# EXISTENCE OF NONNEGATIVE SOLUTIONS FOR SINGULAR ELLIPTIC PROBLEMS 

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#### Abstract

We prove the existence of nonnegative nontrivial weak solutions to the problem $$
\begin{gathered} -\Delta u=a u^{-\alpha} \chi_{\{u>0\}}-b u^{p} \quad \text { in } \Omega, \\ u=0 \quad \text { on } \partial \Omega, \end{gathered}
$$ where $\Omega$ is a bounded domain in $\mathbb{R}^{n}$. A sufficient condition for the existence of a continuous and strictly positive weak solution is also given, and the uniqueness of such a solution is proved. We also prove a maximality property for solutions that are positive a.e. in $\Omega$.


## 1. Introduction and statement of the problem

Let $\Omega$ be a bounded domain in $\mathbb{R}^{n}$ with $C^{1,1}$ boundary, let $a$ and $b$ be nonnegative functions on $\Omega$, and let $\alpha$ and $p$ be positive real numbers. Consider the following singular elliptic problem

$$
\begin{gather*}
-\Delta u=a u^{-\alpha}-b u^{p} \quad \text { in } \Omega, \\
u=0 \quad \text { on } \partial \Omega  \tag{1.1}\\
u>0 \quad \text { in } \Omega
\end{gather*}
$$

Problems like (1.1) appear in chemical catalysts process, non-Newtonian fluids, and in models for the temperature of electrical devices (see e.g., [10, 7, 16, 19]).

Several works can be found concerning the existence of positive solutions to 1.1 for the case $b=0$, i.e., for the problem $-\Delta u=a u^{-\alpha}$ in $\Omega, u=0$ on $\partial \Omega, u>0$ in $\Omega$; let us mention a few: Classical solutions $u \in C^{2}(\Omega) \cap C(\bar{\Omega})$ satisfying $u(x)>0$ for all $x \in \Omega$ were obtained by Crandall, Rabinowitz and Tartar 11 under the following hypothesis: $a \in C^{1}(\bar{\Omega})$ and $\min _{\bar{\Omega}} a>0$. Lazer and McKenna 24 proved the existence of positive weak solutions $u \in H_{0}^{1}(\Omega)$ to (1.1) assuming that $a \in C^{\gamma}(\bar{\Omega})$, $\gamma \in(0,1)$, and, again, $a$ strictly positive on $\bar{\Omega}$. The case $0 \leq a \in L^{\infty}(\Omega), a \not \equiv 0$ (that is: $|\{x \in \Omega: a(x)>0\}|>0$ ) was studied by Del Pino [12]. Situations where $a$ is singular on the boundary $\partial \Omega$ were considered by Bougherara, Giacomoni and Hernández [5].

The existence of classical solutions to problem (1.1) was proved by Coclite and Palmieri [9] for $a$ and $b$ in $C^{1}(\bar{\Omega}), 0<p<1$, and $a$ strictly positive on $\bar{\Omega}$ (see [9,

[^0]Theorem 1]). Related singular elliptic problems were treated by Shi and Yao [29], and by Aranda and Godoy [3], [2]. Elliptic problems with singular terms and free boundaries were considered by Dávila and Montenegro [13, [14].

Ghergu and Rădulescu [22] studied multi-parameter singular bifurcation problems of the form $-\Delta u=g(u)+\lambda|\nabla u|^{p}+\mu f(., u)$ in $\Omega, u=0$ on $\partial \Omega, u>0$ in $\Omega$, where $\Omega$ is a smooth bounded domain in $\mathbb{R}^{n}, \lambda, \mu \geq 0,0<p \leq 2, f: \bar{\Omega} \times[0, \infty) \rightarrow$ $[0, \infty)$ is a Hölder continuous function such that $f(., s)$ is nondecreasing with respect to $s$, and $g:(0, \infty) \rightarrow(0, \infty)$ is a nonincreasing Hölder continuous function such that $\lim _{s \rightarrow 0^{+}} g(s)=\infty$. When $g(s)$ behaves like $s^{-\alpha}$ near the origin, with $0<\alpha<1$, the asymptotic behavior of the solution around the bifurcation point is established.

Dupaigne, Ghergu and Rădulescu [18] obtained various existence and nonexistence results for Lane-Emden-Fowler equations with convection and singular potential of the form $-\Delta u \pm p\left(d_{\Omega}(x)\right) g(u)=\lambda f(x, u)+\mu|\nabla u|^{\beta}$ in $\Omega, u=0$ on $\partial \Omega, u>0$ in $\Omega$, where $\Omega$ is a smooth bounded domain in $\mathbb{R}^{n}, d_{\Omega}(x)=\operatorname{dist}(x, \partial \Omega), \lambda>0, \mu \in \mathbb{R}$, $0<\beta \leq 2, p\left(d_{\Omega}(x)\right)$ is a positive weight possibly singular at $\partial \Omega, g \in C^{1}(0, \infty)$ is a positive decreasing function such that $\lim _{s \rightarrow 0^{+}} g(s)=\infty, f: \bar{\Omega} \times[0, \infty) \rightarrow[0, \infty)$ is a Hölder continuous function which is positive on $\Omega \times(0, \infty)$ and satisfies that $s \rightarrow f(x, s)$ is nondecreasing and also that $f(x, s)$ is either linear or sublinear with respect to $s$.

Rădulescu [28] states existence, nonexistence and uniqueness results for blow-up boundary solutions of logistic equations and for Lane-Emden-Fowler equations with singular nonlinearities and subquadratic convection term.

Existence and nonexistence results for solutions to the inequality $L u \geq K(x) u^{p}$ in $\Omega, u>0$ in $\Omega$ were obtained by Ghergu, Liskevich and Sobol [20] for the case where $\Omega$ is a punctured ball $B_{R}(0) \backslash\{0\}, p \in \mathbb{R}, K \in L_{\mathrm{loc}}^{\infty}\left(B_{R}(0) \backslash\{0\}\right)$, ess $\inf K>$ 0 , and $L u:=\sum_{1 \leq i, j \leq n} a_{i j}(x) \frac{\partial^{2} u}{\partial x_{i} \partial x_{j}}+\sum_{1 \leq j \leq n} b_{j}(x) \frac{\partial u}{\partial x_{j}}$, where the matrix $\mathbf{a}=$ $\left\{a_{i j}(x)\right\}_{1 \leq i, j \leq n}$ is symmetric, uniformly elliptic on $\Omega$, with each $a_{i j} \in L^{\infty}\left(B_{R}(0)\right)$, and each $\bar{b}_{j}$ is a measurable function and satisfies ess $\sup _{x \in B_{R}(0) \backslash\{0\}}|x| b_{j}(x)<\infty$.

Existence and uniqueness results were obtained by Bougherara and Giacomoni [4] for mild solutions to singular initial value parabolic problems involving the pLaplacian operator of the form $u_{t}-\Delta_{p} u=u^{-\alpha}+f(x, u)$ in $Q_{T}:=(0, T) \times \Omega$, $u=0$ on $(0, T) \times \partial \Omega, u>0$ in $Q_{T}, u(0, x)=u_{0}(x)$ in $\Omega$ where $\Omega$ is a regular bounded domain in $\mathbb{R}^{n}$, $f: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a bounded below Carathéodory function and nonincreasing with respect to the second variable, $\Delta_{p} u:=\operatorname{div}\left(|\nabla u|^{p-2} \nabla u\right)$, $1<p<\infty, \alpha>0, T>0$, and $u_{0}$ in a suitable functional space.

Singularly perturbed elliptic problems on an annulus whose solutions concentrate in a circle were studied by Manna and Srikanth 27.

Let us mention also that Loc and Schmitt [26, 25, extended the method of sub and supersolutions to deal with singular elliptic problems. A comprehensive treatment of the subject can be found in Ghergu and Rădulescu's book [21] (see also [28]), and in the survey article [15], by Díaz and Hernández.

Let us state the problem that we will consider from now on: Let $\Omega$ be a bounded domain in $\mathbb{R}^{n}$ with $C^{1,1}$ boundary, $\alpha \in(0,1)$, and $p \in\left(0,2^{*}-1\right)$, where $2^{*}$ is defined by $\frac{1}{2^{*}}=\frac{1}{2}-\frac{1}{n}$ if $n>2$ and $2^{*}=\infty$ if $n \leq 2$. Let $a$ and $b$ be nonnegative functions such that $a$ belongs to $L^{\infty}(\Omega), a \not \equiv 0$, and $b$ is in $L^{r}(\Omega)$, with $r=\frac{2}{1-p}$ if $p<1$, and $r=\infty$ otherwise.

We are concerned with weak solutions to the problem

$$
\begin{gather*}
-\Delta u=a u^{-\alpha} \chi_{\{u>0\}}-b u^{p} \quad \text { in } \Omega, \\
u=0 \quad \text { on } \partial \Omega  \tag{1.2}\\
u \geq 0 \quad \text { in } \Omega
\end{gather*}
$$

where $a u^{-\alpha} \chi_{\{u>0\}}$ stands for the function defined by $a u^{-\alpha} \chi_{\{u>0\}}(x)=a(x) u(x)^{-\alpha}$ if $u(x) \neq 0$, and $a u^{-\alpha} \chi_{\{u>0\}}(x)=0$ if $u(x)=0$.

By a weak solution to 1.2 we mean a nonnegative function $u \in H_{0}^{1}(\Omega)$ such that, for all $\varphi$ in $H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega),\left(a u^{-\alpha} \chi_{\{u>0\}}-b u^{p}\right) \varphi \in L^{1}(\Omega)$, and the following holds

$$
\begin{equation*}
\int_{\Omega}\langle\nabla u, \nabla \varphi\rangle=\int_{\Omega}\left(a u^{-\alpha} \chi_{\{u>0\}}-b u^{p}\right) \varphi . \tag{1.3}
\end{equation*}
$$

The main aim of this work is to prove the existence of at least one nonnegative weak solution $u \not \equiv 0$ to the stated problem (see Theorem 3.1). Additionally, we give a condition on $a, b$ that guarantees the existence of a strictly positive weak solution to 1.2 (see Theorem 3.5). In Theorem 3.8 we prove that there is at most one solution that is positive a.e. in $\Omega$, and give a maximality property for such a solution. Examples of non-existence of strictly positive solutions, and of non-uniqueness of the nonnegative solutions, are also provided.

