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# (p,q)-LAPLACIAN ELLIPTIC SYSTEMS AT RESONANCE

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ABSTRACT. We show the existence of weak solutions for a class of (p,q)-Laplacian elliptic systems at resonance, under certain Landesman-Lazer-type conditions by using critical point theorem.

### 1. INTRODUCTION AND STATEMENT OF MAIN RESULTS

Let  $\Omega$  be a bounded domain with smooth boundary in  $\mathbb{R}^N$  and  $\Delta_p$  be the *p*-Laplacian operator. In this paper, we study the existence of solutions for the problem

$$-\Delta_p u = \lambda_1 |u|^{p-2} u + \frac{\lambda_1}{\beta+1} |u|^{\alpha} |v|^{\beta} v + G_s(x, u, v) - h_1(x) \quad \text{in } \Omega,$$
  

$$-\Delta_q v = \lambda_1 |v|^{q-2} v + \frac{\lambda_1}{\alpha+1} |u|^{\alpha} |v|^{\beta} u + G_t(x, u, v) - h_2(x) \quad \text{in } \Omega,$$
  

$$u = v = 0 \quad \text{on } \partial\Omega,$$
(1.1)

where  $1 < p, q < +\infty$  and  $\alpha \ge 0, \beta \ge 0$  satisfy

$$\frac{\alpha+1}{p} + \frac{\beta+1}{q} = 1.$$
 (1.2)

The nonlinearity  $G: \Omega \times \mathbb{R}^2 \to R$  is a Caratheodory function which has continuous derivatives  $G_s(x, s, t), G_t(x, s, t)$  with respect to s and t for almost any  $x \in \Omega$ , and  $h_1 \in L^{p/(p-1)}(\Omega), h_2 \in L^{q/(q-1)}(\Omega)$ .

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$$||u||_{L^p} \le C ||u||_p, \ \forall u \in W_0^{1,p}(\Omega), \quad \text{and} \quad ||v||_{L^q} \le C ||v||_q, \ \forall v \in W_0^{1,q}(\Omega), \quad (1.3)$$

where  $\|\cdot\|_{L^p}$  denotes the norm of  $L^p(\Omega)$  and throughout this paper, let *C* always denote an embedding constant with relation to (1.3). For the following nonlinear

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eigenvalue problem

$$-\Delta_{p}u = \lambda |u|^{p-2}u + \frac{\lambda}{\beta+1} |u|^{\alpha} |v|^{\beta}v \quad \text{in } \Omega,$$
  

$$-\Delta_{q}v = \lambda |v|^{q-2}v + \frac{\lambda}{\alpha+1} |u|^{\alpha} |v|^{\beta}u \quad \text{in } \Omega,$$
  

$$u = v = 0 \quad \text{on } \partial\Omega,$$
  
(1.4)

consider the functionals  $\phi, \varphi$  on W defined by

$$\phi(u,v) = \frac{1}{p} \int_{\Omega} |\nabla u|^p dx + \frac{1}{q} \int_{\Omega} |\nabla v|^q dx,$$
  
$$\varphi(u,v) = \frac{1}{p} \int_{\Omega} |u|^p dx + \frac{1}{q} \int_{\Omega} |v|^q dx + \frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega} |u|^{\alpha} |v|^{\beta} uv dx,$$
  
we manifold

and the manifold

$$\Sigma = \{(u, v) \in W : \varphi(u, v) = 1\}.$$

It is easy to prove that  $\phi(u, v), \varphi(u, v)$  are (p, q)-homogeneous, namely

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$$\phi(t^{1/p}u, t^{1/q}v) = t\phi(u, v), \quad \varphi(t^{1/p}u, t^{1/q}v) = t\varphi(u, v)$$

for any t > 0 and  $(u, v) \in W$ , and  $\Sigma$  is a symmetric nonempty manifold in W. By an argument similar to the ones in [3, 7], problem (1.4) has a sequence of eigenvalues with the variational characterization

$$\lambda_k = \inf_{\Lambda \in \Sigma_k} \sup_{(u,v) \in \Lambda} \phi(u,v),$$

where  $\Sigma_k = \{\Lambda \subset \Sigma : \text{ there is an odd, continuous and surjective } \gamma : S^{k-1} \to \Lambda \}$ and  $S^{k-1}$  denotes the unit sphere in  $\mathbb{R}^k$ .

