Electronic Journal of Differential Equations, Vol. 2016 (2016), No. 163, pp. 1-10.
ISSN: 1072-6691. URL: http://ejde.math.txstate.edu or http://ejde.math.unt.edu

## $(p, q)$-LAPLACIAN ELLIPTIC SYSTEMS AT RESONANCE

## ZENG-QI OU


#### Abstract

We show the existence of weak solutions for a class of $(p, q)$ Laplacian elliptic systems at resonance, under certain Landesman-Lazer-type conditions by using critical point theorem.


## 1. Introduction and statement of main results

Let $\Omega$ be a bounded domain with smooth boundary in $\mathbb{R}^{N}$ and $\Delta_{p}$ be the $p$ Laplacian operator. In this paper, we study the existence of solutions for the problem

$$
\begin{gather*}
-\Delta_{p} u=\lambda_{1}|u|^{p-2} u+\frac{\lambda_{1}}{\beta+1}|u|^{\alpha}|v|^{\beta} v+G_{s}(x, u, v)-h_{1}(x) \quad \text { in } \Omega \\
-\Delta_{q} v=\lambda_{1}|v|^{q-2} v+\frac{\lambda_{1}}{\alpha+1}|u|^{\alpha}|v|^{\beta} u+G_{t}(x, u, v)-h_{2}(x) \quad \text { in } \Omega  \tag{1.1}\\
u=v=0 \quad \text { on } \partial \Omega
\end{gather*}
$$

where $1<p, q<+\infty$ and $\alpha \geq 0, \beta \geq 0$ satisfy

$$
\begin{equation*}
\frac{\alpha+1}{p}+\frac{\beta+1}{q}=1 . \tag{1.2}
\end{equation*}
$$

The nonlinearity $G: \Omega \times \mathbb{R}^{2} \rightarrow R$ is a Caratheodory function which has continuous derivatives $G_{s}(x, s, t), G_{t}(x, s, t)$ with respect to $s$ and $t$ for almost any $x \in \Omega$, and $h_{1} \in L^{p /(p-1)}(\Omega), h_{2} \in L^{q /(q-1)}(\Omega)$.

Let $W=W_{0}^{1, p}(\Omega) \times W_{0}^{1, q}(\Omega)$ with the norm $\|(u, v)\|=\|u\|_{p}+\|v\|_{q}$ for all $(u, v) \in W$, where $W_{0}^{1, p}(\Omega)$ is the usual Banach space with the norm $\|u\|_{p}=$ $\left(\int_{\Omega}|\nabla u|^{p} d x\right)^{1 / p}$ for any $u \in W_{0}^{1, p}(\Omega)$. From Sobolev embedding Theorem, the embedding $W_{0}^{1, p}(\Omega) \hookrightarrow L^{p}(\Omega)$ is continuous and compact, and there is constant $C>0$ such that

$$
\begin{equation*}
\|u\|_{L^{p}} \leq C\|u\|_{p}, \forall u \in W_{0}^{1, p}(\Omega), \quad \text { and } \quad\|v\|_{L^{q}} \leq C\|v\|_{q}, \forall v \in W_{0}^{1, q}(\Omega) \tag{1.3}
\end{equation*}
$$

where $\|\cdot\|_{L^{p}}$ denotes the norm of $L^{p}(\Omega)$ and throughout this paper, let $C$ always denote an embedding constant with relation to 1.3 . For the following nonlinear

[^0]eigenvalue problem
\[

$$
\begin{gather*}
-\Delta_{p} u=\lambda|u|^{p-2} u+\frac{\lambda}{\beta+1}|u|^{\alpha}|v|^{\beta} v \quad \text { in } \Omega \\
-\Delta_{q} v=\lambda|v|^{q-2} v+\frac{\lambda}{\alpha+1}|u|^{\alpha}|v|^{\beta} u \quad \text { in } \Omega  \tag{1.4}\\
u=v=0 \quad \text { on } \partial \Omega
\end{gather*}
$$
\]

consider the functionals $\phi, \varphi$ on $W$ defined by

$$
\begin{gathered}
\phi(u, v)=\frac{1}{p} \int_{\Omega}|\nabla u|^{p} d x+\frac{1}{q} \int_{\Omega}|\nabla v|^{q} d x \\
\varphi(u, v)=\frac{1}{p} \int_{\Omega}|u|^{p} d x+\frac{1}{q} \int_{\Omega}|v|^{q} d x+\frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega}|u|^{\alpha}|v|^{\beta} u v d x
\end{gathered}
$$

and the manifold

$$
\Sigma=\{(u, v) \in W: \varphi(u, v)=1\}
$$

It is easy to prove that $\phi(u, v), \varphi(u, v)$ are $(p, q)$-homogeneous, namely

$$
\phi\left(t^{1 / p} u, t^{1 / q} v\right)=t \phi(u, v), \quad \varphi\left(t^{1 / p} u, t^{1 / q} v\right)=t \varphi(u, v)
$$

for any $t>0$ and $(u, v) \in W$, and $\Sigma$ is a symmetric nonempty manifold in $W$. By an argument similar to the ones in (3, 7, problem $\sqrt[1.4]{ }$ has a sequence of eigenvalues with the variational characterization

$$
\lambda_{k}=\inf _{\Lambda \in \Sigma_{k}} \sup _{(u, v) \in \Lambda} \phi(u, v)
$$

where $\Sigma_{k}=\left\{\Lambda \subset \Sigma:\right.$ there is an odd, continuous and surjective $\left.\gamma: S^{k-1} \rightarrow \Lambda\right\}$ and $S^{k-1}$ denotes the unit sphere in $\mathbb{R}^{k}$.

