F}^{0}_{\infty,\infty} = \dot{B}^{0}_{\infty,\infty}.$$

We refer to [20] for more properties.

For $1 < q \leq p < \infty$, the homogeneous Morrey-Campanato space in \mathbb{R}^3 is

$$\dot{M}^{p,q} = \left\{ f \in L^q_{\text{loc}}(\mathbb{R}^3); \|f\|_{\dot{M}^{p,q}} = \sup_{x \in \mathbb{R}^3} \sup_{R>0} R^{\frac{3}{p} - \frac{3}{q}} \|f\|_{q(B(x,R))} < \infty \right\},$$

For $1 \le p' \le q' < \infty$, we define the homogeneous space

$$\dot{N}^{p',q'} = \Big\{ f \in L^{q'} | f = \sum_{k \in N} g_k, \text{ where } g_k \in L^{q'}_{comp}(R^3) \text{ and} \\ \sum_{k \in N} d_k^{3(\frac{1}{p'} - \frac{1}{q'})} \| g_k \|_{q'} < \infty, \text{ where } d_k = \text{diam}(\text{supp } g_k) < \infty \Big\}.$$

For $0 < \alpha < 3/2$, we say that a function belongs to the multiplier spaces $M(\dot{H}^{\alpha}, L^2)$ if it maps, by pointwise multiplication, \dot{H}^{α} to L^2 :

$$\dot{X}_{\alpha} := M(\dot{H}^{\alpha}, L^2) := \left\{ f \in S'; \|f \cdot g\|_{L^2} \le C \|g\|_{\dot{H}^{\alpha}}, \forall g \in \dot{H}^{\alpha} \right\}.$$

Here, \dot{H}^{α} is the homogeneous Sobolev space of order α ,

$$\dot{H}^{\alpha} = \left\{ f \in L^{1}_{\text{loc}}; \|f\|_{L^{2}} \equiv \left(\int_{R^{3}} |\xi|^{2\alpha} |\hat{u}(\xi)|^{2} \right)^{1/2} < \infty \right\}.$$

where $L^p(1 \le p \le \infty)$ is the Lebsgue space endowed with norm $\|\cdot\|_p$.

Lemma 2.2 ([4, 12]). Let $1 \le p' \le q' < \infty$, and p, q such that $\frac{1}{p} + \frac{1}{p'} = 1, \frac{1}{q} + \frac{1}{q'} = 1$. Then $\dot{M}^{p,q}$ is the dual space of $\dot{N}^{p',q'}$.

Lemma 2.3 ([4, 7, 12]). Let $1 < p' \le q' < 2, m \ge 2$, and $\frac{1}{p} + \frac{1}{p'} = 1$. Denote $\alpha = -\frac{n}{2} + \frac{n}{p} + \frac{n}{m} \in (0, 1]$ Then there exists a constant C > 0, such that for any $u \in L^m(\mathbb{R}^n), v \in \dot{H}^{\alpha}(\mathbb{R}^n)$,

$$\|u \cdot v\|_{\dot{N}^{p',q'}} \le C \|u\|_{L^m} \|v\|_{\dot{H}^{\alpha}}.$$

Lemma 2.4 ([13]). For $0 \le r \le \frac{3}{2}$, let the space $\mathscr{M}(\dot{B}_2^{r,1} \to L^2)$ be the space of functions which are locally square integrable on \mathbb{R}^3 and such that pointwise multiplication with these functions maps boundedly the Besov space $\dot{B}_2^{r,1}(\mathbb{R}^3)$ to $L^2(\mathbb{R}^3)$. The norm in $\mathscr{M}(\dot{B}_2^{r,1} \to L^2)$ is given by the operator norm of pointwise multiplication:

$$\|f\|_{\mathscr{M}(\dot{B}_{2}^{r,1}\to L^{2})} = \sup\{\|fg\|_{2} : \|g\|_{\dot{B}_{2}^{r,1}} \le 1\}.$$

Then, f belongs to $\mathcal{M}(\dot{B}_2^{r,1} \to L^2)$ if and only if f belongs to $\dot{M}^{2,\frac{3}{r}}$ (with equivalence of norms).

