## NON-TRIVIAL SOLUTIONS OF FRACTIONAL SCHRÖDINGER-POISSON SYSTEMS WITH SUM OF PERIODIC AND VANISHING POTENTIALS

MINGZHU YU, HAIBO CHEN

$$
\begin{aligned}
& \text { AbSTRACT. We consider the fractional Schrödinger-Poisson system } \\
& \qquad \begin{array}{c}
(-\Delta)^{\alpha} u+V(x) u+K(x) \Phi(x) u=f(x, u)-\Gamma(x)|u|^{q-2} u \quad \text { in } \mathbb{R}^{3}, \\
(-\Delta)^{\beta} \Phi=K(x) u^{2} \quad \text { in } \mathbb{R}^{3}
\end{array}
\end{aligned}
$$

where $\alpha, \beta \in(0,1], 4 \alpha+2 \beta>3,4 \leq q<2_{\alpha}^{*}, K(x), \Gamma(x)$ and $f(x, u)$ are periodic in $x, V$ is coercive or $V=V_{\text {per }}+V_{\text {loc }}$ is a sum of a periodic potential $V_{\text {per }}$ and a localized potential $V_{\text {loc }}$. If $f$ has the subcritical growth, but higher than $\Gamma(x)|u|^{q-2} u$, we establish the existence and nonexistence of ground state solutions are dependent on the sign of $V_{\text {loc }}$. Moreover, we prove that such a problem admits infinitely many pairs of geometrically distinct solutions provided that $V$ is periodic and $f$ is odd in $u$. Finally, we investigate the existence of ground state solutions in the case of coercive potential $V$.

## 1. Introduction

This article concerns the nonlinear fractional Schöding-Poisson system

$$
\begin{gather*}
(-\Delta)^{\alpha} u+V(x) u+K(x) \phi(x) u=f(x, u)-\Gamma(x)|u|^{q-2} u, \quad \text { in } \mathbb{R}^{3}  \tag{1.1}\\
(-\Delta)^{\beta} \phi=K(x) u^{2}, \quad \text { in } \mathbb{R}^{3}
\end{gather*}
$$

where $\alpha, \beta \in(0,1], 4 \alpha+2 \beta>3$, the fractional Laplacian $(-\Delta)^{\alpha}(\alpha \in(0,1))$ can be defined by the Fourier transform $(-\Delta)^{\alpha}=\mathcal{F}^{-1}\left(|\xi|^{2 \alpha} \mathcal{F} u\right), \mathcal{F} u$ being the usual Fourier transform in $\mathbb{R}^{3}$. System (1.1) stems from not only an expansion of the Feynman path integral from Brownian-like to Levy-like quantum mechanical paths [17, 18], but also a system of identically charged particles interacting with each other in the case when magnetic effects can be neglected [3].

When $\alpha=\beta=1$, system (1.1) reduces to the Schrödinger-Poisson system

$$
\begin{gather*}
-\Delta u+V(x) u+K(x) \phi(x) u=f(x, u)-\Gamma(x)|u|^{q-2} u, \quad \text { in } \mathbb{R}^{3} \\
-\Delta \phi=K(x) u^{2}, \quad \text { in } \mathbb{R}^{3} \tag{1.2}
\end{gather*}
$$

which has important physical meanings because it appears in quantum mechanical models (see [19]) and in semiconductor theory [4, 20. During the past few years, there has been increasing attention to problems like 1.2 on the existence of ground

[^0]state solutions, positive solutions, multiple solutions and so on. For more details in studying on these kinds of problems, we refer the readers to [1, 2, 7, 15, 22, 27, 32, (33, 39 and the references therein.

When $\phi \equiv 0$, system (1.1) reduces to the fractional Schrödinger problem

$$
\begin{equation*}
(-\Delta)^{\alpha} u+V(x) u=f(x, u)-\Gamma(x)|u|^{q-2} u \quad \text { in } \mathbb{R}^{N} \tag{1.3}
\end{equation*}
$$

which appears in different areas of mathematical physics. For fractional Schrödinger problems like (1.3), pioneered from [12] and [14] via variational methods, there have been many works in the existence and multiplicity of solutions, such as 66, 30, 36. For other related problems, we can refer to [9, 10, 23, 31, 24, 28, 36, 11, At least in our knowledge, however, the only paper considering problem 1.3 is Bieganowski [6], where the authors assumed that
(A1) $f: \mathbb{R}^{N} \times \mathbb{R} \rightarrow \mathbb{R}$ is measurable, $\mathbb{Z}^{N}$ - periodic in $x \in \mathbb{R}^{N}$ and continuous in $u \in \mathbb{R}^{N}$ a.e. $x \in \mathbb{R}^{N}$, and there are $c>0,2 \leq q<p<2_{\alpha}^{*}=\frac{2 N}{N-2 \alpha}$ such that

$$
|f(x, u)| \leq c\left(1+|u|^{p-1}\right) \quad \forall u \in \mathbb{R}, x \in \mathbb{R}^{N}
$$

(A2) $f(x, u)=o(|u|)$ uniformly in $x$ as $|u| \rightarrow 0$;
(A3) $F(x, u) /|u|^{q} \rightarrow \infty$ uniformly in $x$ as $|u| \rightarrow \infty$, where $F(x, u)=\int_{0}^{u} f(x, s) d s$ is the primitive of $f$ with respect to $u$;
(A4) $u \rightarrow f(x, u) /|u|^{q-1}$ is strictly increasing on $(-\infty, 0)$ and $(0,+\infty)$;
(A5) $\Gamma \in L^{\infty}\left(\mathbb{R}^{N}\right)$ is $\mathbb{Z}^{N}$-periodic in $x \in \mathbb{R}^{N}, \Gamma(x) \geq 0$ for a.e. $x \in \mathbb{R}^{N}$;
(A6) $V_{0}:=\operatorname{essinf}_{x \in \mathbb{R}^{N}} V(x)>0$ for $0<\alpha<1$, $\inf \sigma(-\Delta+V(x))>0$ for $\alpha=1$;
(A7) $V \in C\left(\mathbb{R}^{N}, \mathbb{R}\right)$ is such that $\lim _{|x| \rightarrow \infty} V(x)=\infty$ and $V_{0}:=\inf _{x \in \mathbb{R}^{N}} V(x)>$ $0 ;$
(A8) $V=V_{\text {per }}+V_{\text {loc }}, V_{\text {per }} \in L^{\infty}\left(\mathbb{R}^{3}\right)$ is $\mathbb{Z}^{3}$-periodic, $V_{\text {loc }} \in L^{\infty}\left(\mathbb{R}^{3}\right)$ and $V_{\text {loc }}(x) \rightarrow 0$ as $|x| \rightarrow \infty$.
To obtain the ground state solutions of 1.3, the authors first adopted the techniques in 34 to show that the corresponding Nehari manifold is a topological manifold homeomorphic with the unit sphere in the work space, where a minimizing sequence can be found. Then they use the abstract setting introduced [5] which extends some results from [24, [25] to indicate that such a minimizing sequence is bounded. Finally, they decomposed the bounded (PS) sequences including the above minimizing sequence to achieve their aim. The similar reduction techniques are successful used by other authors to study semilinear elliptic systems in 9 and Schrödinger-Poisson system in 37. Bieganowski also applied the methods in 34 to obtain infinitely many pairs of geometrically distinct solutions of 1.3 ).

Presently, more attention has been paid to the study of problems concerning the fractional Schöding-Poisson system (1.1) without $\Gamma(x)|u|^{q-2} u$. And there has been a few results about these kinds of problems, we can refer to [8, 13, 16, 21, 35, 38]. Among them, Ji [16] considered the fractional Schrödinger-Poisson system

$$
\begin{gather*}
(-\Delta)^{s} u+V(x) u+\lambda \phi u=f(x, u), \quad \text { in } \mathbb{R}^{3}  \tag{1.4}\\
(-\Delta)^{t} \phi=u^{2}, \quad \text { in } \mathbb{R}^{3}
\end{gather*}
$$

where $\lambda \in \mathbb{R}^{+}$is a parameter, $s, t \in(0,1)$ and $4 s+2 t>3, V(x)$ satisfies $\left(V_{\alpha} 2\right)$ and $f$ fulfills subcritical growth, 4-suplinear at infinity and monotonicity assumption. By constraint variational method and quantitative deformation lemma, Ji proved that the system (1.4) admits a least an energy sign-changing solution.

Motivated by the results mentioned previously, we study the non-trivial solutions of 1.1). In what follows, we assume that $N=3$ in conditions (A5)-(A8) for convenience. To state our main results, we make the following assumptions:
(A9) $f: \mathbb{R}^{3} \times \mathbb{R} \rightarrow \mathbb{R}$ is measurable, $\mathbb{Z}^{3}$-periodic in $x \in \mathbb{R}^{3}$ and continuous in $u \in \mathbb{R}^{3}$ a.e. $x \in \mathbb{R}^{3}$, and there are $c>0,4 \leq q<p<2_{\alpha}^{*}=\frac{6}{3-2 \alpha}$ such that

$$
|f(x, u)| \leq c\left(1+|u|^{p-1}\right) \quad \forall u \in \mathbb{R}, x \in \mathbb{R}^{3} ;
$$

(A10) $K \in L^{\infty}\left(\mathbb{R}^{3}\right)$ is $\mathbb{Z}^{3}$-periodic, $0 \leq K(x) \leq K_{\infty}$, for all $x \in \mathbb{R}^{3}$ and $K(x) \not \equiv 0$; Now we state our main results.

Theorem 1.1. Let $\alpha \in(0,1]$. Suppose that (A2)-(A6), (A8)-(A10) are satisfied and either $V_{\mathrm{loc}} \equiv 0$ or $V_{\mathrm{loc}}(x)<0$ a.e. $x \in \mathbb{R}^{3}$. Then (1.1) has a ground state of Nehari type.