To prove Theorem 3.1, we show that the energy functional $J$ associated with 1.2 attains its minimum at some nonnegative nontrivial $u \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$. Note that $J$ may fail to be Gateaux differentiable at $u$; despite this fact, we manage to prove that the said minimizer is indeed a weak solution of problem 1.2 . Theorem 3.5 is proved using the sub and supersolutions method for singular elliptic problems developed in 26 .

## 2. Preliminary lemmas

Let $J: H_{0}^{1}(\Omega) \rightarrow \mathbb{R}$ be the energy functional associated with 1.2 ,

$$
\begin{equation*}
J(u):=\frac{1}{2} \int_{\Omega}|\nabla u|^{2}-\frac{1}{1-\alpha} \int_{\Omega} a|u|^{1-\alpha}+\frac{1}{1+p} \int_{\Omega} b|u|^{1+p} . \tag{2.1}
\end{equation*}
$$

Let us start with the following lemma.
Lemma 2.1. The following statements hold:
(i) $J$ is coercive on $H_{0}^{1}(\Omega)$.
(ii) $\inf _{u \in H_{0}^{1}(\Omega)} J(u)>-\infty$.
((iii) $\inf _{u \in H_{0}^{1}(\Omega)} J(u)$ is achieved at some $u \in H_{0}^{1}(\Omega)$.

Proof. Let $u \in H_{0}^{1}(\Omega)$. Since $0<1-\alpha<1$, the Hölder's and Poincare's inequalities give

$$
\frac{1}{1-\alpha} \int_{\Omega} a|u|^{1-\alpha} \leq c\|\nabla u\|_{2}^{1-\alpha}
$$

for some positive constant $c$ independent of $u$, and so $J(u) \geq \frac{1}{2}\|\nabla u\|_{2}^{2}-c\|\nabla u\|_{2}^{1-\alpha}$, which clearly implies (i) and (ii).

To prove (iii), let $\beta=\inf _{u \in H_{0}^{1}(\Omega)} J(u)$, and consider a sequence $\left\{u_{j}\right\}_{j \in \mathbb{N}} \subset H_{0}^{1}(\Omega)$ such that $\lim _{j \rightarrow \infty} J\left(u_{j}\right)=\beta$. Then, by i), $\left\{u_{j}\right\}_{j \in \mathbb{N}}$ is bounded in $H_{0}^{1}(\Omega)$. Let $q$ be in $\left(p+1,2^{*}\right)$. Since the inclusion $H_{0}^{1}(\Omega) \hookrightarrow L^{q}(\Omega)$ is a compact map, we can assume (taking a subsequence if necessary) that $\left\{u_{j}\right\}_{j \in \mathbb{N}}$ converges strongly to some $u \in L^{q}(\Omega)$. Since $\left\{u_{j}\right\}_{k \in \mathbb{N}}$ is bounded in $H_{0}^{1}(\Omega)$, there exists $v \in H_{0}^{1}(\Omega)$,
and a subsequence $\left\{u_{j_{k}}\right\}_{k \in \mathbb{N}}$, such that the subsequence converges strongly to $v$ in $L^{2}(\Omega)$, and $\left\{\nabla u_{j_{k}}\right\}_{k \in \mathbb{N}}$ converges weakly to $\nabla v$ in $L^{2}\left(\Omega, \mathbb{R}^{n}\right)$. Thus $v=u,\left\{u_{j_{k}}\right\}_{k \in \mathbb{N}}$ converges to $u$ in $L^{q}(\Omega)$, and

$$
\begin{equation*}
\|\nabla u\|_{2} \leq \lim \inf _{k \rightarrow \infty}\left\|\nabla u_{j_{k}}\right\|_{2} \tag{2.2}
\end{equation*}
$$

On the other hand, the Nemytskii operators $f(u):=|u|^{1-\alpha}$ and $g(u):=|u|^{1+p}$ are continuous from $L^{2}(\Omega)$ into $L^{\frac{2}{1-\alpha}}(\Omega)$, and from $L^{q}(\Omega)$ into $L^{\frac{q}{1+p}}(\Omega)$, respectively [1. Theorem 1.2.1] and so, since $a \in L^{\infty}(\Omega)$ and $b \in L^{r}(\Omega)$,

$$
\begin{align*}
& \lim _{j \rightarrow \infty} \int_{\Omega}\left(\frac{1}{1-\alpha} a\left|u_{j_{k}}\right|^{1-\alpha}-\frac{1}{1+p} b\left|u_{j_{k}}\right|^{1+p}\right) \\
& =\int_{\Omega}\left(\frac{1}{1-\alpha} a|u|^{1-\alpha}-\frac{1}{1+p} b|u|^{1+p}\right) \tag{2.3}
\end{align*}
$$

which, combined with 2.2 , gives $J(u) \leq \liminf _{k \rightarrow \infty} J\left(u_{j_{k}}\right)=\beta$, therefore (iii) holds (since $\beta \leq J(u)$ ).

Corollary 2.2. $\inf _{u \in H_{0}^{1}(\Omega)} J(u)$ is achieved at some nonnegative $u \in H_{0}^{1}(\Omega)$.
Proof. Lemma 2.1 states that $J$ attains its minimum at some $u \in H_{0}^{1}(\Omega)$. Since $J(u)=J(|u|)$, a nonnegative minimizer exists.

For the rest of this article, we fix a nonnegative minimizer for $J$ on $H_{0}^{1}(\Omega)$, and denote it by $\mathbf{u}$.

Lemma 2.3. The equality

$$
\begin{equation*}
\int_{\Omega}\langle\nabla \mathbf{u}, \nabla(\mathbf{u} \varphi)\rangle=\int_{\Omega}\left(a \mathbf{u}^{1-\alpha}-b \mathbf{u}^{1+p}\right) \varphi \tag{2.4}
\end{equation*}
$$

holds for any $\varphi \in H^{1}(\Omega) \cap L^{\infty}(\Omega)$ such that $\varphi \mathbf{u} \in H_{0}^{1}(\Omega)$.
Proof. Let $\varphi \in H^{1}(\Omega) \cap L^{\infty}(\Omega)$ be such that $\varphi \mathbf{u} \in H_{0}^{1}(\Omega)$; satisfying, in addition, $\|\varphi\|_{\infty} \leq \frac{1}{2}$. Let $\tau \in R$ such that $|\tau|<1$. Then $\mathbf{u}+\tau \mathbf{u} \varphi \geq 0$, and $J(\mathbf{u}) \leq J(\mathbf{u}+\tau \mathbf{u} \varphi)$. A computation shows that this inequality can be written as

$$
\begin{align*}
& \tau \int_{\Omega}\langle\nabla \mathbf{u}, \nabla(\mathbf{u} \varphi)\rangle \\
& \geq \frac{1}{1-\alpha} \int_{\Omega} a \mathbf{u}^{1-\alpha}\left((1+\tau \varphi)^{1-\alpha}-1\right)-\frac{1}{1+p} \int_{\Omega} b \mathbf{u}^{1+p}\left((1+\tau \varphi)^{1+p}-1\right)  \tag{2.5}\\
& \quad-\frac{\tau^{2}}{2} \int_{\Omega} \mathbf{u}^{2}|\nabla \varphi|^{2}-\frac{\tau^{2}}{2} \int_{\Omega} \varphi^{2}|\nabla \mathbf{u}|^{2}-\tau^{2} \int_{\Omega} \mathbf{u} \varphi\langle\nabla \mathbf{u}, \nabla \varphi\rangle
\end{align*}
$$

Note that, for $\gamma>0$, the second-order Taylor expansion of the function $h(t)=$ $(1+t)^{\gamma}-1$ gives

$$
\begin{equation*}
(1+\tau \varphi)^{\gamma}-1=\gamma \tau \varphi-\frac{\tau^{2}}{2} \gamma(\gamma-1)\left(1+\zeta_{\tau, \gamma}\right)^{\gamma-2} \varphi^{2} \tag{2.6}
\end{equation*}
$$

for some measurable function $\zeta_{\tau, \gamma}: \Omega \rightarrow \mathbb{R}$ satisfying $\left|\zeta_{\tau, \gamma}\right| \leq|\tau \varphi| \leq \frac{1}{2}$. Inserting (2.6) (used with $\gamma=1-\alpha$ and $\gamma=1+p$ ) in 2.5), we obtain

$$
\begin{align*}
& \tau \int_{\Omega}\langle\nabla \mathbf{u}, \nabla(\mathbf{u} \varphi)\rangle \\
& \geq \\
& \tau \int_{\Omega} a \mathbf{u}^{1-\alpha} \varphi-\frac{\tau^{2}}{2} \alpha \int_{\Omega} a \mathbf{u}^{1-\alpha}\left(1+\zeta_{\tau, 1-\alpha}\right)^{-\alpha-1} \varphi^{2}  \tag{2.7}\\
& \quad-\left(\tau \int_{\Omega} b \mathbf{u}^{1+p} \varphi+\frac{\tau^{2}}{2} p \int_{\Omega} b \mathbf{u}^{1+p}\left(1+\zeta_{\tau, 1+p}\right)^{p-1} \varphi^{2}\right) \\
& \quad-\frac{\tau^{2}}{2} \int_{\Omega} \mathbf{u}^{2}|\nabla \varphi|^{2}-\frac{\tau^{2}}{2} \int_{\Omega} \varphi^{2}|\nabla \mathbf{u}|^{2}-\tau^{2} \int_{\Omega} \mathbf{u} \varphi\langle\nabla \mathbf{u}, \nabla \varphi\rangle
\end{align*}
$$

Also, $1+\zeta_{\tau, 1-\alpha} \geq \frac{1}{2}$ and $1+\zeta_{\tau, 1+p} \geq \frac{1}{2}$, and thus

$$
\begin{gathered}
\left|\int_{\Omega} a \mathbf{u}^{1-\alpha}\left(1+\zeta_{\tau, 1-\alpha}\right)^{-\alpha-1} \varphi^{2}\right| \leq c \\
\left|\int_{\Omega} b \mathbf{u}^{1+p}\left(1+\zeta_{\tau, 1+p}\right)^{p-1} \varphi^{2}\right| \leq c
\end{gathered}
$$

for some positive constant $c$ independent of $\tau$. Now we take $\tau$ positive in (2.7). Dividing by $\tau$, and then letting $\tau \rightarrow 0^{+}$, from (2.7) we obtain

$$
\int_{\Omega}\langle\nabla \mathbf{u}, \nabla(\mathbf{u} \varphi)\rangle \geq \int_{\Omega} a \mathbf{u}^{1-\alpha} \varphi-\int_{\Omega} b \mathbf{u}^{1+p} \varphi
$$

We note that this inequality holds if we put $-\varphi$ instead of $\varphi$; therefore we obtain also the reverse inequality, and we conclude that 2.4 is valid for $\|\varphi\|_{\infty} \leq \frac{1}{2}$. Finally, since both sides in (2.4) are linear on $\varphi$, the assumption $\|\varphi\|_{\infty} \leq \frac{1}{2}$ can be removed.