3. The proof of main results

Proof of Theorem 1.1. Multiplying $(1.1)_1$ by $-\Delta u$, integrating by parts, noting that $\nabla \cdot u = 0$, we have

$$\frac{1}{2}\frac{d}{dt}\int |\nabla u|^2 \, dx + \int |\Delta u|^2 \, dx = \int [(u \cdot \nabla)u] \cdot \Delta u \, dx =: I. \tag{3.1}$$

Next we estimate the right-hand side of (3.1), with the help of integration by parts and $-\partial_3 u_3 = \partial_1 u_1 + \partial_2 u_2$, one shows that

$$\begin{split} I &= -\sum_{i,j,k=1}^{3} \int \partial_{k} u_{i} \partial_{i} u_{j} \partial_{k} u_{j} \, dx \\ &= -\sum_{i,j=1}^{2} \sum_{k=1}^{3} \int \partial_{k} u_{i} \partial_{i} u_{j} \partial_{k} u_{j} \, dx - \sum_{i,k=1}^{2} \int \partial_{k} u_{i} \partial_{i} u_{3} \partial_{k} u_{3} \, dx \\ &- \sum_{j,k=1}^{2} \int \partial_{k} u_{3} \partial_{3} u_{j} \partial_{k} u_{j} \, dx - \sum_{k=1}^{3} \int \partial_{k} u_{3} \partial_{3} u_{3} \partial_{k} u_{3} \, dx \\ &- \sum_{i=1}^{2} \int \partial_{3} u_{i} \partial_{i} u_{3} \partial_{3} u_{3} \, dx - \sum_{j=1}^{2} \int \partial_{3} u_{3} \partial_{3} u_{j} \partial_{3} u_{j} \, dx \\ &\leq C \int |\nabla_{h} \widetilde{u}| |\nabla u|^{2} \, dx. \end{split}$$

Thus, the above inequality implies

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$$\frac{1}{2}\frac{d}{dt}\|\nabla u\|_2^2 + \|\Delta u\|_2^2 \le C \int |\nabla_h \widetilde{u}| |\nabla u|^2 \, dx. \tag{3.2}$$

Using the Littlewood-Paley decomposition (2.1), $\nabla_h \tilde{u}$ can be written as

$$\nabla_h \widetilde{u} = \sum_{j < -N} \triangle_j (\nabla_h \widetilde{u}) + \sum_{j = -N}^N \triangle_j (\nabla_h \widetilde{u}) + \sum_{j > N} \triangle_j (\nabla_h \widetilde{u}).$$

where N is a positive integer to be chosen later. Substituting this into (3.2), one obtains

$$\frac{1}{2} \frac{d}{dt} \|\nabla u\|_{2}^{2} + \|\Delta u\|_{2}^{2}$$

$$\leq C \sum_{j < -N} \int |\Delta_{j}(\nabla_{h} \widetilde{u})| |\nabla u|^{2} dx + C \sum_{j = -N}^{N} \int |\Delta_{j}(\nabla_{h} \widetilde{u})| |\nabla u|^{2} dx$$

$$+ C \sum_{j > N} \int |\Delta_{j}(\nabla_{h} \widetilde{u})| |\nabla u|^{2} dx$$

$$=: K_{1} + K_{2} + K_{3}.$$
(3.3)

For K_i (i = 1, 2, 3), we now give the estimates one by one. For K_1 , using the Hölder inequality, the Young inequality and Lemma 2.1, it follows that

$$K_{1} \leq C \sum_{j < -N} \|\Delta_{j}(\nabla_{h}\widetilde{u})\|_{\infty} \|\nabla u\|_{2}^{2}$$

$$\leq C \sum_{j < -N} 2^{3j/2} \|\Delta_{j}(\nabla_{h}\widetilde{u})\|_{2} \|\nabla u\|_{2}^{2}$$

$$\leq C \Big(\sum_{j < -N} 2^{3j}\Big)^{1/2} \Big(\sum_{j < -N} \|\Delta_{j}(\nabla_{h}\widetilde{u})\|_{2}^{2}\Big)^{1/2} \|\nabla u\|_{2}^{2}$$

$$\leq C 2^{-3N/2} \|\nabla u\|_{2}^{3}.$$

(3.4)