On the other hand, let

$$
\lambda_{1}^{\prime}=\inf _{(u, v) \in \Sigma} \phi(u, v),
$$

we can see that $\lambda_{1}=\lambda_{1}^{\prime}$. Moreover, $\lambda_{1}$ is a simple, isolated and positive principal eigenvalue of $\sqrt{1.4}$ and has a positive normalized eigenvalue ( $\mu_{0}, \nu_{0}$ ), namely, $\left\|\mu_{0}\right\|_{p}+\left\|\nu_{0}\right\|_{q}=1$. By a simple computation, there exists a positive constant $t_{0}$ such that

$$
\left\|t_{0}^{1 / p} \mu_{0}\right\|_{p}^{p}+\left\|t_{0}^{1 / q} \nu_{0}\right\|_{q}^{q}=1
$$

Let $\mu_{1}=t_{0}^{1 / p} \mu_{0}, \nu_{1}=t_{0}^{1 / q} \nu_{0}$, since $\phi, \varphi$ are $(p, q)$-homogeneous, hence the set of all eigenfunctions corresponding to $\lambda_{1}$ is

$$
E_{1}:=\left\{\left(t^{1 / p} \mu_{1}, t^{1 / q} \nu_{1}\right): t \geq 0\right\} \cup\left\{\left(-t^{1 / p} \mu_{1},-t^{1 / q} \nu_{1}\right): t \geq 0\right\}
$$

The set $E_{1}$ is not an one-dimensional linear subspace of $W$ and the corresponding orthogonal decomposition on $W$ does not hold with respect to the the first eigenvalue $\lambda_{1}$.

In many papers, existence of weak solutions for the resonant elliptic problems were investigated under the well-known Landesman-Lazer-type conditions, which were introduced by Landesman and Lazer in [5] and were extended by Tang in [12]. Since then they were used widely for the different types of equations, for example, in [1, 3, 9] for the quasilinear elliptic equations, in 4] for asymptotically linear noncooperative elliptic systems, in [13] for the forced duffing equations, in [11] for Kirchhoff type equations. Especially, in [2] the case $p=q=2$ (the semilinear elliptic systems) was considered and the case $p=q \geq 2$ (the quasilinear elliptic systems)
was discussed in [6, 7, 14] where $G_{s}(x, s, t)=g_{1}(x, s)$ and $G_{t}(x, s, t)=g_{2}(x, t)$. As far as we know, when $p \neq q>1$, the similar results are not discussed under the Landesman-Lazer-type conditions due to Landesman and Lazer. Motivated by these finding, we consider the existence of solutions for problem (1.1) at resonance with the first eigenvalue under the Landesman-Lazer-type conditions. We first state the following fundamental hypotheses.
(H1) There is $h \in C\left(\bar{\Omega}, \mathbb{R}^{+}\right)$such that $\left|G_{s}(x, s, t)\right| \leq h(x)$ and $\left|G_{t}(x, s, t)\right| \leq h(x)$ for all $(x, s, t) \in \Omega \times \mathbb{R}^{2}$.
(H2) There exist two functions $g_{1}^{++}, g_{1}^{--} \in C(\Omega, R)$ such that

$$
g_{1}^{++}(x)=\liminf _{s \rightarrow+\infty, t \rightarrow+\infty} G_{s}(x, s, t), \quad g_{1}^{--}(x)=\limsup _{s \rightarrow-\infty, t \rightarrow-\infty} G_{s}(x, s, t)
$$

uniformly a.e. $x \in \Omega$.
(H3) There is two functions $g_{2}^{++}, g_{2}^{--} \in C(\Omega, R)$ such that

$$
g_{2}^{++}(x)=\liminf _{s \rightarrow+\infty, t \rightarrow+\infty} G_{t}(x, s, t), \quad g_{2}^{--}(x)=\limsup _{s \rightarrow-\infty, t \rightarrow-\infty} G_{t}(x, s, t)
$$

uniformly a.e. $x \in \Omega$.
The Landesman-Lazer-type conditions for problem 1.1 are read either

$$
\begin{align*}
\int_{\Omega} g_{1}^{--} \mu_{1} d x+\int_{\Omega} g_{2}^{--} \nu_{1} d x & <\int_{\Omega} h_{1} \mu_{1} d x+\int_{\Omega} h_{2} \nu_{1} d x  \tag{1.5}\\
& <\int_{\Omega} g_{1}^{++} \mu_{1} d x+\int_{\Omega} g_{2}^{++} \nu_{1} d x
\end{align*}
$$

or

$$
\begin{align*}
\int_{\Omega} g_{1}^{++} \mu_{1} d x+\int_{\Omega} g_{2}^{++} \nu_{1} d x & <\int_{\Omega} h_{1} \mu_{1} d x+\int_{\Omega} h_{2} \nu_{1} d x \\
& <\int_{\Omega} g_{1}^{--} \mu_{1} d x+\int_{\Omega} g_{2}^{--} \nu_{1} d x \tag{1.6}
\end{align*}
$$

We are ready to state the main results.
Theorem 1.1. Let $h_{1} \in L^{p /(p-1)}(\Omega), h_{2} \in L^{q /(q-1)}(\Omega)$, and 1.2), (H1), (H2), (H3) and (1.5) be satisfied. If $1<p<q$ and the following inequalities hold:

$$
\begin{equation*}
\int_{\Omega} h_{1} \mu_{1} d x-\int_{\Omega} g_{1}^{++} \mu_{1} d x<0, \quad \int_{\Omega} h_{1} \mu_{1} d x-\int_{\Omega} g_{1}^{--} \mu_{1} d x>0 \tag{1.7}
\end{equation*}
$$

then problem (1.1) has at least one solution.
In the other case $1<q<p$, the following result holds.
Theorem 1.2. Let $h_{1} \in L^{p /(p-1)}(\Omega), h_{2} \in L^{q /(q-1)}(\Omega)$, and 1.2 , (H1), (H2), (H3) and 1.5 be satisfied. If $1<q<p$ and the following inequalities hold:

$$
\begin{equation*}
\int_{\Omega} h_{2} \nu_{1} d x-\int_{\Omega} g_{2}^{++} \nu_{1} d x<0, \quad \int_{\Omega} h_{2} \nu_{1} d x-\int_{\Omega} g_{2}^{--} \nu_{1} d x>0 \tag{1.8}
\end{equation*}
$$

then problem (1.1) has at least one solution.
Theorem 1.3. Let $h_{1} \in L^{p /(p-1)}(\Omega), h_{2} \in L^{q /(q-1)}(\Omega)$. If 1.2$)$, (H1), (H2), and (1.6) are satisfied, then problem (1.1) has at least one solution.