Lemma 2.4. There exists $v \in H_{0}^{1}(\Omega)$ such that $J(v)<0$.
Proof. It is sufficient to show that there exists a function $\Phi \in H_{0}^{1}(\Omega)$ such that $\int_{\Omega} a|\Phi|^{1-\alpha}>0$. Indeed, if such a $\Phi$ exists, then, for $t>0$, we have

$$
\begin{aligned}
J(t \Phi) & =\frac{t^{2}}{2}\|\nabla \Phi\|_{2}^{2}-\frac{t^{1-\alpha}}{1-\alpha} \int_{\Omega} a|\Phi|^{1-\alpha}+\frac{t^{1+p}}{1+p} \int_{\Omega} b|\Phi|^{1+p} \\
& =t^{1-\alpha}\left(\frac{t^{1+\alpha}}{2}\|\nabla \Phi\|_{2}^{2}-\frac{1}{1-\alpha} \int_{\Omega} a|\Phi|^{1-\alpha}+\frac{t^{p+\alpha}}{1+p} \int_{\Omega} b|\Phi|^{1+p}\right)
\end{aligned}
$$

which gives that $J(t \Phi)$ is negative for $t$ positive and small enough. Such a $\Phi$ can be constructed as follows: Let $h \in C_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ be a nonnegative radial function with support in the unit ball $B=\left\{x \in \mathbb{R}^{n}:|x|<1\right\}$, and such that $\int_{B} h=1$. For $\varepsilon>0$ let $h_{\varepsilon}(x):=\frac{1}{\varepsilon^{n}} h\left(\frac{x}{\varepsilon}\right)$. For $\delta>0$ let $\Omega_{\delta}:=\{x \in \Omega: \operatorname{dist}(x, \partial \Omega)>\delta\}$. Since $|\{x \in \Omega: a(x)>0\}|>0$, we have $\left|\{x \in \Omega: a(x)>0\} \cap \Omega_{\delta}\right|>0$ for $\delta$ positive and small enough. We fix such a $\delta$, and set $E=\{x \in \Omega: a(x)>0\} \cap \Omega_{\delta}$. For $\varepsilon>0$ we define $\Phi_{\varepsilon}:=h_{\varepsilon} * \chi_{E}$. Then $\Phi_{\varepsilon} \in C^{\infty}\left(\mathbb{R}^{n}\right)$ and $\operatorname{supp}\left(\Phi_{\varepsilon}\right) \subset \Omega$ for $\varepsilon<\delta$. Thus $\Phi_{\varepsilon} \in C_{c}^{\infty}(\Omega)$ for $\varepsilon<\delta$. Also, $\lim _{\varepsilon \rightarrow 0^{+}} \Phi_{\varepsilon}=\chi_{E}$ with convergence in $L^{2}(\Omega)$ (see [6, Theorem 4.22]). Then $\lim _{\varepsilon \rightarrow 0^{+}} a \Phi_{\varepsilon}^{1-\alpha}=a \chi_{E}$ with convergence in $L^{1}(\Omega)$ (see [1, Theorem 1.2.1]), therefore

$$
\lim _{\varepsilon \rightarrow 0^{+}} \int_{\Omega} a \Phi_{\varepsilon}^{1-\alpha}=\int_{\Omega} a\left(\chi_{E}\right)^{1-\alpha}=\int_{\Omega} a \chi_{E}>0
$$

Then $\int_{\Omega} a\left|\Phi_{\varepsilon}\right|^{1-\alpha}>0$ for $\varepsilon$ small enough.
Corollary 2.5. u $\quad \not \equiv 0$.
Remark 2.6. Let us observe that $\nabla\left(v^{2}\right)=2 v \nabla(v)$ for any (possibly unbounded) $v \in H^{1}(\Omega)$. Indeed, for $k \in \mathbb{N}$, let $v_{k}$ be the truncation of $v$, defined by $v_{k}(x)=v(x)$ if $|v(x)| \leq k$, and by $v_{k}(x)=k \operatorname{sign}(v(x))$ otherwise. Then $\left\{v_{k}\right\}_{k \in \mathbb{N}}$ converges to $v$ in $H^{1}(\Omega)$ as $k$ tends to $\infty$, and, since each $v_{k}$ is bounded, it follows from the chain rule (as stated e.g. in [23, Lemma 7.5]) that $\frac{\partial}{\partial x_{i}}\left(v_{k}^{2}\right)=2 v_{k} \frac{\partial v_{k}}{\partial x_{i}}, i=1,2, \ldots, n$. Since $\left\{v_{k}\right\}_{k \in \mathbb{N}}$ converges to $v$ in $L^{2}(\Omega)$, we have that $\left\{v_{k}^{2}\right\}_{k \in \mathbb{N}}$ converges to $v^{2}$ in $L^{1}(\Omega)$, and so also in $D^{\prime}(\Omega)$. Then $\left\{\frac{\partial}{\partial x_{i}}\left(v_{k}^{2}\right)\right\}_{k \in \mathbb{N}}$ converges to $\frac{\partial}{\partial x_{i}}\left(v^{2}\right)$ in $D^{\prime}(\Omega)$. Since $\left\{2 v_{k} \frac{\partial v_{k}}{\partial x_{i}}\right\}_{k \in \mathbb{N}}$ converges to $2 v \frac{\partial v}{\partial x_{i}}$ in $L^{1}(\Omega)$, and therefore in $D^{\prime}(\Omega)$, we obtain that, for each $i, \frac{\partial}{\partial x_{i}}\left(v^{2}\right)=2 v \frac{\partial v}{\partial x_{i}}$.

Lemma 2.7. $\mathbf{u} \in L^{\infty}(\Omega)$.
Proof. Let $\Omega^{\prime}$ be a bounded $C^{0,1}$ domain such that $\bar{\Omega} \subset \Omega^{\prime}$, and let $\widetilde{\mathbf{u}}, \widetilde{a}: \mathbb{R}^{n} \rightarrow \mathbb{R}$ be the extensions by zero of $\mathbf{u}$ and $a$ respectively. We consider first the case $n>2$. Let $r=\frac{1-\alpha}{2}, \eta=\frac{2^{*}}{1-\alpha}$. Then $0<r<1, \eta>1$, and $a \mathbf{u}^{2 r} \in L^{\eta}(\Omega)$. Let $z \in W^{2, \eta}\left(\Omega^{\prime}\right) \cap W_{0}^{1, \eta}\left(\Omega^{\prime}\right)$ be the solution of

$$
\begin{align*}
-\Delta z & =2 \widetilde{a} \widetilde{\mathbf{u}}^{2 r} \quad \text { in } \Omega^{\prime} \\
z & =0 \quad \text { on } \partial \Omega^{\prime} \tag{2.8}
\end{align*}
$$

Let $\widetilde{z}: \mathbb{R}^{n} \rightarrow \mathbb{R}$ be the extension by zero of $z$ and let $\varphi$ be a nonnegative function in $C_{c}^{\infty}\left(\Omega^{\prime}\right)$ By Remark 2.6 and Lemma 2.3 we have

$$
\begin{align*}
\int_{\Omega^{\prime}}\left\langle\nabla\left(\widetilde{\mathbf{u}}^{2}\right), \nabla \varphi\right\rangle & =\int_{\Omega^{\prime}}\langle 2 \widetilde{\mathbf{u}} \nabla \widetilde{\mathbf{u}}, \nabla \varphi\rangle \\
& =\int_{\Omega} 2 \mathbf{u}\langle\nabla \mathbf{u}, \nabla \varphi\rangle \leq \int_{\Omega} 2\langle\nabla(\mathbf{u} \varphi), \nabla \mathbf{u}\rangle \\
& =2 \int_{\Omega}\left(a \mathbf{u}^{1-\alpha}-b \mathbf{u}^{p+1}\right) \varphi  \tag{2.9}\\
& \leq 2 \int_{\Omega^{\prime}} \widetilde{a} \widetilde{\mathbf{u}}^{2 r} \varphi=\int_{\Omega^{\prime}}\langle\nabla z, \nabla \varphi\rangle
\end{align*}
$$