Where in the last inequality, we use the fact that for all $s \in \mathbb{R}$, $\dot{H}^s = \dot{B}_{2,2}^s$. For K_2 , by the Hölder inequality and the Young inequality, one has

$$K_{2} = C \int \sum_{j=-N}^{N} |\Delta_{j}(\nabla_{h}\widetilde{u})| |\nabla u|^{2} dx$$

$$\leq C N^{\frac{2q-3}{2q}} \int \left(\sum_{j=-N}^{N} |\Delta_{j}(\nabla_{h}\widetilde{u})|^{2q/3} \right)^{3/(2q)} |\nabla u|^{2} dx$$

$$\leq C N^{\frac{2q-3}{2q}} \|\nabla_{h}\widetilde{u}\|_{\dot{F}_{q,\frac{2q}{3}}^{0}} \|\nabla u\|_{2}^{\frac{2q}{q-1}}$$

$$\leq C N^{\frac{2q-3}{2q}} \|\nabla_{h}\widetilde{u}\|_{\dot{F}_{q,\frac{2q}{3}}^{0}} \|\nabla u\|_{2}^{\frac{2q-3}{q}} \|\Delta u\|_{2}^{3/q}$$

$$\leq \frac{1}{2} \|\Delta u\|_{2}^{2} + C N \|\nabla_{h}\widetilde{u}\|_{\dot{F}_{q,\frac{2q}{3}}^{0}}^{\frac{2q}{q-3}} \|\nabla u\|_{2}^{2}.$$
(3.5)

where we used the interpolation inequality

$$||u||_{s} \le C ||u||_{2}^{\frac{3}{s}-\frac{1}{2}} ||u||_{\dot{H}^{1}}^{\frac{3}{2}-\frac{3}{s}},$$

for $2 \leq s \leq 6$.

Finally, using Hölder inequality and Lemma 2.1, K_3 can be estimated as

$$K_{3} = C \sum_{j>N} \int |\Delta_{j}(\nabla_{h}\widetilde{u})| |\nabla u|^{2} dx$$

$$\leq C \sum_{j>N} \|\Delta_{j}(\nabla_{h}\widetilde{u})\|_{3} \|\nabla u\|_{3}^{2}$$

$$\leq C \sum_{j>N} 2^{\frac{j}{2}} \|\Delta_{j}(\nabla_{h}\widetilde{u})\|_{2} \|\nabla u\|_{3}^{2}$$

$$\leq C (\sum_{j>N} 2^{-j})^{1/2} (\sum_{j>N} 2^{2j} \|\Delta_{j}(\nabla_{h}\widetilde{u})\|_{2}^{2})^{1/2} \|\nabla u\|_{2} \|\Delta u\|_{2}$$

$$\leq C 2^{-N/2} \|\nabla u\|_{2} \|\Delta u\|_{2}^{2}.$$
(3.6)

Substituting (3.4), (3.5) and (3.6) in (3.3), we obtain

$$\frac{d}{dt} \|\nabla u\|_{2}^{2} + \|\bigtriangleup u\|_{2}^{2}
\leq C2^{-\frac{3}{2}N} \|\nabla u\|_{2}^{3} + CN \|\nabla_{h}\widetilde{u}\|_{\dot{F}_{q,\frac{2q}{3}}}^{\frac{2q}{2q-3}} \|\nabla u\|_{2}^{2} + C2^{-N/2} \|\nabla u\|_{2} \|\bigtriangleup u\|_{2}^{2}.$$
(3.7)

Now we choose N such that $C2^{-N/2} \|\nabla u\|_2 \leq \frac{1}{2}$; that is

$$N \ge \frac{\ln(\|\nabla u\|_2^2 + e) + \ln C}{\ln 2} + 2.$$

Thus (3.7) implies

$$\frac{d}{dt} \|\nabla u\|_2^2 \le C + C \|\nabla_h \widetilde{u}\|_{\dot{F}^0_{q,\frac{2q}{3}}}^p \ln(\|\nabla u\|_2^2 + e) \|\nabla u\|_2^2.$$