On the other hand, let

$$\lambda_1' = \inf_{(u,v)\in\Sigma} \phi(u,v),$$

we can see that  $\lambda_1 = \lambda'_1$ . Moreover,  $\lambda_1$  is a simple, isolated and positive principal eigenvalue of (1.4) and has a positive normalized eigenvalue  $(\mu_0, \nu_0)$ , namely,  $\|\mu_0\|_p + \|\nu_0\|_q = 1$ . By a simple computation, there exists a positive constant  $t_0$  such that

$$||t_0^{1/p}\mu_0||_p^p + ||t_0^{1/q}\nu_0||_q^q = 1.$$

Let  $\mu_1 = t_0^{1/p} \mu_0$ ,  $\nu_1 = t_0^{1/q} \nu_0$ , since  $\phi, \varphi$  are (p, q)-homogeneous, hence the set of all eigenfunctions corresponding to  $\lambda_1$  is

$$E_1 := \{ (t^{1/p}\mu_1, t^{1/q}\nu_1) : t \ge 0 \} \cup \{ (-t^{1/p}\mu_1, -t^{1/q}\nu_1) : t \ge 0 \}.$$

The set  $E_1$  is not an one-dimensional linear subspace of W and the corresponding orthogonal decomposition on W does not hold with respect to the first eigenvalue  $\lambda_1$ .

In many papers, existence of weak solutions for the resonant elliptic problems were investigated under the well-known Landesman-Lazer-type conditions, which were introduced by Landesman and Lazer in [5] and were extended by Tang in [12]. Since then they were used widely for the different types of equations, for example, in [1, 3, 9] for the quasilinear elliptic equations, in [4] for asymptotically linear noncooperative elliptic systems, in [13] for the forced duffing equations, in [11] for Kirchhoff type equations. Especially, in [2] the case p = q = 2 (the semilinear elliptic systems) was considered and the case  $p = q \ge 2$  (the quasilinear elliptic systems)

was discussed in [6, 7, 14] where  $G_s(x, s, t) = g_1(x, s)$  and  $G_t(x, s, t) = g_2(x, t)$ . As far as we know, when  $p \neq q > 1$ , the similar results are not discussed under the Landesman-Lazer-type conditions due to Landesman and Lazer. Motivated by these finding, we consider the existence of solutions for problem (1.1) at resonance with the first eigenvalue under the Landesman-Lazer-type conditions. We first state the following fundamental hypotheses.

- (H1) There is  $h \in C(\overline{\Omega}, \mathbb{R}^+)$  such that  $|G_s(x, s, t)| \le h(x)$  and  $|G_t(x, s, t)| \le h(x)$  for all  $(x, s, t) \in \Omega \times \mathbb{R}^2$ .
- (H2) There exist two functions  $g_1^{++}, g_1^{--} \in C(\Omega, R)$  such that

$$g_1^{++}(x) = \liminf_{s \to +\infty, t \to +\infty} G_s(x, s, t), \quad g_1^{--}(x) = \limsup_{s \to -\infty, t \to -\infty} G_s(x, s, t)$$

uniformly a.e.  $x \in \Omega$ .

(H3) There is two functions  $g_2^{++}, g_2^{--} \in C(\Omega, R)$  such that

$$g_2^{++}(x) = \liminf_{s \to +\infty, t \to +\infty} G_t(x, s, t), \quad g_2^{--}(x) = \limsup_{s \to -\infty, t \to -\infty} G_t(x, s, t)$$

uniformly a.e.  $x \in \Omega$ .

The Landesman-Lazer-type conditions for problem (1.1) are read either

$$\int_{\Omega} g_1^{--} \mu_1 dx + \int_{\Omega} g_2^{--} \nu_1 dx < \int_{\Omega} h_1 \mu_1 dx + \int_{\Omega} h_2 \nu_1 dx < \int_{\Omega} g_1^{++} \mu_1 dx + \int_{\Omega} g_2^{++} \nu_1 dx;$$
(1.5)

or

$$\int_{\Omega} g_1^{++} \mu_1 dx + \int_{\Omega} g_2^{++} \nu_1 dx < \int_{\Omega} h_1 \mu_1 dx + \int_{\Omega} h_2 \nu_1 dx < \int_{\Omega} g_1^{--} \mu_1 dx + \int_{\Omega} g_2^{--} \nu_1 dx.$$
(1.6)

We are ready to state the main results.

**Theorem 1.1.** Let  $h_1 \in L^{p/(p-1)}(\Omega)$ ,  $h_2 \in L^{q/(q-1)}(\Omega)$ , and (1.2), (H1), (H2), (H3) and (1.5) be satisfied. If 1 and the following inequalities hold:

$$\int_{\Omega} h_1 \mu_1 dx - \int_{\Omega} g_1^{++} \mu_1 dx < 0, \quad \int_{\Omega} h_1 \mu_1 dx - \int_{\Omega} g_1^{--} \mu_1 dx > 0, \qquad (1.7)$$

then problem (1.1) has at least one solution.

In the other case 1 < q < p, the following result holds.

