# LOW REGULARITY OF NON- $L^{2}\left(\mathbb{R}^{n}\right)$ LOCAL SOLUTIONS TO GMHD- $\alpha$ SYSTEMS 

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#### Abstract

The Magneto-Hydrodynamic (MHD) system of equations governs viscous fluids subject to a magnetic field and is derived via a coupling of the Navier-Stokes equations and Maxwell's equations. Recently it has become common to study generalizations of fluids-based differential equations. Here we consider the generalized Magneto-Hydrodynamic alpha (gMHD- $\alpha$ ) system, which differs from the original MHD system by including an additional nonlinear terms (indexed by $\alpha$ ), and replacing the Laplace operators by more general Fourier multipliers with symbols of the form $-|\xi|^{\gamma} / g(|\xi|)$. In [8], the problem was considered with initial data in the Sobolev space $H^{s, 2}\left(\mathbb{R}^{n}\right)$ with $n \geq 3$. Here we consider the problem with initial data in $H^{s, p}\left(\mathbb{R}^{n}\right)$ with $n \geq 3$ and $p>2$. Our goal is to minimizing the regularity required for obtaining uniqueness of a solution.


## 1. Introduction

This article concerns the generalized Magneto-Hydrodynamic alpha (gMHD- $\alpha$ ) system of equations, reported in its full generality as,

$$
\begin{gather*}
\partial_{t} v+(u \cdot \nabla) v+\sum_{i=1}^{n} v_{i} \nabla u_{i}-\nu_{1} \mathcal{L}_{1} v+\frac{1}{2} \nabla|B|^{2}=-\nabla p+(B \cdot \nabla) B,  \tag{1.1}\\
\partial_{t} B+(u \cdot \nabla) B-(B \cdot \nabla) u-\nu_{2} \mathcal{L}_{2} B=0  \tag{1.2}\\
v=\left(1-\alpha^{2} \mathcal{L}_{3}\right) u  \tag{1.3}\\
\operatorname{div} u=\operatorname{div} B=0  \tag{1.4}\\
u(0, x)=u_{0}(x), \quad B(0, x)=B_{0}(x), \quad x \in \mathbb{R}^{n} . \tag{1.5}
\end{gather*}
$$

Since these equations govern the motion of fluids subject to a magnetic field, its terms have specific physical meaning: $u$ the fluid velocity, $B$ the magnetic field, and $p$ the scalar-valued pressure of the fluid, $\nu_{1}>0$ is the fluid viscosity, $\nu_{2}>0$ the magnetic diffusion, and $\alpha>0$ a constant coming from varying the Hamiltonian that originally gave rise to the standard MHD equations (see 4]). Finally, the $\mathcal{L}_{i}$ terms are Fourier multipliers with symbol $-|\xi|^{\gamma_{i}} / g_{i}(|\xi|)$, where $g_{i}$ is a positive scalar function and $\gamma_{i}>0$.

The standard MHD system is the special case obtained when setting $\alpha=0$, $g_{1}=g_{2}=1$, and $\gamma_{1}=\gamma_{2}=2$, so that $v=u$ and $\mathcal{L}_{1}=\mathcal{L}_{2}=\Delta$. The existence of

[^0]a global solution up to initial conditions is a classic result, at least in two dimensions. Unfortunately, the presence of nonlinear terms makes the MHD equations particularly complex to solve in arbitrary dimensions, and so a common strategy has been to study modified versions of them.

One modification, termed Lagrangian Averaged MHD- $\alpha$ after the Lagrangian Averaged Navier-Stokes equation, is obtained from Equations 1.1-1.5 by setting $\mathcal{L}_{i}=\Delta$ for $i=1,2,3\left(\gamma_{i}=2\right.$ and $\left.g_{i}=1\right)$. Linshiz and Titi proved the existence of a global solution for smooth initial data in three dimension [4]. Another version is obtained by setting $\alpha=0$ and $g_{1}=g_{2}=1$ and leaving the $\gamma_{i}$ 's unspecified. Zhao and Zhu used these generalized operators to guarantee a global solution to Equations 1.1)-1.5 in the case of $n=3$, provided that $g_{1}=g_{2}=g_{3}=1$, $\gamma_{1}=\gamma_{2}=n / 2$, and $\gamma_{3}=2$ [14].

The first incorporation of a non-constant value for any $g_{i}$ appeared in 9, where Tao proved the existence of a unique global solution to the generalized Navier-Stokes equation ( $B_{0}=\alpha=0$ ) when $\gamma_{1}=n / 2+1$ and when $g_{1}$ is a radial non-decreasing function bounded below satisfying

$$
\begin{equation*}
\int_{1}^{\infty} \frac{d s}{s g_{1}(s)^{4}}=\infty \tag{1.6}
\end{equation*}
$$

the prototypical example of which is essentially a logarithm.
Wu obtained a similar result for the generalized MHD system in [12, specifically showing that there is a unique global solution provided that $u_{0}, B_{0} \in H^{r, 2}\left(\mathbb{R}^{n}\right)$ with $r>n / 2+1 ; \gamma_{1} \geq n / 2+1, \gamma_{2}>0$, and $\gamma_{1}+\gamma_{2} \geq 1$; and $g_{1}, g_{2}$ are non-decreasing, bounded below by 1 , and satisfy

$$
\begin{equation*}
\int_{1}^{\infty} \frac{d s}{s\left(g_{1}(s)+g_{2}(s)\right)^{2}}=\infty \tag{1.7}
\end{equation*}
$$

This work was ultimately extended to the gMHD- $\alpha$ system in [13], where Yamazaki obtained a unique global solution in three dimensions provided that $\gamma_{1}+\gamma_{2}+\gamma_{3} \geq 5$, $\min \left\{\gamma_{1}, \gamma_{3}\right\}>\gamma_{2}>0, \gamma_{3}+2 \gamma_{1}>3$, and the $g_{i}$ satisfy

$$
\begin{equation*}
\int_{1}^{\infty} \frac{d s}{s g_{1}(s)^{2} g_{2}(s) g_{3}(s)^{2}}=\infty \tag{1.8}
\end{equation*}
$$

In [8], one of the authors considered a generalization of the equations in [14] with the incorporation of non-constant $g_{i}, i=1,2,3$, while still leaving $\mathcal{L}_{3}=\Delta$, and guaranteed a unique global solution.

In this article we will extend those results for the case of $\gamma_{3} \neq 2$ and nonconstant $g_{3}$. We will particularly focus on the case of low-regularity initial data in a non- $L^{2}\left(\mathbb{R}^{n}\right)$ setting to then obtain, in the future, global $L^{p}\left(\mathbb{R}^{n}\right)$ solutions using an interpolation technique, the details of which can be found in [5] and [1.

The rest of this article is organized as follows. Section 2 is devoted to explaining the notation we will use and some supporting results necessary for the algorithm. Section 3 contains the main result (Theorem 3.1) of this paper and its proof. We end this section with two important spacial cases of Theorem 3.1.

Theorem 1.1. Let $g_{1}, g_{2}, g_{3}:[0, \infty) \rightarrow \mathbb{R}$ be non-decreasing functions bounded below by 1, satisfying

$$
\begin{equation*}
g_{i}^{(k)}(s) \leq C s^{-k} \tag{1.9}
\end{equation*}
$$

for $i=1,2,3$ and $0 \leq k \leq n / 2+1$. Moreover, assume $0 \leq \gamma_{3} \leq 1$ and $p, q \geq n$ with $2 p>q$. Then, for any divergence-free $u_{0} \in L^{p}\left(\mathbb{R}^{n}\right)$ and $B_{0} \in L^{q}\left(\mathbb{R}^{n}\right)$, there
exists a unique local solution $(u, B)$ to the generalized MHD- $\alpha$ system (1.1)-(1.5) provided that

$$
\begin{aligned}
& \gamma_{1}^{-}>6-\gamma_{3} \\
& \gamma_{2}^{-}>1+\frac{n}{p}
\end{aligned}
$$

Condition 1.9 is a modification of the condition in the Mikhlin multiplier theorem that is necessary for supporting estimates in Proposition (2.4) and (2.5). The functions that satisfy it are still essentially logarithms, the same type of functions that satisfy 1.6 - 1.8 .

Theorem 1.2. Let $\gamma_{3}^{-}-1 \leq \frac{n}{2 p} \leq \gamma_{3}^{-}, \frac{n}{2 q}-1+\gamma_{3}^{-} \leq \frac{n}{2 p}$, and let $p, q \geq n$ with $q<3 p / 2$. Moreover, assume that $g_{1}, g_{2}, g_{3}$ satisfy the inequality (1.9). Then for each divergence-free $u_{0} \in H^{n / 2 p, p}\left(\mathbb{R}^{n}\right)$ and $B_{0} \in H^{n / 2 q, q}\left(\mathbb{R}^{n}\right)$, there exists a unique local solution $(u, B)$ to the generalized $M H D$ system from Equations (1.1)-1.5 provided that

$$
\begin{gathered}
\gamma_{1}^{-}>6-\gamma_{3}^{-}-\frac{n}{p} \\
\gamma_{2}^{-}>1+\frac{n}{2 p}
\end{gathered}
$$

Note that in the statement of Theorem 1.1 and 1.2, and in what follows, we use $x^{-}=x-\varepsilon$ for some positive $\varepsilon$, i.e. $x^{-}$denotes a number arbitrarily close to, but strictly smaller than, $x$.