Theorem 1.2. Let $\alpha \in(0,1]$. Suppose that (A2)-(A6), (A8)-(A10) are satisfied and $\inf \sigma\left(-\Delta+V_{\text {per }}(x)\right)>0$. If $V_{\text {loc }}>0$ a.e. $x \in \mathbb{R}^{3}$, then 1.1) has no ground state solutions.

Theorem 1.3. Let $\alpha \in(0,1]$. Suppose that (A2)-(A6), (A8)-(A10) are satisfied, $V_{\text {loc }} \equiv 0$ and $f$ is odd in $u$. Then 1.1 admits infinitely many pairs $\pm u$ of geometrically distinct solutions.
Theorem 1.4. Let $\alpha \in(0,1]$. Suppose that (A2)-(A5), (A7), (A9), (A10) are satisfied. Then (1.1) has a ground state solution of Nehari type.

The remainder of this articleis organized as follows. In Section 2, some preliminary results are presented. In Section $3-6$, we give the proof of our main results.

## 2. Variational setting and preliminaries

In this section, we introduce some necessary information to be used in this paper. We will denote by $\mathcal{F} u$ the usual Fourier transform of $u$.

Fractional Sobolev spaces are the convenient setting for our equations. A very detailed introduction about it can be found in [26], we offer a review below.

We see that the fractional Sobolev space $W^{\alpha, p}\left(\mathbb{R}^{3}\right)$ is defined for any $p \in[1,+\infty)$ and $\alpha \in(0,1)$ as

$$
W^{\alpha, p}\left(\mathbb{R}^{3}\right)=\left\{u \in L^{p}\left(\mathbb{R}^{3}\right): \int_{\mathbb{R}^{3}} \frac{|u(x)-u(y)|^{p}}{|x-y|^{\alpha p+3}} d x d y<\infty\right\}
$$

This space is endowed with the Gagliardo norm

$$
\|u\|_{W^{\alpha, p}}=\left(\int_{\mathbb{R}^{3}}|u|^{p} d x+\int_{\mathbb{R}^{3}} \frac{|u(x)-u(y)|^{p}}{|x-y|^{\alpha p+3}} d x d y\right)^{1 / p}
$$

When $p=2$, these spaces are also denoted by $H^{\alpha}\left(\mathbb{R}^{3}\right)$.
If $p=2$, an equivalent definition of fractional Sobolev spaces is possible, based on Fourier analysis. Indeed, it turns out that

$$
H^{\alpha}\left(\mathbb{R}^{3}\right)=\left\{u \in L^{2}\left(\mathbb{R}^{3}\right): \int_{\mathbb{R}^{3}}\left(|\xi|^{2 \alpha}|\hat{u}|^{2}+|u|^{2}\right) d \xi<\infty\right\},
$$

and the norm can be equivalently written as

$$
\|u\|_{H^{\alpha}}=\left(\|u\|_{2}^{2}+\int_{\mathbb{R}^{3}}|\xi|^{2 \alpha}|\hat{u}|^{2} d \xi\right)^{1 / 2}
$$

Furthermore, this norm is equivalent to the norm

$$
\|u\|_{\alpha}=\left(\int_{\mathbb{R}^{3}}\left(\left|(-\Delta)^{\alpha / 2} u\right|^{2}+u^{2}\right) d x\right)^{1 / 2} .
$$

The spaces $D^{\alpha, 2}\left(\mathbb{R}^{3}\right)$ is defined as the completion of $C_{0}^{\infty}\left(\mathbb{R}^{3}\right)$ under the norm

$$
\|u\|_{D^{\alpha, 2}}=\left(\int_{\mathbb{R}^{3}}\left|(-\Delta)^{\alpha / 2} u\right|^{2} d x\right)^{1 / 2}=\left(\int_{\mathbb{R}^{3}}|\xi|^{2 \alpha}|\hat{u}(\xi)|^{2} d \xi\right)^{1 / 2}
$$

In this article, in view of the potential $V(x)$, we consider its subspace

$$
\begin{equation*}
E^{\alpha}=\left\{u \in L^{2}\left(\mathbb{R}^{3}\right): \int_{\mathbb{R}^{3}}\left(|\xi|^{2 \alpha}|\hat{u}(\xi)|^{2}+V(x) u^{2}\right) d x<\infty\right\} \tag{2.1}
\end{equation*}
$$

endowed with the norm

$$
\begin{equation*}
\|u\|=\left(\int_{\mathbb{R}^{3}}\left(|\xi|^{2 \alpha}|\hat{u}(\xi)|^{2}+V(x) u^{2}\right) d x\right)^{1 / 2} \tag{2.2}
\end{equation*}
$$

and the scalar product

$$
(u, v)=\int_{\mathbb{R}^{3}}\left(|\xi|^{2 \alpha} \hat{u}(\xi) \hat{v}+V(x) u(x) v(x)\right) d x .
$$

Furthermore, we know that $\|\cdot\|_{E^{\alpha}}$ is equivalent to the norm

$$
\|u\|=\left(\int_{\mathbb{R}^{3}}\left(\left|(-\Delta)^{\alpha / 2} u\right|^{2}+V(x) u^{2}\right) d x\right)^{1 / 2}
$$

The corresponding inner product is

$$
(u, v)=\int_{\mathbb{R}^{3}}\left((-\Delta)^{\alpha / 2} u(-\Delta)^{\alpha / 2} v+V(x) u v\right) d x
$$

First of all, let us study the variational setting of problem (1.1). Under (A10), for any $u \in E^{\alpha}\left(\mathbb{R}^{3}\right)$, the linear operator $T: D^{\beta, 2}\left(\mathbb{R}^{3}\right) \rightarrow \mathbb{R}$ defined as

$$
T(v)=\int_{\mathbb{R}^{3}} K(x) u^{2} v d x
$$

By Hölder's inequality and fractional Sobolev inequality, one has

$$
\begin{align*}
T(v) & =\int_{\mathbb{R}^{3}} K(x) u^{2} v d x \\
& \leq K_{\infty}\left(\int_{\mathbb{R}^{3}}|u(x)|^{\frac{12}{3+2 \beta}} d x\right)^{\frac{3+2 \beta}{6}}\left(\int_{\mathbb{R}^{3}}|v(x)|^{2_{\beta}^{*}} d x\right)^{1 / 2_{\beta}^{*}}  \tag{2.3}\\
& \leq C\|u\|^{2}\|v\|_{D^{\beta, 2}}
\end{align*}
$$

is well defined on $D^{\beta, 2}\left(\mathbb{R}^{3}\right)$ and is continuous, where $2 \leq \frac{12}{3+2 \beta}<2_{\alpha}^{*}$ because of $4 \alpha+2 \beta>3$. Then, by the Lax-Milgram theorem, there exists $\phi_{u}^{\beta} \in D^{\beta, 2}\left(\mathbb{R}^{3}\right)$ such that

$$
\begin{equation*}
\int_{\mathbb{R}^{3}}(-\Delta)^{\beta / 2} \phi_{u}^{\beta}(-\Delta)^{\beta / 2} v d x=\int_{\mathbb{R}^{3}} K(x) u^{2} v d x, \quad \forall v \in D^{\beta, 2}\left(\mathbb{R}^{3}\right) \tag{2.4}
\end{equation*}
$$

Therefore, $(-\Delta)^{\beta} \phi_{u}^{\beta}=K(x) u^{2}$ in a weak sense, and the representation formula holds

$$
\phi_{u}^{\beta}(x)=c_{\beta} \int_{\mathbb{R}^{3}} \frac{K(y) u^{2}(y)}{|x-y|^{3-2 \beta}} d y, \forall x \in \mathbb{R}^{3} \quad \text { where } c_{\beta}=\pi^{-\frac{3}{2}} 2^{-2 \beta} \frac{\Gamma(3-2 \beta)}{\Gamma(\beta)} .
$$

This formula is called $t$-Riesz potential. Substituting $\phi_{u}^{\beta}$ into 1.1), then 1.1) can be reduced to the fractional Schrödinger equation

$$
\begin{equation*}
(-\Delta)^{\alpha} u+V(x) u+K(x) \phi_{u}^{\beta}(x) u=f(x, u)-\Gamma(x)|u|^{q-2} u \quad \text { in } \mathbb{R}^{3} \tag{2.5}
\end{equation*}
$$

whose solutions are the critical points of the function $J(u): E^{\alpha}\left(\mathbb{R}^{3}\right) \rightarrow \mathbb{R}$ defined by

$$
\begin{equation*}
J(u)=\frac{1}{2}\|u\|^{2}+\frac{1}{4} \int_{\mathbb{R}^{3}} K(x) \phi_{u}^{\beta} u^{2} d x-I(u) \tag{2.6}
\end{equation*}
$$

where

$$
I(u)=\int_{\mathbb{R}^{3}}\left(F(x, u(x))-\frac{1}{q} \Gamma(x)|u|^{q}\right) d x
$$

Obviously, the assumptions of our theorems imply that $J$ is a well-defined of class $C^{1}$ functional and that

$$
\begin{align*}
\left(J^{\prime}(u), v\right)= & \int_{\mathbb{R}^{3}}\left((-\Delta)^{\alpha / 2} u(-\Delta)^{\alpha / 2} v+V(x) u v\right) d x+\int_{\mathbb{R}^{3}} K(x) \phi_{u}^{\beta}(x) u v d x  \tag{2.7}\\
& -\int_{\mathbb{R}^{3}} f(x, u) v-\Gamma(x)|u|^{q-1} u d x
\end{align*}
$$

Hence, if $u \in E^{\alpha}\left(\mathbb{R}^{3}\right)$ is a critical point of $J$, then the pair $\left(u, \phi_{u}^{\beta}\right)$ is a solution of (1.1).