**Theorem 1.2.** Let  $h_1 \in L^{p/(p-1)}(\Omega)$ ,  $h_2 \in L^{q/(q-1)}(\Omega)$ , and (1.2), (H1), (H2), (H3) and (1.5) be satisfied. If 1 < q < p and the following inequalities hold:

$$\int_{\Omega} h_2 \nu_1 dx - \int_{\Omega} g_2^{++} \nu_1 dx < 0, \quad \int_{\Omega} h_2 \nu_1 dx - \int_{\Omega} g_2^{--} \nu_1 dx > 0, \qquad (1.8)$$

then problem (1.1) has at least one solution.

**Theorem 1.3.** Let  $h_1 \in L^{p/(p-1)}(\Omega), h_2 \in L^{q/(q-1)}(\Omega)$ . If (1.2), (H1), (H2), (H3) and (1.6) are satisfied, then problem (1.1) has at least one solution.

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Our results extends the ones in [2] from the semilinear elliptic systems to (p, q)-Laplacian elliptic systems, and are also the generalizations of [14], where they considered the case  $p = q \ge 2$  and  $G_s(x, s, t) = g_1(s), G_t(x, s, t) = g_2(t)$ . Moreover, the conditions (1.7) and (1.8) are the technical assumptions. Theorem 1.2 is similar to Theorem 1.1, and we will prove Theorem 1.1 and Theorem 1.3.

#### 2. Proofs of Theorems

Now consider the functionals  $J, J_1, J_2$  on W defined by

$$\begin{split} J(u,v) &= \frac{1}{p} \int_{\Omega} |\nabla u|^{p} dx + \frac{1}{q} \int_{\Omega} |\nabla v|^{q} dx - \frac{\lambda_{1}}{p} \int_{\Omega} |u|^{p} dx \\ &- \frac{\lambda_{1}}{q} \int_{\Omega} |v|^{q} dx - \frac{\lambda_{1}}{(\alpha+1)(\beta+1)} \int_{\Omega} |u|^{\alpha} |v|^{\beta} uv dx \\ &- \int_{\Omega} G(x,u,v) dx + \int_{\Omega} h_{1} u dx + \int_{\Omega} h_{2} v dx, \end{split}$$
$$J_{1}(u,v) &= \frac{1}{p} \int_{\Omega} |\nabla u|^{p} dx - \frac{\lambda_{1}}{p} \int_{\Omega} |u|^{p} dx - \frac{\lambda_{1}}{p(\beta+1)} \int_{\Omega} |u|^{\alpha} |v|^{\beta} uv dx \\ &- \int_{\Omega} \int_{0}^{1} G_{s}(x,ru,rv) u \, dr \, dx + \int_{\Omega} h_{1} u dx, \end{aligned}$$
$$J_{2}(u,v) &= \frac{1}{q} \int_{\Omega} |\nabla v|^{q} dx - \frac{\lambda_{1}}{q} \int_{\Omega} |v|^{q} dx - \frac{\lambda_{1}}{q(\alpha+1)} \int_{\Omega} |u|^{\alpha} |v|^{\beta} uv dx \\ &- \int_{\Omega} \int_{0}^{1} G_{t}(x,ru,rv) v \, dr \, dx + \int_{\Omega} h_{2} v dx. \end{split}$$

Noting that

$$G(x,s,t) = \int_0^1 G_s(x,rs,rt)sdr + \int_0^1 G_t(x,rs,rt)tdr,$$
 (2.1)

from (1.2) and (2.1), it follows that

$$J(u, v) = J_1(u, v) + J_2(u, v)$$
 for all  $(u, v) \in W$ .

From (H1), it is easy to prove that the functional J is well defined and  $J \in C^1(W, R)$ . Moreover, from the variational view of point, a weak solution of problem (1.1) is equivalent to a critical point of the functional J in W. In this paper, we will prove Theorem 1.1 and Theorem 1.2 by using the following G-linking Theorem due to Drábek and Robinson (see [3, 9]) and Theorem 1.3 by using Ekeland's Variational Principle (see [8, 10]). In these abstract theorems, a compact condition, i.e., (PS) condition, is needed.

**Definition 2.1.** Let X be a real Banach space, if for any sequence  $\{u_n\} \subset X$  such that  $f(u_n)$  is bounded and  $f'(u_n) \to 0$  as  $n \to \infty$ ,  $\{u_n\}$  has a convergent subsequence, the functional f satisfies the (PS) condition.

**Definition 2.2** ([3, 9]). Let Q be a submanifold of a Banach space X with relative boundary  $\partial Q, S$  be a closed subset of a Banach space Y and G be a subset of  $C(\partial Q, Y \setminus S)$ . S and  $\partial Q$  are G-linking if for any map  $h \in C(Q, Y)$  such that  $h|_{\partial Q} \in G$  there holds  $h(Q) \cap S \neq \emptyset$ .

