# EXISTENCE AND CONCENTRATION OF POSITIVE BOUND STATES FOR SCHRÖDINGER-POISSON SYSTEMS WITH POTENTIAL FUNCTIONS 

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#### Abstract

In this article we study the existence and concentration behavior of bound states for a nonlinear Schrödinger-Poisson system with a parameter $\varepsilon>0$. Under suitable conditions on the potential functions, we prove that for $\varepsilon$ small the system has a positive solution that concentrates at a point which is a global minimum of the minimax function associated to the related autonomous problem.


## 1. Introduction

In this article we study the Schrödinger-Poisson system

$$
\begin{gather*}
-\varepsilon^{2} \Delta v+V(x) v+K(x) \phi(x) v=|v|^{q-2} v \quad \text { in } \mathbb{R}^{3} \\
-\Delta \phi=K(x) v^{2} \quad \text { in } \mathbb{R}^{3} \tag{1.1}
\end{gather*}
$$

where $\varepsilon>0$ is a parameter, $q \in(4,6)$ and $V, K: \mathbb{R}^{3} \rightarrow \mathbb{R}$ are, respectively, an external potential and a charge density. The unknowns of the system are the field $u$ associated with the particles and the electric potential $\phi$. We are interested in the existence and concentration behavior of solutions of $\sqrt{1.1}$ in the semiclassical limit $\varepsilon \rightarrow 0$.

The first equation of 1.1 is a nonlinear equation in which the potential $\phi$ satisfies a nonlinear Poisson equation. For this reason, 1.1 is called a SchrödingerPoisson system, also known as Schrödinger-Maxwell system. For more information about physical aspects, we refer the reader to [6, 10] and references therein.

We observe that when $\phi \equiv 0,1.1$ reduces to the well known Schrödinger equation

$$
\begin{equation*}
-\varepsilon^{2} \Delta u+V(x) u=f(x, u) \quad x \in \mathbb{R}^{N} \tag{1.2}
\end{equation*}
$$

In the previous years, the nonlinear stationary Schrödinger equation has been widely investigated, mainly in the semiclassical limit as $\varepsilon \rightarrow 0$ (see e.g. [21, 23, 24] and its references). Rabinowitz [21] studied problem (1.2) using mountain pass arguments to find least energy solutions, for $\varepsilon>0$ sufficiently small. Then, Wang [23] proved that the solution in [21] concentrates around the global minimal of $V$ when $\varepsilon$ tends to 0 .

[^0]Wang and Zeng [24] considered the Schrödinger equation

$$
\begin{equation*}
-\varepsilon^{2} \Delta u+V(x) u=K(x)|u|^{p-1} u+Q(x)|u|^{q-1} u, \quad x \in \mathbb{R}^{N} \tag{1.3}
\end{equation*}
$$

where $1<q<p<(n+2) /(n-2)^{+}$. They proved the existence of least energy solutions and their concentration around a point in the semiclassical limit. The authors used the energy function $C(s)$ defined as the minimal energy of the functional associated with $\Delta u+V(s) u=K(s)|u|^{p-1} u+Q(s)|u|^{q-1} u$, where $s \in \mathbb{R}^{N}$ acts as a parameter instead of an independent variable. For each $\varepsilon>0$ sufficiently small, they proved the existence of a solution $u_{\varepsilon}$ for (1.3), whose global maximum approaches to a point $y^{*}$ when $\varepsilon$ tends to 0 . Moreover, under suitable hypothesis on the potentials $V$ and $W$, the function $\xi \mapsto C(\xi)$ assumes a minimum at $y^{*}$.

Motivated by these results, Alves and Soares [2] investigated the same phenomenon for the gradient system

$$
\begin{gather*}
-\varepsilon^{2} \Delta u+V(x) u=Q_{u}(u, v) \quad \text { in } \mathbb{R}^{N} \\
-\varepsilon^{2} \Delta v+W(x) v=Q_{v}(u, v) \quad \text { in } \mathbb{R}^{N}  \tag{1.4}\\
u(x), v(x) \rightarrow 0, \quad \text { as }|x| \rightarrow \infty \\
u, v>0 \quad \mathbb{R}^{N}
\end{gather*}
$$

In this system is natural to expect some competition between the potentials $V$ and $W$, each one trying to attract the local maximum points of the solutions to its minimum points. In fact, in [2] the authors proved that functions $u_{\varepsilon}$ and $v_{\varepsilon}$ satisfies 1.4 and concentrate around the same point which is the minimum of the respective function $C(s)$.

Ianni and Vaira 17 studied the Schrödinger-Poisson system (1.1) proving that if $V$ has a non-degenerated critical point $x_{0}$, then there exists a solution that concentrates around this point. Moreover, they also proved that if $x_{0}$ is degenerated for $V$ and a local minimum for $K$, then there exist a solution concentrating around $x_{0}$. The proof was based in the Lyapunov-Schmidt reduction.

The double parameter perturbation was also considered for system (1.1) by [15, 16. He and Zhou [16] studied the existence and behavior of a ground state solution which concentrates around the global minimum of the potential $V$. They considered $K \equiv 1$ and the presence of the nonlinear term $f(x, u)$.

Yang and Han [26] studied the Schrödinger-Poisson system

$$
\begin{gather*}
-\Delta v+V(x) v+K(x) \phi(x) v=|v|^{q-2} v \quad \text { in } \mathbb{R}^{3} \\
-\Delta \phi=K(x) v^{2} \quad \text { in } \mathbb{R}^{3} \tag{1.5}
\end{gather*}
$$

Under suitable assumptions on $V, K$ and $f$ they proved existence and multiplicity results by using the mountain pass theorem and the fountain theorem. Later, Zhao, Liu and Zhao [27], using variational methods, proved the existence and concentration of solutions for the system

$$
\begin{gather*}
-\Delta v+\lambda V(x) v+K(x) \phi(x) v=|v|^{q-2} v \quad \text { in } \mathbb{R}^{3} \\
-\Delta \phi=K(x) v^{2} \quad \text { in } \mathbb{R}^{3} \tag{1.6}
\end{gather*}
$$

when $\lambda>0$ is a parameter and $2<p<6$.
Several papers dealt with system (1.5) under variety assumptions on potentials $V$ and $K$. Most part of the literature focuses on the study of the system with $V$ or $K$ constant or radially symmetric, mainly studying existence, nonexistence and multiplicity of solutions see e.g. [4, 9, 10, 11, 12, 18, 20, 22].

Using variational methods as in [2, 21, 24, we prove that there exists a solution $u_{\varepsilon}$ for the Schrödinger-Poisson system 1.1 which concentrates around a point, without any additional assumption on the degenerability of such point related with the potentials $V$ and $K$, as used in [17.

More precisely, denote $C_{\infty}$ as the minimax value related to

$$
\begin{gathered}
-\Delta v+V_{\infty} v+K_{\infty} \phi v=|v|^{q-2} v \quad \text { in } \mathbb{R}^{3} \\
-\Delta \phi=K_{\infty} v^{2} \quad \text { in } \mathbb{R}^{3}
\end{gathered}
$$

where the following conditions hold
(H0) There exists $\alpha>0$ such that $V(x), K(x) \geq \alpha>0$ for all $x \in \mathbb{R}^{3}$,
(H1) $V(x)$ and $K(x)$ are continuous functions and $V_{\infty}, K_{\infty}$ are defined by

$$
\begin{aligned}
V_{\infty} & =\liminf _{|x| \rightarrow \infty} V(x)>\inf _{x \in \mathbb{R}^{3}} V(x) \\
K_{\infty} & =\liminf _{|x| \rightarrow \infty} K(x)>\inf _{x \in \mathbb{R}^{3}} K(x) .
\end{aligned}
$$

We prove that if

$$
C_{\infty}>\inf _{\xi \in \mathbb{R}^{3}} C(\xi)
$$

then (1.1) has a positive solution $v_{\varepsilon}$ as $\varepsilon$ tends to zero. After passing to a subsequence, $v_{\varepsilon}$ concentrates at a global minimum point of $C(\xi)$ for $\xi \in \mathbb{R}^{3}$, where the energy function $C(\xi)$ is defined to be the minimax function associated with the problem

$$
\begin{gather*}
-\Delta u+V(\xi) u+K(\xi) \phi(\xi) u=|u|^{q-2} u \quad \text { in } \mathbb{R}^{3} \\
-\Delta \phi=K(\xi) u^{2} \quad \text { in } \mathbb{R}^{3} \tag{1.7}
\end{gather*}
$$

Therefore, $C(\xi)$ plays a central role in our study. The main result for 1.1) reads as follows.