Our results extends the ones in [2] from the semilinear elliptic systems to $(p, q)$ Laplacian elliptic systems, and are also the generalizations of [14], where they considered the case $p=q \geq 2$ and $G_{s}(x, s, t)=g_{1}(s), G_{t}(x, s, t)=g_{2}(t)$. Moreover, the conditions 1.7 and 1.8 are the technical assumptions. Theorem 1.2 is similar to Theorem 1.1. and we will prove Theorem 1.1 and Theorem 1.3 .

## 2. Proofs of Theorems

Now consider the functionals $J, J_{1}, J_{2}$ on $W$ defined by

$$
\begin{aligned}
& J(u, v)= \frac{1}{p} \int_{\Omega}|\nabla u|^{p} d x+\frac{1}{q} \int_{\Omega}|\nabla v|^{q} d x-\frac{\lambda_{1}}{p} \int_{\Omega}|u|^{p} d x \\
&-\frac{\lambda_{1}}{q} \int_{\Omega}|v|^{q} d x-\frac{\lambda_{1}}{(\alpha+1)(\beta+1)} \int_{\Omega}|u|^{\alpha}|v|^{\beta} u v d x \\
&-\int_{\Omega} G(x, u, v) d x+\int_{\Omega} h_{1} u d x+\int_{\Omega} h_{2} v d x \\
& J_{1}(u, v)=\frac{1}{p} \int_{\Omega}|\nabla u|^{p} d x-\frac{\lambda_{1}}{p} \int_{\Omega}|u|^{p} d x-\frac{\lambda_{1}}{p(\beta+1)} \int_{\Omega}|u|^{\alpha}|v|^{\beta} u v d x \\
&- \int_{\Omega} \int_{0}^{1} G_{s}(x, r u, r v) u d r d x+\int_{\Omega} h_{1} u d x \\
& J_{2}(u, v)= \frac{1}{q} \int_{\Omega}|\nabla v|^{q} d x-\frac{\lambda_{1}}{q} \int_{\Omega}|v|^{q} d x-\frac{\lambda_{1}}{q(\alpha+1)} \int_{\Omega}|u|^{\alpha}|v|^{\beta} u v d x \\
&- \int_{\Omega} \int_{0}^{1} G_{t}(x, r u, r v) v d r d x+\int_{\Omega} h_{2} v d x .
\end{aligned}
$$

Noting that

$$
\begin{equation*}
G(x, s, t)=\int_{0}^{1} G_{s}(x, r s, r t) s d r+\int_{0}^{1} G_{t}(x, r s, r t) t d r \tag{2.1}
\end{equation*}
$$

from $\sqrt{1.2}$ ) and (2.1), it follows that

$$
J(u, v)=J_{1}(u, v)+J_{2}(u, v) \quad \text { for all }(u, v) \in W
$$

From (H1), it is easy to prove that the functional $J$ is well defined and $J \in$ $C^{1}(W, R)$. Moreover, from the variational view of point, a weak solution of problem 1.1) is equivalent to a critical point of the functional $J$ in $W$. In this paper, we will prove Theorem 1.1 and Theorem 1.2 by using the following G-linking Theorem due to Drábek and Robinson (see [3, 9]) and Theorem 1.3 by using Ekeland's Variational Principle (see [8, 10]). In these abstract theorems, a compact condition, i.e., $(P S)$ condition, is needed.

Definition 2.1. Let $X$ be a real Banach space, if for any sequence $\left\{u_{n}\right\} \subset X$ such that $f\left(u_{n}\right)$ is bounded and $f^{\prime}\left(u_{n}\right) \rightarrow 0$ as $n \rightarrow \infty,\left\{u_{n}\right\}$ has a convergent subsequence, the functional $f$ satisfies the $(P S)$ condition.

Definition 2.2 ( 3,9 ). Let $Q$ be a submanifold of a Banach space $X$ with relative boundary $\partial Q, S$ be a closed subset of a Banach space $Y$ and $G$ be a subset of $C(\partial Q, Y \backslash S) . \quad S$ and $\partial Q$ are G-linking if for any map $h \in C(Q, Y)$ such that $\left.h\right|_{\partial Q} \in G$ there holds $h(Q) \cap S \neq \emptyset$.

Theorem 2.3 ([3, 9]). Let $X, Y$ be Banach spaces, $S$ be a closed subset of $Y, Q$ be a submanifold of $X$ with relative boundary $\partial Q$ and $G$ be a subset of $C(\partial Q, Y \backslash S)$. Let $\Gamma=\left\{h \in C(Q, Y):\left.h\right|_{\partial Q} \in G\right\}$, assume that $S$ and $\partial Q$ are $G$-linking and $f \in C^{1}(Y, R)$ satisfies
(a) There is $\tilde{h} \in \Gamma$ such that $\sup _{x \in Q} f(\tilde{h}(x))<+\infty$;
(b) There is $\beta_{0}>\alpha_{0}$ such that

$$
\inf _{y \in S} f(y) \geq \beta_{0} \quad \text { and } \quad \sup _{x \in \partial Q} f(h(x)) \leq \alpha_{0}, \quad \forall h \in \Gamma
$$

(c) The (PS) condition holds.