Our goal is to find a critical point being a minimizer of $J$ on the Nehari manifold $\mathcal{N}:=\left\{u \in E^{\alpha}\left(\mathbb{R}^{3}\right) \backslash\{0\}: J^{\prime}(u)(u)=0\right\}$. Obviously, $\mathcal{N}$ contains all nontrivial critical points, hence a ground state is the least energy solution.

Lemma 2.1 ([35]). If $\alpha, \beta \in(0,1)$ and $4 \beta+2 \alpha \geq 3$, then for any $u \in H^{\alpha}\left(\mathbb{R}^{3}\right)$, we have
(i) $\phi_{u}^{t}: H^{\alpha}\left(\mathbb{R}^{3}\right) \rightarrow D^{\beta, 2}\left(\mathbb{R}^{3}\right)$ is continuous, and maps bounded sets into bounded sets;
(ii) $\phi_{\tau u}^{\beta}(x)=\tau^{2} \phi_{u}^{\beta}(x)$ for all $\tau \in \mathbb{R}, \phi_{u(\cdot+y)}^{\beta}=\phi_{u}^{\beta}(x+y)$;
(iii) if $u_{n} \rightharpoonup u$ in $H^{\alpha}\left(\mathbb{R}^{3}\right)$, then $\phi_{u_{n}}^{\beta} \rightharpoonup \phi_{u}^{\beta}$ in $D^{\beta, 2}\left(\mathbb{R}^{3}\right)$;
(iv) if $u_{n} \rightharpoonup u$ in $H^{\alpha}\left(\mathbb{R}^{3}\right)$, then $\int_{\mathbb{R}^{3}} \phi_{u_{n}}^{\beta} u_{n}^{2}=\int_{\mathbb{R}^{3}} \phi_{\left(u_{n}-u\right)}^{\beta}\left(u_{n}-u\right)^{2}+\int_{\mathbb{R}^{3}} \phi_{u}^{\beta} u^{2}+$ $o(1)$.

## 3. Proof of Theorem 1.1

First, we consider the norm in the working space $H^{\alpha / 2}\left(\mathbb{R}^{3}\right)$.
Lemma 3.1 ( 6$])$. Let $\alpha \in(0,2)$ and denote by

$$
\|u\|_{H^{\alpha / 2}}^{2}=\int_{\mathbb{R}^{3}}|\xi|^{\alpha}|\hat{u}(\xi)|^{2} d \xi+\int_{\mathbb{R}^{3}}|u(x)|^{2} d x
$$

the classical norm on the fractional Sobolev space $H^{\alpha / 2}\left(\mathbb{R}^{3}\right)$. Suppose that (A6) and (A8) hold. Then the norm

$$
\|u\|^{2}=\int_{\mathbb{R}^{3}}|\xi|^{\alpha}|\hat{u}(\xi)|^{2} d \xi+\int_{\mathbb{R}^{3}} V(x)|u(x)|^{2} d x
$$

is equivalent to $\|\cdot\|_{H^{\alpha / 2}}$.

Remark 3.2. Obviously, the norm equivalence is also true when (A7) or (A8) holds. And the above lemma implies that $E^{\alpha}\left(\mathbb{R}^{3}\right)$ coincides with $H^{\alpha}\left(\mathbb{R}^{3}\right)$. Thus our functional $J: H^{\alpha}\left(\mathbb{R}^{3}\right) \rightarrow \mathbb{R}$ has the form

$$
J(u)=\frac{1}{2}\|u\|^{2}+\frac{1}{4} \int_{\mathbb{R}^{3}} K(x) \phi_{u}^{\beta}(x) u^{2}-I(u) .
$$

The Nehari manifold has the form

$$
\mathcal{N}=\left\{u \in H^{\alpha}\left(\mathbb{R}^{3}\right) \backslash\{0\}:\|u\|^{2}=I^{\prime}(u)(u)-\int_{\mathbb{R}^{3}} K(x) \phi_{u}^{\beta}(x) u^{2}\right\}
$$

The following lemmas are crucial to the proof of the Theorem 1.1 .
Lemma 3.3. Suppose that the following conditions hold:
(A11) there exists $r \in(0,1)$ such that $a:=\inf _{\|u\|=r} J(u)>J(0)=0$;
(A12) for any $t_{n} \rightarrow \infty$ and $u_{n} \rightarrow u \neq 0$ as $n \rightarrow \infty, I\left(t_{n} u_{n}\right) / t_{n}^{q} \rightarrow \infty$;
(A13) for $t \in(0, \infty) \backslash\{1\}$ and $u \in \mathcal{N}, \frac{t^{4}-1}{4} I^{\prime}(u)(u)-I(t u)+I(u)<0$;
(A14) $J$ is coercive on $\mathcal{N}$.
Then $\inf _{\mathcal{N}} J>0$ and there exists a bounded minimizing sequence for $J$ on $\mathcal{N}$, i.e. there is a sequence $\left\{u_{n}\right\} \subset \mathcal{N}$ such that $J\left(u_{n}\right) \rightarrow \inf _{\mathcal{N}} J$ and $J^{\prime}\left(u_{n}\right) \rightarrow 0$.

Since the proof of Lemma 3.3 is similar to that of [37, Theorem 3.2], we omit it here.

Lemma 3.4. Suppose that (A2)-(A10) are satisfied. Then (A11)-(A14) hold.
Proof. To prove (A11) we fix $\epsilon>0$. Observe that (A2) and (A9) imply that $F(x, u) \leq \epsilon|u|^{2}+C_{\epsilon}|u|^{p}$ for some $C_{\epsilon}>0$. Therefore,

$$
\int_{\mathbb{R}^{3}}\left(F(x, u(x))-\frac{1}{q} \Gamma(x)|u|^{q}\right) d x \leq \int_{\mathbb{R}^{3}} F\left(x, u(x) d x \leq C\left(\epsilon\|u\|^{2}+C_{\epsilon}\|u\|^{p}\right)\right.
$$

for a positive constant $C$ provided by the Sobolev embedding theorem. Thus there is $r>0$ such that

$$
\int_{\mathbb{R}^{3}}\left(F(x, u(x))-\frac{1}{q} \Gamma(x)|u|^{q}\right) d x \leq \frac{1}{4}\|u\|^{2}
$$

for $\|u\| \leq r$. Therefore

$$
J(u)=\frac{1}{2}\|u\|^{2}+\frac{1}{4} \int_{\mathbb{R}^{3}} K(x) \phi_{u}^{\beta}(x) u^{2}-I(u) \geq \frac{1}{4}\|u\|^{2}=\frac{1}{4} r^{2}>0
$$

for all $u \in H^{\alpha}\left(\mathbb{R}^{3}\right), p>4$. So there exists $r \in(0,1)$ such that $\|u\|=r$ and $J(u)>0$, then $a=\inf _{\|u\|=r} J(u)>0$.

To prove (A12), by (A4) and Fatou's lemma we obtain

$$
I\left(t_{n} u_{n}\right) / t_{n}^{q}=\int_{\mathbb{R}^{3}} \frac{F\left(x, t_{n} u_{n}\right)}{t_{n}^{q}} d x-\frac{1}{q} \int_{\mathbb{R}^{3}} \Gamma(x)\left|u_{n}\right|^{q} d x \rightarrow \infty
$$

To prove (A13) we fix $u \in \mathcal{N}$ and consider

$$
\psi(t)=\frac{t^{4}-1}{2} I^{\prime}(u)(u)-I(t u)+I(u) \quad \forall t \geq 0
$$

Then $\psi(1)=0$ and

$$
\frac{d \psi(t)}{d t}=2 t^{3} I^{\prime}(u)(u)-u \frac{d I(t u)}{d t}
$$

$$
\begin{aligned}
& =2 t^{3} \int_{\mathbb{R}^{3}} f(x, u) u-\Gamma(x)|u|^{q} d x-\int_{\mathbb{R}^{3}} f(x, t u) u+\Gamma(x) t^{q-1}|u|^{q-1} d x \\
& =\int_{\mathbb{R}^{3}} 2 t^{3} f(x, u) u-f(x, t u) u d x-\left(2 t^{3}-t^{q-1}\right) \int_{\mathbb{R}^{3}} \Gamma(x)|u|^{q} d x
\end{aligned}
$$

Since $u \in \mathcal{N}$,

$$
I^{\prime}(u)(u)=\int_{\mathbb{R}^{3}} f(x, u) u-\Gamma(x)|u|^{q} d x=\|u\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{u}^{\beta} u^{2}>0 .
$$

Therefore, for $t>1$, we have $\frac{d \psi(t)}{d t}<0$ by (A4). Similarly, $\frac{d \psi(t)}{d t}>0$, for $t<1$. Therefore $\psi(t)<\psi(1)=0$ for $t=1$, that is

$$
\frac{t^{4}-1}{2} I^{\prime}(u)(u)-I(t u)+I(u)<0
$$

To prove (A14), for $u \in \mathcal{N}$, (A4) implies that $f(x, u) u \geq q F(x, u)$. Therefore

$$
\begin{aligned}
J(u) & =J(u)-\frac{1}{q} J^{\prime}(u)(u) \\
& =\left(\frac{1}{2}-\frac{1}{q}\right)\|u\|^{2}+\left(\frac{1}{4}-\frac{1}{q}\right) \int_{\mathbb{R}^{3}} K(x) \phi_{u}^{\beta}(x) u^{2} d x+\int_{\mathbb{R}^{3}}\left(\frac{1}{q} f(x, u) u-F(x, u)\right) d x \\
& \geq\left(\frac{1}{2}-\frac{1}{q}\right)\|u\|^{2}
\end{aligned}
$$

which implies that $J$ is coercive on $\mathcal{N}$.
Now, we will introduce a decomposition result of bounded $(P S)$ sequences which is important to the proof of Theorem 1.1. The following lemma generalizes the decomposition result from [5]. We will denote

$$
J_{\mathrm{per}}=J(u)-\frac{1}{2} \int_{\mathbb{R}^{3}} V_{\mathrm{loc}}(x) u^{2} d x
$$

