# GROUND STATE AND MULTIPLE SOLUTIONS FOR CRITICAL FRACTIONAL SCHRÖDINGER-POISSON EQUATIONS WITH PERTURBATION TERMS 

LINTAO LIU, KAIMIN TENG


#### Abstract

In this article, we study a class of critical fractional SchrödingerPoisson system with two perturbation terms. By using variational methods and Lusternik-Schnirelman category theory, the existence of ground state and two nontrivial solutions are established.


## 1. Introduction

In this article, we consider the nonlinear fractional Schrödinger-Poisson system with critical nonlinearity

$$
\begin{gather*}
(-\Delta)^{s} u+u+K(x) \phi u=a(x)|u|^{p-2} u+\mu b(x)|u|^{q-2} u+|u|^{2_{s}^{*}-2} u, \quad \text { in } \mathbb{R}^{3}, \\
(-\Delta)^{t} \phi=K(x) u^{2}, \quad \text { in } \mathbb{R}^{3}, \tag{1.1}
\end{gather*}
$$

where $(-\Delta)^{\alpha}$ is the fractional Laplacian operator for $\alpha=s, t . \quad p, q \in\left(4,2_{s}^{*}\right)$, $s \in\left(\frac{3}{4}, 1\right), 2 s+2 t>3, \mu>0$ is a parameter, $K(x), a(x)$ and $b(x)$ satisfy the following conditions:
(A1) $K(x) \in C\left(\mathbb{R}^{3}\right), K(x) \geq 0$ and $\lim _{|x| \rightarrow \infty} K(x)=K_{\infty}>0$;
(A2) there exist $C_{0}>0$ and $k>0$ such that $K(x) \leq K_{\infty}-\frac{C_{0}}{(1+|x|)^{k}}$ for all $x \in \mathbb{R}^{3}$;
(A3) there exist $C_{1}>0$ and $d>0$ such that $K(x) \leq K_{\infty}+\frac{C_{1}}{(1+|x|)^{d}}$ for all $x \in \mathbb{R}^{3}$;
(A4) $a(x) \in C\left(\mathbb{R}^{3}\right), a(x) \geq 0$ and $\lim _{|x| \rightarrow \infty} a(x)=a_{\infty}>0$;
(A5) there exist $C_{2}>0$ and $a>0$ such that $a(x) \geq a_{\infty}-\frac{C_{2}}{(1+|x|)^{a}}$ for all $x \in \mathbb{R}^{3}$;
(A6) $b(x) \in C\left(\mathbb{R}^{3}\right), b(x) \geq 0$ and $\lim _{|x| \rightarrow \infty} b(x)=0$;
(A7) there exist $C_{3}>0$ and $b>0$ such that $b(x) \geq \frac{C_{3}}{(1+|x|)^{b}}$ for all $x \in \mathbb{R}^{3}$.
Since the first equation in 1.1) is of fractional Schrödinger equation with a potential $\phi$ satisfying the fractional Poisson equation, we call system (1.1) a fractional Schrödinger-Poisson system. In recent years, equations or systems with fractional Laplace operators have been studied extensively because they are widely used in fractional quantum mechanics, physics, chemistry, obstacle problems, optimization and finance, we refer to see [12, 16, 20, 21, 23] and so on. It is also well applied in the mathematical theory of conformal geometry and minimal surface, see 9 .

[^0]As far as we know, there are a few papers considering (1.1) after it was introduced in [15]. In [15], the author studied the local and global well-posedness of the Cauchy problem

$$
\begin{gathered}
i \partial_{t} \Psi+\frac{1}{2} \Delta_{x} \Psi=A_{0} \Psi+\alpha|\Psi|^{\gamma-1} \Psi, \quad(t, x) \in \mathbb{R} \times \mathbb{R} \\
(-\Delta)_{x}^{\sigma / 2} A_{0}=|\Psi|^{2} \\
\Psi(\cdot, x)=f
\end{gathered}
$$

where $\sigma \in(0,1), \alpha= \pm 1,1<\gamma \leq 5$. Recently, Zhang, Do ó and Squassina 38] established the existence of radial ground state solution to the following fractional Schrödinger-Poisson system with a general subcritical or critical nonlinearity

$$
\begin{gathered}
(-\Delta)^{s} u+\lambda \phi u=f(u), \quad \text { in } \mathbb{R}^{3} \\
(-\Delta)^{t} \phi=\lambda u^{2}, \quad \text { in } \mathbb{R}^{3} .
\end{gathered}
$$

Teng [33] studied the existence of a nontrivial ground state solution through using the method of Pohozaev-Nehari manifold, the monotonic trick and global compactness Lemma for the system

$$
\begin{gathered}
(-\Delta)^{s} u+V(x) u+\phi u=|u|^{p-1} u, \quad \text { in } \mathbb{R}^{3} \\
(-\Delta)^{s} \phi=u^{2}, \quad \text { in } \mathbb{R}^{3}
\end{gathered}
$$

Using a similar argument, Teng in 32 also studied the existence of ground state solutions for the critical problem with a perturbation term

$$
\begin{gathered}
(-\Delta)^{s} u+V(x) u+\phi u=\mu|u|^{q-1} u+|u|^{2_{s}^{*}-2} u, \quad \text { in } \mathbb{R}^{3}, \\
(-\Delta)^{t} \phi=u^{2}, \quad \text { in } \mathbb{R}^{3} .
\end{gathered}
$$

For other related works, see [22, 28] and their references.
On the other hand, when $s=t=1$, system (1.1) reduces to classical SchrödingerPoisson system written by a more general form

$$
\begin{gather*}
-\Delta u+V(x) u+K(x) \phi u=f(x, u), \quad \text { in } \mathbb{R}^{3}  \tag{1.2}\\
-\Delta \phi=K(x) u^{2}, \quad \text { in } \mathbb{R}^{3}
\end{gather*}
$$