Then, the number

$$
c=\inf _{h \in \Gamma} \sup _{x \in Q} f(h(x))
$$

is a critical value of $f$ with $c \geq \beta_{0}$.
Proof. The proof is divided into two steps.
Step 1. The $(P S)$ condition for the functional $J$ is satisfied. Let $\left(u_{n}, v_{n}\right)$ be a $(P S)$ sequence for the functional $J$; that is,

$$
\begin{equation*}
J\left(u_{n}, v_{n}\right) \text { is bouned and } J^{\prime}\left(u_{n}, v_{n}\right) \rightarrow 0 \text { as } n \rightarrow \infty \tag{2.2}
\end{equation*}
$$

From (H1) and by a standard argument, it is sufficient to prove that $\left(u_{n}, v_{n}\right)$ is bounded in $W$. If this does not hold, assume that $\left\|\left(u_{n}, v_{n}\right)\right\|=\left\|u_{n}\right\|_{p}+\left\|v_{n}\right\|_{q} \rightarrow \infty$ as $n \rightarrow \infty$. Define $K_{n}:=\left\|u_{n}\right\|_{p}^{p}+\left\|v_{n}\right\|_{q}^{q}$, hence it follows that $K_{n} \rightarrow \infty$ as $n \rightarrow \infty$. Let $\bar{u}_{n}=u_{n} \backslash K_{n}^{1 / p}, \bar{v}_{n}=v_{n} \backslash K_{n}^{1 / q}$, then $\left(\bar{u}_{n}, \bar{v}_{n}\right)$ is bounded in $W$, i.e.,

$$
\left\|\bar{u}_{n}\right\|_{p}^{p}+\left\|\bar{v}_{n}\right\|_{q}^{q}=1 \quad \text { for all } n
$$

Extracting subsequences if necessary, we can assume that there exists $(\bar{u}, \bar{v}) \in W$ such that

$$
\begin{gather*}
\left(\bar{u}_{n}, \bar{v}_{n}\right) \rightharpoonup(\bar{u}, \bar{v}) \quad \text { weakly in } W,  \tag{2.3}\\
\left(\bar{u}_{n}, \bar{v}_{n}\right) \rightarrow(\bar{u}, \bar{v}) \quad \text { strongly in } L^{p}(\Omega) \times L^{q}(\Omega),  \tag{2.4}\\
\left(\bar{u}_{n}(x), \bar{v}_{n}(x)\right) \rightarrow(\bar{u}(x), \bar{v}(x)) \quad \text { for a.e. } x \in \Omega . \tag{2.5}
\end{gather*}
$$

From 2.2, it follows that

$$
\begin{equation*}
\limsup _{n \rightarrow \infty} \frac{J\left(u_{n}, v_{n}\right)}{K_{n}} \leq 0 \tag{2.6}
\end{equation*}
$$

From (2.1), (H1), the Hölder's inequality and (1.3), we have

$$
\begin{align*}
\left|\int_{\Omega} G(x, u, v) d x\right| & \leq \int_{\Omega}\left|\int_{0}^{1}\left(G_{s}(x, \tau u, \tau v) u+G_{t}(x, \tau u, \tau v) v\right) d \tau\right| d x \\
& \leq \int_{\Omega} h(x)(|u|+|v|) d x  \tag{2.7}\\
& \leq\|h\|_{L^{\infty}}\left(|\Omega|^{\frac{p-1}{p}}\|u\|_{L^{p}}+|\Omega|^{\frac{q-1}{q}}\|v\|_{L^{q}}\right) \\
& \leq C_{1}\left(\|u\|_{p}+\|v\|_{q}\right)
\end{align*}
$$

for all $(u, v) \in W$, where $C_{1}$ is a positive constant, hence it follows that

$$
\begin{equation*}
\frac{1}{K_{n}} \int_{\Omega} G\left(x, u_{n}, v_{n}\right) d x \rightarrow 0 \quad \text { as } n \rightarrow \infty \tag{2.8}
\end{equation*}
$$

From $h_{1} \in L^{p /(p-1)}(\Omega), h_{2} \in L^{q /(q-1)}(\Omega)$ and the Hölder's inequality, we obtain

$$
\begin{equation*}
\frac{1}{K_{n}} \int_{\Omega}\left(h_{1} u_{n}+h_{2} v_{n}\right) d x \rightarrow 0 \quad \text { as } n \rightarrow \infty \tag{2.9}
\end{equation*}
$$

From (2.4) and (2.5), it follows that $\left|\bar{u}_{n}\right|^{\alpha} \bar{u}_{n} \rightarrow|\bar{u}|^{\alpha} \bar{u}$ strongly in $L^{\frac{p}{\alpha+1}}(\Omega)$ and $\left|\bar{v}_{n}\right|^{\beta} \bar{v}_{n} \rightarrow|\bar{v}|^{\alpha} \bar{v}$ strongly in $L^{\frac{q}{\beta+1}}(\Omega)$, hence from Hölder's inequality, we obtain

$$
\begin{align*}
& \left|\int_{\Omega}\left(\left|\bar{u}_{n}\right|^{\alpha}\left|\bar{v}_{n}\right|^{\beta} \bar{u}_{n} \bar{v}_{n}-|\bar{u}|^{\alpha}|\bar{v}|^{\beta} \bar{u} \bar{v}\right) d x\right| \\
& \leq\left.\int_{\Omega}| | \bar{u}_{n}\right|^{\alpha}\left|\bar{v}_{n}\right|^{\beta} \bar{u}_{n} \bar{v}_{n}-\left|\bar{u}_{n}\right|^{\alpha}|\bar{v}|^{\beta} \bar{u}_{n} \bar{v}\left|d x+\int_{\Omega}\right|\left|\bar{u}_{n}\right|^{\alpha}|\bar{v}|^{\beta} \bar{u}_{n} \bar{v}-|\bar{u}|^{\alpha}|\bar{v}|^{\beta} \bar{u} \bar{v} \mid d x \\
& \leq\left.\int_{\Omega}\left|\bar{u}_{n}\right|^{\alpha+1} \cdot| | \bar{v}_{n}\right|^{\beta} \bar{v}_{n}-|\bar{v}|^{\beta} \bar{v}\left|d x+\int_{\Omega}\right|\left|\bar{u}_{n}\right|^{\alpha} \bar{u}_{n}-\left.|\bar{u}|^{\alpha} \bar{u}|\cdot| \bar{v}\right|^{\beta+1} d x \\
& \leq\left\|\bar{u}_{n}\right\|_{L^{p}}^{\alpha+1} \cdot\left\|\left.| | \bar{v}_{n}\right|^{\beta} \bar{v}_{n}-|\bar{v}|^{\beta} \bar{v}\right\|_{L^{\frac{q}{\beta+1}}}+\left\|\bar{v}_{n}\right\|_{L^{q}}^{\beta+1} \cdot\left\|\left.| | \bar{u}_{n}\right|^{\alpha} \bar{u}_{n}-|\bar{u}|^{\alpha} \bar{u}\right\|_{L^{\frac{p}{\alpha+1}}} \\
& \rightarrow 0 \text { as } n \rightarrow \infty . \tag{2.10}
\end{align*}
$$