Lemma 3.5. Suppose that the assumptions of Theorem 1.1 are satisfied. Let $\left(u_{n}\right)$ be a bounded Palais-sequence for $J$. Then passing to a subsequence of $\left(u_{n}\right)$, there exist an integer $l \geq 0$ and sequences $\left(y_{n}^{k}\right) \subset \mathbb{Z}^{3}$, $w_{k} \in H^{\alpha}\left(\mathbb{R}^{3}\right), k=1, \ldots, l$ such that
(a) $u_{n} \rightharpoonup u_{0}$ and $J^{\prime}\left(u_{0}\right)=0$;
(b) $\left|y_{n}^{k}\right| \rightarrow \infty$ and $\left|y_{n}^{k}-y_{n}^{k^{\prime}}\right| \rightarrow \infty$ for $k \neq k^{\prime}$;
(c) $w^{k} \neq 0$ and $J_{\text {per }}^{\prime}\left(w^{k}\right)=0$ for each $1 \leq k \leq l$;
(d) $u_{n}-u_{0}-\sum_{k=1}^{l} w^{k}\left(\cdot-y_{n}^{k}\right) \rightarrow 0$ in $H^{\alpha}\left(\mathbb{R}^{3}\right)$ as $n \rightarrow \infty$;
(e) $J\left(u_{n}\right) \rightarrow J\left(u_{0}\right)+\sum_{k=1}^{l} J_{\text {per }}\left(w^{k}\right)$.

The proof of Lemma 3.5 resembles the proof of [5, Theorem 4.1]. So we omit it.
Proof of Theorem 1.1. By Lemmas 3.3 and 3.4 there exists a bounded minimizing sequence $\left(u_{n}\right) \in \mathcal{N}$ for $J$, that is $J^{\prime}\left(u_{n}\right) \rightarrow 0$ and $J\left(u_{n}\right) \rightarrow c:=\inf _{\mathcal{N}} J>0$. By Lemma 3.5 we have that

$$
J\left(u_{n}\right) \rightarrow J\left(u_{0}\right)+\sum_{k=1}^{l} J_{\mathrm{per}}\left(w^{k}\right)
$$

where $w^{k}$ satisfies

$$
w^{k} \neq 0 \quad \text { and } \quad J_{\text {per }}^{\prime}\left(w_{k}\right)=0 \quad \text { for } 1 \leq k \leq l
$$

that is $w^{k}$ are critical points of the periodic part of the functional $J$. Suppose that $V_{\text {loc }} \equiv 0$, so $J=J_{\text {per }}$. If $u_{0}=0$, we have

$$
c+o(1)=J\left(u_{n}\right) \rightarrow \sum_{k=1}^{l} J_{\mathrm{per}}\left(w^{k}\right) \geq l c
$$

Then $l=1$ and $w^{1} \neq 0$ is a ground state solution. If $u_{0} \neq 0$, we have

$$
c+o(1)=J\left(u_{n}\right) \rightarrow J\left(u_{0}\right)+\sum_{k=1}^{l} J_{\text {per }}\left(w^{k}\right) \geq(l+1) c .
$$

Therefore, $l=0$ and $J\left(u_{n}\right) \rightarrow J\left(u_{0}\right)=c$, so $u_{0}$ is a ground state solution. Suppose that $V_{\text {loc }}<0$ for a.e. $x \in \mathbb{R}^{3}$. We denote $c_{\text {per }}:=\inf _{\mathcal{N}_{\text {per }}} J_{\text {per }}$, where $\mathcal{N}_{\text {per }}=\left\{u \neq 0: J_{\text {per }}^{\prime}(u)(u)=0\right\}$. Because $V_{\text {loc }}<0$, we have $J(u)<J_{\text {per }}(u)$. From the above arguments, there exists a ground state solution $u_{\text {per }}$ of $J_{\text {per }}$ such that $J_{\text {per }}\left(u_{\text {per }}\right)=c_{\text {per }}$. Let $t>0$ be such that $t u_{\text {per }} \in \mathcal{N}$. Then

$$
c \leq J\left(t u_{\mathrm{per}}\right)<J_{\mathrm{per}}\left(t u_{\mathrm{per}}\right) \leq J_{\mathrm{per}}\left(u_{\mathrm{per}}\right)=c_{\mathrm{per}} .
$$

If $u_{0}=0$, we have

$$
c+o(1)=J\left(u_{n}\right)=\sum_{k=1}^{l} J_{\mathrm{per}}\left(w^{k}\right) \geq l c_{\mathrm{per}}>l c
$$

which implies $l=0$ and $J\left(u_{n}\right) \rightarrow c=0$. Obviously, it is a contraction. Therefore, $u_{0} \neq 0$ and note that

$$
c+o(1)=J\left(u_{n}\right) \rightarrow J\left(u_{0}\right)+\sum_{k=1}^{l} J_{\mathrm{per}}\left(w^{k}\right) \geq c+l c_{\mathrm{per}}
$$

So $l=0$ which means that $J\left(u_{n}\right) \rightarrow J\left(u_{0}\right)=c$ and $u_{0}$ is a ground solution.

## 4. Proof of Theorem 1.2

Observe that the condition $\inf \sigma\left(-\Delta+V_{\text {per }}\right)>0$ implies that the norm $\|u\|_{0}^{2}=$ $\left.\left.\int_{\mathbb{R}^{3}}| | \xi\right|^{2 \alpha}|\hat{u}(\xi)|^{2}+V_{\text {per }}(x) u^{2}\right) d x$ is equivalent to $\|\cdot\|$ in $H^{\alpha}\left(\mathbb{R}^{3}\right)$.

Proof of Theorem 1.2. By contradiction, suppose that there exists a ground state solution $u_{0} \in \mathcal{N}$ of $J$. Similar to the proof of [37, Theorem 3.2], there exists $t_{p}>0$ such that $t_{p} u_{0} \in \mathcal{N}_{\text {per }}$. Since $V_{\text {loc }}>0$ for a.e. $x \in \mathbb{R}^{3}$, one has $\int_{\mathbb{R}^{3}} V_{\text {loc }}(x) u_{0}^{2}>0$. Then by Lemma 3.4, we obtain that (A13) holds; that is, for $t \in(0, \infty) \backslash\{1\}$ and $u \in \mathcal{N}$,

$$
\frac{t^{4}-1}{4} I^{\prime}(u)(u)-I(t u)+I(u)<0
$$

Since $u_{0} \in \mathcal{N}$, one has

$$
\begin{aligned}
& J\left(u_{0}\right)-J\left(t_{p} u_{0}\right) \\
& =\frac{1-t_{p}^{2}}{2}\left\|u_{0}\right\|^{2}-\frac{t_{p}^{4}-1}{4} \int_{\mathbb{R}^{3}} K(x) \phi_{u}^{\beta} u_{0}^{2}-I\left(u_{0}\right)+I\left(t_{p} u_{0}\right) \\
& =\left(\frac{1-t_{p}^{2}}{2}\right)^{2}\left\|u_{0}\right\|^{2}-\frac{t_{p}^{4}-1}{4} J^{\prime}\left(u_{0}\right) u_{0}-\frac{t_{p}^{4}-1}{4} I^{\prime}\left(u_{0}\right)\left(u_{0}\right)+I\left(t_{p} u_{0}\right)-I\left(u_{0}\right) \\
& =\frac{\left(1-t_{p}^{2}\right)^{2}}{4}\left\|u_{0}\right\|^{2}-\frac{t_{p}^{4}-1}{4} I^{\prime}\left(u_{0}\right)\left(u_{0}\right)+I\left(t_{p} u_{0}\right)-I\left(u_{0}\right)
\end{aligned}
$$

$$
>\frac{\left(t_{p}^{2}-1\right)^{2}}{4}\left\|u_{0}\right\|^{2}
$$

Then $J\left(u_{0}\right)>J\left(t_{p} u_{0}\right)$. Therefore,

$$
\begin{equation*}
c_{\text {per }}:=\inf _{\mathcal{N}_{\text {per }}} J_{\text {per }} \leq J_{\text {per }}\left(t_{p} u_{0}\right)<J\left(t_{p} u_{0}\right)<J\left(u_{0}\right)=c \tag{4.1}
\end{equation*}
$$

On the other hand, let $u \in \mathcal{N}_{\text {per }}$ and $u_{y}(\cdot)=u(\cdot-y)$ for $y \in \mathbb{Z}^{3}$. Then there exist $t_{u}>0$ such that $t_{u} u_{y} \in \mathcal{N}$. Observe that

$$
\begin{align*}
J_{\mathrm{per}}(u) & =J_{\mathrm{per}}\left(u_{y}\right) \geq J_{\mathrm{per}}\left(t_{u} u_{y}\right) \\
& =J\left(t_{u} u_{y}\right)-\frac{1}{2} \int_{\mathbb{R}^{3}} V_{\mathrm{loc}}(x)\left(t_{u} u_{y}\right)^{2}  \tag{4.2}\\
& \geq c-\int_{\mathbb{R}^{3}} V_{\mathrm{loc}}(x)\left(t_{u} u_{y}\right)^{2} .
\end{align*}
$$

We are going to show that

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} V_{\mathrm{loc}}(x)\left(t_{u} u_{y}\right)^{2} \rightarrow 0 \tag{4.3}
\end{equation*}
$$

Indeed,

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} V_{\mathrm{loc}}(x)\left(t_{u} u_{y}\right)^{2}=t_{u}^{2} \int_{\mathbb{R}^{3}} V_{\mathrm{loc}}(x+y) u^{2} \tag{4.4}
\end{equation*}
$$