This is called the system of Schrödinger-Poisson equations because it consists of a Schrödinger equation coupled with a Poisson term. In the previous decades, there has been a lot of work dealing with the system 1.2 under different assumptions on $V, K$ and $f$, see [2, 3, 4, 8, 10, 11, 15, 17, 19, 24, 27, 29, 35, 37, 39, 40] and the references therein. For example, in [3], the authors proved the existence of ground state solutions for the subcritical $3<p<6$ and the critical case $f=|u|^{p-2} u+u^{5}$ with $4<p<6$. For the case $p \leq 2$ or $p \geq 6$, the reader may see [11] and for the case $2<p<6$, can see [2, 3, 8, 10, 24]. In the case of $V$ being non-radial, $K \equiv 1$ and $f=|u|^{p-2} u$, the existence of ground state solution for system 1.2 was obtained in [3, 40] for $4<p<6$ and $3<p \leq 4$; In [5], the authors proved the existence of ground state and bound states for the case when $V \equiv 1$ and $f=a(x)|u|^{p-2} u$ with $4<p<6$. In [37, the author considered a general critical situation with two perturbation term and obtained the existence and multiplicity of solutions via using Lusternik-Schnirelman category due to [1, 6].

To the best of our knowledge, there are few papers on the multiplicity solutions for system (1.1). Inspired by [1, 6, 37, we construct two mappings:

$$
\begin{gathered}
F_{R}: S^{2}=\left\{y \in \mathbb{R}^{3}:|y|=1\right\} \rightarrow\left\{u \in M: I(u) \leq m_{\infty}-\varepsilon(R)\right\} \\
G:\left\{u \in M: I(u)<m_{\infty}\right\} \rightarrow S^{2}
\end{gathered}
$$

so that $G \circ F_{R}$ homotopic to the identity. Using the theory of Lusternik-Schnirelman category, we will establish the existence of two nontrivial solutions for system (1.1).

Our main results are stated as follows.
Theorem 1.1. Assume that $K, a$ and $b$ satisfy (A1), (A2), (A4)-(A6) with $0<$ $k<\alpha$, where $\alpha=\min \{a,(3+2 s) p\}$. Then problem 1.1) admits a positive ground state solution.

Theorem 1.2. Suppose that (A1), (A3)-(A7) hold with $b<\min \{\alpha, \beta\}$, where $\alpha=\min \{a,(3+2 s) p\}$ and $\beta=\min \{d, 6+4 s\}$. Then problem 1.1 admits $a$ positive ground state solution.

Theorem 1.3. Assume that $K \in C^{1}\left(\mathbb{R}^{3}\right)$, $a \in C^{1}\left(\mathbb{R}^{3}\right)$ and $b(x)$ satisfy (A1), (A3)(A7) with $K(x) \geq K_{\infty}, a(x) \leq a_{\infty}$ and meas $\left\{x \in \mathbb{R}^{3}: K(x) \geq K_{\infty}\right\}>0$. Then there exists $\mu_{0}>0$ small such that for any $\mu \in\left(0, \mu_{0}\right)$, problem 1.1 admits at least two nontrivial solutions.

The rest of the paper is organized as follows: In Section 2, we give some preliminaries. In Section 3, we prove Theorem 1.1 and Theorem 1.2, Section 4 devotes to proving Theorem 1.3 .

## 2. Preliminary lemmas

In the sequel, we use the following notation:

- $H^{s}\left(\mathbb{R}^{3}\right)$ denotes the fractional sobolev space with norm

$$
\|u\|^{2}:=\int_{\mathbb{R}^{3}}\left(\left|(-\Delta)^{\frac{s}{2}} u\right|^{2}+u^{2}\right) d x
$$

and

$$
D^{s, 2}\left(\mathbb{R}^{3}\right):=\left\{u \in L^{2_{s}^{*}}\left(\mathbb{R}^{3}\right):(-\Delta)^{\frac{s}{2}} u \in L^{2}\left(\mathbb{R}^{3}\right)\right\}
$$

denotes the homogeneous fractional sobolev space with the norm

$$
\|u\|_{D^{s, 2}}^{2}:=\int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{s}{2}} u\right|^{2} d x
$$

- $C$ denotes a universal positive constant (possibly different).
- It is well known that $H^{\alpha}\left(\mathbb{R}^{3}\right)$ is continuously embedded into $L^{p}\left(\mathbb{R}^{3}\right)$ for $2 \leq p \leq 2_{\alpha}^{*}\left(2_{\alpha}^{*}=\frac{6}{3-2 \alpha}\right)$, and for any $\alpha \in(0,1)$, there exists a best constant $S_{\alpha}>0$ such that

$$
S_{\alpha}=\inf _{u \in D^{\alpha, 2}} \frac{\int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u\right|^{2} d x}{\left(\int_{\mathbb{R}^{3}}|u(x)|^{2_{\alpha}^{*}} d x\right)^{\frac{2}{2 \alpha}}}
$$

- For simplicity, we assume $K_{\infty}=1$ and $a_{\infty}=1$. Denote $H=H^{s}\left(\mathbb{R}^{3}\right)$ and $D^{s, 2}=D^{s, 2}\left(\mathbb{R}^{3}\right)$.
In this section, we assume (A1), (A4) and (A6) hold. Similar to the argument in [24], we know the function $\phi_{u}^{t}$ has the following properties.
Lemma 2.1. For any $u \in H$, we have
(i) $\phi_{u}^{t} \geq 0$;
(ii) $\phi_{\hbar u}^{t}=\hbar^{2} \phi_{u}^{t}, \forall \hbar>0$;
(iii) $\left\|\phi_{u}^{t}\right\|_{D^{t, 2}} \leq C\|u\|_{\frac{12}{3+2 t}}^{2} \leq C\|u\|^{2}, \int_{\mathbb{R}^{3}} \phi_{u}^{t} u^{2} d x \leq C\|u\|_{\frac{12}{3+2 t}}^{4} \leq C\|u\|^{4}$.