## 2. Notation and supporting facts

We let $H^{r, p}\left(\mathbb{R}^{n}\right)$ be the usual Sobolev space, and we write $\|f\|_{r, p}$ to mean $\|f\|_{H^{r, p}\left(\mathbb{R}^{n}\right)}$ and $\|f\|_{p}$ for $\|f\|_{L^{p}\left(\mathbb{R}^{n}\right)}$. Because of the nature of the procedure we use, we require that the solutions live in an auxiliary continuous-in-time space $C_{a ; r, p}^{T}\left(\mathbb{R}^{n}\right)$ defined by

$$
C_{a ; r, p}^{T}\left(\mathbb{R}^{n}\right):=\left\{f \in C\left((0, T), H^{r, p}\left(\mathbb{R}^{n}\right)\right):\|f\|_{a ; r, p}<\infty\right\}
$$

where $T>0, a \geq 0, C(X, Y)$ is the space of continuous maps $X \rightarrow Y$, and

$$
\|f\|_{a ; r, p}:=\sup _{(0, T)} t^{a}\|f(t)\|_{r, p}
$$

Finally, we denote by $\dot{C}_{a ; r, p}^{T}\left(\mathbb{R}^{n}\right)$ the subspace of $C_{a ; r, p}^{T}\left(\mathbb{R}^{n}\right)$ consisting of functions $f$ such that $\lim _{t \rightarrow 0^{+}} t^{a} f(t)=0$ and by $B C(X, Y) \subset C(X, Y)$ the subspace of bounded continuous maps $X \rightarrow Y$.

The following are supporting propositions that we will use throughout this paper. The first proposition is a product estimate, its proof can be found in [10, Chapter $2]$.

Proposition 2.1. If $r \geq 0$ and $1<p \leq \infty$, then

$$
\|f g\|_{r, p} \leq C\left(\|f\|_{p_{1}}\|g\|_{r, p_{2}}+\|f\|_{r, q_{1}}\|g\|_{q_{2}}\right)
$$

where

$$
\frac{1}{p}=\frac{1}{p_{1}}+\frac{1}{p_{2}}=\frac{1}{q_{1}}+\frac{1}{q_{2}}
$$

and $p_{1}, p_{2}, q_{1}, q_{2} \in[1, \infty]$.

The following is a useful Sobolev embedding which is a straightforward extension of a result from [11, Chapter 13].
Proposition 2.2. Let $s \geq r$ and $(s-r) p<n$. Then

$$
\|f\|_{r, q} \leq C\|f\|_{s, p}
$$

provided that

$$
\frac{1}{q}-\frac{r}{n}=\frac{1}{p}-\frac{s}{n}
$$

Our next result follows from a simple calculus exercise.
Proposition 2.3. If $0<a, b \in \mathbb{R}$, then

$$
\sup _{t \in[0, T]} \int_{0}^{t}(t-s)^{-a} s^{-b} d s \leq C T^{1-a-b}
$$

provided that $a+b<1$.
Our final two propositions consist of an estimate for the semigroup $e^{t \mathcal{L}_{i}}$ analogous to similar results for the heat kernel $e^{t \Delta}$ and an estimate for the operator $\left(1-\mathcal{L}_{i}\right)^{-1}$. The proofs of both propositions can be found in [7]. We recall that $x^{-}$is a number arbitrarily close to, but strictly smaller than, $x$.

Proposition 2.4. Let $1<p_{1} \leq p_{2}<\infty, r_{1} \leq r_{2}, g(x)$ be a non-decreasing function bounded below by 1, satisfying $\left|g^{(k)}(x)\right| \leq C|x|^{-k}$ for $1 \leq k \leq n / 2+1$. Then $e^{t \mathcal{L}_{i}}: H^{r_{1}, p_{1}}\left(\mathbb{R}^{n}\right) \rightarrow H^{r_{2}, p_{2}}\left(\mathbb{R}^{n}\right)$ and

$$
\left\|e^{t \mathcal{L}_{i}} f\right\|_{r_{2}, p_{2}} \leq t^{-\left(r_{2}-r_{1}+n / p_{1}-n / p_{2}\right) / \gamma_{i}^{-}}\|f\|_{r_{1}, p_{1}}
$$

Note that this proposition necessitates the requirements on the $g_{i}$ 's.
Proposition 2.5. Let $1<p<\infty, r \in \mathbb{R}, g(x)$ be a non-decreasing function bounded below by 1, satisfying $\left|g^{(k)}(x)\right| \leq C|x|^{-k}$ for all $1 \leq k \leq n / 2+1$. Then

$$
\left\|\left(1-\mathcal{L}_{i}\right)^{-1} f\right\|_{r, p} \leq C\|f\|_{r-\gamma_{i}^{-}, p}
$$

## 3. Main Result and its proof

In this section we state the most general form of the theorem and then proceed with its proof.

Theorem 3.1. Let $g_{1}, g_{2}, g_{3}:[0, \infty) \rightarrow \mathbb{R}$ be non-decreasing functions bounded below by 1, satisfying

$$
g_{i}^{(k)}(s) \leq C s^{-k}
$$

for $i=1,2,3$ and $0 \leq k \leq n / 2+1$. Let $r_{0}, r_{1}, r_{2} \geq 0$ and let $p_{0}, p_{1}, p_{2} \geq n$ with $p_{0} \leq p_{1}$ and $p_{2}<2 p_{0}$. Moreover, assume that

$$
\begin{gathered}
\gamma_{3}^{-}-1 \leq r_{0} \leq \gamma_{3}^{-} \leq r_{1}, \\
r_{2}-1+\gamma_{3}^{-} \leq r_{0}, \\
r_{2} \leq r_{0}<\frac{n}{p_{1}}, \\
2 r_{1} \geq \max \left\{2,1+\gamma_{3}^{-}-\frac{n}{p_{0}}+\frac{2 n}{p_{1}}\right\}, \\
r_{2}<\min \left\{\frac{n}{p_{2}}, \frac{2 n}{p_{2}}-\frac{n}{p_{0}}\right\} .
\end{gathered}
$$

Then, for any divergence-free $u_{0} \in H^{r_{0}, p_{0}}\left(\mathbb{R}^{n}\right)$ and $B_{0} \in H^{r_{2}, p_{2}}\left(\mathbb{R}^{n}\right)$, there exists a unique local solution $(u, B)$ to the generalized MHD- $\alpha$ system (1.1)-(1.5) provided that

$$
\begin{gathered}
\gamma_{1}^{-}>3 r_{1}-2 r_{0}-\gamma_{3}^{-}+\frac{3 n}{p_{0}}-\frac{3 n}{p_{1}} \\
\gamma_{1}^{-}>1-2 r_{2}+r_{1}-\gamma_{3}^{-}-\frac{n}{p_{1}}+\frac{2 n}{p_{2}} \\
\gamma_{2}^{-}>1-r_{0}+\frac{n}{p_{0}}
\end{gathered}
$$

Note that Theorem 1.1 may be recovered by setting $r_{0}=r_{2}=0, r_{1}=2, p:=$ $p_{0}=p_{1}$, and $q:=p_{2}$, while Theorem 1.2 may be recovered by setting $r_{0}=n /(2 p)$, $r_{1}=2, r_{2}=n /(2 q), p:=p_{0}=p_{1}$, and $q:=p_{2}$.