From the definition of $J,(2.4,, 2.6,2.2,2.20$ and 2.10 , we have

$$
\begin{aligned}
& \limsup _{n \rightarrow \infty}\left(\frac{1}{p} \int_{\Omega}\left|\nabla \bar{u}_{n}\right|^{p} d x+\frac{1}{q} \int_{\Omega}\left|\nabla \bar{v}_{n}\right|^{q} d x\right) \\
& \leq \lambda_{1}\left(\frac{1}{p} \int_{\Omega}|\bar{u}|^{p} d x+\frac{1}{q} \int_{\Omega}|\bar{v}|^{q} d x+\frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega}|\bar{u}|^{\alpha}|\bar{v}|^{\beta} \bar{u} \bar{v} d x\right)
\end{aligned}
$$

From 2.3, it follows that

$$
\int_{\Omega}|\nabla \bar{u}|^{p} d x \leq \liminf _{n \rightarrow \infty} \int_{\Omega}\left|\nabla \bar{u}_{n}\right|^{p} d x, \quad \int_{\Omega}|\nabla \bar{v}|^{q} d x \leq \liminf _{n \rightarrow \infty} \int_{\Omega}\left|\nabla \bar{v}_{n}\right|^{q} d x
$$

hence, combining this with the definition of $\lambda_{1}$, we obtain

$$
\begin{aligned}
& \lambda_{1}\left(\frac{1}{p} \int_{\Omega}|\bar{u}|^{p} d x+\frac{1}{q} \int_{\Omega}|\bar{v}|^{q} d x+\frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega}|\bar{u}|^{\alpha}|\bar{v}|^{\beta} \bar{u} \bar{v} d x\right) \\
& \leq \frac{1}{p} \int_{\Omega}|\nabla \bar{u}|^{p} d x+\frac{1}{q} \int_{\Omega}|\nabla \bar{v}|^{q} d x \\
& \leq \liminf _{n \rightarrow \infty}\left(\frac{1}{p} \int_{\Omega}\left|\nabla \bar{u}_{n}\right|^{p} d x+\frac{1}{q} \int_{\Omega}\left|\nabla \bar{v}_{n}\right|^{q} d x\right) \\
& \leq \limsup _{n \rightarrow \infty}\left(\frac{1}{p} \int_{\Omega}\left|\nabla \bar{u}_{n}\right|^{p} d x+\frac{1}{q} \int_{\Omega}\left|\nabla \bar{v}_{n}\right|^{q} d x\right) \\
& \leq \lambda_{1}\left(\frac{1}{p} \int_{\Omega}|\bar{u}|^{p} d x+\frac{1}{q} \int_{\Omega}|\bar{v}|^{q} d x+\frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega}|\bar{u}|^{\alpha}|\bar{v}|^{\beta} \bar{u} \bar{v} d x\right)
\end{aligned}
$$

hence it follows that

$$
\begin{aligned}
& \frac{1}{p} \int_{\Omega}|\nabla \bar{u}|^{p} d x+\frac{1}{q} \int_{\Omega}|\nabla \bar{v}|^{q} d x \\
& =\lambda_{1}\left(\frac{1}{p} \int_{\Omega}|\bar{u}|^{p} d x+\frac{1}{q} \int_{\Omega}|\bar{v}|^{q} d x+\frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega}|\bar{u}|^{\alpha}|\bar{v}|^{\beta} \bar{u} \bar{v} d x\right),
\end{aligned}
$$

and by the uniform convexity of $W$, we have that $\left(\bar{u}_{n}, \bar{v}_{n}\right)$ converges strongly to $(\bar{u}, \bar{v})$ in $W$, and from the definition of $\left(\mu_{1}, \nu_{1}\right)$, it follows that $(\bar{u}, \bar{v})= \pm\left(\mu_{1}, \nu_{1}\right)$.

In the following, let $(\bar{u}, \bar{v})=\left(\mu_{1}, \nu_{1}\right)$, and the other case where $(\bar{u}, \bar{v})=-\left(\mu_{1}, \nu_{1}\right)$ may be considered similarly. Hence from the definition of $J$, we have

$$
\begin{align*}
& \frac{p J_{1}\left(u_{n}, v_{n}\right)}{(p-1) K_{n}^{1 / p}}+\frac{q J_{2}\left(u_{n}, v_{n}\right)}{(q-1) K_{n}^{1 / q}}-\left\langle J^{\prime}\left(u_{n}, v_{n}\right),\left(\frac{\bar{u}_{n}}{p-1}, \frac{\bar{v}_{n}}{q-1}\right)\right\rangle \\
& =\frac{1}{p-1}\left(\int_{\Omega} G_{s}\left(x, u_{n}, v_{n}\right) \bar{u}_{n} d x-\frac{p}{K_{n}^{1 / p}} \int_{\Omega} \int_{0}^{1} G_{s}\left(x, r u_{n}, r v_{n}\right) u_{n} d r d x\right)  \tag{2.11}\\
& \quad+\frac{1}{q-1}\left(\int_{\Omega} G_{t}\left(x, u_{n}, v_{n}\right) \bar{v}_{n} d x-\frac{q}{K_{n}^{1 / q}} \int_{\Omega} \int_{0}^{1} G_{t}\left(x, r u_{n}, r v_{n}\right) v_{n} d r d x\right) \\
& \quad+\int_{\Omega} h_{1} \bar{u}_{n} d x+\int_{\Omega} h_{2} \bar{v}_{n} d x .
\end{align*}
$$