Theorem 1.1. Suppose (H0)-(H1) hold. If

$$
\begin{equation*}
C_{\infty}>\inf _{\xi \in \mathbb{R}^{3}} C(\xi) \tag{1.8}
\end{equation*}
$$

then there exists $\varepsilon^{*}>0$ such that system (1.1)) has a positive solution $v_{\varepsilon}$ for $\varepsilon \in\left(0, \varepsilon^{*}\right)$. Moreover, $v_{\varepsilon}$ concentrates at a local (hence global) maximum point $y^{*} \in \mathbb{R}^{3}$ such that

$$
C\left(y^{*}\right)=\min _{\xi \in \mathbb{R}^{3}} C(\xi)
$$

Theorem 1.1 complements the study made in [12, 17, 26, 27] in the following sense: we deal with the perturbation problem 1.1 and study the concentration behavior of positive bound states.

To the best of our knowledge, the only previous article regarding the concentration of solutions for the perturbed Schrödinger-Poisson system with potentials $V$ and $K$ is [17], where the smoothness of such potentials is considered. We only need the boundedness of $V$ and $K$. Moreover, we do not assume that the concentration point of solutions $v_{\varepsilon}$ for the system (1.1) is a local minimum (or maximum) of such potentials, as in the previous paper. In our research we shall consider a different variational approach.

The outline of this paper is as follows: in Section 2 we set the variational framework. In Section 3 we study the autonomous system related to (1.1). In section 4 we establish an existence result for system (1.1) with $\varepsilon=1$. In section 5 , we prove Theorem 1.1 .

## 2. Variational framework and preliminary results

Throughout this article we use the following notation:

- $H^{1}\left(\mathbb{R}^{3}\right)$ is the usual Sobolev space endowed with the standard scalar product and norm

$$
(u, v)=\int_{\mathbb{R}^{3}}(\nabla u \nabla v+u v) d x, \quad\|u\|^{2}=\int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+u^{2}\right) d x
$$

- $\mathcal{D}^{1,2}=\mathcal{D}^{1,2}\left(\mathbb{R}^{3}\right)$ represents the completion of $C_{0}^{\infty}\left(\mathbb{R}^{3}\right)$ with respect to the norm

$$
\|u\|_{\mathcal{D}^{1,2}}^{2}=\int_{\mathbb{R}^{3}}|\nabla u|^{2} d x .
$$

- $L^{p}(\Omega), 1 \leq p \leq \infty, \Omega \subset \mathbb{R}^{3}$, denotes a Lebesgue space; the norm in $L^{p}(\Omega)$ is denoted by $\|u\|_{L^{p}(\Omega)}$, where $\Omega$ is a proper subset of $\mathbb{R}^{3} ;\|u\|_{p}$ is the norm in $L^{p}\left(\mathbb{R}^{3}\right)$.

We recall that by the Lax-Milgram theorem, for every $v \in H^{1}\left(\mathbb{R}^{3}\right)$, the Poisson equation $-\Delta \phi=v^{2}$ has a unique positive solution $\phi=\phi_{v} \in \mathcal{D}^{1,2}\left(\mathbb{R}^{3}\right)$ given by

$$
\begin{equation*}
\phi_{v}(x)=\int_{\mathbb{R}^{3}} \frac{v^{2}(y)}{|x-y|} d y \tag{2.1}
\end{equation*}
$$

The function $\phi: H^{1}\left(\mathbb{R}^{3}\right) \rightarrow \mathcal{D}^{1,2}\left(\mathbb{R}^{3}\right), \phi[v]=\phi_{v}$ has the following properties (see for instance Cerami and Vaira [8]).
Lemma 2.1. For any $v \in H^{1}\left(\mathbb{R}^{3}\right)$, we have
(i) $\phi$ is continuous and maps bounded sets into bounded sets;
(ii) $\phi_{v} \geq 0$;
(iii) there exists $C>0$ such that $\|\phi\|_{D^{1,2}} \leq C\|v\|^{2}$ and

$$
\int_{\mathbb{R}^{3}}|\nabla v|^{2} d x=\int_{\mathbb{R}^{3}} \phi_{v} v^{2} d x \leq C\|v\|^{4} ;
$$

(iv) $\phi_{t v}=t^{2} \phi_{v}, \forall t>0$;
(v) if $v_{n} \rightharpoonup v$ in $H^{1}\left(\mathbb{R}^{3}\right)$, then $\phi_{v_{n}} \rightharpoonup \phi_{v}$ in $\mathcal{D}^{1,2}\left(\mathbb{R}^{3}\right)$.

As in [4], for every $v \in H^{1}\left(\mathbb{R}^{3}\right)$, there exist a unique solution $\phi=\phi_{K, v} \in \mathcal{D}^{1,2}\left(\mathbb{R}^{3}\right)$ of $-\Delta \phi=K(x) v^{2}$ where

$$
\begin{equation*}
\phi_{K, v}(x)=\int_{\mathbb{R}^{3}} \frac{K(y) v^{2}(y)}{|x-y|} d y \tag{2.2}
\end{equation*}
$$

and it is easy to see that $\phi_{K, v}$ satisfies Lemma 2.1 if $K$ satisfies conditions (H0)(H1).

Substituting 2.2 into the first equation of (1.1), we obtain

$$
\begin{equation*}
-\varepsilon^{2} \Delta v+V(x) v+K(x) \phi_{K, v}(x) v=|v|^{q-2} v \tag{2.3}
\end{equation*}
$$

Making the changing of variables $x \mapsto \varepsilon x$ and setting $u(x)=v(\varepsilon x)$, 2.3 becomes

$$
\begin{equation*}
-\Delta u+V(\varepsilon x) u+K(\varepsilon x) \phi_{K, v}(\varepsilon x) u=|u|^{q-2} u \tag{2.4}
\end{equation*}
$$

A simple computation shows that

$$
\phi_{K, v}(\varepsilon x)=\varepsilon^{2} \phi_{\varepsilon, u}(x)
$$

where

$$
\phi_{\varepsilon, u}(x)=\int_{\mathbb{R}^{3}} \frac{K(\varepsilon y) u^{2}(y)}{|x-y|} d y .
$$

Substituting this into $(2.4)$, Equation (1.1) can be rewritten in the equivalent equation

$$
\begin{equation*}
-\Delta u+V(\varepsilon x) u+\varepsilon^{2} K(\varepsilon x) \phi_{\varepsilon, u} u=|u|^{q-2} u \tag{2.5}
\end{equation*}
$$

Note that if $u_{\varepsilon}$ is a solution of 2.5), then $v_{\varepsilon}(x)=u_{\varepsilon}(x / \varepsilon)$ is a solution of (2.3). We denote by $H_{\varepsilon}=\left\{u \in H^{1}\left(\mathbb{R}^{3}\right): \int_{\mathbb{R}^{3}} V(\varepsilon x) u^{2}<\infty\right\}$ is a Sobolev space endowed with the norm

$$
\|u\|_{\varepsilon}^{2}=\int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+V(\varepsilon x) u^{2}\right) d x
$$

At this step, we see that (5.1) is variational and its solutions are critical points of the functional

$$
\mathcal{I}_{\varepsilon}(u)=\frac{1}{2} \int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+V(\varepsilon x) u^{2}\right) d x+\frac{\varepsilon^{2}}{4} \int_{\mathbb{R}^{3}} K(\varepsilon x) \phi_{\varepsilon, u}(x) u^{2} d x-\frac{1}{q} \int_{\mathbb{R}^{3}}|u|^{q} d x .
$$

## 3. Autonomous Case

In this section we study the autonomous system

$$
\begin{gather*}
-\Delta u+V(\xi) u+K(\xi) \phi(x) u=|u|^{q-2} u \quad \text { in } \mathbb{R}^{3} \\
-\Delta \phi=K(\xi) u^{2} \quad \text { in } \mathbb{R}^{3} \tag{3.1}
\end{gather*}
$$

where $\xi \in \mathbb{R}^{3}$. To this system we associate the functional $I_{\xi}: H_{\xi} \mapsto \mathbb{R}$,

$$
\begin{equation*}
I_{\xi}(u)=\frac{1}{2} \int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+V(\xi) u^{2}\right) d x+\frac{1}{4} \int_{\mathbb{R}^{3}} K(\xi) \phi_{u}(x) u^{2} d x-\frac{1}{q} \int_{\mathbb{R}^{3}}|u|^{q} d x . \tag{3.2}
\end{equation*}
$$

Hereafter, the Sobolev space $H_{\xi}=H^{1}\left(\mathbb{R}^{3}\right)$ is endowed with the norm

$$
\|u\|_{\xi}=\int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+V(\xi) u^{2}\right) d x .
$$

By standard arguments, the functional $I_{\xi}$ verifies the Mountain-Pass Geometry, more exactly it satisfies the following lemma.