For $\varepsilon>0$ let $h_{\varepsilon}$ be the mollifiers defined as in the proof of Lemma 2.3. For $\varepsilon$ small enough we have $0 \leq \varphi * h_{\varepsilon} \in C_{c}^{\infty}\left(\Omega^{\prime}\right)$, and so, by 2.9 ,

$$
\begin{aligned}
\int_{\Omega^{\prime}}\left\langle\nabla\left(h_{\varepsilon} * \widetilde{\mathbf{u}}^{2}\right), \nabla \varphi\right\rangle & =\int_{\Omega^{\prime}}\left\langle\nabla\left(\widetilde{\mathbf{u}}^{2}\right), h_{\varepsilon} * \nabla \varphi\right\rangle \\
& =\int_{\Omega^{\prime}}\left\langle\nabla\left(\widetilde{\mathbf{u}}^{2}\right), \nabla\left(h_{\varepsilon} * \varphi\right)\right\rangle \\
& \leq \int_{\Omega^{\prime}}\left\langle\nabla z, \nabla\left(h_{\varepsilon} * \varphi\right)\right\rangle
\end{aligned}
$$

where we have used that, since $h_{\varepsilon}$ is an even function, the convolution operator with kernel $h_{\varepsilon}$ is self-adjoint in $L^{2}\left(\mathbb{R}^{n}\right)$. Recall that $\widetilde{z} \in W^{1, \eta}\left(\mathbb{R}^{n}\right)$ and $\operatorname{supp}(\widetilde{z}) \subset \overline{\Omega^{\prime}}$. Also, $\nabla \widetilde{z}=\nabla z$ a.e. in $\Omega^{\prime}$, and $\nabla \widetilde{z}=0$ a.e. in $\mathbb{R}^{n}-\Omega^{\prime}$. Thus

$$
\int_{\Omega^{\prime}}\left\langle\nabla z, \nabla\left(h_{\varepsilon} * \varphi\right)\right\rangle=\int_{\mathbb{R}^{n}}\left\langle\nabla \widetilde{z}, \nabla\left(h_{\varepsilon} * \varphi\right)\right\rangle
$$

$$
\begin{aligned}
& =\int_{\mathbb{R}^{n}}\left\langle\nabla\left(h_{\varepsilon} * \widetilde{z}\right), \nabla \varphi\right\rangle \\
& =\int_{\Omega^{\prime}}\left\langle\nabla\left(h_{\varepsilon} * \widetilde{z}\right), \nabla \varphi\right\rangle
\end{aligned}
$$

Then

$$
\int_{\Omega^{\prime}}\left\langle\nabla\left(h_{\varepsilon} * \widetilde{\mathbf{u}}^{2}\right), \nabla \varphi\right\rangle \leq \int_{\Omega^{\prime}}\left\langle\nabla\left(h_{\varepsilon} * \widetilde{z}\right), \nabla \varphi\right\rangle
$$

and so the divergence theorem gives

$$
-\int_{\Omega^{\prime}} \varphi \Delta\left(h_{\varepsilon} * \widetilde{\mathbf{u}}^{2}\right) \leq-\int_{\Omega^{\prime}} \varphi \Delta\left(h_{\varepsilon} * \widetilde{z}\right)
$$

Since this inequality holds for all nonnegative $\varphi \in C_{c}^{\infty}\left(\Omega^{\prime}\right)$ we obtain

$$
-\Delta\left(h_{\varepsilon} * \widetilde{\mathbf{u}}^{2}\right) \leq-\Delta\left(h_{\varepsilon} * \widetilde{z}\right) \text { in } \Omega^{\prime}
$$

We have also $h_{\varepsilon} * \widetilde{\mathbf{u}}^{2}=0 \leq h_{\varepsilon} * \widetilde{z}$ on $\partial \Omega^{\prime}$. Thus, the classical maximum principle gives $h_{\varepsilon} * \widetilde{\mathbf{u}}^{2} \leq h_{\varepsilon} * \widetilde{z}$ in $\Omega^{\prime}$. Now, $\widetilde{\mathbf{u}}^{2}$ and $\widetilde{z}$ belong to $L^{\frac{2^{*}}{2}}\left(\mathbb{R}^{n}\right)$, and so $\lim _{\varepsilon \rightarrow 0^{+}}\left(h_{\varepsilon} *\right.$ $\left.\widetilde{\mathbf{u}}^{2}\right)=\widetilde{\mathbf{u}}^{2}$, and $\lim _{\varepsilon \rightarrow 0^{+}}\left(h_{\varepsilon} * \widetilde{z}\right)=\widetilde{z}$, in both cases with convergence in $L^{\frac{2}{}^{*}}\left(\mathbb{R}^{n}\right)$. Then, $\left.\lim _{\varepsilon \rightarrow 0^{+}}\left(h_{\varepsilon} * \widetilde{\mathbf{u}}^{2}\right)\right|_{\Omega}=\widetilde{\mathbf{u}}_{\left.\right|_{\Omega}}^{2}$; and $\left.\lim _{\varepsilon \rightarrow 0^{+}}\left(h_{\varepsilon} * \widetilde{z}\right)\right|_{\Omega}=\left.\widetilde{z}\right|_{\Omega}$, in each case with convergence in $L^{\frac{2^{*}}{2}}(\Omega)$. Then $\mathbf{u}^{2} \leq z$ in $\Omega$.

Now the lemma follows from the following standard bootstrap argument: Let $\left\{\eta_{j}\right\}_{j \in \mathbb{N}}$ be recursively defined by $\eta_{1}=\eta^{*}$ and by $\eta_{j+1}=\eta_{j}^{*}$. We can see inductively that $\mathbf{u} \in L^{2 \eta_{j}}(\Omega)$ for all $j$. Indeed, $z \in W^{2, \eta}\left(\Omega^{\prime}\right)$, and so $z \in L^{\eta^{*}}\left(\Omega^{\prime}\right)$. Then $\mathbf{u}^{2} \in$ $L^{\eta^{*}}(\Omega)$, and thus $\mathbf{u} \in L^{2 \eta^{*}}(\Omega)=L^{2 \eta_{1}}(\Omega)$. Suppose now that $\mathbf{u} \in L^{2 \eta_{j}}(\Omega)$, then $2 \widetilde{a} \widetilde{\mathbf{u}}^{2 r} \in L^{\frac{\eta_{j}}{p}}\left(\Omega^{\prime}\right) \subset L^{\eta_{j}}\left(\Omega^{\prime}\right)$, and so $z \in W^{2, \eta_{j}}\left(\Omega^{\prime}\right) \subset L^{\eta_{j}^{*}}\left(\Omega^{\prime}\right)=L^{\eta_{j}^{*}}\left(\Omega^{\prime}\right)$, which gives $\mathbf{u} \in L^{2 \eta_{j}^{*}}(\Omega)=L^{2 \eta_{j+1}}(\Omega)$. Thus $\mathbf{u} \in L^{2 \eta_{j}}(\Omega)$ for all $j$, and so, taking $j$ large enough, we obtain $\mathbf{u} \in L^{s}(\Omega)$ for some $s>2 n$, then $2 \widetilde{a} \widetilde{\mathbf{u}}^{2 r} \in L^{\frac{s}{2 r}}\left(\Omega^{\prime}\right) \subset L^{\frac{s}{2}}\left(\Omega^{\prime}\right)$. Thus $z \in W^{2, \frac{s}{2}}\left(\Omega^{\prime}\right) \subset L^{\infty}\left(\Omega^{\prime}\right)$. Since $\mathbf{u}^{2} \leq z$ in $\Omega$, we obtain $\mathbf{u} \in L^{\infty}(\Omega)$.

Finally, if $n \leq 2$, we have $\mathbf{u} \in L^{s}(\Omega)$ for all $s \in[1, \infty)$. We take $\eta>n$ and, for $r, z, \bar{z}$ and $\widetilde{\mathbf{u}}$ defined as above, we have $a \mathbf{u}^{2 r} \in L^{\eta}(\Omega)$. Thus $\widetilde{z} \in W^{2, \eta}\left(\Omega^{\prime}\right) \subset C\left(\overline{\Omega^{\prime}}\right)$ and, as before, $\mathbf{u}^{2} \leq z$ in $\Omega$. Then $\mathbf{u} \in L^{\infty}(\Omega)$ also in this case.
Lemma 2.8.

$$
\begin{equation*}
\int_{\Omega}\langle\nabla \mathbf{u}, \nabla(\mathbf{u} \varphi)\rangle=\int_{\Omega}\left(a \mathbf{u}^{1-\alpha}-b \mathbf{u}^{1+p}\right) \varphi \tag{2.10}
\end{equation*}
$$

for all $\varphi \in H^{1}(\Omega) \cap L^{\infty}(\Omega)$.
Proof. Let $\varphi \in H^{1}(\Omega) \cap L^{\infty}(\Omega)$. By Lemma 2.7 we have $\mathbf{u} \in L^{\infty}(\Omega)$ and so $\mathbf{u} \varphi \in H_{0}^{1}(\Omega)$. Thus Lemma 2.3 gives 2.10 .

## 3. Main Results

Theorem 3.1. There exists a nonnegative weak solution $0 \not \equiv u \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$ of problem 1.2.

Proof. Let $\mathbf{u}$ be the nonnegative minimizer of $J$ considered in the previous section. Let $\psi$ be a nonnegative function in $H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$, and let $\varepsilon>0$. Note that $\frac{\psi}{\mathbf{u}+\varepsilon} \in H^{1}(\Omega) \cap L^{\infty}(\Omega)$, and that $\nabla\left(\mathbf{u} \frac{\psi}{\mathbf{u}+\varepsilon}\right)=\varepsilon \frac{\nabla \mathbf{u}}{(\mathbf{u}+\varepsilon)^{2}} \psi+\frac{\mathbf{u}}{\mathbf{u}+\varepsilon} \nabla \psi$, and so Lemma 2.8 gives

$$
\begin{equation*}
\varepsilon \int_{\Omega} \psi \frac{|\nabla \mathbf{u}|^{2}}{(\mathbf{u}+\varepsilon)^{2}}+\int_{\Omega} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon}\langle\nabla \mathbf{u}, \nabla \psi\rangle=\int_{\Omega}\left(a \mathbf{u}^{1-\alpha}-b \mathbf{u}^{1+p}\right) \frac{1}{\mathbf{u}+\varepsilon} \psi \tag{3.1}
\end{equation*}
$$

Since $\nabla \mathbf{u}=0$ a.e. on the set $\{x \in \Omega: \mathbf{u}(x)=0\}$, and since $a \mathbf{u}^{1-\alpha}=b \mathbf{u}^{1+p}=0$ on the same set, (3.1) can be written as

$$
\begin{align*}
& \varepsilon \int_{\{\mathbf{u}>0\}} \psi \frac{|\nabla \mathbf{u}|^{2}}{(\mathbf{u}+\varepsilon)^{2}}+\int_{\{\mathbf{u}>0\}} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon}\langle\nabla \mathbf{u}, \nabla \psi\rangle+\int_{\{\mathbf{u}>0\}} b \mathbf{u}^{p} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon} \psi  \tag{3.2}\\
& =\int_{\{\mathbf{u}>0\}} a \mathbf{u}^{-\alpha} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon} \psi
\end{align*}
$$