Taking the Gronwall inequality into consideration, we obtain

$$\ln(\|\nabla u\|_{2}^{2}+e) \leq C \Big[1 + \int_{0}^{T} \|\nabla_{h} \widetilde{u}\|_{\dot{F}_{q}^{0},\frac{2q}{3}}^{p}(\tau) d\tau \cdot e^{\int_{0}^{T} \|\nabla_{h} \widetilde{u}\|_{\dot{F}_{q}^{0},\frac{2q}{3}}^{p}(\tau) d\tau} \Big].$$

The proof of Theorem 1.1 is complete under the condition (1.3).

Proof of Theorem 1.2. Multiplying $(1.1)_1$ by $-\Delta_h u$, integrating by parts, noting that $\nabla \cdot u = 0$, we have

$$\frac{1}{2}\frac{d}{dt}\int |\nabla_h u|^2 \,dx + \int |\nabla\nabla_h u|^2 \,dx = \int [(u \cdot \nabla)u] \cdot \triangle_h u \,dx =: J.$$
(3.8)

Next we estimate the right-hand side of (3.8), with the help of integration by parts and $-\partial_3 u_3 = \partial_1 u_1 + \partial_2 u_2$, one shows that

$$J = -\sum_{i,j=1}^{3} \sum_{k=1}^{2} \int \partial_{k} u_{i} \partial_{i} u_{j} \partial_{k} u_{j} dx$$

$$= -\sum_{i,j,k=1}^{2} \int \partial_{k} u_{i} \partial_{i} u_{j} \partial_{k} u_{j} dx - \sum_{i,k=1}^{2} \int \partial_{k} u_{i} \partial_{i} u_{3} \partial_{k} u_{3} dx$$

$$-\sum_{j=1}^{3} \sum_{k=1}^{2} \int \partial_{k} u_{3} \partial_{3} u_{j} \partial_{k} u_{j} dx$$

$$=: J_{1} + J_{2} + J_{3}.$$

(3.9)

For J_2 and J_3 , we obtain

$$|J_2 + J_3| \le C \int |\nabla u_3| |\nabla_h u| |\nabla u| \, dx. \tag{3.10}$$

 J_1 is a sum of eight terms, using the fact $-\partial_3 u_3 = \partial_1 u_1 + \partial_2 u_2$, we can estimate it as

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$$J_{1} = -\int (\partial_{1}u_{1} + \partial_{2}u_{2})[(\partial_{1}u_{1})^{2} - \partial_{1}u_{1}\partial_{2}u_{2} + (\partial_{2}u_{2})^{2}] dx$$

$$-\int (\partial_{1}u_{1} + \partial_{2}u_{2})[(\partial_{2}u_{1})^{2} + \partial_{1}u_{2}\partial_{2}u_{1} + (\partial_{1}u_{2})^{2}] dx$$

$$= \int \partial_{3}u_{3}[(\partial_{1}u_{1})^{2} - \partial_{1}u_{1}\partial_{2}u_{2} + (\partial_{2}u_{2})^{2} + (\partial_{2}u_{1})^{2} + \partial_{1}u_{2}\partial_{2}u_{1} + (\partial_{1}u_{2})^{2}] dx$$

$$\leq C \int |\nabla u_{3}| |\nabla_{h}u| |\nabla u| dx.$$
(3.11)