Lemma 3.1. The functional $I_{\xi}$ satisfies
(i) There exist positive constants $\beta, \rho$ such that $I_{\xi}(u) \geq \beta$ for $\|u\|_{\xi}=\rho$,
(ii) There exists $u_{1} \in H^{1}\left(\mathbb{R}^{3}\right)$ with $\left\|u_{1}\right\|_{\xi}>\rho$ such that $I_{\xi}\left(u_{1}\right)<0$.

Applying a variant of the Mountain Pass Theorem (see [25]), we obtain a sequence $\left(u_{n}\right) \subset H^{1}\left(\mathbb{R}^{3}\right)$ such that

$$
I_{\xi}\left(u_{n}\right) \rightarrow C(\xi) \quad \text { and } \quad I_{\xi}^{\prime}\left(u_{n}\right) \rightarrow 0
$$

where

$$
\begin{gather*}
C(\xi)=\inf _{\gamma \in \Gamma} \max _{0 \leq t \leq 1} I_{\xi}(\gamma(t)), \quad C(\xi) \geq \alpha,  \tag{3.3}\\
\Gamma=\left\{\gamma \in \mathcal{C}\left([0,1], H^{1}\left(\mathbb{R}^{3}\right)\right) \mid \gamma(0)=0, \gamma(1)=u_{1}\right\} . \tag{3.4}
\end{gather*}
$$

We observe that $C(\xi)$ can be also characterized as

$$
C(\xi)=\inf _{u \neq 0} \max _{t>0} I_{\xi}(t u)
$$

Proposition 3.2. Let $\xi \in \mathbb{R}^{3}$. Then system (3.1) has a positive solution $u \in$ $H^{1}\left(\mathbb{R}^{3}\right)$ such that $I_{\xi}^{\prime}(u)=0$ and $I_{\xi}(u)=C(\xi)$, for any $q \in(4,6)$.

The proof of the above propostion is an easy adaptation of Azzollini and Pomponio [5, Theorem 1.1] and we omit it.

Lemma 3.3. The function $\xi \mapsto C(\xi)$ is continuous.
Proof. The proof consists in proving that there exist sequences $\left(\zeta_{n}\right)$ and $\left(\lambda_{n}\right)$ in $\mathbb{R}^{3}$ such that $C\left(\zeta_{n}\right), C\left(\lambda_{n}\right) \rightarrow C(\xi)$ as $n \rightarrow 0$, where

- $\zeta_{n} \rightarrow \xi$ and $C\left(\zeta_{n}\right) \geq C(\xi)$ for all $n$,
- $\lambda_{n} \rightarrow \xi$ and $C\left(\lambda_{n}\right) \geq C(\xi)$ for all $n$,
as we know by Alves and Soares 2.

$$
\text { 4. System } 1.1 \text { with } \varepsilon=1
$$

Setting $\varepsilon=1$, in this section we consider the system

$$
\begin{gather*}
-\Delta u+V(x) u+K(x) \phi(x) u=|u|^{q-2} u \quad \text { in } \mathbb{R}^{3} \\
-\Delta \phi=K(x) u^{2} \quad \text { in } \mathbb{R}^{3} \tag{4.1}
\end{gather*}
$$

whose solutions are critical points of the corresponding functional

$$
I(u)=\frac{1}{2} \int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+V(x) u^{2}\right) d x+\frac{1}{4} \int_{\mathbb{R}^{3}} K(x) \phi_{u}(x) u^{2} d x-\frac{1}{p} \int_{\mathbb{R}^{3}}|u|^{q} d x
$$

which is well defined for $u \in H_{1}$, where

$$
H_{1}=\left\{u \in H^{1}\left(\mathbb{R}^{3}\right): \int_{\mathbb{R}^{3}} V(x) u^{2} d x<\infty\right\}
$$

with the same norm notation of the Sobolev space $H^{1}\left(\mathbb{R}^{3}\right)$.
Similar to the autonomous case, the functional $I$ satisfies the mountain pass geometry, then there exists a sequence $\left(u_{n}\right) \subset H_{1}$ such that

$$
\begin{equation*}
I\left(u_{n}\right) \rightarrow c \quad \text { and } \quad I^{\prime}\left(u_{n}\right) \rightarrow 0 \tag{4.2}
\end{equation*}
$$

where

$$
\begin{gathered}
c=\inf _{\gamma \in \Gamma} \max _{0 \leq t \leq 1} I(\gamma(t)) \\
\Gamma=\left\{\gamma \in \mathcal{C}\left([0,1], H_{1}\left(\mathbb{R}^{3}\right)\right) \mid \gamma(0)=0, I(\gamma(1))<0\right\}
\end{gathered}
$$

Remark 4.1. The function $(\mu, \nu) \mapsto c_{\mu, \nu}$ is continuous, where $c_{\mu, \nu}$ is the minimax level of

$$
\begin{equation*}
I_{\mu, \nu}(u)=\frac{1}{2} \int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+\mu u^{2}\right) d x+\frac{1}{4} \int_{\mathbb{R}^{3}} \nu \phi_{u}(x) u^{2} d x-\frac{1}{q} \int_{\mathbb{R}^{3}}|u|^{q} d x . \tag{4.3}
\end{equation*}
$$

Remark 4.2. We denote by $C_{\infty}$ the minimax value related to the functional

$$
I_{\infty}(u)=\frac{1}{2} \int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+V_{\infty} u^{2}\right) d x+\frac{1}{4} \int_{\mathbb{R}^{3}} K_{\infty} \phi_{u} u^{2} d x-\frac{1}{q} \int_{\mathbb{R}^{3}}|u|^{q} d x
$$

where $V_{\infty}$ and $K_{\infty}$, given by condition $\left(H_{1}\right)$, belong to $(0, \infty)$. Otherwise, define $C_{\infty}=\infty . I_{\infty}(u)$ is well defined for $u \in E_{\infty}$, where $E_{\infty}$ is a Sobolev space endowed with the norm

$$
\|u\|_{\infty}=\int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+V_{\infty} u^{2}\right) d x
$$

equivalent to the usual Sobolev norm on $H^{1}\left(\mathbb{R}^{3}\right)$.
An important tool in our analysis is the following theorem.
Theorem 4.3. If $c<C_{\infty}$, then $c$ is a nontrivial critical value for $I$.

Proof. From 4.2, $\left(u_{n}\right)$ is bounded in $H_{1}$. As a consequence, passing to a subsequence if necessary, $u_{n} \rightharpoonup u$ in $H_{1}$. From Proposition 2.1 (v), $\phi_{u_{n}} \rightharpoonup \phi_{u}$ in $\mathcal{D}^{1,2}\left(\mathbb{R}^{3}\right)$, as $n \rightarrow \infty$. Then, $\left(u, \phi_{u}\right)$ is a weak solution of 4.1). Similar to the proof of Proposition 3.1, $I(u)=c$. It remains to show that $u \neq 0$.