Proof of Theorem 3.1. For the sake of clarity and to highlight some technical details, the proof will be divided in subsections. We first write the generalized MHD- $\alpha$ system in a more helpful form. Without loss of generality, we set $\alpha=\nu_{1}=\nu_{2}=1$. We pass to divergence-free vector fields by applying the Hodge operator $P$ to equations (1.1) and (1.2) (more information about the Hodge operator can be found in [3. Chapter 11]), and then we apply $\left(1-\mathcal{L}_{3}\right)^{-1}$ to equation (1.1). By noting that $P,\left(1-\mathcal{L}_{3}\right)^{-1}$, and $\partial_{t}$ all commute since they are Fourier multipliers, we obtain

$$
\begin{aligned}
& \partial_{t} u+P\left(1-\mathcal{L}_{3}\right)^{-1}\left((u \cdot \nabla) v+\sum_{i=1}^{n} v_{i} \nabla u_{i}-(B \cdot \nabla) B\right)-\mathcal{L}_{1} u \\
& \quad=P\left(-\nabla p-\frac{1}{2} \nabla|B|^{2}\right)=0
\end{aligned}
$$

An application of the divergence-free condition in (1.4) allows us to rewrite the terms of the form $(x \cdot \nabla) y$ as $\operatorname{div}(x \otimes y)$. Note that $x \otimes y$ is the matrix whose $(i, j)$ entry is $x_{i} y_{j}$, so that the product estimate in Proposition 2.1 applies to $x \otimes y$. We then have the system

$$
\begin{gathered}
\partial_{t} u+P\left(1-\mathcal{L}_{3}\right)^{-1}\left(\operatorname{div}(u \otimes v)+\sum_{i=1}^{n} v_{i} \nabla u_{i}-\operatorname{div}(B \otimes B)\right)-\mathcal{L}_{1} u=0 \\
\partial_{t} B+P(\operatorname{div}(u \otimes B)-\operatorname{div}(B \otimes u))-\mathcal{L}_{2} B=0 \\
v=\left(1-\mathcal{L}_{3}\right) u \\
\operatorname{div} u=\operatorname{div} B=0 \\
u(x, 0)=u_{0}(x), \quad B(0, x)=B_{0}(x), \quad x \in \mathbb{R}^{n}
\end{gathered}
$$

An application of Duhamel's principle shows that $(u, B)$ is a solution to the system if and only if $(u, B)$ is a fixed point of the map $\Phi(u, B):=\left(\Phi_{1}(u, B), \Phi_{2}(u, B)\right)$ defined by

$$
\begin{aligned}
& \Phi_{1}(u, B):=e^{t \mathcal{L}_{1}} u_{0}-\int_{0}^{t} e^{(t-s) \mathcal{L}_{1}}\left(W_{1}(u, v)+W_{2}(u, v)-W_{1}(B, B)\right) d s \\
& \Phi_{2}(u, B):=e^{t \mathcal{L}_{2}} B_{0}-\int_{0}^{t} e^{(t-s) \mathcal{L}_{2}}\left(W_{3}(u, B)-W_{3}(B, u)\right) d s
\end{aligned}
$$

where

$$
W_{1}(x, y)=P\left(1-\mathcal{L}_{3}\right)^{-1} \operatorname{div}(x \otimes y)
$$

$$
\begin{aligned}
& W_{2}(x, y)=P\left(1-\mathcal{L}_{3}\right)^{-1}\left(\sum_{i=1}^{n} y_{i} \nabla x_{i}\right), \\
& W_{3}(x, y)=P \operatorname{div}(x \otimes y)
\end{aligned}
$$

By the contraction mapping theorem, it suffices to show that $\Phi$ is a contraction on the space $X_{T, M} \times Y_{T, M}$, where

$$
\begin{aligned}
X_{T, M}:=\{ & f \in B C\left([0, T), H^{r_{0}, p_{0}}\left(\mathbb{R}^{n}\right)\right) \cap \dot{C}_{a_{1}, r_{1}, p_{1}}\left(\mathbb{R}^{n}\right): \\
& \left.\sup _{(0, T)}\left(\left\|f(t)-e^{t \mathcal{L}_{1}} u_{0}\right\|_{r_{0}, p_{0}}+\|f(t)\|_{a_{1} ; r_{1}, p_{1}}\right)<M\right\}
\end{aligned}
$$

and

$$
Y_{T, M}:=\left\{f \in B C\left([0, T), H^{r_{2}, p_{2}}\left(\mathbb{R}^{n}\right)\right): \sup _{(0, T)}\left\|f(t)-e^{t \mathcal{L}_{2}} B_{0}\right\|_{r_{2}, p_{2}}<M\right\}
$$

for some $0<T<1$ and $M>0$.
Following the methods in [6, 2], we complete the proof by showing that

$$
\begin{aligned}
& I_{1}=\sup _{(0, T)} t^{a_{1}}\left\|e^{t \mathcal{L}_{1}} u_{0}\right\|_{r_{1}, p_{1}}<M / 4 \\
& I_{2}=\sup _{(0, T)}\left\|\int_{0}^{t} e^{(t-s) \mathcal{L}_{1}}\left(W_{1}(u, v)+W_{2}(u, v)-W_{1}(B, B)\right) d s\right\|_{r_{0}, p_{0}}<M / 4, \\
& I_{3}=\sup _{(0, T)} t^{a_{1}}\left\|\int_{0}^{t} e^{(t-s) \mathcal{L}_{1}}\left(W_{1}(u, v)+W_{2}(u, v)-W_{1}(B, B)\right) d s\right\|_{r_{1}, p_{1}}<M / 4, \\
& I_{4}=\underset{(0, T)}{ } \sup _{0}\left\|\int_{0}^{t} e^{(t-s) \mathcal{L}_{2}}\left(W_{3}(u, B)-W_{3}(B, u)\right) d s\right\|_{r_{2}, p_{2}}<M / 4 .
\end{aligned}
$$

We start with $I_{1}$. If $\varphi$ is in the Schwartz space, we have

$$
\begin{aligned}
I_{1} & =\sup _{(0, T)} t^{a_{1}}\left\|e^{t \mathcal{L}_{1}}\left(u_{0}-\varphi+\varphi\right)\right\|_{r_{1}, p_{1}} \\
& \leq \sup _{(0, T)} t^{a_{1}}\left\|e^{t \mathcal{L}_{1}}\left(u_{0}-\varphi\right)\right\|_{r_{1}, p_{1}}+\sup _{(0, T)} t^{a_{1}}\left\|e^{t \mathcal{L}_{1}} \varphi\right\|_{r_{1}, p_{1}} \\
& \leq \sup _{(0, T)} t^{a_{1}} t^{-a_{1}}\left\|u_{0}-\varphi\right\|_{r_{0}, p_{0}}+\sup _{(0, T)} t^{a_{1}}\|\varphi\|_{r_{1}, p_{1}} \\
& \leq\left\|u_{0}-\varphi\right\|_{r_{0}, p_{0}}+T^{a_{1}}\|\varphi\|_{r_{1}, p_{1}},
\end{aligned}
$$

provided that (by Proposition 2.4)

$$
0 \leq a_{1}=\frac{r_{1}-r_{0}+\frac{n}{p_{0}}-\frac{n}{p_{1}}}{\gamma_{1}^{-}}<1
$$

and $p_{0} \leq p_{1}$. We can choose $\varphi$ so that $\left\|u_{0}-\varphi\right\|_{r_{0}, p_{0}}$ is arbitrarily small, and then we can choose $T$ small enough to reduce $T^{a_{1}}\|\varphi\|_{r_{1}, p_{1}}$ so that the sum of the two is bounded by $M / 4$.
3.1. $I_{2}$ and $I_{3}$. Minkowski's inequality gives us

$$
\begin{gathered}
I_{2} \leq J_{1}+J_{2}+J_{3} \\
I_{3} \leq K_{1}+K_{2}+K_{3}
\end{gathered}
$$

where

$$
\begin{aligned}
J_{1} & :=\sup _{(0, T)} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{1}} W_{1}(u, v)\right\|_{r_{0}, p_{0}} d s \\
K_{1} & :=\sup _{(0, T)} t^{a_{1}} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{1}} W_{1}(u, v)\right\|_{r_{1}, p_{1}} d s \\
J_{2} & :=\sup _{(0, T)} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{1}} W_{2}(u, v)\right\|_{r_{0}, p_{0}} d s \\
K_{2} & :=\sup _{(0, T)} t^{a_{1}} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{1}} W_{2}(u, v)\right\|_{r_{1}, p_{1}} d s \\
J_{3} & :=\sup _{(0, T)} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{1}} W_{1}(B, B)\right\|_{r_{0}, p_{0}} d s \\
K_{3} & :=\sup _{(0, T)} t^{a_{1}} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{1}} W_{1}(B, B)\right\|_{r_{1}, p_{1}} d s
\end{aligned}
$$

We will show that each term is bounded above by $C M^{2} T^{k}$ for various values of $k>0$, which, since $T<1$, will imply $I_{2}, I_{3}<M / 4$ provided that $C M^{2}<M / 4$.