**Theorem 2.3** ([3, 9]). Let X, Y be Banach spaces, S be a closed subset of Y, Q be a submanifold of X with relative boundary  $\partial Q$  and G be a subset of  $C(\partial Q, Y \setminus S)$ . Let  $\Gamma = \{h \in C(Q, Y) : h|_{\partial Q} \in G\}$ , assume that S and  $\partial Q$  are G-linking and  $f \in C^1(Y, R)$  satisfies

- (a) There is  $\tilde{h} \in \Gamma$  such that  $\sup_{x \in Q} f(\tilde{h}(x)) < +\infty$ ;
- (b) There is  $\beta_0 > \alpha_0$  such that

$$\inf_{y\in S} f(y) \geq \beta_0 \quad and \quad \sup_{x\in \partial Q} f(h(x)) \leq \alpha_0, \ \forall h\in \Gamma;$$

(c) The (PS) condition holds.

Then, the number

$$c = \inf_{h \in \Gamma} \sup_{x \in Q} f(h(x))$$

is a critical value of f with  $c \geq \beta_0$ .

*Proof.* The proof is divided into two steps.

**Step 1.** The (PS) condition for the functional J is satisfied. Let  $(u_n, v_n)$  be a (PS) sequence for the functional J; that is,

$$J(u_n, v_n)$$
 is bound and  $J'(u_n, v_n) \to 0$  as  $n \to \infty$ . (2.2)

From (H1) and by a standard argument, it is sufficient to prove that  $(u_n, v_n)$  is bounded in W. If this does not hold, assume that  $||(u_n, v_n)|| = ||u_n||_p + ||v_n||_q \to \infty$ as  $n \to \infty$ . Define  $K_n := ||u_n||_p^p + ||v_n||_q^q$ , hence it follows that  $K_n \to \infty$  as  $n \to \infty$ . Let  $\bar{u}_n = u_n \setminus K_n^{1/p}, \bar{v}_n = v_n \setminus K_n^{1/q}$ , then  $(\bar{u}_n, \bar{v}_n)$  is bounded in W, i.e.,

$$\|\bar{u}_n\|_p^p + \|\bar{v}_n\|_q^q = 1$$
 for all  $n$ .

Extracting subsequences if necessary, we can assume that there exists  $(\bar{u}, \bar{v}) \in W$  such that

$$(\bar{u}_n, \bar{v}_n) \rightharpoonup (\bar{u}, \bar{v})$$
 weakly in  $W$ , (2.3)

$$(\bar{u}_n, \bar{v}_n) \to (\bar{u}, \bar{v})$$
 strongly in  $L^p(\Omega) \times L^q(\Omega)$ , (2.4)

$$(\bar{u}_n(x), \bar{v}_n(x)) \to (\bar{u}(x), \bar{v}(x))$$
 for a.e.  $x \in \Omega$ . (2.5)

From (2.2), it follows that

$$\limsup_{n \to \infty} \frac{J(u_n, v_n)}{K_n} \le 0, \tag{2.6}$$

From (2.1), (H1), the Hölder's inequality and (1.3), we have

$$\begin{split} \left| \int_{\Omega} G(x, u, v) dx \right| &\leq \int_{\Omega} \left| \int_{0}^{1} (G_{s}(x, \tau u, \tau v) u + G_{t}(x, \tau u, \tau v) v) d\tau \right| dx \\ &\leq \int_{\Omega} h(x) (|u| + |v|) dx \\ &\leq \|h\|_{L^{\infty}} (|\Omega|^{\frac{p-1}{p}} \|u\|_{L^{p}} + |\Omega|^{\frac{q-1}{q}} \|v\|_{L^{q}}) \\ &\leq C_{1}(\|u\|_{p} + \|v\|_{q}) \end{split}$$
(2.7)

for all  $(u, v) \in W$ , where  $C_1$  is a positive constant, hence it follows that

$$\frac{1}{K_n} \int_{\Omega} G(x, u_n, v_n) dx \to 0 \quad \text{as } n \to \infty.$$
(2.8)

From  $h_1 \in L^{p/(p-1)}(\Omega), h_2 \in L^{q/(q-1)}(\Omega)$  and the Hölder's inequality, we obtain

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$$\frac{1}{K_n} \int_{\Omega} (h_1 u_n + h_2 v_n) dx \to 0 \quad \text{as } n \to \infty.$$
(2.9)