From Alves, Souto and Soares [3], if there exist constants $\eta, R$ such that

$$
\liminf _{n \rightarrow+\infty} \int_{B_{R}(0)} u_{n}^{2} d x \geq \eta>0
$$

then $u \neq 0$.
By contradiction, consider $u \equiv 0$. Hence, there exists a subsequence of $\left(u_{n}\right)$, still denoted by $\left(u_{n}\right)$, such that

$$
\lim _{n \rightarrow+\infty} \int_{B_{R}(0)} u_{n}^{2} d x=0
$$

Let $\mu$ and $\nu$ be such that

$$
\inf _{x \in \mathbb{R}^{3}} V(x)<\mu<\liminf _{|x| \rightarrow \infty} V(x)=V_{\infty} \inf _{x \in \mathbb{R}^{3}} K(x)<\nu<\liminf _{|x| \rightarrow \infty} K(x)=K_{\infty}
$$

and take $R>0$ such that

$$
V(x)>\mu, \quad \forall x \in \mathbb{R}^{3} \backslash B_{R}(0) K(x)>\nu, \quad \forall x \in \mathbb{R}^{3} \backslash B_{R}(0)
$$

For each $n \in \mathbb{N}$, there exist $t_{n}>0, t_{n} \rightarrow 1$ such that $I\left(t_{n} u_{n}\right)=\max _{t \geq 0} I\left(t u_{n}\right)$. The convergence of $\left(t_{n}\right)$ follows from (4.2). In fact, since $I^{\prime}\left(u_{n}\right) u_{n}=\bar{o}_{n}(1)$ and $I^{\prime}\left(t_{n} u_{n}\right) t_{n} u_{n}=o_{n}(1)$, we have

$$
\left\|u_{n}\right\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}} u_{n}^{2} d x=\int_{\mathbb{R}^{3}}\left|u_{n}\right|^{q} d x+o_{n}(1)
$$

we have

$$
t_{n}^{2}\left\|u_{n}\right\|^{2}+t_{n}^{4} \int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}} u_{n}^{2} d x=t_{n}^{q} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{q} d x+o_{n}(1)
$$

Then

$$
\left(1-t_{n}^{2}\right)\left\|u_{n}\right\|^{2}=\left(t_{n}^{q-2}-t_{n}^{2}\right) \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{q} d x+o_{n}(1)
$$

Observe that $t_{n}$ neither converge to 0 nor to $\infty$, otherwise we would have $\left\|u_{n}\right\| \rightarrow \infty$ as $n \rightarrow \infty$, which is impossible since $c>0$. See e.g. [1].

Suppose $t_{n} \rightarrow t_{0}$. Letting $n \rightarrow+\infty$,

$$
0=\left(t_{0}^{2}-1\right) \ell_{1}+t_{0}^{2}\left(t_{0}^{q-4}-1\right) \ell_{2}
$$

where $\ell_{1}, \ell_{2}>0$. Hence, $t_{0}=1$. Consequently, we have

$$
\begin{aligned}
& I\left(u_{n}\right)-I\left(t_{n} u_{n}\right) \\
& =\frac{1-t_{n}^{2}}{2}\left\|u_{n}\right\|^{2}+\frac{1}{4}\left(1-t_{n}^{4}\right) \int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}} u_{n}^{2} d x+\frac{t_{n}^{q}-1}{q} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{q} d x=o_{n}(1)
\end{aligned}
$$

which implies, for every $t \geq 0$,

$$
\begin{align*}
I\left(u_{n}\right) \geq & I\left(t u_{n}\right)+o_{n}(1) \\
= & \frac{t^{2}}{2} \int_{\mathbb{R}^{3}}\left|\nabla u_{n}\right|^{2}+V(x) u_{n}^{2} d x+\frac{t^{4}}{4} \int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}} u_{n}^{2} d x \\
& -\frac{t^{q}}{q} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{q} d x+I_{\mu, \nu}\left(t u_{n}\right)-I_{\mu, \nu}\left(t u_{n}\right)+o_{n}(1)  \tag{4.4}\\
\geq & \frac{t^{2}}{2} \int_{B_{R}(0)}(V(x)-\mu) u_{n}^{2} d x+\frac{t^{4}}{4} \int_{B_{R}(0)}(K(x)-\nu) \phi_{u_{n}} u_{n}^{2} d x \\
& +I_{\mu, \nu}\left(t u_{n}\right)+o_{n}(1),
\end{align*}
$$

where $I_{\mu, \nu}(u)$ is given by 4.3).
Consider $\tau_{n}$ such that $I_{\mu, \nu}\left(\tau_{n} u_{n}\right)=\max _{t \geq 0} I_{\mu, \nu}\left(t u_{n}\right)$. As in the above arguments, $\tau_{n} \rightarrow 1$. Letting $t=\tau_{n}$ in 4.4, we have

$$
\begin{aligned}
I\left(u_{n}\right) \geq & \frac{\tau_{n}^{2}}{2} \int_{B_{R}(0)}(V(x)-\mu) u_{n}^{2} d x+\frac{\tau_{n}^{4}}{4} \int_{B_{R}(0)}(K(x)-\nu) \phi_{u_{n}} u_{n}^{2} d x \\
& +c_{\mu, \nu}+o_{n}(1)
\end{aligned}
$$

Taking the limit $n \rightarrow+\infty$, we have $c \geq c_{\mu, \nu}$. Next, taking $\mu \rightarrow V_{\infty}$ and $\nu \rightarrow K_{\infty}$, we obtain $c \geq C_{\infty}$, proving Theorem 4.3 .

## 5. Proof of Theorem 1.1

This section is devoted to study the existence, regularity and the asymptotic behavior of solutions for the system (1.1)), which is equivalent to

$$
\begin{equation*}
-\Delta u+V(\varepsilon x) u+\varepsilon^{2} K(\varepsilon x) \phi_{\varepsilon, u} u=|u|^{q-2} u \tag{5.1}
\end{equation*}
$$

where

$$
\mathcal{I}_{\varepsilon}(u)=\frac{1}{2} \int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+V(\varepsilon x) u^{2}\right) d x+\frac{\varepsilon^{2}}{4} \int_{\mathbb{R}^{3}} K(\varepsilon x) \phi_{\varepsilon, u}(x) u^{2} d x-\frac{1}{q} \int_{\mathbb{R}^{3}}|u|^{q} d x .
$$

is the Euler-Lagrange functional related to (5.1).
The proof of Theorem 1.1 is divided into three subsections as follows:

### 5.1. Existence of a solution.

Theorem 5.1. Suppose (H0)-(H1) hold and consider

$$
\begin{equation*}
C_{\infty}>\inf _{\xi \in \mathbb{R}^{3}} C(\xi) . \tag{5.2}
\end{equation*}
$$

Then, there exists $\varepsilon^{*}>0$ such that system (5.1) has a positive solution for every $0<\varepsilon<\varepsilon^{*}$.

Proof. By hypothesis (5.2), there exists $b \in \mathbb{R}^{3}$ and $\delta>0$ such that

$$
\begin{equation*}
C(b)+\delta<C_{\infty} \tag{5.3}
\end{equation*}
$$

Define $u_{\varepsilon}(x)=u\left(x-\frac{b}{\varepsilon}\right)$, where, from Proposition 3.2, $u$ is a solution of the autonomous Schrödinger-Poisson system

$$
\begin{gather*}
-\Delta u+V(b) u+K(b) \phi(x) u=|u|^{q-2} u \quad \text { in } \mathbb{R}^{3} \\
-\Delta \phi=K(b) u^{2} \quad \text { in } \mathbb{R}^{3} \tag{5.4}
\end{gather*}
$$

with $I_{b}(u)=C(b)$.

Let $t_{\varepsilon}$ be such that $\mathcal{I}_{\varepsilon}\left(t_{\varepsilon} u_{\varepsilon}\right)=\max _{t \geq 0} \mathcal{I}_{\varepsilon}\left(t u_{\varepsilon}\right)$. Similar to the proof of Theorem 4.3. we have $\lim _{\varepsilon \rightarrow 0} t_{\varepsilon}=1$.

Since

$$
c_{\varepsilon}=\inf _{\gamma \in \Gamma} \max _{0 \leq t \leq 1} \mathcal{I}_{\varepsilon}(\gamma(t))=\inf _{\substack{u \in H^{1} \\ u \neq 0}} \max _{t \geq 0} \mathcal{I}_{\varepsilon}(t u) \leq \max _{t \geq 0} \mathcal{I}_{\varepsilon}\left(t u_{\varepsilon}\right)=\mathcal{I}_{\varepsilon}\left(t_{\varepsilon} u_{\varepsilon}\right),
$$

we have

$$
\limsup _{\varepsilon \rightarrow 0} c_{\varepsilon} \leq \limsup _{\varepsilon \rightarrow 0} \mathcal{I}_{\varepsilon}\left(t_{\varepsilon} u_{\varepsilon}\right)=I_{b}(u)=C(b)<C(b)+\delta
$$

which, from 5.3), implies

$$
\limsup _{\varepsilon \rightarrow 0} c_{\varepsilon}<C_{\infty}
$$

Therefore, there exists $\varepsilon^{*}>0$ such that $c_{\varepsilon}<C_{\infty}$ for every $0<\varepsilon<\varepsilon^{*}$. In view of Theorem 4.3, system (5.1) has a positive solution for every $0<\varepsilon<\varepsilon^{*}$.
5.2. Regularity of the solution. The first result is a suitable version of Brezis and Kato [7] and the second one is a particular version of from Gilbarg and Trudinger [14, Theorem 8.17].
Proposition 5.2. Consider $u \in H^{1}\left(\mathbb{R}^{3}\right)$ satisfying

$$
-\Delta u+b(x) u=f(x, u) \quad \text { in } \mathbb{R}^{3}
$$

where $b: \mathbb{R}^{3} \rightarrow \mathbb{R}$ is a $L_{\mathrm{loc}}^{\infty}\left(\mathbb{R}^{3}\right)$ function and $f: \mathbb{R}^{3} \rightarrow \mathbb{R}$ is a Caratheodory function such that

$$
0 \leq f(x, s) \leq C_{f}\left(s^{r}+s\right), \quad \forall s>0, x \in \mathbb{R}^{3}
$$