From $h_{1} \in L^{p /(p-1)}(\Omega), h_{2} \in L^{q /(q-1)}(\Omega)$, we have

$$
\begin{equation*}
\int_{\Omega} h_{1} \bar{u}_{n} d x \rightarrow \int_{\Omega} h_{1} \mu_{1} d x \text { and } \int_{\Omega} h_{2} \bar{v}_{n} d x \rightarrow \int_{\Omega} h_{2} \nu_{1} d x \quad \text { as } n \rightarrow \infty . \tag{2.12}
\end{equation*}
$$

From (H2) and (H3), it is easy to know that

$$
\begin{align*}
\int_{\Omega} G_{s}\left(x, u_{n}, v_{n}\right) \bar{u}_{n} d x & \rightarrow \int_{\Omega} g_{1}^{++} \mu_{1} d x \\
\int_{\Omega} G_{t}\left(x, u_{n}, v_{n}\right) \bar{v}_{n} d x & \rightarrow \int_{\Omega} g_{2}^{++} \nu_{1} d x \tag{2.13}
\end{align*}
$$

as $n \rightarrow \infty$. Finally, from (H2) and Lebesgue dominated convergence theorem, we have

$$
\begin{align*}
\frac{1}{K_{n}^{1 / p}} \int_{\Omega} \int_{0}^{1} G_{s}\left(x, r u_{n}, r v_{n}\right) u_{n} d r d x & =\int_{\Omega} \int_{0}^{1} G_{s}\left(x, r u_{n}, r v_{n}\right) \frac{u_{n}}{K_{n}^{1 / p}} d r d x  \tag{2.14}\\
& \rightarrow \int_{\Omega} g_{1}^{++} \mu_{1} d x \quad \text { as } n \rightarrow \infty
\end{align*}
$$

Similarly, we obtain

$$
\begin{equation*}
\frac{1}{K_{n}^{1 / q}} \int_{\Omega} \int_{0}^{1} G_{t}\left(x, r u_{n}, r v_{n}\right) v_{n} d r d x \rightarrow \int_{\Omega} g_{2}^{++} \nu_{1} d x \quad \text { as } n \rightarrow \infty \tag{2.15}
\end{equation*}
$$

Therefore, letting $n \rightarrow \infty$ in (2.11) and from (2.2), 2.12, 2.13), 2.14) and 2.15), we obtain

$$
\int_{\Omega} h_{1} \mu_{1} d x+\int_{\Omega} h_{2} \nu_{1} d x=\int_{\Omega} g_{1}^{++} \mu_{1} d x+\int_{\Omega} g_{2}^{++} \nu_{1} d x
$$

which contradicts with 1.5 . Hence, $\left(u_{n}, v_{n}\right)$ is bounded in $W$.
Step 2. The functional $J$ satisfies the geometries of Theorem 2.3. For any

$$
(u, v) \in E_{1}=\left\{\left(t^{1 / p} \mu_{1}, t^{1 / q} \nu_{1}\right): t \geq 0\right\} \cup\left\{\left(-t^{1 / p} \mu_{1},-t^{1 / q} \nu_{1}\right): t \geq 0\right\}
$$

we have

$$
\begin{aligned}
& \frac{1}{p} \int_{\Omega}|\nabla u|^{p} d x+\frac{1}{q} \int_{\Omega}|\nabla v|^{q} d x \\
& =\lambda_{1}\left(\frac{1}{p} \int_{\Omega}|u|^{p} d x+\frac{1}{q} \int_{\Omega}|v|^{q} d x+\frac{1}{(\alpha+1)(\beta+1)} \int_{\Omega}|u|^{\alpha}|v|^{\beta} u v d x\right) .
\end{aligned}
$$

From the above equality and the definition of $J$, for any $\left(t^{1 / p} \mu_{1}, t^{1 / q} \nu_{1}\right) \in E_{1}$, we obtain

$$
\begin{align*}
& J\left(t^{1 / p} \mu_{1}, t^{1 / q} \nu_{1}\right) \\
& =t^{1 / p} \int_{\Omega} h_{1} \mu_{1} d x+t^{1 / q} \int_{\Omega} h_{2} \nu_{1} d x-\int_{\Omega} G\left(x, t^{1 / p} \mu_{1}, t^{1 / q} \nu_{1}\right) d x \\
& =t^{1 / p}\left(\int_{\Omega} h_{1} \mu_{1} d x-\int_{\Omega} \int_{0}^{1} G_{s}\left(x, r t^{1 / p} \mu_{1}, r t^{1 / q} \nu_{1}\right) \mu_{1} d r d x\right)  \tag{2.16}\\
& \quad+t^{1 / q}\left(\int_{\Omega} h_{2} \nu_{1} d x-\int_{\Omega} \int_{0}^{1} G_{t}\left(x, r t^{1 / p} \mu_{1}, r t^{1 / q} \nu_{1}\right) \nu_{1} d r d x\right)
\end{align*}
$$

From (H1), (H2) and Lebesgue dominated convergence theorem, it follows that

$$
\begin{align*}
& \lim _{t \rightarrow+\infty} \int_{\Omega} \int_{0}^{1} G_{s}\left(x, r t^{1 / p} \mu_{1}, r t^{1 / q} \nu_{1}\right) \mu_{1} d r d x=\int_{\Omega} g_{1}^{++} \mu_{1} d x  \tag{2.17}\\
& \lim _{t \rightarrow+\infty} \int_{\Omega} \int_{0}^{1} G_{t}\left(x, r t^{1 / q} \mu_{1}, r t^{1 / q} \nu_{1}\right) \nu_{1} d r d x=\int_{\Omega} g_{2}^{++} \nu_{1} d x \tag{2.18}
\end{align*}
$$

Hence, from 1.7, 2.16, 2.17 and 2.18, we obtain

$$
J\left(t^{1 / p} \mu_{1}, t^{1 / q} \nu_{1}\right) \rightarrow-\infty \quad \text { as } t \rightarrow+\infty
$$

Similarly, the following result can be obtained with $g_{1}^{++}$and $g_{2}^{++}$exchanged with $g_{1}^{--}$and $g_{2}^{--}$respectively,

$$
J\left(-t^{1 / p} \mu_{1},-t^{1 / q} \nu_{1}\right) \rightarrow-\infty \quad \text { as } t \rightarrow+\infty
$$