We begin our algorithm with $J_{1}$ and $K_{1}$, showing the details of the calculations and highlighting the choices of parameters. The argument for the other two pairs of integrals is very similar, and the details will be omitted.
3.2. $J_{1}$ and $K_{1}$. By Proposition $2.4, J_{1}$ is bounded by

$$
J_{1} \leq \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(r_{0}-\left(\gamma_{3}^{-}-1\right)+n / \pi_{1}-n / p_{0}\right) / \gamma_{1}^{-}}\left\|W_{1}(u, v)\right\|_{\gamma_{3}^{-}-1, \pi_{1}} d s
$$

provided that $\gamma_{3}^{-}-1 \leq r_{0}$ and where $\pi_{1}$ is an intermediate parameter that will be specified later. Now we work towards bounding $W_{1}(u, v)$, and an application of Proposition 2.5 gives us

$$
\left\|W_{1}(u, v)\right\|_{\gamma_{3}^{-}-1, \pi_{1}}=\left\|P\left(1-\mathcal{L}_{3}\right)^{-1} \operatorname{div}(u \otimes v)\right\|_{\gamma_{3}^{-}-1, \pi_{1}} \leq C\|u \otimes v\|_{\pi_{1}}
$$

We chose a regularity of $\gamma_{3}^{-}-1$ when performing the semigroup estimate in order to end up with a product in zero regularity, so that we can apply Holder's inequality: by Proposition 2.1. if

$$
\begin{equation*}
\frac{1}{\pi_{1}}=\frac{1}{p^{\prime}}+\frac{1}{p_{1}}, \tag{3.1}
\end{equation*}
$$

we have

$$
\|u \otimes v\|_{\pi_{1}} \leq C\|u\|_{p^{\prime}}\|v\|_{p_{1}} \leq C\|u\|_{p^{\prime}}\|u\|_{\gamma_{3}^{-}, p_{1}}
$$

Note that (3.1) specifies the required value of $\pi_{1}$. We obtain $\|u\|_{p^{\prime}} \leq\|u\|_{r_{0}, p_{0}}$ by Proposition 2.2 if

$$
\begin{equation*}
r_{0}<\frac{n}{p_{0}} \quad \text { and } \quad \frac{1}{p^{\prime}}=\frac{1}{p_{0}}-\frac{r_{0}}{n} \tag{3.2}
\end{equation*}
$$

we also get $\|u\|_{\gamma_{3}^{-}, p_{1}} \leq\|u\|_{r_{1}, p_{1}}$ by requiring that $r_{1} \geq \gamma_{3}^{-}$. Combining the two bounds gives us

$$
\begin{equation*}
\left\|W_{1}(u, v)\right\|_{\gamma_{3}^{-}-1, \pi_{1}} \leq C\|u \otimes v\|_{\pi_{1}} \leq C\|u\|_{r_{0}, p_{0}}\|u\|_{r_{1}, p_{1}} \tag{3.3}
\end{equation*}
$$

Note that (3.1) and (3.2) give us

$$
\frac{1}{\pi_{1}}=\frac{1}{p_{0}}+\frac{1}{p_{1}}-\frac{r_{0}}{n}
$$

With this new bound on $W_{1}(u, v)$, we come back to $J_{1}$ and see that

$$
\begin{aligned}
J_{1} & \leq \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(r_{0}-\left(\gamma_{3}^{-}-1\right)+n / \pi_{1}-n / p_{0}\right) / \gamma_{1}^{-}}\left\|W_{1}(u, v)\right\|_{\gamma_{3}^{-}-1, \pi_{1}} d s \\
& \leq C \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(r_{0}-\left(\gamma_{3}^{-}-1\right)+n / \pi_{1}-n / p_{0}\right) / \gamma_{1}^{-}}\|u\|_{r_{0}, p_{0}}\|u\|_{r_{1}, p_{1}} d s \\
& =C \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(r_{0}-\left(\gamma_{3}^{-}-1\right)+n / \pi_{1}-n / p_{0}\right) / \gamma_{1}^{-}} s^{-a_{1}}\|u\|_{r_{0}, p_{0}} s^{a_{1}}\|u\|_{r_{1}, p_{1}} d s \\
& \leq C\|u\|_{0 ; r_{0}, p_{0}}\|u\|_{a_{1} ; r_{1}, p_{1}} \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(r_{0}-\left(\gamma_{3}^{-}-1\right)+n / \pi_{1}-n / p_{0}\right) / \gamma_{1}^{-}} s^{-a_{1}} d s \\
& <C M^{2} T^{1-\left(r_{0}-\left(\gamma_{3}^{-}-1\right)+n / \pi_{1}-n / p_{0}\right) / \gamma_{1}^{-}-a_{1}},
\end{aligned}
$$

where the last inequality holds by Proposition 2.3. if

$$
\begin{align*}
\gamma_{1}^{-} & >r_{0}-\left(\gamma_{3}^{-}-1\right)+\frac{n}{\pi_{1}}-\frac{n}{p_{0}}+\gamma_{1} a_{1} \\
& =r_{0}-\left(\gamma_{3}^{-}-1\right)+n\left(\frac{1}{p_{0}}+\frac{1}{p_{1}}-\frac{r_{0}}{n}\right)-\frac{n}{p_{0}}+r_{1}-r_{0}+\frac{n}{p_{0}}-\frac{n}{p_{1}}  \tag{3.4}\\
& =1-r_{0}+r_{1}-\gamma_{3}^{-}+\frac{n}{p_{0}}
\end{align*}
$$

We further note that the requirement in (3.4) also guarantees that the exponent on $T$ is positive, as desired.

We now turn our attention to $K_{1}$. Proposition 2.4 guarantees that, if $p_{1} \geq \pi_{1}^{\prime}$, then

$$
K_{1} \leq \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(r_{1}-\left(r_{1}-r_{0}+\gamma_{3}^{-}-1\right)+n / \pi_{1}^{\prime}-n / p_{1}\right) / \gamma_{1}^{-}}\left\|W_{1}(u, v)\right\|_{r_{1}-r_{0}+\gamma_{3}^{-}-1, \pi_{1}^{\prime}} d s
$$

This time, we chose $r_{1}-r_{0}+\gamma_{3}^{-}-1$ in order to match the previous "jump" in regularity from $r_{0}$ to $\gamma_{3}^{-}-1$. Propositions 2.1 and 2.5 give us

$$
\begin{aligned}
\left\|W_{1}(u, v)\right\|_{r_{1}-r_{0}+\gamma_{3}^{-}-1, \pi_{1}^{\prime}} & =\left\|P\left(1-\mathcal{L}_{3}\right)^{-1} \operatorname{div}(u \otimes v)\right\|_{r_{1}-r_{0}+\gamma_{3}^{-}-1, \pi_{1}^{\prime}} \\
& \leq C\|u \otimes v\|_{r_{1}-r_{0}, \pi_{1}^{\prime}} \\
& \leq C\left(\|u\|_{r_{1}-r_{0}, p^{\prime}}\|v\|_{p^{\prime \prime}}+\|v\|_{r_{1}-r_{0}, q^{\prime}}\|u\|_{q^{\prime \prime}}\right)
\end{aligned}
$$

provided that

$$
\frac{1}{\pi_{1}^{\prime}}=\frac{1}{p^{\prime}}+\frac{1}{p^{\prime \prime}}=\frac{1}{q^{\prime}}+\frac{1}{q^{\prime \prime}}
$$

Four applications of Proposition 2.2 lead us to the following bounds:

$$
\begin{array}{ll}
\|u\|_{r_{1}-r_{0}, p^{\prime}} \leq\|u\|_{r_{1}, p_{1}} & \text { if } r_{0}<\frac{n}{p_{1}} \text { and } \frac{1}{p^{\prime}}=\frac{1}{p_{1}}-\frac{r_{0}}{n} \\
\|v\|_{p^{\prime \prime}} \leq\|v\|_{r_{0}-\gamma_{3}^{-}, p_{0}} & \text { if } r_{0}<\frac{n}{p_{0}}+\gamma_{3}^{-} \text {and } \frac{1}{p^{\prime \prime}}=\frac{1}{p_{0}}-\frac{r_{0}-\gamma_{3}^{-}}{n}, \\
\|v\|_{r_{1}-r_{0}, q^{\prime}} \leq\|v\|_{r_{1}-\gamma_{3}^{-}, p_{1}} & \text { if } r_{0}<\frac{n}{p_{0}}+\gamma_{3}^{-} \text {and } \frac{1}{q^{\prime}}=\frac{1}{p_{1}}-\frac{r_{0}-\gamma_{3}^{-}}{n} \tag{3.7}
\end{array}
$$

$$
\begin{equation*}
\|u\|_{q^{\prime \prime}} \leq\|u\|_{r_{0}, p_{0}} \quad \text { if } r_{0}<\frac{n}{p_{0}} \text { and } \frac{1}{q^{\prime \prime}}=\frac{1}{p_{0}}-\frac{r_{0}}{n} \tag{3.8}
\end{equation*}
$$

Combining the parameters specified by (3.5)-3.8), we obtain

$$
\frac{1}{\pi_{1}^{\prime}}=\frac{1}{p_{0}}+\frac{1}{p_{1}}-\frac{2 r_{0}-\gamma_{3}^{-}}{n}
$$

and

$$
\left\|W_{1}(u, v)\right\|_{r_{1}-r_{0}+\gamma_{3}^{-}-1, \pi_{1}^{\prime}} \leq C\|u \otimes v\|_{r_{1}-r_{0}, \pi_{1}^{\prime}} \leq C\|u\|_{r_{0}, p_{0}}\|u\|_{r_{1}, p_{1}} .
$$