Then, $u \in L^{t}\left(\mathbb{R}^{3}\right)$ for every $t \geq 2$. Moreover, there exists a positive constant $C=C\left(t, C_{f}\right)$ such that

$$
\|u\|_{L^{t}\left(\mathbb{R}^{3}\right)} \leq C\|u\|_{H^{1}\left(\mathbb{R}^{3}\right)}
$$

Proposition 5.3. Consider $t>3$ and $g \in L^{1 / 2}(\Omega)$, where $\Omega$ is an open subset of $\mathbb{R}^{3}$. Then, if $u \in H^{1}(\Omega)$ is a subsolution of

$$
\Delta u=g \quad \text { in } \Omega
$$

we have that for any $y \in \mathbb{R}^{3}$ and $B_{2 R}(y) \subset \Omega, R>0$ and

$$
\sup _{B_{R}(y)} u \leq C\left(\left\|u^{+}\right\|_{L^{2}\left(B_{2 R}(y)\right)}+\|g\|_{L^{1 / 2}\left(B_{2 R}(y)\right)}\right)
$$

where $C=C(t, R)$.
In view of Propositions 5.2 and 5.3 , the positive solutions of 1.1 are in $C^{2}\left(\mathbb{R}^{3}\right) \cap$ $L^{\infty}\left(\mathbb{R}^{3}\right)$ for all $\varepsilon>0$. Similar arguments was employed by He and Zou [16.

### 5.3. Concentration of solutions.

Lemma 5.4. Suppose $(\mathrm{H} 0)-(\mathrm{H} 1)$ hold. Then, there exists $\beta_{0}>0$ such that

$$
c_{\varepsilon} \geq \beta_{0}
$$

for every $\varepsilon>0$. Moreover,

$$
\limsup _{\varepsilon \rightarrow 0} c_{\varepsilon} \leq \inf _{\xi \in \mathbb{R}^{3}} C(\xi)
$$

Proof. Let $w_{\varepsilon} \in H_{\varepsilon}$ be such that $c_{\varepsilon}=\mathcal{I}_{\varepsilon}\left(w_{\varepsilon}\right)$. Then, from condition $\left(H_{0}\right)$

$$
c_{\varepsilon}=\mathcal{I}_{\varepsilon}\left(w_{\varepsilon}\right) \geq \inf _{\substack{u \in H_{1}^{1} \\ u \neq 0}} \sup _{t \geq 0} J(t u)=\beta_{0}, \quad \forall \varepsilon>0
$$

where

$$
J(u)=\frac{1}{2} \int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+\alpha u^{2}\right) d x+\frac{1}{4} \int_{\mathbb{R}^{3}} \alpha \phi_{u} u^{2} d x-\frac{1}{q} \int_{\mathbb{R}^{3}}|u|^{q} d x
$$

Let $\xi \in \mathbb{R}^{3}$ and consider $w \in H^{1}\left(\mathbb{R}^{3}\right)$ a least energy solution for system (1.7), that is, $I_{\xi}(w)=C(\xi)$ and $I_{\xi}^{\prime}(w)=0$. Let $w_{\varepsilon}(x)=w\left(x-\frac{\xi}{\varepsilon}\right)$ and take $t_{\varepsilon}>0$ such that

$$
c_{\varepsilon} \leq \mathcal{I}_{\varepsilon}\left(t_{\varepsilon} w_{\varepsilon}\right)=\max _{t \geq 0} \mathcal{I}_{\varepsilon}\left(t w_{\varepsilon}\right)
$$

Similar to the proof of Theorem 4.3, $t_{\varepsilon} \rightarrow 1$ as $\varepsilon \rightarrow 0$, then

$$
c_{\varepsilon} \leq \mathcal{I}_{\varepsilon}\left(t_{\varepsilon} w_{\varepsilon}\right) \rightarrow I_{\xi}(w)=C(\xi), \quad \text { as } \varepsilon \rightarrow 0
$$

which implies that $\lim \sup _{\varepsilon \rightarrow 0} c_{\varepsilon} \leq C(\xi)$ for all $\xi \in \mathbb{R}^{3}$. Therefore,

$$
\limsup _{\varepsilon \rightarrow 0} c_{\varepsilon} \leq \inf _{\xi \in \mathbb{R}^{3}} C(\xi)
$$

Lemma 5.5. There exist a family $\left(y_{\varepsilon}\right) \subset \mathbb{R}^{3}$ and constants $R, \beta>0$ such that

$$
\liminf _{\varepsilon \rightarrow 0} \int_{B_{R}\left(y_{\varepsilon}\right)} u_{\varepsilon}^{2} d x \geq \beta, \quad \text { for each } \varepsilon>0
$$

Proof. By contradiction, suppose that there exists a sequence $\varepsilon_{n} \rightarrow 0$ such that

$$
\lim _{n \rightarrow \infty} \sup _{y \in \mathbb{R}^{3}} \int_{B_{R}(y)} u_{n}^{2} d x=0, \quad \text { for all } R>0
$$

where, for the sake of simplicity, we denote $u_{n}(x)=u_{\varepsilon_{n}}(x)$. Hereafter, denote $\phi_{\varepsilon_{n}, u_{n}}(x)=\phi_{u_{n}}(x)$. From [19, Lemma I.1], we have

$$
\int_{\mathbb{R}^{3}}\left|u_{n}\right|^{q} d x \rightarrow 0, \quad \text { as } n \rightarrow \infty
$$

But, since

$$
\int_{\mathbb{R}^{3}}\left(\left|\nabla u_{n}\right|^{2}+V\left(\varepsilon_{n} x\right) u_{n}^{2}\right) d x+\int_{\mathbb{R}^{3}} \varepsilon_{n}^{2} K\left(\varepsilon_{n} x\right) \phi_{u_{n}} u_{n}^{2} d x=\int_{\mathbb{R}^{3}}\left|u_{n}\right|^{q} d x
$$

we have

$$
\int_{\mathbb{R}^{3}}\left(\left|\nabla u_{n}\right|^{2}+V\left(\varepsilon_{n} x\right) u_{n}^{2}\right) d x \rightarrow 0, \quad \text { as } n \rightarrow \infty
$$

Therefore,

$$
\lim _{n \rightarrow \infty} c_{\varepsilon_{n}}=\lim _{n \rightarrow \infty} I_{\varepsilon_{n}}\left(u_{n}\right)=0
$$

which is an absurd, since for some $\beta_{0}>0, c_{\varepsilon} \geq \beta_{0}$, from Lemma 5.4.
Lemma 5.6. The family $\left(\varepsilon y_{\varepsilon}\right)$ is bounded. Moreover, if $y^{*}$ is the limit of the sequence $\left(\varepsilon_{n} y_{\varepsilon_{n}}\right)$ in the family $\left(\varepsilon y_{\varepsilon}\right)$, then we have

$$
C\left(y^{*}\right)=\inf _{\xi \in \mathbb{R}^{3}} C(\xi)
$$

Proof. Consider $u_{n}(x)=u_{\varepsilon_{n}}\left(x+y_{\varepsilon_{n}}\right)$. Suppose by contradiction that $\left(\varepsilon_{n} y_{\varepsilon_{n}}\right)$ approaches infinity. It follows from Lemma 5.5 that there exists constants $R, \beta>0$ such that

$$
\begin{equation*}
\int_{B_{R}(0)} u_{n}^{2}(x) d x \geq \beta>0, \quad \text { for all } n \in \mathbb{N} . \tag{5.5}
\end{equation*}
$$

Since $u_{n}(x)$ satisfies

$$
\begin{equation*}
-\Delta u_{n}+V\left(\varepsilon_{n} x+\varepsilon_{n} y_{\varepsilon_{n}}\right) u_{n}+\varepsilon_{n}^{2} K\left(\varepsilon_{n} x+\varepsilon_{n} y_{\varepsilon_{n}}\right) \phi_{\varepsilon_{n}, u_{n}} u_{n}=\left|u_{n}\right|^{q-2} u_{n} \tag{5.6}
\end{equation*}
$$

it follows that $u_{n}(x)$ is bounded in $H_{\varepsilon}$. Hence, passing to a subsequence if necessary, $u_{n} \rightarrow \hat{u} \geq 0$ weakly in $H_{\varepsilon}$, strongly in $L_{\mathrm{loc}}^{p}\left(\mathbb{R}^{3}\right)$ for $p \in(2,6)$ and a.e. in $\mathbb{R}^{3}$. From (5.5), $\hat{u} \neq 0$.