By (A8), one has

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} V_{\mathrm{loc}}(x+y) u^{2} \rightarrow 0 \quad \text { as }|y| \rightarrow \infty \tag{4.5}
\end{equation*}
$$

If $t_{u} \rightarrow \infty$, then by (A4) one has

$$
\begin{aligned}
o(1)= & \frac{c}{t_{u}^{q}} \leq \frac{J\left({ }_{u} u_{y}\right)}{t_{u}^{q}} \\
= & \frac{1}{2}\left\|u_{y}\right\|^{2} t_{u}^{2-q}+\frac{1}{4} t_{u}^{4-q} \int_{\mathbb{R}^{3}} K(x) \phi_{u_{y}}^{\beta} u_{y}^{2} \\
& -\int_{\mathbb{R}^{3}} \frac{F\left(x, t_{u} u_{y}\right)}{\left|t_{u} u_{y}\right|^{q}} u_{y}^{q} d x+\int_{\mathbb{R}^{3}} \frac{1}{q} \Gamma(x)\left|u_{y}\right|^{q} d x \rightarrow-\infty .
\end{aligned}
$$

Thus $t_{u}<\infty$, which, together with 4.5, implies that 4.3) holds. By (4.2), we have

$$
J_{\mathrm{per}}(u) \geq c-\int_{\mathbb{R}^{3}} V_{\mathrm{loc}}(x)\left(t_{u} u_{y}\right)^{2} \rightarrow c
$$

Taking infimum over all $u \in \mathcal{N}_{\text {per }}$ we have $c_{\text {per }} \geq c$, which contradicts with 4.1). The proof is complete.

## 5. Proof of Theorem 1.3

Let $\beta:=\inf _{\mathcal{N}}\|u\|>0$, where $\|\cdot\|$ is the norm defined by 2.2 . Theorem 1.1 provides that $c$ is attained at some function in $\mathcal{N}$. By $\tau_{u}$ we denote the $\mathbb{Z}^{3}$-action on $H^{\alpha}\left(\mathbb{R}^{3}\right)$, i.e. $\tau_{k} u=u(\cdot-k)$. Obviously, $\tau_{k} \tau_{-k} u=u$. For given $k \in \mathbb{Z}^{3}$, let us consider $\tau_{k}$ as an operator $\tau_{k}: H^{\alpha}\left(\mathbb{R}^{3}\right) \rightarrow H^{\alpha}\left(\mathbb{R}^{3}\right)$. It is easy to prove that $\tau_{k}$ is linear.

Lemma $5.1(\boxed{6})$. . For every $u, v \in H^{\alpha}\left(\mathbb{R}^{N}\right)$ and $k \in \mathbb{Z}^{N}$, we have

$$
\left(\tau_{k} u, v\right)=\left(u, \tau_{-k} v\right)
$$

Remark 5.2. When $N=3$, the conclusion in Lemma 5.1 still holds. Obviously, by the definition of $\tau_{k}$, we have $\left\|\tau_{k}\right\|=\|u\|$. Thus $\tau_{k}$ is a bounded operator and $\left\|\tau_{k}\right\|=1$. Then we may consider an adjoint operator $\tau_{k}^{*}: H^{\alpha}\left(\mathbb{R}^{3}\right) \rightarrow H^{\alpha}\left(\mathbb{R}^{3}\right)$. From Lemma5.1 we have $\tau_{k}^{*}=\tau_{-k}$. Moreover, $\tau_{k}$ is an isomorphism and $\tau_{k}^{-1}=\tau_{-k}=\tau_{k}^{*}$. So $\tau_{k}$ is an orthogonal operator.

The proof of Lemmas 5.3, 5.5, 5.6 are similar to those in [6, Lemmas 5.3, 5.4 5.6, 5.7]. So we omit them here.

Lemma 5.3. Let $\alpha \in(0,1]$.
(1) The functional $J$ is $\mathbb{Z}^{3}$-invariant;
(2) $\mathcal{N}$ is $\mathbb{Z}^{3}$-invariant.

Remark 5.4. Lemma 5.1 implies that the unit sphere $S^{1}$ is $\mathbb{Z}^{3}$-invariant.
From [5] we know for each $u \in H^{\alpha}\left(\mathbb{R}^{3}\right)$ there is a unique number $t(u)>0$ such that $t(u) u \in \mathcal{N}$ and the function $m: S^{1} \rightarrow \mathcal{N}$ given by $m(u)=t(u) u$ is a homeomorphism. The inverse $m^{-1}: \mathcal{N} \rightarrow S^{1}$ is given by $m^{-1}(u)=\frac{u}{\|u\|}$. Assume that $A, B \subset H^{\alpha}\left(\mathbb{R}^{3}\right)$ are $\mathbb{Z}^{3}$-invariant subsets. We say that the function $h: A \rightarrow B$ is $\mathbb{Z}^{3}$-equivariant if

$$
h\left(\tau_{k} u\right)=\tau_{k} h(u)
$$

for any $u \in A$ and $k \in \mathbb{Z}^{3}$.
Lemma 5.5. The following four functions are $\mathbb{Z}^{3}$-equivariant:
(1) $m: S^{1} \rightarrow \mathcal{N}$,
(2) $m^{-1}: \mathcal{N} \rightarrow S^{1}$,
(3) $\nabla J: H^{\alpha}\left(\mathbb{R}^{3}\right) \rightarrow H^{\alpha}\left(\mathbb{R}^{3}\right)$,
(4) $\nabla(J \circ m): S^{1} \rightarrow H^{\alpha}\left(\mathbb{R}^{3}\right)$

Lemma 5.6. The function $m^{-1}: \mathcal{N} \rightarrow S^{1}$ is Lipschitz continuous.
We denote $\mathcal{O}(u)=\left\{u(\cdot-k): k \in \mathbb{Z}^{3}\right\}$. If $u$ is a solution of 1.1) and $k \in \mathbb{Z}^{3}$, then $u(\cdot-k)$ is also a solution, provided that $V_{\text {loc }} \equiv 0$. Therefore, all elements of the orbit $\mathcal{O}(u)$ of $u$ under the $\mathbb{Z}^{3}$-action are solutions. We define that $u_{1}$ and $u_{2}$ are geometrically distinct if their orbits satisfy $\mathcal{O}\left(u_{1}\right) \cap \mathcal{O}\left(u_{2}\right)=\emptyset$.

To prove Theorem 1.3 we use the method introduced by Szulkin and Weth 34. Let $\mathcal{L}=\left\{u \in S^{1}:(J \circ m)^{\prime}(u)=0\right\}$. Take a set $F \subset \mathcal{L}$ such that $F=-F$ and for each orbit $\mathcal{O}(w)$ there is a unique representative $v \in F$. Observe the [24, Theorem 3.1(b)], we know that to prove Theorem 1.3 it is sufficient to show $J \circ m$ has infinitely many geometrically distinct critical points. Therefore, we only need to show that $F$ is infinite.

As in [34, we put

$$
\kappa=\inf _{v, w \in \mathcal{L}, v \neq w}\|v-w\|>0
$$

Next, we show a lemma which is crucial to the proof of the Theorem 1.3 .
Lemma 5.7. Let $d \geq c$. If $\left(v_{n}^{1}\right),\left(v_{n}^{2}\right) \subset S^{1}$ are two $(P S)$ sequences for $J \circ m$ such that $J\left(m\left(v_{n}^{i}\right)\right) \leq d, i=1,2$, then

$$
\left\|v_{n}^{1}-v_{n}^{2}\right\| \rightarrow 0
$$

or

$$
\liminf _{n \rightarrow \infty}\left\|v_{n}^{1}-v_{n}^{2}\right\| \geq \rho(d)>0
$$

where $\rho(d)$ depends only on $d$, but not on the particular choice of the sequences.

Proof. Note that $u_{n}^{i}=m\left(v_{n}^{i}\right), i=1,2$ are $(P S)$ sequences for $J$. Moreover, they are bounded in $H^{\alpha}\left(\mathbb{R}^{3}\right)$, since $J$ is coercive on $\mathcal{N}$. Therefore $\left(u_{n}^{1}\right)$ and $\left(u_{n}^{2}\right)$ are bounded in $L^{2}\left(\mathbb{R}^{3}\right)$, say $\left|u_{n}^{1}\right|_{2}+\left|u_{n}^{2}\right|_{2} \leq M$ for some $M>0$.
Case 1: Assume that $\left|u_{n}^{1}-u_{n}^{2}\right|_{p} \rightarrow 0$. Fix $\epsilon>0$, by (A2) and (A9) we have