By the Lax-Milgram theorem, there exists a unique $\phi_{u}^{t} \in D^{t, 2}\left(\mathbb{R}^{3}\right)$ such that $(-\Delta)^{t} \phi_{u}^{t}=K(x) u^{2}$. Thus, we can rewrite (1.1) as

$$
\begin{equation*}
(-\Delta)^{s} u+u+K(x) \phi_{u}^{t} u=a(x)|u|^{p-2} u+\mu b(x)|u|^{q-2} u+|u|^{2_{s}^{*}-2} u \tag{2.1}
\end{equation*}
$$

To find weak solutions to $(2.1)$, we look for critical points of the functional $I(u)$ : $H \rightarrow \mathbb{R}$ associated with 2.1 which is defined by

$$
\begin{aligned}
I(u)= & \frac{1}{2} \int_{\mathbb{R}^{3}}\left(\left|(-\Delta)^{\frac{s}{2}} u\right|^{2}+u^{2}\right) d x+\frac{1}{4} \int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x \\
& -\frac{1}{p} \int_{\mathbb{R}^{3}} a(x)|u|^{p} d x-\frac{\mu}{q} \int_{\mathbb{R}^{3}} b(x)|u|^{q} d x-\frac{1}{2_{s}^{*}} \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x .
\end{aligned}
$$

To prove the compactness, we need to consider the following problem at infinity associated with (2.1):

$$
\begin{equation*}
(-\Delta)^{s} u+u+\hat{\phi}_{u}^{t} u=|u|^{p-2} u+|u|^{2_{s}^{*}-2} u, \quad u>0 \tag{2.2}
\end{equation*}
$$

where $\hat{\phi}_{u}^{t} \in \mathcal{D}^{t, 2}\left(\mathbb{R}^{3}\right)$ is the unique solution to problem

$$
(-\Delta)^{t} \phi=u^{2}
$$

The functional associated with 2.2 is

$$
\begin{aligned}
I_{\infty}(u)= & \frac{1}{2} \int_{\mathbb{R}^{3}}\left(\left|(-\Delta)^{\frac{s}{2}} u\right|^{2}+u^{2}\right) d x+\frac{1}{4} \int_{\mathbb{R}^{3}} \hat{\phi}_{u}^{t} u^{2} d x \\
& -\frac{1}{p} \int_{\mathbb{R}^{3}}|u|^{p} d x-\frac{1}{2_{s}^{*}} \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x .
\end{aligned}
$$

Let

$$
m=\inf _{u \in M} I(u), \quad m_{\infty}=\inf _{u \in M_{\infty}} I_{\infty}(u)
$$

where

$$
\begin{gathered}
M=\left\{u \in H^{s}\left(\mathbb{R}^{3}\right) \backslash\{0\}:\left[I^{\prime}(u), u\right]=0\right\} \\
M_{\infty}=\left\{u \in H^{s}\left(\mathbb{R}^{3}\right) \backslash\{0\}:\left[I_{\infty}^{\prime}(u), u\right]=0\right\}
\end{gathered}
$$

are Nehari manifolds correspond to the functionals $I$ and $I_{\infty}$, respectively. Similar argument as [22, Proposition 3.4], we can obtain the following Lemma.
Lemma 2.2. By using [22, Proposition 3.4], problem (2.2) has a positive ground state solution $u_{\infty} \in C^{1,2 s+\sigma-1}\left(\mathbb{R}^{3}\right) \cap L^{\infty}\left(\mathbb{R}^{3}\right)$, where $\sigma \in(0,1)$ and $2 s+\sigma>1$.

From $u \in H \cap C^{1,2 s+\sigma-1}\left(\mathbb{R}^{3}\right)$, we see that $\lim _{|x| \rightarrow \infty} u_{\infty}(x)=0$. Similar as the proof of [32, Proposition 3.8], we conclude that there exists $C>0$ such that

$$
0<u_{\infty}(x) \leq \frac{C}{(1+|x|)^{3+2 s}}, \quad \forall x \in \mathbb{R}^{3}
$$

Moreover, in [22], the authors showed that $m_{\infty}=c_{\infty}=\inf _{\gamma \in \Gamma} \max _{t \in[0,1]} I_{\infty}(\gamma(t))$, where $\Gamma=\left\{\gamma \in C([0,1], H): \gamma(0)=0, I_{\infty}(\gamma(1))<0\right\}$ and

$$
\begin{equation*}
m_{\infty}=c_{\infty}=\inf _{u \in H \backslash\{0\}} \max _{t \geq 0} I_{\infty}(t u) \tag{2.3}
\end{equation*}
$$

To prove the $(P S)_{c}$ condition, we need the following function and its estimates (see [32])

$$
v_{\varepsilon}(x)=\psi(x) U_{\varepsilon}(x), \quad x \in \mathbb{R}^{3}
$$

where $U_{\varepsilon}(x)=\varepsilon^{-\frac{3-2 s}{2}} u^{*}(x / \varepsilon)$,

$$
u^{*}\left(\frac{x}{\varepsilon}\right)=\frac{\tilde{u}\left(x / S_{s}^{\frac{1}{2 s}}\right)}{\|\tilde{u}\|_{2_{s}^{*}}}
$$