Using $\hat{u}$ as a test function in (5.6) and taking the limit, we obtain

$$
\begin{equation*}
\int_{\mathbb{R}^{3}}\left(|\nabla \hat{u}|^{2}+\mu \hat{u}^{2}\right) d x \leq \int_{\mathbb{R}^{3}}\left(|\nabla \hat{u}|^{2}+\mu \hat{u}^{2}\right) d x+\int_{\mathbb{R}^{3}} \nu \phi_{\hat{u}} \hat{u}^{2} d x \leq \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x \tag{5.7}
\end{equation*}
$$

where, $\mu$ and $\nu$ are positive constantes such that

$$
\mu<\liminf _{|x| \rightarrow \infty} V(x) \quad \text { and } \quad \nu<\liminf _{|x| \rightarrow \infty} K(x)
$$

Consider the functional $I_{\mu, \nu}: H^{1}\left(\mathbb{R}^{3}\right) \rightarrow \mathbb{R}$ given by

$$
I_{\mu, \nu}(u)=\frac{1}{2} \int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+\mu u^{2}\right) d x+\frac{1}{4} \int_{\mathbb{R}^{3}} \nu \phi_{u}(x) u^{2} d x-\frac{1}{q} \int_{\mathbb{R}^{3}}|u|^{q} d x .
$$

Let $\sigma>0$ be such that $I_{\mu, \nu}(\sigma \hat{u})=\max _{t>0} I_{\mu, \nu}(t \hat{u})$. We claim that

$$
\begin{equation*}
\sigma^{2} \int_{\mathbb{R}^{3}}\left(|\nabla \hat{u}|^{2}+\mu \hat{u}^{2}\right) d x+\sigma^{4} \int_{\mathbb{R}^{3}} \nu \phi_{\hat{u}} \hat{u}^{2} d x=\sigma^{q} \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x . \tag{5.8}
\end{equation*}
$$

In fact, from 5.7

$$
\begin{aligned}
I_{\mu, \nu}(\sigma \hat{u}) & =\frac{\sigma^{2}}{2} \int_{\mathbb{R}^{3}}\left(|\nabla \hat{u}|^{2}+\mu \hat{u}^{2}\right) d x+\frac{\sigma^{4}}{4} \int_{\mathbb{R}^{3}} \nu \phi \hat{u} \hat{u}^{2} d x-\frac{\sigma^{q}}{q} \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x \\
& \leq \frac{\sigma^{2}}{2} \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x+\frac{\sigma^{4}}{4} \int_{\mathbb{R}^{3}} \nu \phi_{\hat{u}} \hat{u}^{2} d x-\frac{\sigma^{q}}{q} \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x
\end{aligned}
$$

it follows that $\sigma \leq 1$, and since $\left.\frac{d}{d t} I_{\mu, \nu}(t \hat{u})\right|_{t=\sigma}=0$, we obtain

$$
\left.\frac{d}{d t} I_{\mu, \nu}(t \hat{u})\right|_{t=\sigma}=\sigma \int_{\mathbb{R}^{3}}\left(|\nabla \hat{u}|^{2}+\mu \hat{u}^{2}\right) d x+\sigma^{3} \int_{\mathbb{R}^{3}} \nu \phi_{\hat{u}} \hat{u}^{2} d x-\sigma^{q-1} \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x=0
$$

proving (5.8).
From Lemma 5.4, equation 5.8 and the fact that $\sigma \leq 1$, we have

$$
\begin{aligned}
c_{\mu, \nu} & =\inf _{u \neq 0} \max _{t>0} I_{\mu, \nu}(t u)=\inf _{u \neq 0} I_{\mu, \nu}(\sigma u) \leq I_{\mu, \nu}(\sigma \hat{u}) \\
& =\frac{\sigma^{2}}{2} \int_{\mathbb{R}^{3}}\left(|\nabla \hat{u}|^{2}+\mu \hat{u}^{2}\right) d x+\frac{\sigma^{4}}{4} \int_{\mathbb{R}^{3}} \nu \phi_{\hat{u}} \hat{u}^{2} d x-\frac{\sigma^{q}}{q} \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x \\
& =\frac{\sigma^{2}}{4} \int_{\mathbb{R}^{3}}\left(|\nabla \hat{u}|^{2}+\mu \hat{u}^{2}\right) d x+\frac{\sigma^{q}(q-4)}{4 q} \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x \\
& \leq \frac{1}{4} \int_{\mathbb{R}^{3}}\left(|\nabla \hat{u}|^{2}+\mu \hat{u}^{2}\right) d x+\frac{q-4}{4 q} \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x \\
& \leq \liminf _{n \rightarrow \infty}\left(\mathcal{I}_{\varepsilon_{n}}\left(u_{n}\right)-\frac{1}{4} \mathcal{I}_{\varepsilon_{n}}^{\prime}\left(u_{n}\right) u_{n}\right)
\end{aligned}
$$

$$
=\liminf _{n \rightarrow \infty} c_{\varepsilon_{n}} \leq \limsup _{n \rightarrow \infty} c_{\varepsilon_{n}} \leq \inf _{\xi \in \mathbb{R}^{3}} C(\xi)
$$

hence, $c_{\mu, \nu} \leq \inf _{\xi \in \mathbb{R}^{3}} C(\xi)$.
If we consider

$$
\mu \rightarrow \liminf _{|x| \rightarrow \infty} V(x)=V_{\infty} \quad \text { and } \quad \nu \rightarrow \liminf _{|x| \rightarrow \infty} K(x)=K_{\infty}
$$

then by the continuity of the function $(\mu, \nu) \mapsto c_{\mu \nu}$ we obtain $C_{\infty} \leq \inf _{\xi \in \mathbb{R}^{3}} C(\xi)$, which contradicts condition $\left(C^{\infty}\right)$. Therefore, $\left(\varepsilon y_{\varepsilon}\right)$ is bounded and there exists a subsequence of $\left(\varepsilon y_{\varepsilon}\right)$ such that $\varepsilon_{n} y_{\varepsilon_{n}} \rightarrow y^{*}$.

Now we proceed to prove that $C\left(y^{*}\right)=\inf _{\xi \in \mathbb{R}^{3}} C(\xi)$. Recalling that $u_{n}(x)=$ $u_{\varepsilon_{n}}\left(x+y_{\varepsilon_{n}}\right)$ and from the arguments above, $\hat{u}$ satisfies the equation

$$
\begin{equation*}
-\Delta u+V\left(y^{*}\right) u+K\left(y^{*}\right) \phi_{u} u=|u|^{q-2} u \tag{5.9}
\end{equation*}
$$

The Euler-Lagrange functional associated to this equation is $I_{y^{*}}: H_{y^{*}}\left(\mathbb{R}^{3}\right)$, defined as in (3.2) with $\xi=y^{*}$.