Moreover, the integrability requirement from Proposition 2.4 necessitates

$$
\frac{1}{p_{1}} \leq \frac{1}{\pi_{1}^{\prime}}=\frac{1}{p_{0}}+\frac{1}{p_{1}}-\frac{2 r_{0}-\gamma_{3}^{-}}{n}
$$

and so

$$
\begin{equation*}
r_{0} \leq \frac{1}{2}\left(\frac{n}{p_{0}}+\gamma_{3}^{-}\right) \tag{3.9}
\end{equation*}
$$

We can finally plug this bound into $K_{1}$ :

$$
\begin{aligned}
K_{1} \leq & \sup _{(0, T)} t^{a_{1}} \int_{0}^{t}(t-s)^{-\left(r_{1}-\left(r_{1}-r_{0}+\gamma_{3}^{-}-1\right)+n / \pi_{1}^{\prime}-n / p_{1}\right) / \gamma_{1}^{-}} \\
& \times\left\|W_{1}(u, v)\right\|_{r_{1}-r_{0}+\gamma_{3}^{-}-1, \pi_{1}^{\prime}} d s \\
\leq & C \sup _{(0, T)} t^{a_{1}} \int_{0}^{t}(t-s)^{-\left(r_{1}-\left(r_{1}-r_{0}+\gamma_{3}^{-}-1\right)+n / \pi_{1}^{\prime}-n / p_{1}\right) / \gamma_{1}^{-}}\|u\|_{r_{0}, p_{0}}\|u\|_{r_{1}, p_{1}} d s \\
= & C \sup _{(0, T)} t^{a_{1}} \int_{0}^{t}(t-s)^{-\left(r_{1}-\left(r_{1}-r_{0}+\gamma_{3}^{-}-1\right)+n / \pi_{1}^{\prime}-n / p_{1}\right) / \gamma_{1}^{-}} s^{-a_{1}} \\
& \times\|u\|_{r_{0}, p_{0}} s^{a_{1}}\|u\|_{r_{1}, p_{1}} d s \\
\leq & C\|u\|_{0 ; r_{0}, p_{0}}\|u\|_{a_{1} ; r_{1}, p_{1}} \sup _{(0, T)} t^{a_{1}} \int_{0}^{t}(t-s)^{-\left(r_{1}-\left(r_{1}-r_{0}+\gamma_{3}^{-}-1\right)+n / \pi_{1}^{\prime}-n / p_{1}\right) / \gamma_{1}^{-}} \\
& \times s^{-a_{1}} d s \\
< & C M^{2} T^{1-\left(r_{1}-\left(r_{1}-r_{0}+\gamma_{3}^{-}-1\right)+n / \pi_{1}^{\prime}-n / p_{0}\right) / \gamma_{1}^{-}},
\end{aligned}
$$

where the last inequality follows by Proposition 2.3 if

$$
\begin{align*}
\gamma_{1}^{-} & >r_{1}-\left(r_{1}-r_{0}+\gamma_{3}^{-}-1\right)+\frac{n}{\pi_{1}^{\prime}}-\frac{n}{p_{1}}+\gamma_{1}^{-} a_{1} \\
& =r_{0}-\gamma_{3}^{-}+1+n\left(\frac{1}{p_{0}}+\frac{1}{p_{1}}-\frac{2 r_{0}-\gamma_{3}^{-}}{n}\right)-\frac{n}{p_{1}}+r_{1}-r_{0}+\frac{n}{p_{0}}-\frac{n}{p_{1}}  \tag{3.10}\\
& =1-2 r_{0}+r_{1}+\frac{2 n}{p_{0}}-\frac{n}{p_{1}} .
\end{align*}
$$

Note that, once again, the requirement that Proposition 2.3 hold is sufficient to guarantee that the exponent on $T$ be positive.

To summarize, here is the list of inequalities needed to obtain the desired bounds on $J_{1}$ and $K_{1}$ :

$$
\begin{array}{ll}
r_{0} \leq \gamma_{3}^{-} & (\text {assumption) }, \\
\gamma_{3}^{-}-1 \leq r_{0} & \left(\text { semigroup estimate for } J_{1}\right), \\
r_{1}-r_{0}+\gamma_{3}^{-}-1 \leq r_{1} & \left(\text { semigroup estimate for } K_{1}\right)
\end{array}
$$

$$
\begin{array}{ll}
r_{0}<\frac{n}{p_{0}} & \\
r_{0}<\frac{n}{p_{1}} & (\text { by } \sqrt[3.2)]{ }), \\
r_{0} \leq \frac{1}{2}\left(\frac{n}{p_{0}}+\gamma_{3}^{-}\right) & (\text {by }(3.5), \\
r_{0}<\frac{n}{p_{0}}+\gamma_{3}^{-} & \\
r_{1} \geq \gamma_{3}^{-} & (\text {by } 3.12), \\
\gamma_{1}^{-}>1-r_{0}+r_{1}-\gamma_{3}^{-}+\frac{n}{p_{0}} & (\text { by }, \\
\gamma_{1}^{-}>1-2 r_{0}+r_{1}+\frac{2 n}{p_{0}}-\frac{n}{p_{1}} & (\text { by }), \\
3.10) .
\end{array}
$$

After some obvious simplifications and after noting that (3.10) implies (3.4) since

$$
\begin{aligned}
& \underbrace{\left(1-2 r_{0}+r_{1}+\frac{2 n}{p_{0}}-\frac{n}{p_{1}}\right)}_{\text {RHS of }}-\underbrace{\left(1-r_{0}+r_{1}-\gamma_{3}^{-}+\frac{n}{p_{0}}\right)}_{\text {RHS of } \sqrt{3.4}} \\
& =\gamma_{3}^{-}-r_{0}+\frac{n}{p_{0}}-\frac{n}{p_{1}} \geq 0,
\end{aligned}
$$

the list reduces to

$$
\begin{gathered}
\gamma_{3}^{-}-1 \leq r_{0} \leq \gamma_{3}^{-} \leq r_{1} \\
r_{0}<\frac{n}{p_{1}} \\
\gamma_{1}^{-}>1-2 r_{0}+r_{1}+\frac{2 n}{p_{0}}-\frac{n}{p_{1}} .
\end{gathered}
$$

3.3. $J_{2}$ and $K_{2}$. We have

$$
J_{2}=\sup _{(0, T)} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{1}} W_{2}(u, v)\right\|_{r_{0}, p_{0}} d s \leq \sup _{(0, T)} \int_{0}^{t}\left\|W_{2}(u, v)\right\|_{r_{0}, p_{0}} d s
$$

We now work towards bounding $W_{2}(u, v)$. We immediately see, thanks to Proposition 2.5, that

$$
\begin{aligned}
\left\|W_{2}(u, v)\right\|_{r_{0}, p_{0}} & =\left\|P\left(1-\mathcal{L}_{3}\right)^{-1} \sum_{i=1}^{n} v_{i} \nabla u_{i}\right\|_{r_{0}, p_{0}} \\
& \leq C\left\|\sum_{i=1}^{n} v_{i} \nabla u_{i}\right\|_{r_{0}-\gamma_{3}^{-}, p_{0}} \\
& \leq C \sum_{i=1}^{n}\left\|v_{i} \nabla u_{i}\right\|_{p_{0}}
\end{aligned}
$$

since $r_{0} \leq \gamma_{3}^{-}$. Now the product estimate is nothing more than Holder's inequality, so if

$$
\frac{1}{p_{0}}=\frac{1}{p^{\prime}}+\frac{1}{p^{\prime \prime}}
$$

we obtain

$$
\begin{aligned}
\left\|W_{2}(u, v)\right\|_{r_{0}, p_{0}} & \leq C \sum_{i=1}^{n}\left\|v_{i}\right\|_{p^{\prime}}\left\|\nabla u_{i}\right\|_{p^{\prime \prime}} \\
& \leq C \sum_{i=1}^{n}\|v\|_{p^{\prime}}\|\nabla u\|_{p^{\prime \prime}} \\
& \leq C\|u\|_{\gamma_{3}^{-}, p^{\prime}}\|u\|_{1, p^{\prime \prime}} .
\end{aligned}
$$

By Proposition 2.2 we have

$$
\|u\|_{\gamma_{3}^{-}, p^{\prime}} \leq C\|u\|_{\gamma_{3}^{-}+\beta, p_{1}} \leq C\|u\|_{r_{1}, p_{1}} \quad \text { and } \quad\|u\|_{1, p^{\prime \prime}} \leq C\|u\|_{r_{1}, p_{1}},
$$

where the first set of inequalities requires that

$$
\begin{equation*}
0 \leq \beta<\frac{n}{p_{1}}, \quad \frac{1}{p^{\prime}}=\frac{1}{p_{1}}-\frac{\beta}{n}, \quad r_{1} \geq \gamma_{3}^{-}+\beta \tag{3.11}
\end{equation*}
$$

and the second inequality requires that

$$
\begin{equation*}
\frac{1}{p^{\prime \prime}}=\frac{1}{p_{1}}-\frac{r_{1}-1}{n} \quad \text { and } \quad r_{1} \geq 1 \tag{3.12}
\end{equation*}
$$

We finally obtain

$$
\left\|W_{2}(u, v)\right\|_{r_{0}, p_{0}} \leq C\|u\|_{\gamma_{3}^{-}, p^{\prime}}\|u\|_{1, p^{\prime \prime}} \leq C\|u\|_{r_{1}, p_{1}}^{2}
$$

We pause here to note that, without the presence of the space $\dot{C}_{a_{1} ; r_{1}, p_{1}}\left(\mathbb{R}^{n}\right)$ in the definition of $X_{T, M}$, we would not be able to bound this $W_{2}$ term.