$$
\begin{aligned}
&\left\|u_{n}^{1}-u_{n}^{2}\right\|^{2} \\
&= J^{\prime}\left(u_{n}^{1}\right)\left(u_{n}^{1}-u_{n}^{2}\right)-J^{\prime}\left(u_{n}^{2}\right)\left(u_{n}^{1}-u_{n}^{2}\right)+\int_{\mathbb{R}^{3}}\left[f\left(x, u_{n}^{1}\right)-f\left(x, u_{n}^{2}\right)\right]\left(u_{n}^{1}-u_{n}^{2}\right) d x \\
&-\int_{\mathbb{R}^{3}} \Gamma(x)\left(\left|u_{n}^{1}\right|^{q-2} u_{n}^{1}-\left|u_{n}^{2}\right|^{q-2} u_{n}^{2}\right)\left(u_{n}^{1}-u_{n}^{2}\right) d x \\
&+\int_{\mathbb{R}^{3}} K(x)\left(\phi_{u_{n}^{1}}^{\beta} u_{n}^{1}-\phi_{u_{n}^{2}}^{\beta} u_{n}^{2}\right)\left(u_{n}^{1}-u_{n}^{2}\right) \\
& \leq \epsilon\left\|u_{n}^{1}-u_{n}^{2}\right\|+\int_{\mathbb{R}^{3}}\left(\epsilon\left(\left|u_{n}^{1}\right|+\left|u_{n}^{2}\right|\right)+C_{\epsilon}\left(\left|u_{n}^{1}\right|^{p-1}+\left|u_{n}^{2}\right|^{p-1}\right)\right)\left|u_{n}^{1}-u_{n}^{2}\right| d x \\
&-\int_{\mathbb{R}^{3}} \Gamma(x)\left(\left|u_{n}^{1}\right|^{q-2} u_{n}^{1}-\left|u_{n}^{2}\right|^{q-2} u_{n}^{2}\right)\left(u_{n}^{1}-u_{n}^{2}\right) d x \\
&+\int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}^{1}}^{\beta} u_{n}^{1}\left(u_{n}^{1}-u_{n}^{2}\right) d x+\int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}^{2}}^{\beta} u_{n}^{2}\left(u_{n}^{1}-u_{n}^{2}\right) d x \\
& \leq\left(1+C_{0}\right) \epsilon\left\|u_{n}^{1}-u_{n}^{2}\right\|+D_{\epsilon}\left|u_{n}^{1}-u_{n}^{2}\right|_{p}+C_{1}|\Gamma|_{\infty}\left|u_{n}^{1}-u_{n}^{2}\right|_{q}^{q} \\
&+\int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}^{1}}^{\beta} u_{n}^{1}\left(u_{n}^{1}-u_{n}^{2}\right) d x+\int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}^{2}}^{\beta} u_{n}^{2}\left(u_{n}^{1}-u_{n}^{2}\right) d x
\end{aligned}
$$

for each $n \geq n_{\epsilon}$ and some constants $C_{0}, C_{1}, D_{\epsilon}>0$. By assumption we have that $D_{\epsilon}\left|u_{n}^{1}-u_{n}^{2}\right|_{p} \rightarrow 0$. Observe that

$$
C_{1}|\Gamma|_{\infty}\left|u_{n}^{1}-u_{n}^{2}\right|_{q}^{q} \leq C_{1}|\Gamma|_{\infty}\left|u_{n}^{1}-u_{n}^{2}\right|_{2}^{\theta q}\left|u_{n}^{1}-u_{n}^{2}\right|_{p}^{(1-\theta) q}
$$

where $\theta \in(0,1)$ is such that $\frac{1}{q}=\frac{\theta}{2}+\frac{1-\theta}{p}$. Thus

$$
C_{1}|\Gamma|_{\infty}\left|u_{n}^{1}-u_{n}^{2}\right|_{q}^{q} \leq C_{1}|\Gamma|_{\infty} M^{\theta q}\left|u_{n}^{1}-u_{n}^{2}\right|_{p}^{(1-\theta) q} \rightarrow 0 .
$$

And by Hölder inequality,

$$
\begin{align*}
& \int_{\mathbb{R}^{N}} K(x) \phi_{u_{n}^{1}}^{\beta} u_{n}^{1}\left(u_{n}^{1}-u_{n}^{2}\right) d x \\
& \leq|K|_{\infty}\left(\int_{\mathbb{R}^{3}}\left|\phi_{u_{n}^{1}}^{\beta}\right|^{2^{*}} d x\right)^{1 / 2_{\beta}^{*}}\left(\int_{\mathbb{R}^{3}}\left|u_{n}^{1}\left(u_{n}^{1}-u_{n}^{2}\right)\right|^{\frac{6}{3+2 \beta}} d x\right)^{\frac{3+2 \beta}{6}} \\
& \leq C_{2}| | \phi_{u_{n}^{1}}^{\beta} \|_{D^{\beta, 2}}\left(\int_{\mathbb{R}^{3}}\left|u_{n}^{1} \frac{6}{3+2 \beta}\right| u_{n}^{1}-\left.u_{n}^{2}\right|^{\frac{6}{3+2 \beta}} d x\right)^{\frac{3+2 \beta}{6}}  \tag{5.1}\\
& \leq C_{3}\left\|\phi_{u_{n}^{\beta}}^{\beta}\right\|_{D^{\beta, 2}}\left|u_{n}^{1}\right|_{\frac{12}{3+2 \beta}}\left|u_{n}^{1}-u_{n}^{2}\right|_{\frac{12}{3+2 \beta}} \\
& \leq C_{4}\left|u_{n}^{1}-u_{n}^{2}\right|_{\frac{12}{3+2 \beta}},
\end{align*}
$$

where $C_{2}, C_{3}, C_{4}>0$ are constants. By $\beta \in(0,1]$, we have $\frac{12}{3+2 \beta}<4<p$. Thus there exists $\theta^{\prime} \in(0,1)$ such that

$$
\frac{1}{\frac{12}{3+2 \beta}}=\frac{\theta^{\prime}}{2}+\frac{1-\theta^{\prime}}{p}
$$

and

$$
\begin{equation*}
\left|u_{n}^{1}-u_{n}^{2}\right|_{\frac{12}{3+2 \beta}}^{3+1} \leq\left|u_{n}^{1}-u_{n}^{2}\right|_{2}^{\theta^{\prime}}\left|u_{n}^{1}-u_{n}^{2}\right|_{p}^{1-\theta^{\prime}} \leq M^{\theta^{\prime}}\left|u_{n}^{1}-u_{n}^{2}\right|_{p}^{1-\theta^{\prime}} \tag{5.2}
\end{equation*}
$$

Then

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}^{1}}^{\beta} u_{n}^{1}\left(u_{n}^{1}-u_{n}^{2}\right) d x \leq C_{4} M^{\theta^{\prime}}\left|u_{n}^{1}-u_{n}^{2}\right|_{p}^{1-\theta^{\prime}} \rightarrow 0 \tag{5.3}
\end{equation*}
$$

Similarly we have

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}^{2}}^{\beta} u_{n}^{2}\left(u_{n}^{1}-u_{n}^{2}\right) d x \leq C_{5} M^{\theta^{\prime \prime}}\left|u_{n}^{1}-u_{n}^{2}\right|_{p}^{1-\theta^{\prime \prime}} \rightarrow 0 \tag{5.4}
\end{equation*}
$$

where $C_{5}>0$ is some constant and $\theta^{\prime \prime} \in(0,1)$. Finally

$$
\begin{align*}
\limsup _{n \rightarrow \infty}\left\|u_{n}^{1}-u_{n}^{2}\right\|^{2} \leq & \limsup _{n \rightarrow \infty}\left(1+C_{0}\right) \epsilon\left\|u_{n}^{1}-u_{n}^{2}\right\| \\
& +\limsup _{n \rightarrow \infty} D_{\epsilon}\left|u_{n}^{1}-u_{n}^{2}\right|_{p}+\limsup _{n \rightarrow \infty} C_{1}|\Gamma|_{\infty}\left|u_{n}^{1}-u_{n}^{2}\right|_{q}^{q} \\
& +\limsup _{n \rightarrow \infty} C_{4} M^{\theta^{\prime}}\left|u_{n}^{1}-u_{n}^{2}\right|_{p}^{1-\theta^{\prime}}+\limsup _{n \rightarrow \infty} C_{5}^{\theta^{\prime \prime}}\left|u_{n}^{1}-u_{n}^{2}\right|_{p}^{1-\theta^{\prime \prime}} \\
= & \left(1+C_{0}\right) \epsilon \limsup _{n \rightarrow \infty}\left\|u_{n}^{1}-u_{n}^{2}\right\| \tag{5.5}
\end{align*}
$$

for every $\epsilon>0$. Thus $\lim _{n \rightarrow \infty}\left\|u_{n}^{1}-u_{n}^{2}\right\|=0$. Finally

$$
\left\|v_{n}^{1}-v_{n}^{2}\right\|=\left\|m^{-1}\left(u_{n}^{1}\right)-m^{-1}\left(u_{n}^{2}\right)\right\| \leq L\left\|u_{n}^{1}-u_{n}^{2}\right\| \rightarrow 0
$$

where $L>0$ is a Lipschitz constant for $m^{-1}$.
Case 2: Assume that $\left|u_{n}^{1}-u_{n}^{2}\right|_{p} \nrightarrow 0$. By the Lions lemma, there are $y_{n} \in \mathbb{R}^{3}$ such that

$$
\int_{B\left(y_{n}, 1\right)}\left|u_{n}^{1}-u_{n}^{2}\right|^{2} d x=\sup _{y \in \mathbb{R}^{3}} \int_{B(y, 1)}\left|u_{n}^{1}-u_{n}^{2}\right|^{2} d x \geq \epsilon
$$

for some $\epsilon>0$. In view of Lemma 5.5 we can assume that $\left(y_{n}\right)$ is bounded. Therefore, up to a subsequence we have

$$
u_{n}^{1} \rightharpoonup u^{1}, u_{n}^{2} \rightharpoonup u^{2}
$$

where $u^{1} \neq u^{2}$ and $J^{\prime}\left(u^{1}\right)=J^{\prime}\left(u^{2}\right)=0$ and

$$
\left\|u_{n}^{1}\right\| \rightarrow \alpha^{1}, \| u_{n}^{2} \rightarrow \alpha^{2}
$$

where $\beta \leq \alpha^{i} \leq \nu(d)=\sup \{\|u\|: u \in \mathcal{N}, J(u) \leq d\}, i=1,2$. Suppose that $u^{1} \neq 0$ and $u^{2} \neq 0$. Then $u^{i} \in \mathcal{N}$ for $i=1,2$. Moreover

$$
v^{i}=m^{-1}\left(u^{i}\right) \in S^{1}, \quad i=1,2 \text { and } v^{1} \neq v^{2} .
$$