Using $\hat{u}$ as a test function in (5.9) and taking the limit, we obtain

$$
\int_{\mathbb{R}^{3}}\left(|\nabla \hat{u}|^{2}+V\left(y^{*}\right) \hat{u}^{2}\right) d x \leq \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x .
$$

Then

$$
I_{y^{*}}(\sigma \hat{u})=\max _{t>0} I_{y^{*}}(t \hat{u})
$$

Finally, from Lemma 5.4 and since $0<\sigma \leq 1$ we have

$$
\begin{aligned}
& \inf _{\xi \in \mathbb{R}^{3}} C(\xi) \\
& \leq C\left(y^{*}\right) \leq I_{y^{*}}(\sigma \hat{u}) \\
& =\frac{\sigma^{2}}{4} \int_{\mathbb{R}^{3}}\left(|\nabla \hat{u}|^{2}+V\left(y^{*}\right) \hat{u}^{2}\right) d x+\frac{\sigma^{q}(q-4)}{4 q} \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x \\
& \leq \frac{1}{4} \int_{\mathbb{R}^{3}}\left(|\nabla \hat{u}|^{2}+V\left(y^{*}\right) \hat{u}^{2}\right) d x+\frac{q-4}{4 q} \int_{\mathbb{R}^{3}}|\hat{u}|^{q} d x \\
& \leq \liminf _{n \rightarrow \infty}\left[\frac{1}{4} \int_{\mathbb{R}^{3}}\left(\left|\nabla u_{n}\right|^{2}+V\left(\varepsilon_{n} x+\varepsilon_{n} y_{\varepsilon_{n}}\right) u_{n}^{2}\right) d x+\frac{q-4}{4 q} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{q} d x\right] \\
& \leq \liminf _{n \rightarrow \infty}\left(\mathcal{I}_{\varepsilon_{n}}\left(u_{n}\right)-\frac{1}{4} \mathcal{I}_{\varepsilon_{n}}^{\prime}\left(u_{n}\right) u_{n}\right) \\
& =\liminf _{n \rightarrow \infty} c_{\varepsilon_{n}} \leq \inf _{\xi \in \mathbb{R}^{3}} C(\xi)
\end{aligned}
$$

which implies that $C\left(y^{*}\right)=\inf _{\xi \in \mathbb{R}^{3}} C(\xi)$.
As a consequence of the previous lemma, there exists a subsequence of $\left(\varepsilon_{n} y_{\varepsilon_{n}}\right)$ such that $\varepsilon_{n} y_{\varepsilon_{n}} \rightarrow y^{*}$.

Let $u_{\varepsilon_{n}}\left(x+y_{\varepsilon_{n}}\right)=u_{n}(x)$ and consider $\tilde{u} \in H^{1}$ such that $u_{n} \rightharpoonup \tilde{u}$.
Lemma 5.7. $u_{n} \rightarrow \tilde{u}$ in $H^{1}\left(\mathbb{R}^{3}\right)$, as $n \rightarrow \infty$. Moreover, there exists $\varepsilon^{*}>0$ such that $\lim _{|x| \rightarrow \infty} u_{\varepsilon}(x)=0$ uniformly on $\varepsilon \in\left(0, \varepsilon^{*}\right)$.

Proof. By Lemmas 5.4 and 5.6, we have

$$
\begin{aligned}
& \inf _{\xi \in \mathbb{R}^{3}} C(\xi) \\
& =C\left(y^{*}\right) \leq I_{y^{*}}(\tilde{u})-\frac{1}{4} I_{y^{*}}^{\prime}(\tilde{u}) \tilde{u}
\end{aligned}
$$

$$
\begin{aligned}
& =\frac{1}{4} \int_{\mathbb{R}^{3}}\left(|\nabla \tilde{u}|^{2}+V\left(y^{*}\right) \tilde{u}^{2}\right) d x+\left(\frac{q-4}{4 q}\right) \int_{\mathbb{R}^{3}}|\tilde{u}|^{q} d x \\
& \leq \liminf _{n \rightarrow \infty}\left\{\frac{1}{4} \int_{\mathbb{R}^{3}}\left(\left|\nabla u_{n}\right|^{2}+V\left(\varepsilon_{n} x+\varepsilon_{n} y_{\varepsilon_{n}}\right) u_{n}^{2}\right) d x+\left(\frac{q-4}{4 q}\right) \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{q} d x\right\} \\
& \leq \limsup _{n \rightarrow \infty}\left\{\frac{1}{4} \int_{\mathbb{R}^{3}}\left(\left|\nabla u_{n}\right|^{2}+V\left(\varepsilon_{n} x+\varepsilon_{n} y_{\varepsilon_{n}}\right) u_{n}^{2}\right) d x+\left(\frac{q-4}{4 q}\right) \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{q} d x\right\} \\
& =\limsup _{n \rightarrow \infty}\left\{\mathcal{I}_{\varepsilon_{n}}\left(u_{\varepsilon_{n}}\right)-\frac{1}{4} \mathcal{I}_{\varepsilon_{n}}^{\prime}\left(u_{\varepsilon_{n}}\right) u_{\varepsilon_{n}}\right\} \\
& =\limsup _{n \rightarrow \infty} c_{\varepsilon_{n}} \leq \inf _{\xi \in \mathbb{R}^{3}} C(\xi) .
\end{aligned}
$$

Then

$$
\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left(\left|\nabla u_{n}\right|^{2}+V\left(\varepsilon_{n} x+\varepsilon_{n} y_{\varepsilon_{n}}\right) u_{n}^{2}\right) d x=\int_{\mathbb{R}^{3}}\left(|\nabla \tilde{u}|^{2}+V\left(y^{*}\right) \tilde{u}^{2}\right) d x
$$

Now observe that

$$
\begin{aligned}
c_{\varepsilon_{n}} & =\mathcal{I}_{\varepsilon_{n}}\left(u_{\varepsilon_{n}}\right)-\frac{1}{4} \mathcal{I}_{\varepsilon_{n}}^{\prime}\left(u_{\varepsilon_{n}}\right) u_{\varepsilon_{n}} \\
& =\frac{1}{4} \int_{\mathbb{R}^{3}}\left(\left|\nabla u_{\varepsilon_{n}}\right|^{2}+V\left(\varepsilon_{n} x\right) u_{\varepsilon_{n}}^{2}\right) d x+\left(\frac{q-4}{4 q}\right) \int_{\mathbb{R}^{3}}\left|u_{\varepsilon_{n}}\right|^{q} d x \\
& =\frac{1}{4} \int_{\mathbb{R}^{3}}\left(\left|\nabla u_{n}\right|^{2}+V\left(\varepsilon_{n} x+\varepsilon_{n} y_{\varepsilon_{n}}\right) u_{n}^{2}\right) d x+\left(\frac{q-4}{4 q}\right) \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{q} d x \\
& :=\alpha_{n}
\end{aligned}
$$

hence,

$$
\limsup _{n \rightarrow \infty} \alpha_{n}=\limsup _{n \rightarrow \infty} c_{\varepsilon_{n}} \leq C\left(y^{*}\right)
$$

On the other hand, using Fatou's Lemma,

$$
\begin{aligned}
\liminf _{n \rightarrow \infty} \alpha_{n} & \geq \frac{1}{4} \int_{\mathbb{R}^{3}}\left(|\nabla \tilde{u}|^{2}+V\left(y^{*}\right) \tilde{u}^{2}\right) d x+\left(\frac{q-4}{4 q}\right) \int_{\mathbb{R}^{3}}|\tilde{u}|^{q} d x \\
& =I_{y^{*}}(\tilde{u})-\frac{1}{4} I_{y^{*}}^{\prime}(\tilde{u}) \tilde{u} \\
& \geq C\left(y^{*}\right)
\end{aligned}
$$

then $\lim _{n \rightarrow \infty} \alpha_{n}=C\left(y^{*}\right)$.
Therefore, since $\tilde{u}$ is the weak limit of $\left(u_{n}\right)$ in $H^{1}\left(\mathbb{R}^{3}\right)$, we conclude that $u_{n} \rightarrow \tilde{u}$ strongly in $H^{1}\left(\mathbb{R}^{3}\right)$. In particular, we have

$$
\begin{equation*}
\lim _{R \rightarrow \infty} \int_{|x| \geq R} u_{n}^{2^{*}} d x=0 \quad \text { uniformly on } n \tag{5.10}
\end{equation*}
$$

Applying Proposition 5.2 with $b(x)=V\left(\varepsilon_{n} x+\varepsilon_{n} y_{\varepsilon_{n}}\right)+\varepsilon_{n}^{2} K\left(\varepsilon_{n} x+\varepsilon_{n} y_{\varepsilon_{n}}\right) \phi_{u_{n}}$, we obtain $u_{n} \in L^{t}\left(\mathbb{R}^{3}\right), t \geq 2$ and

$$
\left\|u_{n}\right\|_{t} \leq C\left\|u_{n}\right\|
$$

where $C$ does not depend on $n$.
Now consider

$$
\begin{aligned}
-\Delta u_{n} & \leq-\Delta u_{n}+V\left(\varepsilon_{n} x+\varepsilon_{n} y_{\varepsilon_{n}}\right) u_{n}+\varepsilon_{n}^{2} K\left(\varepsilon_{n} x+\varepsilon_{n} y_{\varepsilon_{n}}\right) \phi_{u_{n}} u_{n} \\
& =\left|u_{n}\right|^{q-2} u_{n}:=g_{n}(x) .
\end{aligned}
$$