Also

$$
\lim _{\varepsilon \rightarrow 0^{+}}\left(\frac{\mathbf{u}}{\mathbf{u}+\varepsilon}\langle\nabla \mathbf{u}, \nabla \psi\rangle\right)=\chi_{\{\mathbf{u}>0\}}\langle\nabla \mathbf{u}, \nabla \psi\rangle=\langle\nabla \mathbf{u}, \nabla \psi\rangle
$$

a.e. in $\Omega$, and $\left|\frac{\mathbf{u}}{\mathbf{u}+\varepsilon}\langle\nabla \mathbf{u}, \nabla \psi\rangle\right| \leq|\langle\nabla \mathbf{u}, \nabla \psi\rangle| \in L^{1}(\Omega)$, and so Lebesgue's dominated convergence theorem gives

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0^{+}} \int_{\{\mathbf{u}>0\}} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon}\langle\nabla \mathbf{u}, \nabla \psi\rangle=\int_{\Omega}\langle\nabla \mathbf{u}, \nabla \psi\rangle \tag{3.3}
\end{equation*}
$$

On the other hand, $\lim _{\varepsilon \rightarrow 0^{+}} a \mathbf{u}^{-\alpha} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon} \psi=a \mathbf{u}^{-\alpha} \psi$ on the set $\{x \in \Omega: \mathbf{u}(x)>0\}$ and, since $a \mathbf{u}^{-\alpha} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon} \psi$ is non-increasing in $\varepsilon$, the monotone convergence theorem gives

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0^{+}} \int_{\{\mathbf{u}>0\}} a \mathbf{u}^{-\alpha} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon} \psi=\int_{\{\mathbf{u}>0\}} a \mathbf{u}^{-\alpha} \psi=\int_{\Omega} a \mathbf{u}^{-\alpha} \chi_{\{\mathbf{u}>0\}} \psi \tag{3.4}
\end{equation*}
$$

Also

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0^{+}} \int_{\{\mathbf{u}>0\}} b \mathbf{u}^{p} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon} \psi=\int_{\Omega} b \mathbf{u}^{p} \psi \tag{3.5}
\end{equation*}
$$

Then, from (3.2), (3.3), (3.4 and (3.5), we obtain

$$
\begin{aligned}
& \int_{\Omega}\langle\nabla \mathbf{u}, \nabla \psi\rangle+\int_{\Omega} b \mathbf{u}^{p} \psi \\
& =\lim _{\varepsilon \rightarrow 0^{+}}\left(\int_{\{\mathbf{u}>0\}} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon}\langle\nabla \mathbf{u}, \nabla \psi\rangle+\int_{\{\mathbf{u}>0\}} b \mathbf{u}^{p} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon} \psi\right) \\
& \leq \limsup _{\varepsilon \rightarrow 0^{+}}\left(\int_{\{\mathbf{u}>0\}} \frac{\varepsilon \psi|\nabla \mathbf{u}|^{2}}{(\mathbf{u}+\varepsilon)^{2}}+\int_{\{\mathbf{u}>0\}} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon}\langle\nabla \mathbf{u}, \nabla \psi\rangle+\int_{\{\mathbf{u}>0\}} b \mathbf{u}^{p} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon} \psi\right) \\
& =\limsup _{\varepsilon \rightarrow 0^{+}} \int_{\{\mathbf{u}>0\}} a \mathbf{u}^{-\alpha} \frac{\mathbf{u}}{\mathbf{u}+\varepsilon} \psi \\
& =\int_{\Omega} a \mathbf{u}^{-\alpha} \chi_{\{\mathbf{u}>0\}} \psi .
\end{aligned}
$$

Thus

$$
\begin{equation*}
\int_{\Omega}\langle\nabla \mathbf{u}, \nabla \psi\rangle+\int_{\Omega} b \mathbf{u}^{p} \psi \leq \int_{\Omega} a \mathbf{u}^{-\alpha} \chi_{\{\mathbf{u}>0\}} \psi \tag{3.6}
\end{equation*}
$$

Let us see that the reverse inequality in (3.6 holds: A computation gives, for $t>0$,

$$
\begin{aligned}
0 \leq & \frac{1}{t}(J(\mathbf{u}+t \psi)-J(\mathbf{u})) \\
= & \int_{\Omega}\langle\nabla \mathbf{u}, \nabla \psi\rangle+\frac{t}{2} \int_{\Omega}|\nabla \psi|^{2}-\frac{1}{(1-\alpha) t} \int_{\Omega} a\left((\mathbf{u}+t \psi)^{1-\alpha}-\mathbf{u}^{1-\alpha}\right) \\
& +\frac{1}{(1+p) t} \int_{\Omega} b\left((\mathbf{u}+t \psi)^{1+p}-\mathbf{u}^{1+p}\right),
\end{aligned}
$$

and so

$$
\begin{align*}
& \frac{1}{(1-\alpha) t} \int_{\Omega} a\left((\mathbf{u}+t \psi)^{1-\alpha}-\mathbf{u}^{1-\alpha}\right) \\
& \leq \int_{\Omega}\langle\nabla \mathbf{u}, \nabla \psi\rangle+\frac{1}{(1+p) t} \int_{\Omega} b\left((\mathbf{u}+t \psi)^{1+p}-\mathbf{u}^{1+p}\right)+\frac{t}{2} \int_{\Omega}|\nabla \psi|^{2} \tag{3.7}
\end{align*}
$$

The mean value theorem gives $(\mathbf{u}+t \psi)^{1-\alpha}-\mathbf{u}^{1-\alpha}=(1-\alpha)\left(\mathbf{u}+\sigma_{t}\right)^{-\alpha} t \psi$ for some measurable function $\sigma_{t}$ such that $0<\sigma_{t}<t \psi$. Thus

$$
\begin{aligned}
& \frac{1}{(1-\alpha) t} \int_{\Omega} a\left((\mathbf{u}+t \psi)^{1-\alpha}-\mathbf{u}^{1-\alpha}\right) \\
& =\frac{1}{(1-\alpha) t} \int_{\{a>0\} \cap\{\psi>0\}} a\left((\mathbf{u}+t \psi)^{1-\alpha}-\mathbf{u}^{1-\alpha}\right) \\
& =\int_{\{a>0\} \cap\{\psi>0\}} a\left(\mathbf{u}+\sigma_{t}\right)^{-\alpha} \psi
\end{aligned}
$$

Now, $\lim _{t \rightarrow 0^{+}} a\left(\mathbf{u}+\sigma_{t}\right)^{-\alpha} \psi=a \mathbf{u}^{-\alpha} \psi$ a.e on the set $\{a>0\} \cap\{\psi>0\}$. Then, by Fatou's Lemma,

$$
\begin{align*}
& \liminf _{t \rightarrow 0^{+}} \frac{1}{(1-\alpha) t} \int_{\Omega} a\left((\mathbf{u}+t \psi)^{1-\alpha}-\mathbf{u}^{1-\alpha}\right) \\
& =\liminf _{t \rightarrow 0^{+}} \int_{\{a>0\} \cap\{\psi>0\}} a\left(\mathbf{u}+\sigma_{t}\right)^{-\alpha} \psi  \tag{3.8}\\
& \geq \int_{\{a>0\} \cap\{\psi>0\}} a \mathbf{u}^{-\alpha} \psi \geq \int_{\Omega} a \mathbf{u}^{-\alpha} \chi_{\{\mathbf{u}>0\}} \psi
\end{align*}
$$

Again by the mean value theorem, we have

$$
\frac{1}{(1+p) t} \int_{\Omega} b\left((\mathbf{u}+t \psi)^{1+p}-\mathbf{u}^{1+p}\right)=\int_{\Omega} b\left(\mathbf{u}+\sigma_{t}\right)^{p} \psi
$$

Note that, for $0<t<1$, we have $0 \leq b\left(\mathbf{u}+\sigma_{t}\right)^{p} \psi \leq b(\mathbf{u}+\psi)^{p+1} \in L^{1}(\Omega)$. Also, $\lim _{t \rightarrow 0^{+}} b\left(\mathbf{u}+\sigma_{t}\right)^{p} \psi=b \mathbf{u}^{p} \psi$ a.e. in $\Omega$. Thus, by Lebesgue's dominated convergence theorem, we have

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} \frac{1}{(1+p) t} \int_{\Omega} b\left((\mathbf{u}+t \psi)^{1+p}-\mathbf{u}^{1+p}\right)=\int_{\Omega} b \mathbf{u}^{p} \psi \tag{3.9}
\end{equation*}
$$

Now, from 3.7, 3.8), and 3.9, we obtain

$$
\begin{equation*}
\int_{\Omega}\langle\nabla \mathbf{u}, \nabla \psi\rangle+\int_{\Omega} b \mathbf{u}^{p} \psi \geq \int_{\Omega} a \mathbf{u}^{-\alpha} \chi_{\{\mathbf{u}>0\}} \psi \tag{3.10}
\end{equation*}
$$

Since $b \mathbf{u}^{p} \psi \in L^{1}(\Omega)$, 3.10) implies that $a \mathbf{u}^{-\alpha} \chi_{\{\mathbf{u}>0\}} \psi \in L^{1}(\Omega)$. We apply (3.10), combined with (3.6), to complete the proof.