Therefore,

$$
\liminf _{n \rightarrow \infty}\left\|v_{n}^{1}-v_{n}^{2}\right\|=\liminf _{n \rightarrow \infty}\left\|\frac{u_{n}^{1}}{\left\|u_{n}^{1}\right\|}-\frac{u_{n}^{1}}{\left\|u_{n}^{2}\right\|}\right\| \geq\left\|\frac{u^{1}}{\alpha^{1}}-\frac{u^{2}}{\alpha^{2}}\right\|=\left\|\beta_{1} v^{1}-\beta_{2} v^{2}\right\|
$$

where $\beta_{i}=\frac{\left\|u^{i}\right\|}{\alpha^{i}} \geq \frac{\beta}{\nu(d)}, i=1,2$. Of course $\left\|v^{1}\right\|=\left\|v^{2}\right\|=1$. So

$$
\liminf _{n \rightarrow \infty}\left\|v_{n}^{1}-v_{n}^{2}\right\| \geq\left\|\beta_{1} v^{1}-\beta_{2} v^{2}\right\| \geq \min _{i=1,2}\left\{\beta_{i}\right\}\left\|v^{1}-v^{2}\right\| \geq \frac{\beta \kappa}{\nu(d)}
$$

If $u^{2}=0$, then $u^{1} \neq u^{2}=0$. Therefore,

$$
\liminf _{n \rightarrow \infty}\left\|v_{n}^{1}-v_{n}^{2}\right\|=\liminf _{n \rightarrow \infty}\left\|\frac{u_{n}^{1}}{\left\|u_{n}^{1}\right\|}-\frac{u_{n}^{1}}{\left\|u_{n}^{2}\right\|}\right\| \geq\left\|\frac{u^{1}}{\alpha^{1}}-\frac{u^{2}}{\alpha^{2}}\right\|=\left\|\frac{u^{1}}{\alpha^{1}}\right\| \geq \frac{\beta}{\nu(d)}
$$

The case $u^{1}=0$ is similar, the proof is complete.

Proof of Theorem 1.3. Suppose by contraction that $F$ is finite. Because the unit sphere $S^{1} \subset H^{\alpha}\left(\mathbb{R}^{3}\right)$ is a Finsler $C^{1,1}$-manifold and by [29, Lemma II.3.9], we obtain $J \circ m: S^{1} \rightarrow \mathbb{R}$ has a pseudo-gradient vector field. Because of the obtained discreteness of $(P S)$ sequences in Lemma 5.7, we can apply the methods in 34 , Lemma 2.15, Lemma 2.16 and Theorem 1.2] in our case. In fact, for every $k \in \mathbb{N}$ there is $u \in S^{1}$ such that

$$
(J \circ m)^{\prime}(u)=0 \quad \text { and } \quad J(m(u))=c_{k},
$$

where

$$
c_{k}=\inf \left\{d \in \mathbb{R}: \gamma\left(\left\{v \in S^{1}: J(m(v)) \leq d\right\}\right) \geq k\right\}
$$

is the Lusternik-Schnirelmann value and $\gamma$ denotes the Krasnoselskii genus (see [29]). Moreover $c_{k} \leq c_{k+1}$, thus we get contradiction (for the detailed arguments see (34). The proof is complete.

## 6. Proof of Theorem 1.4

Before showing the proof, We state the following form of the Sobolev-GagliardoNirenberg inequality.
Lemma 6.1 ([28, Proposition II.3]). Let $r>1$. Then there is a positive constant $C>0$ such that for every function $u \in H^{\alpha}\left(\mathbb{R}^{N}\right)$ there holds

$$
|u|_{r+1}^{r+1} \leq C\|u\|_{H^{\alpha^{\alpha}}}^{\frac{(r-1) N}{\alpha^{\alpha}}}|u|_{2}^{r+1-\frac{(r-1) N}{\alpha}} .
$$

In [11], it was proven that $E^{\alpha}$ is compactly embedded into $L^{r}\left(\mathbb{R}^{N}\right)$ for $N>\alpha$ and $2 \leq r<2_{\alpha}^{*}$. We will show that the method introduced by Secchi in [28] may be applied in this case.

Proof of Theorem 1.4. In view of Lemma 3.3. we obtain a bounded minimizing sequence $\left(u_{n}\right) \subset \mathcal{N}$, i.e.

$$
J\left(u_{n}\right) \rightarrow \inf _{\mathcal{N}} J=: c, \quad J^{\prime}\left(u_{n}\right) \rightarrow 0
$$

Then we assume that $u_{n} \rightharpoonup u_{0}$ in $E^{\alpha}\left(\mathbb{R}^{3}\right)$ and $u_{n} \rightarrow u_{0}$ in $L_{\text {loc }}^{r}\left(\mathbb{R}^{3}\right)$ for $1 \leq r<2_{\alpha}^{*}$. It is a easy to proof that $u_{0}$ is a weak solution to our problem. Then we just need to check whether $u_{0} \neq 0$. Observe that for $n \geq n_{0}$

$$
\begin{aligned}
\frac{c}{2} \leq & J\left(u_{n}\right)=J\left(u_{n}\right)-\frac{1}{2} J^{\prime}\left(u_{n}\right)\left(u_{n}\right) \\
= & \frac{1}{2} \int_{\mathbb{R}^{3}} f\left(x, u_{n}\right) u_{n}-2 F\left(x, u_{n}\right) d x-\left(\frac{1}{2}-\frac{1}{q}\right) \int_{\mathbb{R}^{3}} \Gamma(x)|u|^{q} d x \\
& -\frac{1}{4} \int_{\mathbb{R}^{3}} K(x) \Phi_{u}^{\beta}(x) u^{2} d x \\
\leq & \frac{1}{2} \int_{\mathbb{R}^{3}} f\left(x, u_{n}\right) u_{n}-2 F\left(x, u_{n}\right) d x \\
\leq & \frac{1}{2} \int_{\mathbb{R}^{3}} f\left(x, u_{n}\right) u_{n} d x \\
\leq & \frac{1}{2} \int_{\mathbb{R}^{3}}\left(\epsilon\left|u_{n}\right|^{2}+C_{\epsilon}\left|u_{n}\right|^{p}\right) d x,
\end{aligned}
$$

where $n_{0} \geq 1$ is large enough. Then

$$
\frac{c}{2} \leq \frac{\epsilon}{2}\left|u_{n}\right|_{2}^{2}+\frac{C_{\epsilon}}{2}\left|u_{n}\right|_{p}^{p}
$$

By Lemma 6.1. we have

$$
\frac{c}{2} \leq \frac{\epsilon}{2}\left|u_{n}\right|_{2}^{2}+\frac{C C_{\epsilon}}{2}\left\|u_{n}\right\|_{H^{\alpha}\left(\mathbb{R}^{3}\right)}^{\frac{3(p-2)}{\alpha}}\left|u_{n}\right|_{2}^{p-\frac{3(p-2)}{\alpha}} .
$$

Observe that the boundedness of $\left(u_{n}\right)$ in $H^{\alpha}$ with respect to the norm $\|\cdot\|$ implies the boundedness of $\left(u_{n}\right)$ in $E^{\alpha}$ with respect to the classical norm $\|\cdot\|_{H^{\alpha}}\left(\mathbb{R}^{3}\right)$. Therefore $\left\|u_{n}\right\|_{H^{\alpha}}\left(\mathbb{R}^{3}\right) \leq D$ for some $D>0$. Thus

$$
\frac{c}{2} \leq \frac{\epsilon}{2}\left|u_{n}\right|_{2}^{2}+\frac{C C_{\epsilon}}{2} D^{\frac{3(p-2)}{\alpha}}\left|u_{n}\right|_{2}^{p-\frac{3(p-2)}{\alpha}}
$$

Denote $\hat{C}_{\epsilon}=\frac{C \cdot C_{\epsilon}}{2} D^{\frac{3(p-2)}{\alpha}}$. Then

$$
\frac{c}{2} \leq \frac{\epsilon}{2}\left|u_{n}\right|_{2}^{2}+\hat{C}_{\epsilon}\left|u_{n}\right|_{2}^{p-\frac{3(p-2)}{\alpha}}
$$

Take any $\epsilon \leq \frac{c}{2\left(\sup _{n}\left\|u_{n}\right\|\right)^{2}}$. Then

$$
\frac{c}{2} \leq \frac{c\left|u_{n}\right|_{2}^{2}}{4\left(\sup _{n}\left\|u_{n}\right\|\right)^{2}}+\hat{C}_{\epsilon}\left|u_{n}\right|_{2}^{p-\frac{3(p-2)}{\alpha}} .
$$

Obviously,

$$
\frac{\left|u_{n}\right|_{2}^{2}}{\left(\sup _{n}\left\|u_{n}\right\|\right)^{2}} \leq 1
$$

and therefore

$$
\frac{c}{2} \leq \frac{c}{4}+\hat{C}_{\epsilon}\left|u_{n}\right|_{2}^{p-\frac{3(p-2)}{\alpha}}
$$

Finally

$$
\frac{c}{4} \leq \hat{C}_{\epsilon}\left|u_{n}\right|_{2}^{p-\frac{3(p-2)}{\alpha}}
$$

So

$$
\ln \frac{c}{4} \leq\left(p-\frac{3(p-2)}{\alpha}\right) \ln \left(C_{1}\left|u_{n}\right|_{2}\right)
$$

where $C_{1}=\hat{C}_{\epsilon}^{p-\frac{3(p-2)}{\alpha}}$. Thus

$$
\frac{\alpha}{\alpha p-3(p-2)} \ln \frac{c}{4} \leq \ln \left(C_{1}\left|u_{n}\right|_{2}\right)
$$

And finally

$$
\left.\left|u_{n}\right|_{2}^{2} \geq\left(\frac{1}{C_{1}} \exp \left(\frac{\alpha}{\alpha p-3(p-2)}\right) \ln \frac{c}{4}\right)\right)^{2}=: \hat{c}>0
$$