$\kappa \in \mathbb{R} \backslash\{0\}, \mu>0$, and $x_{0} \in \mathbb{R}^{3}$ are fixed constants, $\tilde{u}(x)=\kappa\left(\mu^{2}+\left|x-x_{0}\right|^{2}\right)^{-\frac{3-2 s}{2}}$, and $\psi \in C^{\infty}\left(\mathbb{R}^{3}\right)$ such that $0 \leq \psi \leq 1$ in $\mathbb{R}^{3}, \psi(x) \equiv 1$ in $B_{\delta}$ and $\psi(x) \equiv 0$ in $\mathbb{R}^{3} \backslash B_{2 \delta}$. We know that

$$
\begin{gather*}
\int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}(x)\right|^{2} d x \leq S_{s}^{\frac{3}{2 s}}+O\left(\varepsilon^{3-2 s}\right),  \tag{2.4}\\
\int_{\mathbb{R}^{3}}\left|v_{\varepsilon}(x)\right|^{2_{s}^{*}} d x=S_{s}^{\frac{3}{2 s}}+O\left(\varepsilon^{3}\right),  \tag{2.5}\\
\int_{\mathbb{R}^{3}}\left|v_{\varepsilon}(x)\right|^{p} d x= \begin{cases}O\left(\varepsilon^{\frac{(2-p) 3+2 s p}{2}}\right), & p>\frac{3}{3-2 s} ; \\
O\left(\varepsilon^{\frac{(2-p) 3+2 s p}{2}}|\log \varepsilon|\right), & p=\frac{3}{3-2 s} ; \\
O\left(\varepsilon^{\frac{3-2 s}{2} p}\right), & p<\frac{3}{3-2 s} .\end{cases} \tag{2.6}
\end{gather*}
$$

Lemma 2.3. Let $\left\{u_{n}\right\} \subset H$ be a bounded sequence such that $I\left(u_{n}\right) \rightarrow c \in\left(0, m_{\infty}\right)$ and $I^{\prime}\left(u_{n}\right) \rightarrow 0$. Then $\left\{u_{n}\right\}$ admits a strongly convergent subsequence in $H$.

Proof. First we show that $m_{\infty}<\frac{s}{3} S_{s}^{\frac{3}{2 s}}$. By (2.3), we see that $c_{\infty} \leq \sup _{t \geq 0} I_{\infty}\left(t v_{\varepsilon}\right)$. Thus we only need to prove $\sup _{t \geq 0} I_{\infty}\left(t v_{\varepsilon}\right)<\frac{s}{3} S_{s}^{\frac{3}{2 s}}$ for $\varepsilon>0$ small. By Lemma 2.1, we have

$$
\begin{equation*}
I_{\infty}\left(t v_{\varepsilon}\right) \leq \frac{1}{2} t^{2}\left\|v_{\varepsilon}\right\|^{2}+C t^{4}\left\|v_{\varepsilon}\right\|^{4}-\frac{1}{2_{s}^{*}} t^{2_{s}^{*}}\left\|v_{\varepsilon}\right\|_{2_{s}^{*}}^{2_{s}^{*}} . \tag{2.7}
\end{equation*}
$$

Form (2.4)-(2.6), there exists $\varepsilon_{1}>0$ small enough such that

$$
\begin{gather*}
\left\|v_{\varepsilon}\right\|^{2}:=\int_{\mathbb{R}^{3}}\left(\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2}+v_{\varepsilon}^{2}\right) d x \leq S_{s}^{\frac{3}{2 s}}+O\left(\varepsilon^{3-2 s}\right)+O\left(\varepsilon^{3-2 s}\right) \leq \frac{3}{2} S_{s}^{\frac{3}{2 s}}  \tag{2.8}\\
\left\|v_{\varepsilon}\right\|_{2_{s}^{*}}^{2_{s}^{*}}=S_{s}^{\frac{3}{2 s}}+O\left(\varepsilon^{3}\right) \geq \frac{1}{2} S_{s}^{\frac{3}{2 s}} \tag{2.9}
\end{gather*}
$$

for $\varepsilon \in\left(0, \varepsilon_{1}\right)$. Thus, form 2.7)-2.9), we have

$$
\begin{equation*}
I_{\infty}\left(t v_{\varepsilon}\right) \leq \frac{3}{4} t^{2} S_{s}^{\frac{3}{2 s}}+C \frac{9}{4} t^{4} S_{s}^{\frac{3}{s}}-\frac{1}{2_{s}^{*}} t^{2_{s}^{*}} \frac{1}{2} S_{s}^{\frac{3}{2 s}} \tag{2.10}
\end{equation*}
$$

By $2<4<2_{s}^{*}$, there exist a small $t_{1}>0$ and a large $t_{2}>0$ independent of $\varepsilon \in\left(0, \varepsilon_{1}\right)$ such that

$$
\begin{equation*}
\sup _{t \in\left[0, t_{1}\right] \cup\left[t_{2},+\infty\right)} I_{\infty}\left(t v_{\varepsilon}\right)<\frac{s}{3} S_{s}^{\frac{3}{2 s}} . \tag{2.11}
\end{equation*}
$$