Remark 3.2. It is well known (see e.g., 17) that, for $m \in L^{\infty}(\Omega)$ such that $|\{x \in \Omega: m(x)>0\}|>0$, there exists a unique $\lambda=\lambda_{1}(-\Delta, \Omega, m)$ such that the problem

$$
\begin{gathered}
-\Delta \varphi_{1}=\lambda m \varphi_{1} \quad \text { in } \Omega \\
\varphi_{1}=0 \quad \text { on } \partial \Omega \\
\varphi_{1}>0 \quad \text { in } \Omega
\end{gathered}
$$

has a solution $\varphi_{1} \in H_{0}^{1}(\Omega)$. This solution is unique up to a multiplicative constant, belongs to $C^{1 . \gamma}(\bar{\Omega})$ for some $0<\gamma<1$, satisfies that $|\nabla \varphi|(x)>0$ for all $x \in \partial \Omega$, and there are positive constants $c_{1}, c_{2}$ such that $c_{1} d_{\Omega} \leq \varphi \leq c_{2} d_{\Omega}$ in $\Omega$, where $d_{\Omega}: \Omega \rightarrow \mathbb{R}$ is the function defined by

$$
d_{\Omega}(x)=\operatorname{dist}(x, \partial \Omega)
$$

$\lambda_{1}$ and $\varphi_{1}$ are called, respectively, the principal eigenvalue and a positive principal eigenfunction for $-\Delta$ in $\Omega$, with Dirichlet boundary condition and weight $m$.
Remark 3.3. It is well known that, under our assumptions on $\Omega, \alpha$, and $a$, the problem

$$
\begin{gathered}
-\Delta \theta=a \theta^{-\alpha} \text { in } \Omega \\
\theta=0 \text { on } \partial \Omega \\
\theta>0 \text { in } \Omega
\end{gathered}
$$

has a unique weak solution $\theta \in H_{0}^{1}(\Omega)$. Moreover, $\theta$ is in $C(\bar{\Omega})$, and $\theta \geq c^{\prime} d_{\Omega}$ for some positive constant $c^{\prime}$ (see [12, 3]). A computation shows that (in weak sense) $-\Delta\left(\theta^{\alpha+1}\right)=-(\alpha+1) \theta^{\alpha} \Delta \theta-(\alpha+1) \alpha \theta^{\alpha-2}|\nabla \theta|^{2} \leq(\alpha+1)\|a\|_{\infty}$ in $\Omega$, and so we have $\theta \leq c^{\prime \prime} d_{\Omega}^{\frac{1}{\alpha+1}}$ in $\Omega$, for some constant $c^{\prime \prime}>0$.
Remark 3.4. Following [26], we say that $w \in W_{\text {loc }}^{1,2}(\Omega)$ is a subsolution (supersolution) to the problem

$$
\begin{equation*}
-\Delta z=a z^{-\alpha}-b z^{p} \text { in } \Omega \tag{3.11}
\end{equation*}
$$

in the sense of distributions, if, and only if: $w>0$ a.e. in $\Omega, a w^{-\alpha}-b w^{p} \in L_{\mathrm{loc}}^{1}(\Omega)$, and for all nonnegative $\varphi \in C_{c}^{\infty}(\Omega)$, it holds that

$$
\int_{\Omega}\langle\nabla w, \nabla \varphi\rangle \leq(\geq) \int_{\Omega}\left(a w^{-\alpha}-b w^{p}\right) \varphi
$$

We also say that $z \in W_{\text {loc }}^{1,2}(\Omega)$ is a solution, in the sense of distributions, of 3.11 if, and only if, $z>0$ a.e. in $\Omega$, and, for all $\varphi \in C_{c}^{\infty}(\Omega)$ it holds that

$$
\int_{\Omega}\langle\nabla z, \nabla \varphi\rangle=\int_{\Omega}\left(a z^{-\alpha}-b z^{p}\right) \varphi
$$

For subsolutions, supersolutions and solutions defined in the above sense, [26, Theorem 2.4] says that, if (3.11) has a subsolution $\underline{z}$ and a supersolution $\bar{z}$ (in the sense of distributions), both in $L_{\mathrm{loc}}^{\infty}(\Omega)$, and such such that $0<\underline{z}(x) \leq \bar{z}(x)$ a.e. $x \in \Omega$, and if there exists $k \in L_{\mathrm{loc}}^{\infty}(\Omega)$ such that $\left|a s^{-\alpha}-b s^{p}\right| \leq k(x)$ a.e. $x \in \Omega$ for all $s \in[\underline{z}(x), \bar{z}(x)]$; then (3.11) has a solution $z$ in the sense of distributions, and $z$ satisfies $\underline{z} \leq z \leq \bar{z}$ a.e. in $\Omega$.

Theorem 3.5. Suppose that $a \geq \varepsilon b$ for some $\varepsilon>0$. Then there exists a weak solution $v \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$ of $\left(1.2\right.$ such that $v \geq c d_{\Omega}$ in $\Omega$ for some $c>0$, and $v \in C_{\mathrm{loc}}^{1}(\Omega) \cap C(\bar{\Omega})$.
Proof. Suppose that $a \geq \varepsilon b$ for some $\varepsilon>0$. Let $\varphi_{1} \in H_{0}^{1}(\Omega)$ be the positive principal eigenfunction associated to the weight function $a$, normalized by $\left\|\varphi_{1}\right\|_{\infty}=$ 1 (see Remark 3.2). Note that (in weak sense), for $t$ positive and small enough,

$$
\begin{equation*}
-\Delta\left(t \varphi_{1}\right) \leq a\left(t \varphi_{1}\right)^{-\alpha}-b\left(t \varphi_{1}\right)^{p} \text { in } \Omega . \tag{3.12}
\end{equation*}
$$

Indeed, $-\Delta\left(t \varphi_{1}\right)=\lambda_{1} a t \varphi_{1}$, and so 3.12 is equivalent to $\left(1-\lambda_{1}\left(t \varphi_{1}\right)^{1+\alpha}\right) a \geq$ $\left(t \varphi_{1}\right)^{p+\alpha} b$ in $\Omega$. But, for $t$ small enough, we have $b\left(t \varphi_{1}\right)^{p+\alpha} \leq b t^{p+\alpha} \leq \frac{1}{2} \varepsilon b \leq$
$\frac{1}{2} a \leq\left(1-\lambda_{1}\left(t \varphi_{1}\right)^{1+\alpha}\right) a$ in $\Omega$. Since $t \varphi_{1}>0$ in $\Omega$, it follows that, for such a $t$, $t \varphi_{1}$ is a subsolution of (1.2), in the sense of Remark 3.4. On the other hand, let $\theta \in H_{0}^{1}(\Omega) \cap C(\bar{\Omega})$ be the solution of the problem $-\Delta \theta=a \theta^{-\alpha}$ in $\Omega, \theta=0$ on $\partial \Omega$. Since $\theta \geq c^{\prime} d_{\Omega}$ in $\Omega$ for some $c^{\prime}>0$, we have that $\theta$ is strictly positive in $\Omega$, and, by diminishing $t$ if necessary, we can assume, that $t \varphi_{1} \leq \theta$. Clearly (in weak sense) $-\Delta \theta \geq a \theta^{-\alpha}-b \theta^{p}$ in $\Omega$, and so $\theta$ is a supersolution of $\sqrt{1.2}$, again in the sense of Remark 3.4. Since $t \varphi_{1} \geq c_{1} t d_{\Omega}$ in $\Omega$ for some $c_{1}>0$, and since $\theta \leq c^{\prime \prime} d_{\Omega}^{\frac{1}{\alpha+1}}$ in $\Omega$ for some $c^{\prime \prime}>0$, we have $\left[t \varphi_{1}(x), \theta(x)\right] \subset\left[c_{1} t d_{\Omega}(x), c^{\prime \prime} d_{\Omega}^{\frac{1}{\alpha+1}}(x)\right]$ for $x \in \Omega$. Therefore a.e. $x \in \Omega$, for all $s \in\left[t \varphi_{1}(x), \theta(x)\right]$, the following holds

$$
\left|a s^{-\alpha}-b s^{p}\right| \leq\|a\|_{\infty}\left(c_{1} t\right)^{-\alpha} d_{\Omega}(x)^{-\alpha}+\|b\|_{\infty}\left(c^{\prime \prime}\right)^{p} d_{\Omega}^{\frac{p}{\alpha+1}}(x):=k(x)
$$

Since $k \in L_{\text {loc }}^{\infty}(\Omega)$, [26, Theorem 2.4] (see Remark 3.4), says that there exists $v \in W_{\mathrm{loc}}^{1,2}(\Omega)$ such that $t \varphi_{1} \leq v \leq \theta$ in $\Omega$, and such that, for any $\varphi \in C_{c}^{\infty}(\Omega)$,

$$
\begin{equation*}
\int_{\Omega}\langle\nabla v, \nabla \varphi\rangle=\int_{\Omega}\left(a v^{-\alpha}-b v^{p}\right) \varphi . \tag{3.13}
\end{equation*}
$$

Note that $v \in H_{0}^{1}(\Omega)$ : Indeed, let $\Omega^{\prime}$ be a subdomain of $\Omega$ such that $\overline{\Omega^{\prime}} \subset \Omega$. Since $v \geq c^{\prime \prime} d_{\Omega}$ in $\Omega$ for some $c^{\prime \prime}>0$, we have $a v^{-\alpha}-b v^{p} \in L^{\infty}\left(\Omega^{\prime}\right)$. Therefore, from (3.13), a density argument, and Lebesgue's dominated convergence theorem give that, for any $\varphi \in H_{0}^{1}\left(\Omega^{\prime}\right)$, it holds

$$
\begin{equation*}
\int_{\Omega^{\prime}}\langle\nabla v, \nabla \varphi\rangle=\int_{\Omega^{\prime}}\left(a v^{-\alpha}-b v^{p}\right) \varphi \tag{3.14}
\end{equation*}
$$

Let $\varepsilon>0$. Since $v \leq \theta \leq c^{\prime \prime} d_{\Omega}^{\frac{1}{1+\alpha}}$ for some $c^{\prime \prime}>0$, we have that $\operatorname{supp}(v-\varepsilon)^{+} \subset \Omega^{\prime}$ for some subdomain $\Omega^{\prime}$ such that $\overline{\Omega^{\prime}} \subset \Omega$. Also $(v-\varepsilon)^{+} \in H^{1}(\Omega)$ and so $(v-\varepsilon)^{+} \in$ $H_{0}^{1}(\Omega)$. Thus, from (3.14), we obtain

$$
\begin{align*}
\int_{\Omega} \chi_{\{v>\varepsilon\}} \nabla v \cdot \nabla v & =\int_{\Omega^{\prime}} \nabla v \cdot \nabla(v-\varepsilon)^{+} \\
& =\int_{\Omega^{\prime}}\left(a v^{-\alpha}-b v^{p}\right)(v-\varepsilon)^{+}  \tag{3.15}\\
& =\int_{\Omega}\left(a v^{-\alpha}-b v^{p}\right)(v-\varepsilon) \chi_{\{v>\varepsilon\}}
\end{align*}
$$