From (2.4) and (2.5), it follows that  $|\bar{u}_n|^{\alpha}\bar{u}_n \to |\bar{u}|^{\alpha}\bar{u}$  strongly in  $L^{\frac{p}{\alpha+1}}(\Omega)$  and  $|\bar{v}_n|^{\beta}\bar{v}_n \to |\bar{v}|^{\alpha}\bar{v}$  strongly in  $L^{\frac{q}{\beta+1}}(\Omega)$ , hence from Hölder's inequality, we obtain

$$\begin{split} \left| \int_{\Omega} (|\bar{u}_{n}|^{\alpha} |\bar{v}_{n}|^{\beta} \bar{u}_{n} \bar{v}_{n} - |\bar{u}|^{\alpha} |\bar{v}|^{\beta} \bar{u} \bar{v}) dx \right| \\ &\leq \int_{\Omega} ||\bar{u}_{n}|^{\alpha} |\bar{v}_{n}|^{\beta} \bar{u}_{n} \bar{v}_{n} - |\bar{u}_{n}|^{\alpha} |\bar{v}|^{\beta} \bar{u}_{n} \bar{v}| dx + \int_{\Omega} ||\bar{u}_{n}|^{\alpha} |\bar{v}|^{\beta} \bar{u}_{n} \bar{v} - |\bar{u}|^{\alpha} |\bar{v}|^{\beta} \bar{u} \bar{v}| dx \\ &\leq \int_{\Omega} |\bar{u}_{n}|^{\alpha+1} \cdot ||\bar{v}_{n}|^{\beta} \bar{v}_{n} - |\bar{v}|^{\beta} \bar{v}| dx + \int_{\Omega} ||\bar{u}_{n}|^{\alpha} \bar{u}_{n} - |\bar{u}|^{\alpha} \bar{u}| \cdot |\bar{v}|^{\beta+1} dx \\ &\leq ||\bar{u}_{n}||_{L^{p}}^{\alpha+1} \cdot ||\bar{v}_{n}|^{\beta} \bar{v}_{n} - |\bar{v}|^{\beta} \bar{v}||_{L^{\frac{q}{\beta+1}}} + ||\bar{v}_{n}||_{L^{q}}^{\beta+1} \cdot ||\bar{u}_{n}|^{\alpha} \bar{u}_{n} - |\bar{u}|^{\alpha} \bar{u}||_{L^{\frac{p}{\alpha+1}}} \\ &\rightarrow 0 \quad \text{as } n \to \infty. \end{split}$$

$$(2.10)$$

From the definition of J, (2.4), (2.6), (2.8), (2.9) and (2.10), we have

$$\begin{split} &\limsup_{n \to \infty} \left( \frac{1}{p} \int_{\Omega} |\nabla \bar{u}_n|^p dx + \frac{1}{q} \int_{\Omega} |\nabla \bar{v}_n|^q dx \right) \\ &\leq \lambda_1 \left( \frac{1}{p} \int_{\Omega} |\bar{u}|^p dx + \frac{1}{q} \int_{\Omega} |\bar{v}|^q dx + \frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega} |\bar{u}|^{\alpha} |\bar{v}|^{\beta} \bar{u} \bar{v} dx \right). \end{split}$$

From (2.3), it follows that

$$\int_{\Omega} |\nabla \bar{u}|^p dx \leq \liminf_{n \to \infty} \int_{\Omega} |\nabla \bar{u}_n|^p dx, \quad \int_{\Omega} |\nabla \bar{v}|^q dx \leq \liminf_{n \to \infty} \int_{\Omega} |\nabla \bar{v}_n|^q dx,$$

hence, combining this with the definition of  $\lambda_1$ , we obtain

$$\begin{split} \lambda_1 \Big( \frac{1}{p} \int_{\Omega} |\bar{u}|^p dx + \frac{1}{q} \int_{\Omega} |\bar{v}|^q dx + \frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega} |\bar{u}|^{\alpha} |\bar{v}|^{\beta} \bar{u} \bar{v} dx \Big) \\ &\leq \frac{1}{p} \int_{\Omega} |\nabla \bar{u}|^p dx + \frac{1}{q} \int_{\Omega} |\nabla \bar{v}|^q dx \\ &\leq \liminf_{n \to \infty} \Big( \frac{1}{p} \int_{\Omega} |\nabla \bar{u}_n|^p dx + \frac{1}{q} \int_{\Omega} |\nabla \bar{v}_n|^q dx \Big) \\ &\leq \limsup_{n \to \infty} \Big( \frac{1}{p} \int_{\Omega} |\nabla \bar{u}_n|^p dx + \frac{1}{q} \int_{\Omega} |\nabla \bar{v}_n|^q dx \Big) \\ &\leq \lambda_1 \Big( \frac{1}{p} \int_{\Omega} |\bar{u}|^p dx + \frac{1}{q} \int_{\Omega} |\bar{v}|^q dx + \frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega} |\bar{u}|^{\alpha} |\bar{v}|^{\beta} \bar{u} \bar{v} dx \Big). \end{split}$$

hence it follows that

$$\begin{split} &\frac{1}{p} \int_{\Omega} |\nabla \bar{u}|^p dx + \frac{1}{q} \int_{\Omega} |\nabla \bar{v}|^q dx \\ &= \lambda_1 \Big( \frac{1}{p} \int_{\Omega} |\bar{u}|^p dx + \frac{1}{q} \int_{\Omega} |\bar{v}|^q dx + \frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega} |\bar{u}|^{\alpha} |\bar{v}|^{\beta} \bar{u} \bar{v} dx \Big), \end{split}$$

and by the uniform convexity of W, we have that  $(\bar{u}_n, \bar{v}_n)$  converges strongly to  $(\bar{u}, \bar{v})$  in W, and from the definition of  $(\mu_1, \nu_1)$ , it follows that  $(\bar{u}, \bar{v}) = \pm (\mu_1, \nu_1)$ .