Form Lemma 2.1 and 2.4-2.6, we obtain

$$
\begin{align*}
\sup _{t \in\left[t_{1}, t_{2}\right]} I_{\infty}\left(t v_{\varepsilon}\right) \leq & \sup _{t \geq 0}\left[\frac{1}{2} t^{2} \int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{s}{2}} v_{\varepsilon}\right|^{2} d x-\frac{1}{2_{s}^{*}} t^{2_{s}^{*}} \int_{\mathbb{R}^{3}}\left|v_{\varepsilon}(x)\right|^{2_{s}^{*}} d x\right] \\
& +C\left\|v_{\varepsilon}\right\|_{2}^{2}+C\left\|v_{\varepsilon}\right\|_{\frac{12}{3+2 t}}^{4}-C\left\|v_{\varepsilon}\right\|_{p}^{p}  \tag{2.12}\\
= & \frac{s}{3} S_{s}^{\frac{3}{2 s}}+O\left(\varepsilon^{3-2 s}\right)-C \varepsilon^{\frac{(2-p) 3+2 s p}{2}}
\end{align*}
$$

In view of $p \in\left(4,2_{s}^{*}\right), s \in\left(\frac{3}{4}, 1\right)$, so we see that $\frac{(2-p) 3+2 s p}{2}<3-2 s$. By choosing $\varepsilon \in\left(0, \varepsilon_{1}\right)$ small, we obtain

$$
\begin{equation*}
\sup _{t \in\left[t_{1}, t_{2}\right]} I_{\infty}\left(t v_{\varepsilon}\right)<\frac{s}{3} S_{s}^{\frac{3}{2 s}} \tag{2.13}
\end{equation*}
$$

By (2.11) and 2.13), we have

$$
\begin{equation*}
m_{\infty}<\frac{s}{3} S_{s}^{\frac{3}{2 s}} \tag{2.14}
\end{equation*}
$$

Since $\left\{u_{n}\right\}$ is bounded in $H$, up to a subsequence, we may assume that $u_{n} \rightharpoonup u$ weakly in $H, u_{n} \rightarrow u$ in $L_{\text {loc }}^{r}\left(\mathbb{R}^{3}\right)$ for $1 \leq r<2_{s}^{*}$ and $u_{n} \rightarrow u$ a.e. $\mathbb{R}^{3}$. Thus by standard argument, we can show that $I^{\prime}(u)=0$. Set $v_{n}=u_{n}-u$. By the Brezis-Lieb Lemma in 36, we have that

$$
\begin{align*}
\left\|v_{n}\right\|^{2} & =\left\|u_{n}\right\|^{2}-\|u\|^{2}+o(1) \\
\left\|v_{n}\right\|_{2_{s}^{*}}^{2_{s}^{*}} & =\left\|u_{n}\right\|_{2_{s}^{*}}^{2_{s}^{*}}-\|u\|_{2_{s}^{*}}^{2_{s}^{*}}+o(1), \tag{2.15}
\end{align*}
$$

and

$$
\begin{aligned}
\int_{\mathbb{R}^{3}} a(x)\left|v_{n}\right|^{p} d x & =\int_{\mathbb{R}^{3}} a(x)\left|u_{n}\right|^{p}-\int_{\mathbb{R}^{3}} a(x)|u|^{p}+o(1), \\
\int_{\mathbb{R}^{3}} b(x)\left|v_{n}\right|^{q} d x & =\int_{\mathbb{R}^{3}} b(x)\left|u_{n}\right|^{q}-\int_{\mathbb{R}^{3}} b(x)|u|^{q}+o(1)
\end{aligned}
$$

From $\lim _{|x| \rightarrow \infty} a(x)=1, \lim _{|x| \rightarrow \infty} b(x)=0$, and $v_{n} \rightarrow 0$ in $L_{\text {loc }}^{r}\left(\mathbb{R}^{3)}\right.$ for any $r \in\left[1,2_{s}^{*}\right)$, we deduce that

$$
\begin{gather*}
\int_{\mathbb{R}^{3}} a(x)\left|u_{n}\right|^{p}-\int_{\mathbb{R}^{3}} a(x)|u|^{p}=\int_{\mathbb{R}^{3}}\left|v_{n}\right|^{p} d x+o(1) \\
\int_{\mathbb{R}^{3}} b(x)\left|u_{n}\right|^{q}-\int_{\mathbb{R}^{3}} b(x)|u|^{q}=o(1) \tag{2.16}
\end{gather*}
$$

By [33, Lemma 2.5], we can see that

$$
\int_{\mathbb{R}^{3}} K(x) \phi_{v_{n}}^{t} v_{n}^{2} d x=\int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}}^{t} u_{n}^{2} d x-\int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x+o(1)
$$

From $\lim _{|x| \rightarrow \infty} K(x)=1$ and Hölder's inequality, it is easy to deduce that

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} \hat{\phi}_{v_{n}}^{t} v_{n}^{2} \mathrm{~d} x=\int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}}^{t} u_{n}^{2} d x-\int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x+o(1) . \tag{2.17}
\end{equation*}
$$

Thus, from (2.15)-2.17), it follows that

$$
\begin{equation*}
c-I(u)=I_{\infty}\left(v_{n}\right)+o(1) . \tag{2.18}
\end{equation*}
$$

By using [13, Proposition 5.1.1], we see that $u \in L^{\infty}\left(\mathbb{R}^{3}\right)$. Then by [36, Lemmas 8.1 and 8.9], we have that

$$
\begin{gather*}
\left|\int_{\mathbb{R}^{3}}\left(u_{n}^{2_{s}^{*}-1}-u^{2_{s}^{*}-1}-v_{n}^{2_{s}^{*}-1}\right) \varphi d x\right|=o(1)\|\varphi\|, \quad \forall \varphi \in H \\
\left|\int_{\mathbb{R}^{3}} a(x)\left(\left|u_{n}\right|^{p-2} u_{n}-|u|^{p-2} u-\left|v_{n}\right|^{p-2} v_{n}\right) \varphi d x\right|=o(1)\|\varphi\|, \quad \forall \varphi \in H,  \tag{2.19}\\
\left|\int_{\mathbb{R}^{3}} b(x)\left(\left|u_{n}\right|^{q-2} u_{n}-|u|^{q-2} u-\left|v_{n}\right|^{q-2} v_{n}\right) \varphi d x\right|=o(1)\|\varphi\|, \quad \forall \varphi \in H .
\end{gather*}
$$