Substituting the estimates (3.9)-(3.11) in (3.8), we obtain

$$\frac{1}{2}\frac{d}{dt}\|\nabla_{h}u\|_{2}^{2} + \|\nabla\nabla_{h}u\|_{2}^{2} \le C \int |\nabla u_{3}||\nabla_{h}u||\nabla u| \, dx =: L.$$
(3.12)

when 2 , using Lemmas 2.2 and 2.3, and the Young inequality, we obtain

$$L \leq C \|\nabla u_{3}\|_{\dot{M}^{p,\frac{3}{r}}} \||\nabla u| \cdot |\nabla_{h} u|\|_{\dot{N}^{p',\frac{3}{3-r}}} \leq C \|\nabla u_{3}\|_{\dot{M}^{p,\frac{3}{r}}} \|\nabla_{h} u\|_{\dot{H}^{r}} \|\nabla u\|_{2} \leq C \|\nabla u_{3}\|_{\dot{M}^{p,\frac{3}{r}}} \|\nabla u\|_{L^{2}} \|\nabla_{h} u\|_{2}^{1-r} \|\nabla \nabla_{h} u\|_{2}^{r} \leq \frac{1}{2} \|\nabla \nabla_{h} u\|_{2}^{2} + C \|\nabla u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{2}{2-r}} \|\nabla u\|_{2}^{2}.$$

$$(3.13)$$

where we used the inequality

$$\|f\|_{\dot{H}^r} = \||\xi|^r \hat{f}\|_2 = (\int |\xi|^{2r} |\hat{f}|^{2r} |\hat{f}|^{2-2r} d\xi)^{1/2} \le \|f\|_2^{1-r} \|\nabla f\|_2^r,$$

with $0 < r \leq 1$.

In the case p = 2, using Hölders inequality, Lemma 2.4, and the Young inequality, we can estimate L as

$$L \leq C \||\nabla u_{3}| \cdot |\nabla_{h}u|\|_{2} \|\nabla u\|_{2}$$

$$\leq C \|\nabla u_{3}\|_{\dot{M}^{2,\frac{3}{r}}} \|\nabla_{h}u\|_{\dot{B}_{2}^{r,1}} \|\nabla u\|_{2}$$

$$\leq C \|\nabla u_{3}\|_{\dot{M}^{2,\frac{3}{r}}} \|\nabla_{h}u\|_{2}^{1-r} \|\nabla \nabla_{h}u\|_{2}^{r} \|\nabla u\|_{2}$$

$$\leq \frac{1}{2} \|\nabla \nabla_{h}u\|_{2}^{2} + C \|\nabla u_{3}\|_{\dot{M}^{2,\frac{3}{r}}}^{\frac{2}{r-r}} \|\nabla u\|_{2}^{2}.$$
(3.14)

where we used the following interpolation inequality [14]: for $0 \leq r \leq 1$, $||f||_{\dot{B}_{2}^{r,1}} \leq$
$$\begin{split} \|f\|_2^{1-r}\|\nabla f\|_2^r. \\ \text{Now, gathering (3.13) and (3.14) together and substituting into (3.12), we obtain} \end{split}$$

$$\frac{d}{dt} \|\nabla_h u\|_2^2 + \|\nabla\nabla_h u\|_2^2 \le C \|\nabla u_3\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{2}{2-r}} \|\nabla u\|_2^2.$$
(3.15)

Multiplying $(1.1)_1$ by $-\Delta u$, integrating by parts, noting that $\nabla \cdot u = 0$, we have (see [25])

$$\frac{1}{2} \frac{d}{dt} \|\nabla u\|_{2}^{2} + \|\Delta u\|_{2}^{2} = \int [(u \cdot \nabla)u] \cdot \Delta u \, dx$$

$$\leq C \int |\nabla_{h}u| |\nabla u|^{2} \, dx$$

$$\leq C \|\nabla_{h}u\|_{2} \|\nabla u\|_{4}^{2}$$

$$\leq C \|\nabla_{h}u\|_{2} \|\nabla u\|_{2}^{1/2} \|\nabla u\|_{6}^{\frac{3}{2}}$$

$$\leq C \|\nabla_{h}u\|_{2} \|\nabla u\|_{2}^{1/2} \|\nabla \nabla_{h}u\|_{2} \|\Delta u\|_{2}^{1/2}.$$