Finally, it follows that

$$
\begin{equation*}
\lim _{|t| \rightarrow \infty} J\left( \pm t^{1 / p} \mu_{1}, \pm t^{1 / q} \nu_{1}\right)=-\infty \tag{2.19}
\end{equation*}
$$

On the other hand, letting $\Lambda_{2}:=\left\{(u, v) \in W: \phi(u, v) \geq \lambda_{2} \varphi(u, v)\right\}$, from (1.3), 2.7) and the Hölder's inequality, for any $(u, v) \in \Lambda_{2}$, we obtain

$$
\begin{aligned}
J(u, v) \geq & \frac{\lambda_{2}-\lambda_{1}}{p \lambda_{2}}\|u\|_{p}^{p}+\frac{\lambda_{2}-\lambda_{1}}{q \lambda_{2}}\|v\|_{q}^{q}-C_{1}\left(\|u\|_{p}+\|v\|_{q}\right) \\
& -\left(\left\|h_{1}\right\|_{L^{\frac{p}{p-1}}}\|u\|_{L^{p}}+\left\|h_{2}\right\|_{L^{\frac{q}{q-1}}}\|v\|_{L^{q}}\right) \\
\geq & \frac{\lambda_{2}-\lambda_{1}}{p \lambda_{2}}\left(\|u\|_{p}^{p}+\|v\|_{q}^{q}\right)-C_{2}\left(\|u\|_{p}+\|v\|_{q}\right)
\end{aligned}
$$

where $C_{2}=C_{1}+C \max \left\{\left\|h_{1}\right\|_{L^{\frac{p}{p-1}}},\left\|h_{2}\right\|_{L^{\frac{q}{q-1}}}\right\}$. Combining the above expression with 2.19), we obtain that there exists a positive constant $T$ such that

$$
\begin{equation*}
\alpha_{0}:=\sup _{t \geq T} J\left( \pm t^{1 / p} \mu_{1}, \pm t^{1 / q} \nu_{1}\right)<\beta_{0}:=\inf _{(u, v) \in \Lambda_{2}} J(u, v) . \tag{2.20}
\end{equation*}
$$

Let $M=\left\{\left( \pm t^{1 / p} \mu_{1}, \pm t^{1 / q} \nu_{1}\right): t \geq T\right\}$ and

$$
G=\left\{h \in C\left(S^{0}, W\right): h \text { is odd and } h\left(S^{0}\right) \subset M\right\}
$$

where $S^{0}$ is the boundary of the closed unit ball $B^{1}$ in $\mathbb{R}^{1}$, i.e., $S^{0}=\partial B^{1}$. For any $h \in G$, by 2.20 , we have $h\left(S^{0}\right) \cap \Lambda_{2}=\emptyset$, which implies that $G$ is a subset of $C\left(S^{0}, W \backslash \Lambda_{2}\right)$. Let

$$
\Gamma=\left\{h \in C\left(B^{1}, W\right):\left.h\right|_{S^{0}} \in G\right\}
$$

we can claim: $\Gamma$ is nonempty and $\Lambda_{2}$ and $S^{0}$ are G-linking, that is $h\left(B^{1}\right) \cap \Lambda_{2} \neq \emptyset$ for any $h \in \Gamma$. The similar proof of the conclusion may be found in [7, 3, 9], but for the readers convenience and completeness, we write it.

In fact, define $\bar{h}: B^{1} \rightarrow W$ by

$$
\begin{gathered}
\bar{h}(t)=\left((t T)^{1 / p} \mu_{1},(t T)^{1 / q} \nu_{1}\right) \quad \text { for all } t \in[0,1] \\
\bar{h}(-t)=\left(-(t T)^{1 / p} \mu_{1},-(t T)^{1 / q} \nu_{1}\right) \quad \text { for all } t \in[0,1] .
\end{gathered}
$$

Hence, $\bar{h} \in \Gamma$ and $\Gamma$ is nonempty. Now let $h \in \Gamma$, if there is $(u, v) \in h\left(B^{1}\right)$ such that $\varphi(u, v)=0$, we get $h\left(B^{1}\right) \cap \Lambda_{2} \neq \emptyset$. If not, we consider the map $\hat{h}: S^{1} \rightarrow \Sigma$ defined by

$$
\hat{h}\left(x_{1}, x_{2}\right)= \begin{cases}\pi \circ h\left(x_{1}\right), & \text { if } x_{2} \geq 0 \\ -\pi \circ h\left(-x_{1}\right), & \text { if } x_{2} \leq 0\end{cases}
$$

where $\pi(u, v)=\left(u \backslash(\varphi(u, v))^{1 / p}, v \backslash(\varphi(u, v))^{1 / q}\right)$. It is easy to know that $\hat{h}\left(S^{1}\right) \subset \Sigma_{2}$. Therefore, $\phi\left(u_{0}, v_{0}\right) \geq \lambda_{2}$ for some $\left(u_{0}, v_{0}\right) \in \hat{h}\left(S^{1}\right)$, namely, $\left(u_{0}, v_{0}\right) \in \Lambda_{2}$. From $\pi \circ h(x) \in \Lambda_{2}$, we have implies $h(x) \in \Lambda_{2}$, which implies that $h\left(B^{1}\right) \cap \Lambda_{2} \neq \emptyset$. Hence $\Lambda_{2}$ and $S^{0}$ are G-linking.