Together with $\lim _{|x| \rightarrow \infty} a(x)=1, \lim _{|x| \rightarrow \infty} b(x)=0$, we deduce that

$$
\begin{gather*}
\left|\int_{\mathbb{R}^{3}}\left[a(x)\left(\left|u_{n}\right|^{p-2} u_{n}-|u|^{p-2} u\right)-\left|v_{n}\right|^{p-2} v_{n}\right] \varphi d x\right|=o(1)\|\varphi\|, \quad \forall \varphi \in H  \tag{2.20}\\
\left|\int_{\mathbb{R}^{3}} b(x)\left(\left|u_{n}\right|^{q-2} u_{n}-|u|^{q-2} u\right) \varphi d x\right|=o(1)\|\varphi\|, \quad \forall \varphi \in H
\end{gather*}
$$

Using [33, Lemma 2.5], we have

$$
\left|\int_{\mathbb{R}^{3}} K(x)\left(\phi_{u_{n}}^{t} u_{n}-\phi_{u}^{t} u-\phi_{v_{n}}^{t} v_{n}\right) \varphi d x\right|=o(1)\|\varphi\|, \quad \forall \varphi \in H
$$

From $\lim _{|x| \rightarrow \infty} K(x)=1$, and similar to the of proof of (2.17), we obtain

$$
\begin{equation*}
\left|\int_{\mathbb{R}^{3}} K(x)\left(\phi_{u_{n}}^{t} u_{n}-\phi_{u}^{t} u\right) \varphi d x-\int_{\mathbb{R}^{3}} \hat{\phi}_{v_{n}}^{t} v_{n} \varphi d x\right|=o(1)\|\varphi\|, \quad \forall \varphi \in H \tag{2.21}
\end{equation*}
$$

Hence, by 2.19)-2.21, it holds

$$
\begin{equation*}
I_{\infty}^{\prime}\left(v_{n}\right)=o(1) \tag{2.22}
\end{equation*}
$$

We claim $v_{n} \rightarrow 0$ in $H$. Two cases occur: either

$$
\lim _{n \rightarrow \infty} \sup _{y \in \mathbb{R}_{3}} \int_{B_{1}(y)}\left|v_{n}\right|^{2} d x=0
$$

or there exists $\gamma>0$ such that

$$
\lim _{n \rightarrow \infty} \sup _{y \in \mathbb{R}_{3}} \int_{B_{1}(y)}\left|v_{n}\right|^{2} d x \geq \gamma
$$

Thus, either $\left\|v_{n}\right\|_{r} \rightarrow 0$ for any $r \in\left(2,2_{s}^{*}\right)$ through using vanishing Lemma, or there $y_{n} \in \mathbb{R}^{3}$ with $\left|y_{n}\right| \rightarrow \infty$ such that $v_{n}\left(.+y_{n}\right) \rightharpoonup v \neq 0$ weakly in $H$. If $v_{n}\left(.+y_{n}\right) \rightharpoonup v \neq 0$ weakly in $H$, from $\sqrt{2.18}$ and 2.22 , it follows that $c-I(u)=$ $I_{\infty}\left(v_{n}\left(.+y_{n}\right)\right)+o(1)$ and $I_{\infty}^{\prime}\left(v_{n}\left(.+y_{n}\right)\right)=o(1)$. Thus $I_{\infty}^{\prime}(v)=0$ and

$$
\begin{aligned}
c-I(u)= & I_{\infty}\left(v_{n}\left(.+y_{n}\right)\right)-\frac{1}{4}\left[I_{\infty}^{\prime}\left(v_{n}\left(.+y_{n}\right)\right), v_{n}\left(.+y_{n}\right)\right] \\
= & \frac{1}{4} \|\left(\left.v_{n}\left(.+y_{n}\right) \|^{2}+\left(\frac{1}{4}-\frac{1}{p}\right) \int_{\mathbb{R}^{3}} \right\rvert\,\left(\left.v_{n}\left(.+y_{n}\right)\right|^{p} d x\right.\right. \\
& \left.+\left(\frac{1}{4}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}} \right\rvert\,\left(\left.v_{n}\left(\cdot+y_{n}\right)\right|^{2_{s}^{*}} d x+o(1),\right.
\end{aligned}
$$

form which we obtain

$$
\begin{aligned}
c & \geq I(u)+\frac{1}{4}\|v\|^{2}+\left(\frac{1}{4}-\frac{1}{p}\right) \int_{\mathbb{R}^{3}}|v|^{p} d x+\left(\frac{1}{4}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}}|v|^{2_{s}^{*}} d x \\
& \left.=I(u)+I_{\infty}(v)-\frac{1}{4}\left[I_{\infty}^{\prime}(v), v\right)\right]=I(u)+I_{\infty}(v)
\end{aligned}
$$