Returning to $J_{2}$, we have

$$
\begin{aligned}
J_{2} & \leq \sup _{(0, T)} \int_{0}^{t}\left\|W_{2}(u, v)\right\|_{r_{0}, \pi_{2}} d s \\
& \leq C \sup _{(0, T)} \int_{0}^{t}\|u\|_{r_{1}, p_{1}}^{2} d s \\
& =C \sup _{(0, T)} \int_{0}^{t} s^{-2 a_{1}} s^{a_{1}}\|u\|_{r_{1}, p_{1}} s^{a_{1}}\|u\|_{r_{1}, p_{1}} d s \\
& \leq C\|u\|_{a_{1} ; r_{1}, p_{1}}^{2} \sup _{(0, T)} \int_{0}^{t} s^{-2 a_{1}} d s \\
& <C M^{2} T^{1-2 a_{1}},
\end{aligned}
$$

provided $2 a_{1}>1$, which is equivalent to

$$
\begin{equation*}
\gamma_{1}^{-}>2 r_{1}-2 r_{0}+\frac{2 n}{p_{0}}-\frac{2 n}{p_{1}} \tag{3.13}
\end{equation*}
$$

and we recall that

$$
\frac{n}{p_{0}}=\frac{n}{p^{\prime}}+\frac{n}{p^{\prime \prime}}=\frac{2 n}{p_{1}}-\beta-r_{1}+1
$$

We choose $\beta$ to be exactly

$$
\begin{equation*}
\beta=1-r_{1}-\frac{n}{p_{0}}+\frac{2 n}{p_{1}}, \tag{3.14}
\end{equation*}
$$

and so the two requirements in (3.11) become

$$
\begin{equation*}
r_{1} \geq 1-\frac{n}{p_{0}}+\frac{n}{p_{1}} \quad \text { and } \quad 2 r_{1} \geq 1+\gamma_{3}^{-}-\frac{n}{p_{0}}+\frac{2 n}{p_{1}} \tag{3.15}
\end{equation*}
$$

Turning to $K_{2}$, noting that we go down to $\gamma_{3}^{-}$instead of $r_{0}$, we have

$$
\begin{aligned}
K_{2} & =\sup _{(0, T)} t^{a_{1}} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{1}} W_{2}(u, v)\right\|_{r_{1}, p_{1}} d s \\
& \leq \sup _{(0, T)} t^{a_{1}} \int_{0}^{t}(t-s)^{-\left(r_{1}-\gamma_{3}^{-}+n / p_{0}-n / p_{1}\right) / \gamma_{1}^{-}}\left\|W_{2}(u, v)\right\|_{\gamma_{3}^{-}, p_{0}} d s \\
& \leq C \sup _{(0, T)} t^{a_{1}} \int_{0}^{t}(t-s)^{-\left(r_{1}-\gamma_{3}^{-}+n / p_{0}-n / p_{1}\right) / \gamma_{1}^{-}}\|u\|_{r_{1}, p_{1}}\|u\|_{r_{1}, p_{1}} d s \\
& =C \sup _{(0, T)} t^{a_{1}} \int_{0}^{t}(t-s)^{-\left(r_{1}-\gamma_{3}^{-}+n / p_{0}-n / p_{1}\right) / \gamma_{1}^{-}} s^{-2 a_{1}} s^{a_{1}}\|u\|_{r_{1}, p_{1}} s^{a_{1}}\|u\|_{r_{1}, p_{1}} d s \\
& \leq C\|u\|_{a_{1} ; r_{0}, p_{0}}^{2} \sup _{(0, T)} t^{a_{1}} \int_{0}^{t}(t-s)^{-\left(r_{1}-\gamma_{3}^{-}+n / p_{0}-n / p_{1}\right) / \gamma_{1}^{-}} s^{-2 a_{1}} d s \\
& <C M^{2} T^{1-\left(r_{1}-\gamma_{3}^{-}+n / \pi_{2}-n / p_{1}\right) / \gamma_{1}^{-}-a_{1}},
\end{aligned}
$$

where, by Proposition 2.3, the last inequality holds if

$$
\begin{align*}
\gamma_{1}^{-} & >r_{1}-\gamma_{3}^{-}+\frac{n}{p_{0}}-\frac{n}{p_{1}}+2 \gamma_{1}^{-} a_{1} \\
& =3 r_{1}-2 r_{0}-\gamma_{3}^{-}+\frac{3 n}{p_{0}}-\frac{3 n}{p_{1}} \tag{3.16}
\end{align*}
$$

Here is a summary of the inequalities needed to obtain the required bounds on $J_{2}$ and $K_{2}$ :

$$
\begin{array}{ll}
r_{1} \geq 1 & (\text { by }(3.12)), \\
r_{1} \geq 1-\frac{n}{p_{0}}+\frac{n}{p_{1}} & (\text { by }(3.15), \\
2 r_{1} \geq 1+\gamma_{3}^{-}-\frac{n}{p_{0}}+\frac{2 n}{p_{1}} & (\text { by }(3.15), \\
\gamma_{1}^{-}>2 r_{1}-2 r_{0}+\frac{2 n}{p_{0}}-\frac{2 n}{p_{1}} & (\text { by }(3.13), \\
\gamma_{1}^{-}>3 r_{1}-2 r_{0}-\gamma_{3}^{-}+\frac{3 n}{p_{0}}-\frac{3 n}{p_{1}} & (\text { by }(3.16) .
\end{array}
$$

By noting that

$$
\begin{aligned}
& \underbrace{\left(3 r_{1}-2 r_{0}-\gamma_{3}^{-}+\frac{3 n}{p_{0}}-\frac{3 n}{p_{1}}\right)}_{\text {RHS of }(3.16)}-\underbrace{\left(2 r_{1}-2 r_{0}+\frac{2 n}{p_{0}}-\frac{2 n}{p_{1}}\right)}_{\text {RHS of }} \\
& =r_{1}-\gamma_{3}^{-}+\frac{n}{p_{0}}-\frac{n}{p_{1}} \geq 0
\end{aligned}
$$

we conclude that 3.16 implies 3.13 , and so the list reduces to

$$
\begin{aligned}
2 r_{1} & \geq \max \left\{2,1+\gamma_{3}^{-}-\frac{n}{p_{0}}+\frac{2 n}{p_{1}}\right\} \\
\gamma_{1}^{-} & >3 r_{1}-2 r_{0}-\gamma_{3}^{-}+\frac{3 n}{p_{0}}-\frac{3 n}{p_{1}}
\end{aligned}
$$

3.4. $J_{3}$ and $K_{3}$. Provided that $r_{2}-1+\gamma_{3}^{-} \leq r_{0}$ and $\frac{1}{\pi_{3}} \geq \frac{1}{p_{0}}$, Proposition 2.4 gives us

$$
\begin{aligned}
J_{3} & =\sup _{(0, T)} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{1}} W_{1}(B, B)\right\|_{r_{0}, p_{0}} d s \\
& \leq \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(r_{0}-\left(r_{2}-1+\gamma_{3}^{-}\right)+n / \pi_{3}-n / p_{0}\right) / \gamma_{1}^{-}}\left\|W_{1}(B, B)\right\|_{r_{2}-1+\gamma_{3}^{-}, \pi_{3}} d s .
\end{aligned}
$$

Once again, applying Propositions 2.1 and 2.5 gives us

$$
\begin{aligned}
\left\|W_{1}(B, B)\right\|_{r_{2}-1+\gamma_{3}^{-}, \pi_{3}} & =\left\|P\left(1-\mathcal{L}_{3}\right)^{-1} \operatorname{div}(B \otimes B)\right\|_{r_{2}-1+\gamma_{3}^{-}, \pi_{3}} \\
& \leq C\|B \otimes B\|_{r_{2}, \pi_{3}} \\
& \leq C\|B\|_{r_{2}, p_{2}}\|B\|_{p^{\prime}}
\end{aligned}
$$

where the product estimate requires that $r_{2} \geq 0$ and

$$
\frac{1}{\pi_{3}}=\frac{1}{p_{2}}+\frac{1}{p^{\prime}}
$$

Provided that

$$
\begin{equation*}
r_{2}<\frac{n}{p_{2}} \quad \text { and } \quad \frac{1}{p^{\prime}}=\frac{1}{p_{2}}-\frac{r_{2}}{n} \tag{3.17}
\end{equation*}
$$

which combines with the previous equation to give

$$
\frac{1}{\pi_{3}}=\frac{2}{p_{2}}-\frac{r_{2}}{n}
$$

we can bound $\|B\|_{p^{\prime}}$ by $\|B\|_{r_{2}, p_{2}}$ thanks to Proposition 2.2. Thus,

$$
\left\|W_{1}(B, B)\right\|_{r_{2}-1+\gamma_{3}^{-}, \pi_{3}} \leq C\|B\|_{r_{2}, p_{2}}^{2}
$$