In the following, let  $(\bar{u}, \bar{v}) = (\mu_1, \nu_1)$ , and the other case where  $(\bar{u}, \bar{v}) = -(\mu_1, \nu_1)$  may be considered similarly. Hence from the definition of J, we have

$$\frac{pJ_1(u_n, v_n)}{(p-1)K_n^{1/p}} + \frac{qJ_2(u_n, v_n)}{(q-1)K_n^{1/q}} - \langle J'(u_n, v_n), (\frac{\bar{u}_n}{p-1}, \frac{\bar{v}_n}{q-1}) \rangle \\
= \frac{1}{p-1} \Big( \int_{\Omega} G_s(x, u_n, v_n) \bar{u}_n dx - \frac{p}{K_n^{1/p}} \int_{\Omega} \int_0^1 G_s(x, ru_n, rv_n) u_n dr dx \Big) \\
+ \frac{1}{q-1} \Big( \int_{\Omega} G_t(x, u_n, v_n) \bar{v}_n dx - \frac{q}{K_n^{1/q}} \int_{\Omega} \int_0^1 G_t(x, ru_n, rv_n) v_n dr dx \Big) \\
+ \int_{\Omega} h_1 \bar{u}_n dx + \int_{\Omega} h_2 \bar{v}_n dx.$$
(2.11)

From  $h_1 \in L^{p/(p-1)}(\Omega), h_2 \in L^{q/(q-1)}(\Omega)$ , we have

$$\int_{\Omega} h_1 \bar{u}_n dx \to \int_{\Omega} h_1 \mu_1 dx \quad \text{and} \quad \int_{\Omega} h_2 \bar{v}_n dx \to \int_{\Omega} h_2 \nu_1 dx \quad \text{as} \quad n \to \infty.$$
(2.12)

From (H2) and (H3), it is easy to know that

$$\int_{\Omega} G_s(x, u_n, v_n) \bar{u}_n dx \to \int_{\Omega} g_1^{++} \mu_1 dx,$$

$$\int_{\Omega} G_t(x, u_n, v_n) \bar{v}_n dx \to \int_{\Omega} g_2^{++} \nu_1 dx$$
(2.13)

as  $n \to \infty.$  Finally, from (H2) and Lebesgue dominated convergence theorem, we have

$$\frac{1}{K_n^{1/p}} \int_{\Omega} \int_0^1 G_s(x, ru_n, rv_n) u_n \, dr \, dx = \int_{\Omega} \int_0^1 G_s(x, ru_n, rv_n) \frac{u_n}{K_n^{1/p}} \, dr \, dx$$
$$\to \int_{\Omega} g_1^{++} \mu_1 dx \quad \text{as } n \to \infty.$$
(2.14)

Similarly, we obtain

$$\frac{1}{K_n^{1/q}} \int_{\Omega} \int_0^1 G_t(x, ru_n, rv_n) v_n \, dr \, dx \to \int_{\Omega} g_2^{++} \nu_1 dx \quad \text{as } n \to \infty.$$
(2.15)

Therefore, letting  $n \to \infty$  in (2.11) and from (2.2), (2.12), (2.13), (2.14) and (2.15), we obtain

$$\int_{\Omega} h_1 \mu_1 dx + \int_{\Omega} h_2 \nu_1 dx = \int_{\Omega} g_1^{++} \mu_1 dx + \int_{\Omega} g_2^{++} \nu_1 dx,$$

which contradicts with (1.5). Hence,  $(u_n, v_n)$  is bounded in W.

Step 2. The functional J satisfies the geometries of Theorem 2.3. For any

$$(u,v) \in E_1 = \{(t^{1/p}\mu_1, t^{1/q}\nu_1) : t \ge 0\} \cup \{(-t^{1/p}\mu_1, -t^{1/q}\nu_1) : t \ge 0\},\$$

we have

$$\begin{split} &\frac{1}{p} \int_{\Omega} |\nabla u|^p dx + \frac{1}{q} \int_{\Omega} |\nabla v|^q dx \\ &= \lambda_1 \Big( \frac{1}{p} \int_{\Omega} |u|^p dx + \frac{1}{q} \int_{\Omega} |v|^q dx + \frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega} |u|^{\alpha} |v|^{\beta} uv \, dx \Big). \end{split}$$