By the definition of $m_{\infty}$, we have $I_{\infty}(v) \geq m_{\infty}$. Since $I^{\prime}(u)=0$, we have

$$
\begin{aligned}
I(u)= & I(u)-\frac{1}{4}\left[I^{\prime}(u), u\right] \\
= & \frac{1}{4}\|u\|^{2}+\left(\frac{1}{4}-\frac{1}{p}\right) \int_{\mathbb{R}^{3}} a(x)|u|^{p} d x \\
& +\left(\frac{1}{4}-\frac{1}{q}\right) \mu \int_{\mathbb{R}^{3}} b(x)|u|^{q} d x+\left(\frac{1}{4}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x \geq 0
\end{aligned}
$$

which leads to a contradiction with $c<m_{\infty}$. Thus $\left\|v_{n}\right\|_{L^{r}} \rightarrow 0$ for any $r \in\left(2,2_{s}^{*}\right)$. By 2.18 and 2.22, we have

$$
\begin{gathered}
c-I(u)=\frac{1}{2}\left\|v_{n}\right\|^{2}-\frac{1}{2_{s}^{*}}\left\|v_{n}\right\|_{2_{s}^{*}}^{2_{s}^{*}}+o(1) \\
\left\|v_{n}\right\|^{2}-\left\|v_{n}\right\|_{2_{s}^{*}}^{2_{s}^{*}}=o(1)
\end{gathered}
$$

Up to a subsequence, we may assume that $\left\|v_{n}\right\|^{2} \rightarrow l$. Thus $\left\|v_{n}\right\|_{2_{s}^{*}}^{2_{s}^{*}} \rightarrow l$. If $l>0$, by the definition of $S_{s}$, we obtain $l \geq\left(S_{s}\right)^{\frac{3}{2 s}}$. Hence,

$$
c=I(u)+\frac{1}{2}\left\|v_{n}\right\|^{2}-\frac{1}{2_{s}^{*}}\left\|v_{n}\right\|_{2_{s}^{*}}^{2_{s}^{*}}=I(u)+\left(\frac{1}{2}-\frac{1}{2_{s}^{*}}\right) l \geq \frac{s}{3} S_{s}^{\frac{3}{2 s}},
$$

which contradicts with $c<m_{\infty}<\frac{s}{3} S_{s}^{\frac{3}{2 s}}$. Thus $l=0$ and we complete the proof.
Lemma 2.4. Suppose that $\alpha, \beta>n, f, g \in L^{\infty}\left(\mathbb{R}^{n}\right)$ and

$$
f(x) \leq \frac{C_{1}}{(1+|x|)^{\alpha}}, \quad g(x) \leq \frac{C_{2}}{(1+|x|)^{\beta}}
$$

Then there exits $C>0$ such that

$$
|f * g(x)| \leq \frac{C}{(1+|x|)^{\gamma}}
$$

where $\gamma=\min \{\alpha, \beta\}$.
Proof. By direct computations,

$$
\begin{aligned}
& |f * g(x)| \\
& =\left|\int_{\mathbb{R}^{n}} \frac{C_{1}}{(1+|x-y|)^{\alpha}} \frac{C_{2}}{(1+|y|)^{\beta}} \mathrm{d} y\right| \\
& =\int_{|x-y| \geq \frac{|x|}{2}} \frac{C_{1}}{(1+|x-y|)^{\alpha}} \frac{C_{2}}{(1+|y|)^{\beta}} \mathrm{d} y+\int_{|x-y|<\frac{|x|}{2}} \frac{C_{1}}{(1+|x-y|)^{\alpha}} \frac{C_{2}}{(1+|y|)^{\beta}} \mathrm{d} y \\
& \leq \frac{C_{1}}{\left(1+\frac{|x|}{2}\right)^{\alpha}} \int_{|x-y| \geq \frac{|x|}{2}} \frac{C_{2}}{(1+|y|)^{\beta}} \mathrm{d} y+\int_{\frac{|x|}{2}<|y|<\frac{3}{2}|x|} \frac{C_{1}}{(1+|x-y|)^{\alpha}} \frac{C_{2}}{(1+|y|)^{\beta}} \mathrm{d} y \\
& \leq \frac{C_{1}}{\left(1+\frac{|x|}{2}\right)^{\alpha}} \int_{\mathbb{R}^{n}} \frac{C_{2}}{(1+|y|)^{\beta}} \mathrm{d} y+\frac{C_{2}}{\left(1+\frac{|x|}{2}\right)^{\beta}} \int_{\frac{|x|}{2}<|y|<\frac{3}{2}|x|} \frac{C_{1}}{(1+|x-y|)^{\alpha}} \mathrm{d} y \\
& \leq \frac{C_{1} C_{2} 2^{\alpha}}{(2+|x|)^{\alpha}} \int_{\mathbb{R}^{n}} \frac{1}{(1+|y|)^{\beta}} \mathrm{d} y+\frac{C_{1} C_{2} 2^{\beta}}{(2+|x|)^{\beta}} \int_{|x-y|<\frac{5}{2}|x|} \frac{1}{(1+|x-y|)^{\alpha}} \mathrm{d} y \\
& \leq \frac{C_{1} C_{2} 2^{\alpha}}{(2+|x|)^{\alpha}} \int_{\mathbb{R}^{n}} \frac{1}{(1+|y|)^{\beta}} \mathrm{d} y+\frac{C_{1} C_{2} 2^{\beta}}{(2+|x|)^{\beta}} \int_{\mathbb{R}^{n}} \frac{1}{(1+|x-y|)^{\alpha}} \mathrm{d} y \\
& \leq C\left(\frac{1}{(2+|x|)^{\alpha}}+\frac{1}{(2+|x|)^{\beta}}\right) \\
& \leq \frac{C}{(1+|x|)^{\gamma}},
\end{aligned}
$$

where $\gamma=\min \{\alpha, \beta\}$.
Now we recall the definition of Lusternik-Schnirelman category.