Moreover, Proposition (2.4) requires that

$$
\frac{1}{p_{0}} \leq \frac{1}{\pi_{3}}=\frac{2}{p_{2}}-\frac{r_{2}}{n}
$$

which can be restated as

$$
\begin{equation*}
r_{2} \leq \frac{2 n}{p_{2}}-\frac{n}{p_{0}} \tag{3.18}
\end{equation*}
$$

Plugging the bound for $W_{1}(B, B)$ back into the integral gives us

$$
\begin{aligned}
J_{3} & \leq \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(r_{0}-\left(r_{2}-1+\gamma_{3}^{-}\right)+n / \pi_{3}-n / p_{0}\right) / \gamma_{1}^{-}}\left\|W_{1}(B, B)\right\|_{r_{2}-1+\gamma_{3}^{-}, \pi_{3}} d s \\
& \leq C \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(r_{0}-\left(r_{2}-1+\gamma_{3}^{-}\right)+n / \pi_{3}-n / p_{0}\right) / \gamma_{1}^{-}}\|B\|_{r_{2}, p_{2}}^{2} d s \\
& \leq C\|B\|_{0 ; r_{2}, p_{2}}^{2} \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(r_{0}-\left(r_{2}-1+\gamma_{3}^{-}\right)+n / \pi_{3}-n / p_{0}\right) / \gamma_{1}^{-}} d s \\
& <C M^{2} T^{1-\left(r_{0}-\left(r_{2}-1+\gamma_{3}^{-}\right)+n / \pi_{3}-n / p_{0}\right) / \gamma_{1}^{-}},
\end{aligned}
$$

where once again the last inequality holds if

$$
\begin{align*}
\gamma_{1}^{-} & >r_{0}-\left(r_{2}-1+\gamma_{3}^{-}\right)+\frac{n}{\pi_{3}}-\frac{n}{p_{0}} \\
& =r_{0}-2 r_{2}-\gamma_{3}^{-}+1+\frac{2 n}{p_{2}}-\frac{n}{p_{0}} \tag{3.19}
\end{align*}
$$

The same bounds for $W_{1}(B, B)$ work in the case of $K_{3}$, so that

$$
\begin{aligned}
K_{3} & =\sup _{(0, T)} t^{a_{1}} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{1}} W_{1}(B, B)\right\|_{r_{1}, p_{1}} d s \\
& \leq \sup _{(0, T)} t^{a_{1}} \int_{0}^{t}(t-s)^{-\left(r_{1}-\left(r_{2}-1+\gamma_{3}^{-}\right)+n / \pi_{3}-n / p_{1}\right) / \gamma_{1}^{-}}\left\|W_{1}(B, B)\right\|_{r_{2}-1+\gamma_{3}^{-}, \pi_{3}} d s \\
& \leq C \sup _{(0, T)} t^{a_{1}} \int_{0}^{t}(t-s)^{-\left(r_{1}-\left(r_{2}-1+\gamma_{3}^{-}\right)+n / \pi_{3}-n / p_{1}\right) / \gamma_{1}^{-}}\|B\|_{r_{2}, p_{2}}^{2} d s \\
& \leq C\|B\|_{0 ; r_{2}, p_{2}}^{2} \sup _{(0, T)} t^{a_{1}} \int_{0}^{t}(t-s)^{-\left(r_{1}-\left(r_{2}-1+\gamma_{3}^{-}\right)+n / \pi_{3}-n / p_{1}\right) / \gamma_{1}^{-}} d s \\
& <C M^{2} T^{1-\left(r_{1}-\left(r_{2}-1+\gamma_{3}^{-}\right)+n / \pi_{3}-n / p_{1}\right) / \gamma_{1}^{-}+a_{1}},
\end{aligned}
$$

which holds provided that

$$
\begin{align*}
\gamma_{1}^{-} & >r_{1}-\left(r_{2}-1+\gamma_{3}^{-}\right)+\frac{n}{\pi_{3}}-\frac{n}{p_{1}} \\
& =r_{1}-2 r_{2}-\gamma_{3}^{-}+1-\frac{n}{p_{1}}+\frac{2 n}{p_{2}} \tag{3.20}
\end{align*}
$$

What follows is the list of inequalities needed to bound $J_{3}$ and $K_{3}$ as desired:

$$
\begin{array}{ll}
r_{2}-1+\gamma_{3}^{-} \leq r_{0} & \left(\text { semigroup estimate for } J_{3}\right), \\
r_{2}-1+\gamma_{3}^{-} \leq r_{1} & \left(\text { semigroup estimate for } K_{3}\right), \\
r_{2} \geq 0 & (\text { product estimate) }, \\
r_{2}<\min \left\{\frac{n}{p_{2}}, \frac{2 n}{p_{2}}-\frac{n}{p_{0}}\right\} & (\text { by } 3.17)-3.18), \\
\gamma_{1}^{-}>r_{0}-2 r_{2}-\gamma_{3}^{-}+1-\frac{n}{p_{0}}+\frac{2 n}{p_{2}} & (\text { by } 3.19), \\
\gamma_{1}^{-}>r_{1}-2 r_{2}-\gamma_{3}^{-}+1+\frac{2 n}{p_{2}}-\frac{n}{p_{1}} & (\text { by } 3.20) .
\end{array}
$$

We see that (3.20) suffices for 3.19 since

$$
\begin{aligned}
& \underbrace{=r_{1}-r_{0}+\frac{n}{p_{0}}-\frac{n}{p_{1}} \geq 0,}_{\text {RHS of } \sqrt[3.20]{\left(r_{1}-2 r_{2}-\gamma_{3}^{-}+1+\frac{2 n}{p_{2}}-\frac{n}{p_{1}}\right)}-\underbrace{\left(r_{0}-2 r_{2}-\gamma_{3}^{-}+1-\frac{n}{p_{0}}+\frac{2 n}{p_{2}}\right)}_{\text {RHS of } \sqrt{3.19}}}
\end{aligned}
$$

and so the list reduces to

$$
\begin{gathered}
r_{2}-1+\gamma_{3}^{-} \leq r_{0} \\
0 \leq r_{2}<\min \left\{\frac{n}{p_{2}}, \frac{2 n}{p_{2}}-\frac{n}{p_{0}}\right\}, \\
\gamma_{1}^{-}>r_{1}-2 r_{2}-\gamma_{3}^{-}+1+\frac{2 n}{p_{2}}-\frac{n}{p_{1}}
\end{gathered}
$$

3.5. Bounding $I_{4}$. Applying Minkowski's inequality to $I_{4}$ gives

$$
I_{4} \leq L_{1}+L_{2}
$$

where

$$
\begin{aligned}
L_{1} & :=\sup _{(0, T)} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{2}} W_{3}(u, B)\right\|_{r_{2}, p_{2}} d s \\
L_{2} & :=\sup _{(0, T)} \int_{0}^{t}\left\|e^{(t-s) \mathcal{L}_{2}} W_{3}(B, u)\right\|_{r_{2}, p_{2}} d s
\end{aligned}
$$

We can immediately note that, since $W_{3}$ is not symmetric, $L_{1} \neq L_{2}$, but our techniques will give the same bound for each. So, we set $L:=L_{1}$ and proceed to bound only $L_{1}$. Proposition 2.4 gives us

$$
L \leq \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(r_{2}-\left(r_{2}-1\right)+n / \pi_{4}-n / p_{2}\right) / \gamma_{2}^{-}}\left\|W_{3}(u, B)\right\|_{r_{2}-1, \pi_{4}} d s
$$

provided that $1 / \pi_{4} \geq 1 / p_{2}$.
Continuing with $W_{3}(u, B)$, we obtain

$$
\left\|W_{3}(u, B)\right\|_{r_{2}-1, \pi_{4}}=\|P \operatorname{div}(u \otimes B)\|_{r_{2}-1, \pi_{4}} \leq C\|u \otimes B\|_{r_{2}, \pi_{4}}
$$

an application of Proposition 2.1 gives us

$$
\|u \otimes B\|_{r_{2}, \pi_{4}} \leq C\left(\|u\|_{r_{2}, p^{\prime}}\|B\|_{p^{\prime \prime}}+\|B\|_{r_{2}, p_{2}}\|u\|_{q^{\prime \prime}}\right)
$$

as long as

$$
\frac{1}{\pi_{4}}=\frac{1}{p^{\prime}}+\frac{1}{p^{\prime \prime}}=\frac{1}{p_{2}}+\frac{1}{q^{\prime \prime}} .
$$