Integrating, with respect to t, yields

$$\frac{1}{2} \|\nabla u(t)\|_{2}^{2} + \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau
\leq \frac{1}{2} \|\nabla u_{0}\|_{2}^{2} + C \sup_{0 \leq \tau \leq t} \|\nabla_{h} u(\tau)\|_{2} (\int_{0}^{t} \|\nabla u(\tau)\|_{2}^{2} d\tau)^{1/4}
\times \left(\int_{0}^{t} \|\nabla \nabla_{h} u(\tau)\|_{2}^{2} d\tau\right)^{1/2} \left(\int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau\right)^{1/4}.$$
(3.16)

Substituting (3.15) in (3.16), using Hölders inequality and the Young inequality, we obtain

$$\begin{split} &\frac{1}{2} \|\nabla u(t)\|_{2}^{2} + \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau \\ &\leq \frac{1}{2} \|\nabla u_{0}\|_{2}^{2} + (C + C \int_{0}^{t} \|\nabla u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{2}{2-r}} \|\nabla u(\tau)\|_{2}^{2} d\tau) \Big(\int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau\Big)^{1/4} \\ &\leq C + C \Big(\int_{0}^{t} \|\nabla u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{2}{2-r}} \|\nabla u(\tau)\|_{2}^{\frac{3}{2}} \|\nabla u(\tau)\|_{2}^{1/2} d\tau\Big)^{4/3} \\ &\quad + \frac{1}{2} \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau \qquad (3.17) \\ &\leq C + C \Big(\int_{0}^{t} \|\nabla u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{8}{3(2-r)}} \|\nabla u(\tau)\|_{2}^{2} d\tau\Big) \Big(\int_{0}^{t} \|\nabla u(\tau)\|_{2}^{2} d\tau\Big)^{1/3} \\ &\quad + \frac{1}{2} \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau \\ &\leq C + C \int_{0}^{t} \|\nabla u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{8}{3(2-r)}} \|\nabla u(\tau)\|_{2}^{2} d\tau + \frac{1}{2} \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau. \end{split}$$

Absorbing the last term into the left hand side, applying the Gronwall inequality and combining with the standard continuation argument, we conclude that the solutions u can be extended beyond t = T provided that $\nabla u_3 \in L^{\frac{8}{3(2-r)}}(0,T;\dot{M}^{p,\frac{3}{r}}(\mathbb{R}^3))$. This completes the proof of Theorem 1.2.

Proof of Theorem 1.3. We start from (3.9), we can estimate J_2 and J_3 as

$$|J_2 + J_3| \le C \int |u_3| |\nabla u| |\nabla \nabla_h u| \, dx. \tag{3.18}$$

From (3.11), we find that

$$J_1 \le C \int |u_3| |\nabla u| |\nabla \nabla_h u| \, dx. \tag{3.19}$$

Combining (3.9), (3.17), (3.18) with (3.8), we obtain

$$\frac{1}{2}\frac{d}{dt}\|\nabla_h u\|_2^2 + \|\nabla\nabla_h u\|_2^2 \le C \int |u_3||\nabla u||\nabla\nabla_h u| \, dx =: V.$$
(3.20)

When 2 , similarly as in the proof of L in (3.13), we obtain

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$$V \leq C \|u_{3}\|_{\dot{M}^{p,\frac{3}{r}}} \||\nabla u| \cdot |\nabla \nabla_{h} u|\|_{\dot{N}^{p',\frac{3}{3-r}}}$$

$$\leq C \|u_{3}\|_{\dot{M}^{p,\frac{3}{r}}} \|\nabla u\|_{\dot{H}^{r}} \|\nabla \nabla_{h} u\|_{2}$$

$$\leq C \|u_{3}\|_{\dot{M}^{p,\frac{3}{r}}} \|\nabla \nabla_{h} u\|_{L^{2}} \|\nabla u\|_{2}^{1-r} \|\Delta u\|_{2}^{r}$$

$$\leq \frac{1}{2} \|\nabla \nabla_{h} u\|_{2}^{2} + C \|u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{2} \|\nabla u\|_{2}^{2(1-r)} \|\Delta u\|_{2}^{2r}.$$

(3.21)