Now, from the compactness of $B^{1}$, (a) of Theorem 2.3 holds, (b) of Theorem 2.3 is satisfied from 2.20), (c) of Theorem 2.3 comes from (i). Accordingly, Theorem 1.1 holds from the G-linking Theorem with the critical value

$$
c=\inf _{h \in \Gamma} \sup _{x \in B^{1}} J(h(x)) .
$$

Proof of Theorem 1.3. (i) The functional $J$ satisfies the $(P S)$ condition. From (1.6), the claim can be proved with similar to step 1 of Theorem 1.1 .
(ii) Now we will prove that the functional $J$ is coercive, that is,

$$
J(u, v) \rightarrow+\infty \quad \text { as } \quad\|(u, v)\| \rightarrow \infty
$$

If the claim does not hold, there is a constant $c$ and a sequence $\left(u_{n}, v_{n}\right)$ such that $J\left(u_{n}, v_{n}\right) \leq c$ and $\left\|\left(u_{n}, v_{n}\right)\right\| \rightarrow \infty$. From the proof of the (PS) condition of Theorem 1.1. $\left(\bar{u}_{n}, \bar{v}_{n}\right)$ converges strongly to $\pm\left(\mu_{1}, \nu_{1}\right)$, where $\bar{u}_{n}=u_{n} \backslash K_{n}^{1 / p}, \bar{v}_{n}=$ $v_{n} \backslash K_{n}^{1 / q}$. Assume that $\left(\bar{u}_{n}, \bar{v}_{n}\right)$ converges strongly to $\left(\mu_{1}, \nu_{1}\right)$ (the case $\left(\bar{u}_{n}, \bar{v}_{n}\right)$ converges strongly to $\left(-\mu_{1},-\nu_{1}\right)$ may be treated similarly) and $p \geq q>1$ (the case $q \geq p>1$ may also be treated similarly). From the definitions of $J, J_{1}, J_{2}$ and $J\left(u_{n}, v_{n}\right) \leq c$ for all $n, 2.14$ and 2.15, we have

$$
\begin{aligned}
0 \geq & \limsup _{n \rightarrow \infty} \frac{J\left(u_{n}, v_{n}\right)}{K_{n}^{1 / p}}=\limsup _{n \rightarrow \infty}\left(\frac{J_{1}\left(u_{n}, v_{n}\right)}{K_{n}^{1 / p}}+\frac{J_{2}\left(u_{n}, v_{n}\right)}{K_{n}^{1 / p}}\right) \\
\geq & \limsup _{n \rightarrow \infty}\left(\frac{J_{1}\left(u_{n}, v_{n}\right)}{K_{n}^{1 / p}}+\frac{J_{2}\left(u_{n}, v_{n}\right)}{K_{n}^{1 / q}}\right) \\
\geq & \lim _{n \rightarrow \infty}\left(\int_{\Omega}\left(h_{1} \bar{u}_{n}-\frac{1}{K_{n}^{1 / p}} \int_{0}^{1} G_{s}\left(x, r u_{n}, r v_{n}\right) u_{n} d r\right) d x\right. \\
& \left.+\int_{\Omega}\left(h_{2} \bar{v}_{n}-\frac{1}{K_{n}^{1 / q}} \int_{0}^{1} G_{t}\left(x, r u_{n}, r v_{n}\right) v_{n} d r\right) d x\right) \\
= & \int_{\Omega}\left(h_{1} \mu_{1}-g_{1}^{++} \mu_{1}\right) d x+\int_{\Omega}\left(h_{2} \nu_{1}-g_{2}^{++} \nu_{1}\right) d x
\end{aligned}
$$

which is a contradiction to 1.6 . By Ekeland's Variational Principle, the proof is complete.

Acknowledgments. This research was supported by National Natural Science Foundation of China (No. 11471267), by the Fundamental Research Funds for the Central Universities (No. XDJK2011C039), and by the Postdoctoral Research Foundation of Chongqing (No. Xm201319).

## References

[1] D. Arcoya, L. Orsina; Landesman-Lazer conditions and quasilinear elliptic equations. Nonliear Anal. TMA, 28(10)(1997),1623-1632.
[2] E. D. Da Silva; Multiplicity of solutions for gradient systems using Landesman-Lazer conditions. Abstr. Appl. Anal., (2010), Article ID:237826.
[3] P. Drábek, S.B. Robinson; Resonance problems for the p-Laplacian. J. Funct. Anal., 169(1)(1999). 189-200.
[4] X. F. Ke, C. L. Tang; Existence and multiplicity of solutions for asymptotically linear noncooperative elliptic systems. J. Math. Anal. Appl. 375(2011), 631-647.
[5] M. Landesman, A. C. Lazer; Nonlinear perturbations of linear elliptic boundary value problems at resonance. J. Math. Mech., 19(1970), 609-623.
[6] Z. Q. Ou, C. Li; Existence of weak solutions for a class of quasilinear elliptic systems, Boundary Value Problems. (2015) 2015:195.
[7] Z. Q. Ou, C. L. Tang; Resonance problems for the p-Laplacian systems. J. Math. Anal. Appl., 345(1)(2008), 511-521.
[8] P. H. Rabinowitz; Minimax methods in critical point theory with applications to differential equations, CBMS Regional Conference Series in Mathematics, Vol. 65, American Math. Soc., Providence, RI(1986).
[9] S. Z. Song, C. L. Tang; Resonance problems for the p-Laplacian with a nonlinear boundary condition. Nonlinear Anal. TMA, 64(9),(2006) 2007-2021.
[10] M. Struwe; Variational Methods: Applications to Nonlinear Partial Differential Equations and Hamiltonian Systems, Springer-Verlag, New York(1990).
[11] J. J. Sun, C. L. Tang; Resonance problems for Kirchhoff type equations. Discrete Contin. Dyn. Syst. 33(5)(2013), 2139-2154.
[12] C. L. Tang; Solvability for two-point boundary value problems. J. Math. Anal. Appl., 216(1)(1997), 368-374.
[13] C. L. Tang; Solvability of the forced duffing equation at resonance. J. Math. Anal. Appl., 219(1)(1998), 110-124.
[14] N. B. Zographopoulos; p-Laplacian systems on resonance. Appl. Anal., 83(5)(2004), 509-519.
Zeng-Qi Ou
School of Mathematics and Statistics, Southwest University, Chongqing 400715, China
E-mail address: ouzengq707@sina.com, Phone +86 2368253135


[^0]:    2010 Mathematics Subject Classification. 35D30, 35J50, 35J92.
    Key words and phrases. Elliptic systems; Landesman-Lazer-type conditions; resonance; critical point theorem.
    (C) 2016 Texas State University.

    Submitted June 25, 2015. Published June 28, 2016.