We want to bound $\|u \otimes B\|_{r_{2}, \pi_{4}}$ by $\|u\|_{r_{0}, p_{0}}\|B\|_{r_{2}, p_{2}}$, which requires three applications of Proposition 2.2. First, we obtain $\|u\|_{r_{2}, p^{\prime}} \leq\|u\|_{r_{0}, p_{0}}$ if

$$
\begin{equation*}
0 \leq r_{0}-r_{2}<\frac{n}{p_{0}} \quad \text { and } \quad \frac{1}{p^{\prime}}-\frac{r_{2}}{n}=\frac{1}{p_{0}}-\frac{r_{0}}{n} \tag{3.21}
\end{equation*}
$$

We further obtain $\|u\|_{q^{\prime \prime}} \leq\|u\|_{r_{0}, p_{0}}$ provided that

$$
\begin{equation*}
r_{0}<\frac{n}{p_{0}} \quad \text { and } \quad \frac{1}{q^{\prime \prime}}=\frac{1}{p_{0}}-\frac{r_{0}}{n} \tag{3.22}
\end{equation*}
$$

The last embedding, $\|B\|_{p^{\prime \prime}} \leq\|B\|_{r_{2}, p_{2}}$, requires

$$
\begin{equation*}
r_{2}<\frac{n}{p_{2}} \quad \text { and } \quad \frac{1}{p^{\prime \prime}}=\frac{1}{p_{2}}-\frac{r_{2}}{n} . \tag{3.23}
\end{equation*}
$$

Combining Equations (3.21)-(3.23) gives us

$$
\frac{1}{\pi_{4}}=\frac{1}{p_{0}}+\frac{1}{p_{2}}-\frac{r_{0}}{n}
$$

which is required to satisfy

$$
\begin{equation*}
\frac{1}{p_{2}} \leq \frac{1}{\pi_{4}}=\frac{1}{p_{0}}+\frac{1}{p_{2}}-\frac{r_{0}}{n} \Longrightarrow r_{0} \leq \frac{n}{p_{0}} \tag{3.24}
\end{equation*}
$$

This is the bound we were looking for:

$$
\left\|W_{3}(u, B)\right\|_{r_{2}-1, \pi_{4}} \leq C\left(\|u\|_{r_{2}, p^{\prime}}\|B\|_{p^{\prime \prime}}+\|B\|_{r_{2}, q^{\prime}}\|u\|_{q^{\prime \prime}}\right) \leq C\|u\|_{r_{0}, p_{0}}\|B\|_{r_{2}, p_{2}} .
$$

We can plug the above into $L$ and obtain

$$
L \leq \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(1+n / \pi_{4}-n / p_{2}\right) / \gamma_{2}^{-}}\left\|W_{3}(u, B)\right\|_{r_{2}-1, \pi_{4}} d s
$$

$$
\begin{aligned}
& \leq C \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(1+n / \pi_{4}-n / p_{2}\right) / \gamma_{2}^{-}}\|u\|_{r_{0}, p_{0}}\|B\|_{r_{2}, p_{2}} d s \\
& \leq C\|u\|_{0 ; r_{0}, p_{0}}\|B\|_{0 ; r_{2}, p_{2}} \sup _{(0, T)} \int_{0}^{t}(t-s)^{-\left(1+n / \pi_{4}-n / p_{2}\right) / \gamma_{2}^{-}} d s \\
& \leq C M^{2} T^{1-\left(1+n / \pi_{4}-n / p_{2}\right) / \gamma_{2}^{-}}
\end{aligned}
$$

which holds if

$$
\begin{equation*}
\gamma_{2}^{-}>1+\frac{n}{\pi_{4}}-\frac{n}{p_{2}}=1-r_{0}+\frac{n}{p_{0}} \tag{3.25}
\end{equation*}
$$

The list of inequalities necessary to bound $L$ is thus

$$
\begin{array}{ll}
0 \leq r_{0}-r_{2}<\frac{n}{p_{0}} & (\text { by } 3.21), \\
r_{2}<\frac{n}{p_{2}} & (\text { by } 3.23), \\
r_{0} \leq \frac{n}{p_{0}} & (\text { by } 3.24), \\
\gamma_{2}^{-}>1-r_{0}+\frac{n}{p_{0}} & (\text { by } 3.25) .
\end{array}
$$

3.6. Wrapping up. On one final note, we point out that since

$$
\begin{aligned}
& \underbrace{\left(3 r_{1}-2 r_{0}-\gamma_{3}^{-}+\frac{3 n}{p_{0}}-\frac{3 n}{p_{1}}\right)}_{\text {RHS of } 3.16}-\underbrace{\left(1-2 r_{0}+r_{1}+\frac{2 n}{p_{0}}-\frac{n}{p_{1}}\right)}_{\text {RHS of } 3.10} \\
& =2 r_{1}-1-\gamma_{3}^{-}+\frac{n}{p_{0}}-\frac{2 n}{p_{1}} \geq 0
\end{aligned}
$$

we have that (3.16) implies (3.10), and so the following is the definitive list containing all the inequalities needed for $I_{2}, I_{3}$, and $I_{4}$ :

$$
\begin{gathered}
\gamma_{3}^{-}-1 \leq r_{0} \leq \gamma_{3}^{-} \leq r_{1}, \\
r_{2}-1+\gamma_{3}^{-} \leq r_{0} \\
r_{2} \leq r_{0}<\frac{n}{p_{1}}, \\
2 r_{1} \geq \max \left\{2,1+\gamma_{3}^{-}-\frac{n}{p_{0}}+\frac{2 n}{p_{1}}\right\}, \\
r_{2}<\min \left\{\frac{n}{p_{2}}, \frac{2 n}{p_{2}}-\frac{n}{p_{0}}\right\}, \\
\gamma_{1}^{-}>3 r_{1}-2 r_{0}-\gamma_{3}^{-}+\frac{3 n}{p_{0}}-\frac{3 n}{p_{1}}, \\
\gamma_{1}^{-}>1-2 r_{2}+r_{1}-\gamma_{3}^{-}-\frac{n}{p_{1}}+\frac{2 n}{p_{2}}, \\
\gamma_{2}^{-}>1-r_{0}+\frac{n}{p_{0}} .
\end{gathered}
$$

The above inequalities coincide with those in Theorem 3.1, and so the proof is complete.

## References

[1] Gallagher, I.; Planchon, F.; On global infinite energy solutions to the Navier-Stokes equations in two dimensions. Arch. Ration. Mech. Anal., 161(4) (2002), 307-337.
[2] Kato, T.; Ponce, G.; The Navier-Stokes equation with weak initial data. Internat. Math. Res. Notices, 435 (10) (1994). ff., approx. 10 pp.
[3] Lemarié-Rieusset, P. G.; Recent developments in the Navier-Stokes problem, volume 431 of Chapman $\mathcal{E}^{3}$ Hall/CRC Research Notes in Mathematics. Chapman \& Hall/CRC, Boca Raton, FL., 2002.
[4] Linshiz, J. S.; Titi, E. S.; Analytical study of certain magnetohydrodynamic- $\alpha$ models. J. Math. Phys., 48 (6) (2007), 065504, 28.
[5] Pennington, N.; Global solutions to the Lagrangian averaged Navier-Stokes equation in lowregularity Besov spaces. Adv. Differential Equations, 17(7-8) (2012), 697-724.
[6] Pennington, N.; The Lagrangian averaged Navier-Stokes equation with rough data in Sobolev spaces. J. Math. Anal. Appl., 403(1) (2013), 72-88.
[7] Pennington, N.; Local and global low-regularity solutions to generalized Leray-alpha equations. Electron. J. Differential Equations, 2015 No. 170 (2015), 1-24.
[8] Pennington, N.; Low regularity global solutions for a generalized MHD- $\alpha$ system. Nonlinear Anal. Real World Appl., 38 (2017), 171-183.
[9] Tao, T.; Global regularity for a logarithmically supercritical hyperdissipative Navier-Stokes equation. Anal. PDE, 2(3) (2009), 361-366.
[10] Taylor, M. E.; Tools for PDE, volume 81 of Mathematical Surveys and Monographs. American Mathematical Society, Providence, RI. Pseudodifferential operators, paradifferential operators, and layer potentials, 2000.
[11] Taylor, M. E.; Partial differential equations III. Nonlinear equations, volume 117 of Applied Mathematical Sciences. Springer, New York, second edition, 2011.
[12] Wu, J.; Global regularity for a class of generalized magnetohydrodynamic equations. J. Math. Fluid Mech., 13(2) (2011), 295-305.
[13] Yamazaki, K.; On the global regularity of generalized Leray-alpha type models. Nonlinear Anal., 75(2) (2012), 503-515.
[14] Zhao, J.; Zhu, M.; Global regularity for the incompressible MHD- $\alpha$ system with fractional diffusion. Appl. Math. Lett., 29 (2014), 26-29.

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