When p = 2, similarly as in the proof of L in (3.14), we have

$$V \leq C |||u_{3}| \cdot |\nabla u|||_{2} ||\nabla \nabla_{h} u||_{2}$$

$$\leq C ||u_{3}||_{\dot{M}^{2,\frac{3}{r}}} ||\nabla u||_{\dot{B}^{r,1}_{2}} ||\nabla \nabla_{h} u||_{2}$$

$$\leq C ||u_{3}||_{\dot{M}^{2,\frac{3}{r}}} ||\nabla u||_{2}^{1-r} ||\Delta u||_{2}^{r} ||\nabla \nabla_{h} u||_{2}$$

$$\leq \frac{1}{2} ||\nabla \nabla_{h} u||_{2}^{2} + C ||u_{3}||_{\dot{M}^{2,\frac{3}{r}}}^{2} ||\nabla u||_{2}^{2(1-r)} ||\Delta u||_{2}^{2r}.$$
(3.22)

Substituting (3.21) and (3.22) in (3.20), we find that

$$\frac{d}{dt} \|\nabla_h u\|_2^2 + \|\nabla\nabla_h u\|_2^2 \le C \|u_3\|_{\dot{M}^{2,\frac{3}{r}}}^2 \|\nabla u\|_2^{2(1-r)} \|\Delta u\|_2^{2r}.$$
(3.23)

Substituting (3.23) in (3.16), we obtain

$$\begin{split} &\frac{1}{2} \|\nabla u(t)\|_{2}^{2} + \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau \\ &\leq \frac{1}{2} \|\nabla u_{0}\|_{2}^{2} + \left(C + C \int_{0}^{t} \|u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{2} \|\nabla u(\tau)\|_{2}^{2(1-r)} \|\Delta u(\tau)\|_{2}^{2r} d\tau \right) \\ &\times \left(\int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau \right)^{1/4} \\ &\leq C + \frac{1}{4} \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau + C \Big(\int_{0}^{t} \|u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{2} \|\nabla u(\tau)\|_{2}^{2(1-r)} \|\Delta u(\tau)\|_{2}^{2r} d\tau \Big)^{4/3} \\ &\leq C + \frac{1}{4} \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau + C \Big(\int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau \Big)^{4r/3} \\ &\times \Big(\int_{0}^{t} \|u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{2}{1-r}} \|\nabla u(\tau)\|_{2}^{2} d\tau \Big)^{\frac{4(1-r)}{3}} \\ &\leq C + \frac{1}{2} \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau + C \Big(\int_{0}^{t} \|u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{2}{1-r}} \|\nabla u(\tau)\|_{2}^{\frac{3-4r}{2(1-r)}} \|\nabla u(\tau)\|_{2}^{\frac{1}{2(1-r)}} d\tau \Big)^{\frac{4(1-r)}{3-4r}} \\ &\leq C + \frac{1}{2} \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau + C \Big(\int_{0}^{t} \|u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{3-4r}{3}} \|\nabla u(\tau)\|_{2}^{2} d\tau \Big) \Big(\int_{0}^{t} \|\nabla u(\tau)\|_{2}^{2} d\tau \Big)^{\frac{1}{3-4r}} \\ &\leq C + \frac{1}{2} \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau + C \Big(\int_{0}^{t} \|u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{3-4r}{3}} \|\nabla u(\tau)\|_{2}^{2} d\tau \Big) \Big(\int_{0}^{t} \|\nabla u(\tau)\|_{2}^{2} d\tau \Big)^{\frac{1}{3-4r}} \\ &\leq C + \frac{1}{2} \int_{0}^{t} \|\Delta u(\tau)\|_{2}^{2} d\tau + C \int_{0}^{t} \|u_{3}\|_{\dot{M}^{p,\frac{3}{r}}}^{\frac{3-4r}{3}} \|\nabla u(\tau)\|_{2}^{2} d\tau \Big). \end{split}$$

By a similar argument as in the proof of Theorem 1.2, provided that $u_3 \in L^{\frac{8}{3-4r}}(0,T;\dot{M}^{p,\frac{3}{r}}(\mathbb{R}^3))$, we complete the proof of Theorem 1.3.

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