From the above equality and the definition of J, for any  $(t^{1/p}\mu_1, t^{1/q}\nu_1) \in E_1$ , we obtain

$$J(t^{1/p}\mu_{1}, t^{1/q}\nu_{1})$$

$$= t^{1/p} \int_{\Omega} h_{1}\mu_{1}dx + t^{1/q} \int_{\Omega} h_{2}\nu_{1}dx - \int_{\Omega} G(x, t^{1/p}\mu_{1}, t^{1/q}\nu_{1})dx$$

$$= t^{1/p} \Big( \int_{\Omega} h_{1}\mu_{1}dx - \int_{\Omega} \int_{0}^{1} G_{s}(x, rt^{1/p}\mu_{1}, rt^{1/q}\nu_{1})\mu_{1} dr dx \Big)$$

$$+ t^{1/q} \Big( \int_{\Omega} h_{2}\nu_{1}dx - \int_{\Omega} \int_{0}^{1} G_{t}(x, rt^{1/p}\mu_{1}, rt^{1/q}\nu_{1})\nu_{1} dr dx \Big).$$
(2.16)

From (H1), (H2) and Lebesgue dominated convergence theorem, it follows that

$$\lim_{t \to +\infty} \int_{\Omega} \int_{0}^{1} G_s(x, rt^{1/p} \mu_1, rt^{1/q} \nu_1) \mu_1 \, dr \, dx = \int_{\Omega} g_1^{++} \mu_1 dx, \qquad (2.17)$$

$$\lim_{t \to +\infty} \int_{\Omega} \int_{0}^{1} G_t(x, rt^{1/q} \mu_1, rt^{1/q} \nu_1) \nu_1 \, dr \, dx = \int_{\Omega} g_2^{++} \nu_1 dx.$$
(2.18)

Hence, from (1.7), (2.16), (2.17) and (2.18), we obtain

$$J(t^{1/p}\mu_1, t^{1/q}\nu_1) \to -\infty \text{ as } t \to +\infty.$$

Similarly, the following result can be obtained with  $g_1^{++}$  and  $g_2^{++}$  exchanged with  $g_1^{--}$  and  $g_2^{--}$  respectively,

$$J(-t^{1/p}\mu_1, -t^{1/q}\nu_1) \to -\infty \text{ as } t \to +\infty.$$

Finally, it follows that

$$\lim_{t \to \infty} J(\pm t^{1/p} \mu_1, \pm t^{1/q} \nu_1) = -\infty.$$
(2.19)

On the other hand, letting  $\Lambda_2 := \{(u, v) \in W : \phi(u, v) \ge \lambda_2 \varphi(u, v)\}$ , from (1.3), (2.7) and the Hölder's inequality, for any  $(u, v) \in \Lambda_2$ , we obtain

$$J(u,v) \ge \frac{\lambda_2 - \lambda_1}{p\lambda_2} \|u\|_p^p + \frac{\lambda_2 - \lambda_1}{q\lambda_2} \|v\|_q^q - C_1(\|u\|_p + \|v\|_q) - (\|h_1\|_{L^{\frac{p}{p-1}}} \|u\|_{L^p} + \|h_2\|_{L^{\frac{q}{q-1}}} \|v\|_{L^q}) \ge \frac{\lambda_2 - \lambda_1}{p\lambda_2} (\|u\|_p^p + \|v\|_q^q) - C_2(\|u\|_p + \|v\|_q),$$

where  $C_2 = C_1 + C \max\{\|h_1\|_{L^{\frac{p}{p-1}}}, \|h_2\|_{L^{\frac{q}{q-1}}}\}$ . Combining the above expression with (2.19), we obtain that there exists a positive constant T such that

$$\alpha_0 := \sup_{t \ge T} J(\pm t^{1/p} \mu_1, \pm t^{1/q} \nu_1) < \beta_0 := \inf_{(u,v) \in \Lambda_2} J(u,v).$$
(2.20)

Let  $M = \{(\pm t^{1/p}\mu_1, \pm t^{1/q}\nu_1) : t \ge T\}$  and  $G = \{h \in C(S^0, W) : h \text{ is odd and } h(S^0) \subset M\},$ 

where  $S^0$  is the boundary of the closed unit ball  $B^1$  in  $\mathbb{R}^1$ , i.e.,  $S^0 = \partial B^1$ . For any  $h \in G$ , by (2.20), we have  $h(S^0) \cap \Lambda_2 = \emptyset$ , which implies that G is a subset of  $C(S^0, W \setminus \Lambda_2)$ . Let

$$\Gamma = \{ h \in C(B^1, W) : h |_{S^0} \in G \},\$$

we can claim:  $\Gamma$  is nonempty and  $\Lambda_2$  and  $S^0$  are G-linking, that is  $h(B^1) \cap \Lambda_2 \neq \emptyset$ for any  $h \in \Gamma$ . The similar proof of the conclusion may be found in [7, 3, 9], but for the readers convenience and completeness, we write it.