For some $t>3,\left\|g_{n}\right\|_{\frac{t}{2}} \leq C$, for all $n$. Using Proposition 5.3, we have

$$
\sup _{B_{R}(y)} u_{n} \leq C\left(\left\|u_{n}\right\|_{L^{2}\left(B_{2 R}(y)\right)}+\left\|g_{n}\right\|_{L^{1 / 2}\left(B_{2 R}(y)\right)}\right)
$$

for every $y \in \mathbb{R}^{3}$, which implies that $\left\|u_{n}\right\|_{L^{\infty}\left(\mathbb{R}^{3}\right)}$ is uniformly bounded. Then, from 5.10,

$$
\lim _{|x| \rightarrow \infty} u_{n}(x)=0 \quad \text { uniformly on } n \in \mathbb{N} .
$$

Consequently, there exists $\varepsilon^{*}>0$ such that

$$
\lim _{|x| \rightarrow \infty} u_{\varepsilon}(x)=0 \quad \text { uniformly on } \varepsilon \in\left(0, \varepsilon^{*}\right)
$$

To complete the proof of Theorem 1.1, it remains to show that the solutions of (1.1) have at most one local (hence global) maximum point $y^{*}$ such that $C\left(y^{*}\right)=$ $\min _{\xi \in \mathbb{R}^{3}} C(\xi)$.

From the previous Lemma, we can focus our attention only in a fixed ball $B_{R}(0) \subset \mathbb{R}^{3}$. If $w \in L^{\infty}\left(\mathbb{R}^{3}\right)$ is the limit in $C_{\text {loc }}^{2}\left(\mathbb{R}^{3}\right)$ of

$$
w_{n}(x)=u_{n}\left(x+y_{n}\right)
$$

then, from Gidas, Ni and Nirenberg [13], $w$ is radially symmetric and has a unique local maximum at zero which is a non-degenerate global maximum. Therefore, there exists $n_{0} \in \mathbb{N}$ such that $w_{n}$ does not have two critical points in $B_{R}(0)$ for all $n \geq n_{0}$. Consider $p_{\varepsilon} \in \mathbb{R}^{3}$ this local (hence global) maximum of $w_{\varepsilon}$.

Recall that if $u_{\varepsilon}$ is a solution of $\left(S_{\varepsilon}\right)$, then

$$
v_{\varepsilon}(x)=u_{\varepsilon}\left(\frac{x}{\varepsilon}\right)
$$

is a solution of 1.1. Since $p_{\varepsilon}$ is the unique maximum of $w_{\varepsilon}$, then $\hat{y}_{\varepsilon}=p_{\varepsilon}+y_{\varepsilon}$ is the unique maximum of $u_{\varepsilon}$. Hence, $\tilde{y}_{\varepsilon}=\varepsilon p_{\varepsilon}+\varepsilon y_{\varepsilon}$ is the unique maximum of $v_{\varepsilon}$. Once $p_{\varepsilon} \in B_{R}(0)$, that is, it is bounded, and $\varepsilon y_{\varepsilon} \rightarrow y^{*}$, we have

$$
\tilde{y}_{\varepsilon} \rightarrow y^{*}
$$

where $C\left(y^{*}\right)=\inf _{\xi \in \mathbb{R}^{3}} C(\xi)$. Consequently, the concentration of functions $v_{\varepsilon}$ approaches $y^{*}$.

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## References

[1] C. O. Alves, P. C. Carrião, O. H. Miyagaki; Nonlinear perturbations of a periodic elliptic problem with critical growth, J. Math. Anal. Appl., 260 (2001), 133-146.
[2] C. O. Alves, S. H. M. Soares; Existence and concentration of positive solutions for a class of gradient systems, Nonlinear Differ. Equ. Appl., 12 (2005), 437-457.
[3] C. O. Alves, S. H. M. Soares, M. A. S. Souto; Schrödinguer-Poisson equations without Ambrosetti-Rabinowitz condition, J. Math. Anal. Appl., 377 (2011), 584-592.
[4] A. Ambrosetti; On Schrödinger-Poisson systems, Milan J. Math., 76 (2008), 257-274.
[5] A. Azzollini, A. Pomponio; Ground state solutions for the nonlinear Schrödinger-Maxwell equations, J. Math. Anal. Appl., 345 (2008), 90-108.
[6] V. Benci, D. Fortunato; An eigenvalue problem for the Schrödinger-Maxwell equations, Top. Meth. Nonlinear Anal, 11 (1998), 283-293.
[7] H. Brezis, T. Kato; Remarks on the Schrödinger operator with singular complex potentials, J. Math. pures et appl., 58 (1979), 137-151.
[8] G. Cerami, G. Vaira; Positive solutions for some non-autonomous Schrödinger-Poisson systems, J. Differential Equations, 248 (2010), 521-543.
[9] G. M. Coclite; A multiplicity result for the nonlinear Schrödinger-Maxwell equations, Commun. Appl. Anal., 7 (2003), 417-423.
[10] T. D'Aprile, D. Mugnai; Solitary waves for nonlinear Klein-Gordon-Maxwell and Schrödinger-Maxwell equations, Proc. R. Soc. Edinb., Sect. A 134 (2004), 1-14.
[11] T. D'Aprile, D. Mugnai; Non-existence results for the coupled Klein-Gordon-Maxwell equations, Adv. Nonlinear Stud., 4 (2004), 307-322.
[12] Y. Fang, J. Zhang; Multiplicity of solutions for the nonlinear Schrödinger-Maxwell system, Commun. Pure App. Anal., 10 (2011), 1267-1279.
[13] B. Gidas, Wei-Ming Ni, L. Nirenberg; Symmetry and related Properties via the Maximum Principle, Commun. Math. Phys., 68 (1979), 209-243.
[14] D. Gilbarg, N. S. Trudinger; Elliptic Partial Differential Equations of Second Order, Second edition, Grundlehrem der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences], (224). Springer-Verlag, Berlin, 1983.
[15] X. He; Multiplicity and concentration of positive solutions for the Schrödinger-Poisson equations, Z. Angew. Math. Phys., 5 (2011), 869-889.
[16] X. He, W. Zou; Existence and concentration of ground states for Schrödinger-Poisson equations with critical growth, J. Math. Phys., 53 (2012), 023702.
[17] I. Ianni, G. Vaira; On concentration of positive bound states for the Schrödinger-Poisson problem with potentials, Adv. Nonlinear Studies, 8 (2008), 573-595.
[18] H. Kikuchi, On the existence of a solution for a elliptic system related to the MaxwellSchrödinger equations, Nonlinear Anal., 67 (2007), 1445-1456.
[19] P. Lions; The concentration-compactness principle in the calculus of variations. The locally compact case. II, Ann. Inst. H. Poincaré Anal. Non Linéaire, 1 (1984), 223-283.
[20] C. Mercuri; Positive solutions of nonlinear Schrödinger-Poisson systems with radial potentials vanishing at infinity, Atti Accad. Naz. Lincei Cl. Sci. Fis. Mat. Natur. Rend. Lincei Mat. Appl., 19 (2008), 211-227.
[21] P. H. Rabinowitz; On a class of nonlinear Schrödinger equations, Z. Angew. Math. Phys., 43 (1992), 270-291.
[22] D. Ruiz; The Schrödinger-Poisson equation under the effect of a nonlinear local term, J. Funct. Anal., 237 (2006), 665-674.
[23] X. Wang; On concentration of positive solutions bounded states of nonlinear Schrödinger equations, Comm. Math. Phys., 153 (1993), 229-244.
[24] X. Wang, B. Zeng; On concentration of positive bound states of nonlinear Schrödinger equations with competing potential functions, SIAM J. Math. Anal., 28 (1997), 633-655.
[25] M. Willem, Minimax theorems, Progress in Nonlinear Differential Equations and their Applications, 24, Birkhäuser, Boston, 1996.
[26] M.-H. Yang, Z.-Q. Han; Existence and multiplicity results for the nonlinear SchrödingerPoisson systems, Nonlinear Anal., 13 (2012), 1093-1101.
[27] L. Zhao, H. Liu, F. Zhao; Existence and concentration of solutions for the SchrödingerPoisson equations with steep potential, J. Differential Equations, 255 (2013), 1-23.

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