The monotone convergence theorem gives

$$
\lim _{\varepsilon \rightarrow 0^{+}} \int_{\Omega} \chi_{\{v>\varepsilon\}} \nabla v \cdot \nabla v=\int_{\Omega} \nabla v \cdot \nabla v
$$

and, since $a v^{-\alpha}-b v^{p} \in L^{1}(\Omega)$, and $v \in L^{\infty}(\Omega)$, Lebesgue's dominated convergence theorem gives

$$
\lim _{\varepsilon \rightarrow 0^{+}} \int_{\Omega}\left(a v^{-\alpha}-b v^{p}\right)(v-\varepsilon) \chi_{\{v>\varepsilon\}}=\int_{\Omega}\left(a v^{1-\alpha}-b v^{1+p}\right)
$$

Taking limits in 3.15, we obtain

$$
\int_{\Omega} \nabla v \cdot \nabla v=\int_{\Omega}\left(a v^{1-\alpha}-b v^{1+p}\right)<\infty
$$

Thus $v \in H^{1}(\Omega)$ and, since $t \varphi_{1} \leq v \leq \theta$, we have $v \in H_{0}^{1}(\Omega)$. Note also that $a v^{-\alpha}-b v^{p} \in L^{1}(\Omega)$ and so, again by a density argument, and applying Lebesgue's
dominated convergence theorem, we conclude that 3.13 holds for all $\varphi$ in $H_{0}^{1}(\Omega) \cap$ $L^{\infty}(\Omega)$.

Let $\Omega^{\prime}$ be an arbitrary subdomain of $\Omega$ such that $\overline{\Omega^{\prime}} \subset \Omega$, and let $\Omega^{\prime \prime}$ be such that $\overline{\Omega^{\prime}} \subset \Omega^{\prime \prime} \subset \overline{\Omega^{\prime \prime}} \subset \Omega$. Since $v \in L^{\infty}\left(\Omega^{\prime \prime}\right)$ and $\left.\left(a v^{-\alpha}-b v^{p}\right)\right|_{\Omega^{\prime \prime}} \in L^{\infty}\left(\Omega^{\prime \prime}\right)$, we have $\left.v\right|_{\Omega^{\prime}} \in W^{2, s}\left(\Omega^{\prime}\right)$ for all $s \in[1, \infty)$ (see e.g., Proposition 4.1.2 in [8) and so $\left.v\right|_{\Omega^{\prime}} \in C^{1}\left(\overline{\Omega^{\prime}}\right)$. Thus $v \in C_{\text {loc }}^{1}(\Omega)$ and, since $t \varphi_{1} \leq v \leq \theta, v$ is continuous on $\partial \Omega$.

Example 3.6. Let $\Omega=(0,2 \pi), \alpha=1 / 3$, and $p \in(0,1 / 5)$. Let $a$ and $b$ be the functions defined on $\Omega$ by $a=2(1-\cos (2 x)) \sqrt[3]{\sin ^{2}(x)}, b(x)=2\left|\sin ^{2}(x)\right|^{-p}$. Then $a \geq 0, b \geq 0,0 \not \equiv a \in L^{\infty}(\Omega)$ and $b \in L^{\frac{2}{1-p}}(\Omega)$. Consider now the following three functions in $C^{1}(\bar{\Omega}): u(x)=\sin ^{2}(x) \chi_{(0, \pi)}, v(x)=\sin ^{2}(x) \chi_{(0,2 \pi)}$, and $w(x)=$ $\sin ^{2}(x) \chi_{(\pi, 2 \pi)}$. A computation shows that $u, v$, and $w$ are all weak solutions of 1.2 ( $v$ is in fact a classical solution). Therefore (without additional assumptions on $a$ and $b$ ) uniqueness is not to be expected for nonnegative nontrivial weak solutions of 1.2 . Notice that $w \equiv 0$ on $(0, \pi)$. Note also that $v(x)>0$ for $x \in \Omega-\{\pi\}$ and $v(\pi)=0$, therefore, by Theorem 3.8 below, there is no continuous and strictly positive solution to 1.2 .

Example 3.7. Let $\Omega=(0,2)$, let $\alpha \in(0,1), p \in(0,1)$, let $b:=\chi_{(0,1)}$ and let $a:=\chi_{(1,1+\delta)}$, with

$$
0<\delta \leq\left(\frac{1-\alpha}{2}\right)^{\frac{1}{1-\alpha}}\left(\left(\frac{2}{p+1}\right)^{\frac{1}{1-p}}\left(\frac{1-p}{2}\right)^{\frac{1+p}{1-p}}\right)^{\frac{1+\alpha}{1-\alpha}}
$$

Let us show that the problem

$$
\begin{gather*}
-u^{\prime \prime}=a u^{-\alpha}-b u^{p} \quad \text { in } \Omega, \\
u=0 \quad \text { on } \partial \Omega \tag{3.16}
\end{gather*}
$$

has no weak solution $u \in H_{0}^{1}(\Omega)$ such that $u>0$ a.e. in $\Omega$. Let us suppose, for the sake of contradiction, that $u$ is a weak solution such that $u>0$ a.e. in $\Omega$. Since $H_{0}^{1}(\Omega) \subset C^{\gamma}(\bar{\Omega})$ for some $\gamma \in(0,1)$, we have $u \in C^{\gamma}(\bar{\Omega})$ for such a $\gamma$. Throughout this example, unless there is risk of confusion, the restrictions of $u$ to $(0,1),(1,1+\delta)$, and $(1+\delta, 2)$, will be still denoted by $u$. Since $u$ belongs to $C^{\gamma}([0,1])$, and $\left|u^{p}(x)-u^{p}(y)\right| \leq|u(x)-u(y)|^{p}$ for any $x, y \in[0,1]$, we have $u^{p} \in C^{\cdot \gamma p}([0,1])$. Let $A=u(1)$. Since

$$
\begin{gather*}
-u^{\prime \prime}=-u^{p} \quad \text { in }(0,1), \\
u(0)=0  \tag{3.17}\\
u(1)=A
\end{gather*}
$$

we have that $u$ is a classical solution of (3.17) that belongs to $C^{2}([0,1]) \cap C([0,1])$ and so $-u^{\prime \prime}=-u^{p}$ in $[0,1]$. (see, e.g., [23, Theorem 6.14]). Note also that

$$
\begin{equation*}
u(x) \geq\left(\frac{1-p}{2}\right)^{\frac{2}{1-p}}\left(\frac{2}{p+1}\right)^{\frac{1}{1-p}} x^{\frac{2}{1-p}} \quad \text { for all } x \in[0,1] \tag{3.18}
\end{equation*}
$$

Indeed, multiplying 3.17 by $u^{\prime}$ we obtain $\frac{1}{2}\left(\left(u^{\prime}\right)^{2}\right)^{\prime}=\frac{1}{p+1}\left(u^{p+1}\right)^{\prime}$ on $[0,1]$, and so $\frac{1}{2}\left(u^{\prime}(x)\right)^{2}-\frac{1}{p+1} u(x)^{p+1}=\frac{1}{2}\left(u^{\prime}(0)\right)^{2} \geq 0$ for all $x \in[0,1]$. Thus

$$
\begin{equation*}
\left(u^{\prime}\right)^{2} \geq \frac{2}{p+1} u^{p+1} \quad \text { in }[0,1] \tag{3.19}
\end{equation*}
$$

As $u \geq 0$ on $[0,1]$ and $u(0)=0$, we have $u^{\prime}(0) \geq 0$. Observe also that 3.17) implies $u^{\prime \prime} \geq 0$ on $[0,1]$, and so $u$ is a convex function on $[0,1]$. Thus $u^{\prime}$ is nondecreasing on $[0,1]$ and, since $u^{\prime}(0) \geq 0$, we have $u^{\prime} \geq 0$ in $[0,1]$, and so, from 3.19, we conclude

$$
\begin{equation*}
u^{\prime} \geq\left(\frac{2}{p+1}\right)^{1 / 2} u^{\frac{p+1}{2}} \quad \text { in }[0,1] \tag{3.20}
\end{equation*}
$$

If $u(\bar{x})=0$ for some $\bar{x} \in(0,1)$ we would have $u(x)=0$ for all $x \in(0, \bar{x})$, which contradicts the assumption that $u>0$ a.e. in $\Omega$. Thus $u(x)>0$ for all $x \in[0,1]$, therefore 3.20 can be rewritten as $u^{-\frac{p+1}{2}} u^{\prime} \geq\left(\frac{2}{p+1}\right)^{1 / 2}$ on $[0,1]$. By integrating this inequality over $(0, x)$ we obtain $\frac{2}{1-p}(u(x))^{\frac{1-p}{2}} \geq\left(\frac{2}{p+1}\right)^{1 / 2} x$ for all $x \in[0,1]$, and so 3.18 holds. In particular we have

$$
\begin{equation*}
u(1) \geq\left(\frac{1-p}{2}\right)^{\frac{2}{1-p}}\left(\frac{2}{p+1}\right)^{\frac{1}{1-p}} x^{\frac{2}{1-p}} \tag{3.21}
\end{equation*}
$$

and then, by 3.20,

$$
\begin{equation*}
u^{\prime}(1) \geq\left(\frac{2}{p+1}\right)^{\frac{1}{1-p}}\left(\frac{1-p}{2}\right)^{\frac{1+p}{1-p}} . \tag{3.22}
\end{equation*}
$$

Consider now the restriction of $u$ to $(1,1+\delta) ; u \in H^{1}(1,1+\delta) \subset C([1,1+\delta])$, and solves

$$
\begin{aligned}
& -u^{\prime \prime}=u^{-\alpha} \quad \text { in }(1,1+\delta) \\
& u(1) \geq 0, u(1+\delta) \geq 0
\end{aligned}
$$

Let $\zeta \in H_{0}^{1}(1,1+\delta) \subset C([1,1+\delta])$ be the solution to the problem

$$
\begin{aligned}
-\zeta^{\prime \prime} & =\zeta^{-\alpha} \quad \text { in }(1,1+\delta) \\
\zeta & >0 \quad \text { in }(1,1+\delta) \\
\zeta(1) & =0, \quad \zeta(1+\delta)=0
\end{aligned}
$$

Observe that $u \geq \zeta$ on $(1,1+\delta)$. To prove this, suppose, for the sake of contradiction, that $\{x \in(1,1+\delta): u(x)<\zeta(x)\} \neq \emptyset$, and let $U$ be one of its connected components. Note that $U$ is an open interval, since $u$ and $\zeta$ are continuous on $(1,1+\delta)$. Since $-\zeta^{\prime \prime}=\zeta^{-\alpha} \leq u^{-\alpha}=-u^{\prime \prime}$ on $U$, and $\zeta=u$ on $\partial U$, the maximum principle gives $\zeta \leq u$ on $U$, which is a contradiction. Thus $u \geq \zeta$ on $(1,1+\delta)$ as claimed.