Take any $R>0$ and observe that

$$
\left|u_{n}\right|_{2}^{2}=\int_{B(0, R)}\left|u_{n}\right|^{2} d x+\int_{\mathbb{R}^{3} \backslash B(0, R)}\left|u_{n}\right|^{2} d x
$$

Assume by a contraction that $u_{0}=0$, then $u_{n} \rightarrow 0$ in $L_{\text {loc }}^{2}\left(\mathbb{R}^{3}\right)$. Then for every $R>0$ there is $n_{0}$ such that for $n \geq n_{0}$ we have

$$
\int_{B(0, R)}\left|u_{n}\right|^{2} d x \leq \frac{\hat{c}}{2}
$$

Thus

$$
\int_{\mathbb{R}^{3} \backslash B(0, R)}\left|u_{n}\right|^{2} d x \geq \frac{\hat{c}}{2} .
$$

On the other hand

$$
\frac{\hat{c}}{2} \leq \int_{\mathbb{R}^{3} \backslash B(0, R)}\left|u_{n}\right|^{2} d x=\int_{\mathbb{R}^{3} \backslash B(0, R)} \frac{V(x)\left|u_{n}\right|^{2}}{V(x)} d x
$$

$$
\begin{aligned}
& \leq \frac{1}{\inf _{|x| \geq R} V(x)} \int_{\mathbb{R}^{3} \backslash B(0, R)} V(x)\left|u_{n}\right|^{2} d x \\
& \leq \frac{\left\|u_{n}\right\|^{2}}{\inf _{|x| \geq R} V(x)} \leq \frac{\sup _{n}\left\|u_{n}\right\|^{2}}{\inf _{|x| \geq R} V(x)}
\end{aligned}
$$

Taking $R>0$ big enough we obtain a contradiction, since $V(x) \rightarrow \infty$ as $|x| \rightarrow \infty$. Therefore, $u_{0} \neq 0$. The proof is complete.

Acknowledgments. This work was supported by Natural Science Foundation of China 11671403.

## References

[1] Ambrosetti, A.; Ruiz, D.; Multiple bound states for the Schrödinger-Poisson problem, Commun. Contemp. Math., 10 (2008), 391-404.
[2] Azzollini, A.; Pomponio, A.; Ground state solutions for the nonlinear Schrödinger-Mexwell equations, J. Math. Anal. Appl., 345 (2008), 90-108.
[3] Benci, V.; Fortunato, D. F.; An eigenvalue problem for the Schrödinger-Mexwell equations, Topol. Methods. Nonliner Anal., 11 (1998), 283-293.
[4] Benci, V;, Fortunato, D. F.; Solitary waves of the nonlinear Klein-Gordon equation coupled with the Maxwell equations, Rev. Math. Phys., 14 (2002), 409-420.
[5] Bieganowski, B.; Mederski, J.; Nonlinear Schrödinger equations with sum of periodic and vanishing potentials and sign-changing nonlinearities, Com. Pure Appl. Anal., 17 (2018), 143-161.
[6] Bieganowski, B.; Solutions of the fractional Schrödinger equation with a sign-changing nonlinearity, J. Math. Anal. Appl., 450 (2017), 461-479.
[7] Cerami, G., Vaira, G.; Positive solutions for some non-autonomous Schrödinger-Poisson systems, J. Differential Equations 248 (2010), 521-543.
[8] Chen, S. T.; Peng, J. W., Tang, X. H.; Ground state solutions for asymptotically periodic fractional Schrödinger-Poisson problems with saymptotically cubic or super-cubic nonlinearities, Math. Methods Appl. Sci., 40 (2017), 4948-4961.
[9] Che, G. F.; Chen, H. B.; Ground state solutions for a class of semilinear elliptic systems with sum of periodic and vanishing potentials, Topol. Methods. Nonliner Anal., 51 (2017), 215-242.
[10] Che, G. F.; Chen, H. B.; Existence and multiplicity of solutions for semilinear elliptic systems with periodic potential, B. Malylas. Math. Sci. Soc., (2017), 1-20.
[11] Cheng, M.; Bounded state for the fractional Schrödinger equations with unbounded potential, J. Math. Phys., 53 (2012), 043507.
[12] Dipierro, S.; Palatucci, G., Valdinoci, E.; Existence and symmetry results for a Schrödinger type problem involving the fractional Laplacian, Le matematiche, 68 (2013), 201-216.
[13] Duarte, R. C.; Souto, M. A. S. ; Fractional Schrödinger-Poisson equations with general nonlinearity, Electron. J. Differential Equations, 319 (2013), 1-19.
[14] Felmer, P.; Quaas, A.; Tan, J. G.; Positive solutions of nonlinear Schrödinger equation with the fractional Laplacian, Proc. Roy. Soc. Edinburgh A, 142 (2012), 1237-1262.
[15] Huang, L.; Rocha, E. M.; Chen, J.; Two positive solutions of a class of Schrödinger-Poisson system with indefinite nonlinearity, J. Differential Equations, 255 (2013), 2463-2483.
[16] Ji, C.; Ground state sign-changing solutions for a class of nonlinear fractional SchrödingerPoisson system in $\mathbb{R}^{3}$, (2017) arXiv: 1703.03723 v 1 .
[17] Laskin, N.; Fractional quantum mechanics and Lévy path integrals, Phys. Lett. A, 268 (2000) 298-305.
[18] Laskin, N.; Fractional Schrödinger equation, Phys. Rev. E., 66 (2002), 249-264.
[19] Lieb, E. H.; Thomas-Fermi and related theories of atoms and molecules, Rev. Mod. Phys., 53 (1981), 603.
[20] Lions, P. L.; Solutions of Hartree-Fock equations for Coulomb systems, Commun. Math. Phys. 109 (1987) 33-97.
[21] Liu, W.; Existence of multi-bump solutions for the fractional Schrödinger-Poisson system, $J$. Math. Phys., 57 (2016), 779-791.
[22] Liu, Z. S.; Zhang, Z. T.; Huang, S. B.; Existence and nonexistence of positive solutions for a static Schrödinger-Poisson-Slater equation. J. Differential Equations (2018), DOI:10.1016/j.jde.2018.10.048.
[23] Liu, Z. S.; Zhang, J.J.; Multiplicity and concentration of positive solutions for the fractional Schrödinger-Poisson systems with critical growth, ESAIM: Contr. Optim. Calc. Var., 23 (2017), 1515-1542.
[24] Mederski J; Solutions to a nonlinear fractional Schrödinger equation with periodic potential and zero on the boundary of the spectrum, Topol. Methods Nonlinear Anal., 46 (2015), 755-771.
[25] Mederski, J.; Ground states of a system of nonlinear Schrödinger equations with periodic potentials, Comm. Partial Differential Equations, 41 (2016), 1426-1440.
[26] Nezza, E. Di; Palatucci, G.; Valdinoci, E.; Hitchhiker's guide to the fractional Sobolev spaces, Bull. Sci. Math., 136 (2012), 521-573.
[27] Ruiz, D.; The Schrödinger-Poisson equation under the effect of a nonlinear local trem, $J$. Funct. Analysis, 237 (2006), 655-647.
[28] Secchi, S.; Ground state solutions for nonlinear fractional Schrödinger equations in $\mathbb{R}^{N}, J$. Math. Phys., 54 (2013), 031501.
[29] Struwe, M.; Variational Methods, Springer, (2008).
[30] Shi, H. X.; Chen, H. B.; Multiplicity of solutions for a class of fractional Schrödinger equations, Electn. J. Differential Equations, 25 (2015), 1-11.
[31] Shi, H. X.; Chen, H. B.; Existence and multiplicity of solutions for a class of generalized quasilinear Schrodinger equations, J. Math. Anal. Appl., 425 (2017), 578-594.
[32] Sun, J. T., Wu, T. F.; Feng, Z. S.; Multiplicity of positive solutions for nonlinear SchrödingerPoission system, J. Differential Equations, 260 (2016), 586-627.
[33] Sun, J. T.; Wu, T. F.; Feng, Z. S.; Non-autonomous Schrödinger-Poisson system in $\mathbb{R}^{3}$, Discrete Contin. Dyn. Syst., 38 (2018), 1889-1933.
[34] Szulkin, A.; Weth, T.; Ground state solutions for some indefinite variational problems, J. Func. Anal., 257 (2009), 3802-3822.
[35] Teng, K.; Existence of ground state solutions for the nonlinear fractional Schrödinger-Poissom system with critical Sobolev exponent, J. Differential Equations, 261 (2016), 3061-3106.
[36] Teng, K.; He, X. M.; Ground state solutions for fractional Schrödinger equations with critical Sobolev exponent, Com. Pure Appl. Anal., 15 (2016), 991-1008.
[37] Xie, W. H.; Chen, H. B.; Shi, H. X.; Ground state solutions for the nonlinear SchrödingerPoisson systems with sum of periodic and vanishing potentials, Mathe. Methods Appl. Sci., 41 (2018), 144-158.
[38] Zhang, J.; Squassina, M.; Fractional Schrödinger-Poisson systems with a general subcritical or critical nonlinearity, Adv. Non. Stud., 16 (2016), 15-30.
[39] Zhao, L. G.; Zhao, F. K.; On the existence of solutions for the Schrödinger-Poisson equations, J. Math. Anal. Appl., 346 (2008), 155-167.

Mingzhe Yu
School of Mathematics and Statistics, Central South University, Changsha, 410083 Hunan, China

Email address: yumz_math@csu.edu.cn
Haibo Chen (corresponding author)
School of Mathematics and Statistics, Central South University, Changsha, 410083
Hunan, China
Email address: math_chb@163.com


[^0]:    2010 Mathematics Subject Classification. 35B38, 35G99.
    Key words and phrases. Fractional Schrödinger-Poisson system; coercive potential; periodic and localized potential; Nehari manifold; variational method.
    (c) 2019 Texas State University.

    Submitted March 22, 2018. Published September 4, 2019.