In fact, define  $\bar{h}: B^1 \to W$  by

$$\bar{h}(t) = ((tT)^{1/p}\mu_1, (tT)^{1/q}\nu_1) \text{ for all } t \in [0, 1],$$
  
$$\bar{h}(-t) = (-(tT)^{1/p}\mu_1, -(tT)^{1/q}\nu_1) \text{ for all } t \in [0, 1].$$

Hence,  $\bar{h} \in \Gamma$  and  $\Gamma$  is nonempty. Now let  $h \in \Gamma$ , if there is  $(u, v) \in h(B^1)$  such that  $\varphi(u, v) = 0$ , we get  $h(B^1) \cap \Lambda_2 \neq \emptyset$ . If not, we consider the map  $\hat{h} : S^1 \to \Sigma$  defined by

$$\hat{h}(x_1, x_2) = \begin{cases} \pi \circ h(x_1), & \text{if } x_2 \ge 0, \\ -\pi \circ h(-x_1), & \text{if } x_2 \le 0, \end{cases}$$

where  $\pi(u, v) = (u \setminus (\varphi(u, v))^{1/p}, v \setminus (\varphi(u, v))^{1/q})$ . It is easy to know that  $\hat{h}(S^1) \subset \Sigma_2$ . Therefore,  $\phi(u_0, v_0) \geq \lambda_2$  for some  $(u_0, v_0) \in \hat{h}(S^1)$ , namely,  $(u_0, v_0) \in \Lambda_2$ . From  $\pi \circ h(x) \in \Lambda_2$ , we have implies  $h(x) \in \Lambda_2$ , which implies that  $h(B^1) \cap \Lambda_2 \neq \emptyset$ . Hence  $\Lambda_2$  and  $S^0$  are G-linking.

Now, from the compactness of  $B^1$ , (a) of Theorem 2.3 holds, (b) of Theorem 2.3 is satisfied from (2.20), (c) of Theorem 2.3 comes from (i). Accordingly, Theorem 1.1 holds from the G-linking Theorem with the critical value

$$c = \inf_{h \in \Gamma} \sup_{x \in B^1} J(h(x)). \quad \Box$$

*Proof of Theorem 1.3.* (i) The functional J satisfies the (PS) condition. From (1.6), the claim can be proved with similar to step 1 of Theorem 1.1.

(ii) Now we will prove that the functional J is coercive, that is,

$$J(u,v) \to +\infty$$
 as  $||(u,v)|| \to \infty$ .

If the claim does not hold, there is a constant c and a sequence  $(u_n, v_n)$  such that  $J(u_n, v_n) \leq c$  and  $||(u_n, v_n)|| \to \infty$ . From the proof of the (PS) condition of Theorem 1.1,  $(\bar{u}_n, \bar{v}_n)$  converges strongly to  $\pm(\mu_1, \nu_1)$ , where  $\bar{u}_n = u_n \setminus K_n^{1/p}, \bar{v}_n = v_n \setminus K_n^{1/q}$ . Assume that  $(\bar{u}_n, \bar{v}_n)$  converges strongly to  $(\mu_1, \nu_1)$  (the case  $(\bar{u}_n, \bar{v}_n)$  converges strongly to  $(-\mu_1, -\nu_1)$  may be treated similarly) and  $p \geq q > 1$  (the case  $q \geq p > 1$  may also be treated similarly). From the definitions of  $J, J_1, J_2$  and  $J(u_n, v_n) \leq c$  for all n, (2.14) and (2.15), we have

$$\begin{split} 0 &\geq \limsup_{n \to \infty} \frac{J(u_n, v_n)}{K_n^{1/p}} = \limsup_{n \to \infty} \left( \frac{J_1(u_n, v_n)}{K_n^{1/p}} + \frac{J_2(u_n, v_n)}{K_n^{1/p}} \right) \\ &\geq \limsup_{n \to \infty} \left( \frac{J_1(u_n, v_n)}{K_n^{1/p}} + \frac{J_2(u_n, v_n)}{K_n^{1/q}} \right) \\ &\geq \lim_{n \to \infty} \left( \int_{\Omega} (h_1 \bar{u}_n - \frac{1}{K_n^{1/p}} \int_0^1 G_s(x, ru_n, rv_n) u_n dr) dx \right) \\ &+ \int_{\Omega} (h_2 \bar{v}_n - \frac{1}{K_n^{1/q}} \int_0^1 G_t(x, ru_n, rv_n) v_n dr) dx \Big) \\ &= \int_{\Omega} (h_1 \mu_1 - g_1^{++} \mu_1) dx + \int_{\Omega} (h_2 \nu_1 - g_2^{++} \nu_1) dx, \end{split}$$

which is a contradiction to (1.6). By Ekeland's Variational Principle, the proof is complete.  $\hfill \Box$ 

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