Recall that there exists $c>0$ such that $\zeta \geq c d$ on $(1,1+\delta)$, where $d(x)=$ $\operatorname{dist}(x, \partial(1,1+\delta))$ for all $x \in(1,1+\delta)$ (see Remark 3.3); therefore $u \geq c d$ on $(1,1+\delta)$. Note also that $u(1+\delta)>0$. If not, since $u(2)=0$ and $u^{\prime \prime}=0$ in $(1,2)$, we would have $u=0$ in (1,2); which would contradict $u>0$ a.e. in $\Omega$. Since $u(1)>0, u(1+\delta)>0$, and $u \geq c d$ on $(1,1+\delta)$, it follows that $u(x)>0$ for any $x \in[1,1+\delta]$ and, since $u$ is continuous on $[1,1+\delta]$, we have $u \geq$ const $>0$ on $[1,1+\delta]$. Now

$$
\begin{aligned}
\left|u^{-\alpha}(x)-u^{-\alpha}(y)\right| & =(u(x) u(y))^{-\alpha}\left|u(x)^{\alpha}-u(y)^{\alpha}\right| \\
& \leq(u(x) u(y))^{-\alpha}|u(x)-u(y)|^{\alpha}
\end{aligned}
$$

and so, since $u \in C^{\gamma}(\bar{\Omega})$, we have $u^{-\alpha} \in C^{\alpha \gamma}([1,1+\delta])$. Let $A=u(1), B=u(1+\delta)$. Since u solves

$$
\begin{align*}
& -u^{\prime \prime}=u^{-\alpha} \quad \text { in }(1,1+\delta) \\
& u(1)=A, \quad u(1+\delta)=B \tag{3.23}
\end{align*}
$$

it follows that $u$ is a classical solution of (3.23) that belongs to $C^{2}([1,1+\delta]) \cap$ $C([1,1+\delta])($ see $[23$, Theorem 6.14]).

On the other hand, since $u^{\prime \prime}=0$ on $(1+\delta, 2)$ and $u(2)=0$, we have

$$
\begin{equation*}
u(x)=\frac{u(1+\delta)}{1-\delta}(2-x) \quad \text { for all } x \in(1+\delta, 2) \tag{3.24}
\end{equation*}
$$

Since $u^{-\alpha} \in C^{\alpha \gamma}([1,1+\delta])$ and $u \in H_{0}^{1}(\Omega) \subset C(\bar{\Omega})$, we have au ${ }^{-\alpha}-b u^{p} \in L^{2}(\Omega)$, and thus, from (3.16), it follows that $u \in W^{2,2}(\Omega) \subset C^{1}(\bar{\Omega})$. Multiplying 3.23) by $u^{\prime}$ we obtain

$$
\begin{equation*}
\left(\frac{1}{2}\left(u^{\prime}\right)^{2}\right)^{\prime}=-\frac{1}{1-\alpha}\left(u^{1-\alpha}\right)^{\prime} \quad \text { on }(1,1+\delta) \tag{3.25}
\end{equation*}
$$

and so $\frac{1}{2}\left(u^{\prime}\right)^{2}+\frac{1}{1-\alpha} u^{1-\alpha}=\mathrm{const}=\frac{1}{2}\left(u^{\prime}(1)\right)^{2}+\frac{1}{1-\alpha} u(1)^{1-\alpha}$. Therefore, for $x \in$ $(1,1+\delta): u^{\prime}(x)=0$ if, and only if, $\frac{1}{1-\alpha} u^{1-\alpha}(x)=\frac{1}{2}\left(u^{\prime}(1)\right)^{2}+\frac{1}{1-\alpha} u(1)^{1-\alpha}$. If there were no $x$ in $(1,1+\delta)$ such that $\frac{1}{1-\alpha} u^{1-\alpha}(x)=\frac{1}{2}\left(u^{\prime}(1)\right)^{2}+\frac{1}{1-\alpha} u(1)^{1-\alpha}$, we would have $u^{\prime}(x) \neq 0$ for all $x \in(1,1+\delta)$; which would imply that $u^{\prime}(x)>0$ for all $x \in(1,1+\delta)$ (since $u^{\prime}$ is continuous on $[1,1+\delta]$, and since $\left.u^{\prime}(1)>0\right)$. Thus $u^{\prime}(1+\delta) \geq 0$, but, by $3.24, u^{\prime}(1+\delta)=-\frac{u(1+\delta)}{1-\delta}<0$, which is a contradiction. Therefore $\left\{x \in(1,1+\delta): \frac{1}{1-\alpha} u^{1-\alpha}(x)=\frac{1}{2}\left(u^{\prime}(1)\right)^{2}+\frac{1}{1-\alpha} u(1)^{1-\alpha}\right\} \neq \emptyset$; let $x_{1}$ be its infimum. Since $u$ is continuous, $x_{1}$ is a minimum, therefore we have $u\left(x_{1}\right)=$ $\left(\frac{1-\alpha}{2}\left(u^{\prime}(1)\right)^{2}+u(1)^{1-\alpha}\right)^{\frac{1}{1-\alpha}}$. Note that $u^{\prime}(x)>0$ for all $x \in\left[1, x_{1}\right)$. Moreover, (3.23) gives that $u$ is concave on $[1,1+\delta]$, and so $\frac{u\left(x_{1}\right)-u(1)}{x_{1}-1} \leq u^{\prime}(1)$. Then, recalling 3.22,

$$
\begin{aligned}
x_{1}-1 & \geq \frac{u\left(x_{1}\right)-u(1)}{u^{\prime}(1)}=\frac{\left(\frac{1-\alpha}{2}\left(u^{\prime}(1)\right)^{2}+u(1)^{1-\alpha}\right)^{\frac{1}{1-\alpha}}-u(1)}{u^{\prime}(1)} \\
& \geq \frac{\left(\frac{1-\alpha}{2}\left(u^{\prime}(1)\right)^{2}\right)^{\frac{1}{1-\alpha}}+\left(u(1)^{1-\alpha}\right)^{\frac{1}{1-\alpha}}-u(1)}{u^{\prime}(1)} \\
& =\frac{\left(\frac{1-\alpha}{2}\left(u^{\prime}(1)\right)^{2}\right)^{\frac{1}{1-\alpha}}}{u^{\prime}(1)}=\left(\frac{1-\alpha}{2}\right)^{\frac{1}{1-\alpha}}\left(u^{\prime}(1)\right)^{\frac{1+\alpha}{1-\alpha}} \\
& \geq\left(\frac{1-\alpha}{2}\right)^{\frac{1}{1-\alpha}}\left(\left(\frac{2}{p+1}\right)^{\frac{1}{1-p}}\left(\frac{1-p}{2}\right)^{\frac{1+p}{1-p}}\right)^{\frac{1+\alpha}{1-\alpha}} \geq \delta,
\end{aligned}
$$

which contradicts $x_{1}<1+\delta$.
Theorem 3.8. There is at most one weak solution $v \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$ of 1.2 such that $v(x)>0$ a.e. in $\Omega$; and, if it exists, it satisfies $v \geq u$ for any other nonnegative weak solution $u \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$ of (1.2).

Proof. Since $s \rightarrow f(s):=a s^{-\alpha}-b s^{p}$ is nondecreasing, the uniqueness assertion of the theorem follows from a standard argument: If $w$ is another solution which is positive a.e. in $\Omega$, take $\varphi:=v-w$ as a test function in the weak form of the equation

$$
\begin{gathered}
-\Delta(v-w)=f(v)-f(w) \quad \text { in } \Omega \\
v-w=0 \quad \text { on } \partial \Omega
\end{gathered}
$$

to obtain $\int_{\Omega}|\nabla(v-w)|^{2}=\int_{\Omega}(f(v)-f(w))(v-w) \leq 0$, which implies $v=w$.

Let $u \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$ be a nonnegative solution of 1.2 . Therefore, for any $\varphi \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$, we have

$$
\begin{align*}
& \int_{\Omega}\langle\nabla(u-v), \nabla \varphi\rangle \\
& =\int_{\Omega}\left(a u^{-\alpha} \chi_{\{u>0\}}-b u^{p}-\left(a v^{-\alpha}-b v^{p}\right)\right) \varphi  \tag{3.26}\\
& =\int_{\{u>0\}}(f(u)-f(v)) \varphi+\int_{\{u=0\}}\left(-a v^{-\alpha}+b v^{p}\right) \varphi .
\end{align*}
$$

Now, we take $\varphi=(u-v)^{+}$. Since $v>0$ a.e. in $\Omega$, we have

$$
\int_{\{u=0\}}\left(-a v^{-\alpha}+b v^{p}\right)(u-v)^{+}=0 .
$$

Thus, from 3.26), we obtain $\int_{\Omega}\left|\nabla(u-v)^{+}\right|^{2} \leq 0$, and so $u \leq v$ in $\Omega$.

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