Definition 2.5. (i) For a topological space $X$, we say a non-empty, closed subset $A \subset X$ is contractible to a point in $X$ if and only if there exist a continuous mapping $\eta:[0,1] \times A \rightarrow X$ such that for some $x_{0} \in X$,
(a) $\eta(0, x)=x$ for all $x \in A$,
(b) $\eta(1, x)=x_{0}$ for all $x \in A$.
(ii) We define
$\operatorname{cat}(X)=\min \left\{k \in \mathbb{N}\right.$ : there exist closed subsets $A_{1}, \ldots, A_{k} \subset X$ such that
$A_{i}$ is contractible to a point in X for all $i$ and

$$
\left.\cup_{i=1}^{k} A_{i}=X\right\}
$$

We say $\operatorname{cat}(X)=\infty$ if do not exist finitely many closed subsets $A_{1}, \ldots, A_{k} \subset X$ such that $A_{i}$ is contractible to a point in $X$ for all $i$ and $\cup_{i=1}^{k} A_{i}=X$.

We need the following two important lemmas. See [1, Proposition 2.4 and Lemma 2.5].

Lemma 2.6. Suppose that $\mathcal{M}$ is a Hilbert manifold and $\Psi \in C^{1}(\mathcal{M}, \mathbb{R})$. Assume that there exist $c_{0} \in \mathbb{R}$ and $k \in \mathbb{N}$ such that $\Psi(u)$ satisfies the Palais-Smale condition for $c \leq c_{0}$ and $\operatorname{cat}\left(\left\{u \in \mathcal{M}: \Psi(u) \leq c_{0}\right\}\right) \geq k$. Then $\Psi(u)$ has at least $k$ critical points in $\left\{u \in \mathcal{M}: \Psi(u) \leq c_{0}\right\}$.

Lemma 2.7. Let $X$ be a topological space. Suppose that there exist two continuous mappings $F: S^{2}=\left\{y \in \mathbb{R}^{3}:|y|=1\right\} \rightarrow X$ and $G: X \rightarrow S^{2}$, such that $G \circ F$ is homotopic to identity id : $S^{2} \rightarrow S^{2}$, that is, there is a continuous mapping $\zeta:[0,1] \times S^{2} \rightarrow S^{2}$ such that $\zeta(0, x)=(G \circ F)(x)$ for all $x \in S^{2}$ and $\zeta(1, x)=x$ for all $x \in S^{2}$. Then $\operatorname{cat}(X) \geq 2$.

## 3. Proof of main results

Proof of Theorem 1.1. Let $\left\{u_{n}\right\} \subset M$ be a minimizing sequence for functional $I$, that is, $\left\{u_{n}\right\} \subset M$ and $I\left(u_{n}\right) \rightarrow m$, where

$$
M=\left\{u \in H \backslash\{0\}: G(u)=\left[I^{\prime}(u), u\right]=0\right\} .
$$

We claim $I^{\prime}\left(u_{n}\right) \rightarrow 0$. By the Lagrange multiplier Theorem, there exists $\lambda_{n} \in \mathbb{R}$ such that

$$
I^{\prime}\left(u_{n}\right)-\lambda_{n} G^{\prime}\left(u_{n}\right) \rightarrow 0
$$

Since $u_{n} \subset M$, we have

$$
m+o(1)=I\left(u_{n}\right)-\frac{1}{4}\left(I^{\prime}\left(u_{n}\right), u_{n}\right) \geq \frac{1}{4}\left\|u_{n}\right\|^{2}
$$

which implies that $\left\{u_{n}\right\}$ is bounded in $H$. Hence

$$
\begin{equation*}
\lambda_{n}\left[G^{\prime}\left(u_{n}\right), u_{n}\right] \rightarrow 0 \tag{3.1}
\end{equation*}
$$

By (A4) and (A6), for any $\varepsilon>0$, there exists $C_{\varepsilon}>0$ such that

$$
a(x)|u|^{p}+b(x)|u|^{q}+|u|^{2_{s}^{*}} \leq \varepsilon|u|^{2}+C_{\varepsilon}|u|^{2_{s}^{*}} .
$$

Taking $\varepsilon=1 / 2$ and recalling the definition of $S_{s}$, we have

$$
\left\|u_{n}\right\|^{2} \leq \int_{\mathbb{R}^{3}} a(x)\left|u_{n}\right|^{p} d x+\int_{\mathbb{R}^{3}} \mu b(x)\left|u_{n}\right|^{q} d x+\int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{s}^{*}} d x
$$

$$
\begin{aligned}
& \leq \frac{1}{2} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2} d x+C_{1 / 2} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{s}^{*}} d x \\
& \leq \frac{1}{2} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2} d x+C_{1 / 2} \frac{\left\|u_{n}\right\|^{2_{s}^{*}}}{S_{s}^{\frac{3}{3-2 s}}}
\end{aligned}
$$

which implies that

$$
\begin{equation*}
\left\|u_{n}\right\|^{2} \geq \frac{S_{s}^{\frac{3}{2 s}}}{\left(2 C_{1 / 2}\right)^{\frac{3-2 s}{2 s}}} \tag{3.2}
\end{equation*}
$$

By (3.2), we obtain

$$
\begin{aligned}
& {\left[G^{\prime}\left(u_{n}\right), u_{n}\right] } \\
&= {\left[G^{\prime}\left(u_{n}\right), u_{n}\right]-4\left[I^{\prime}\left(u_{n}\right), u_{n}\right] } \\
&= 2\left\|u_{n}\right\|^{2}+4 \int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}}^{t} u_{n}^{2} d x-p \int_{\mathbb{R}^{3}} a(x)\left|u_{n}\right|^{p} d x-q \int_{\mathbb{R}^{3}} \mu b(x)\left|u_{n}\right|^{q} d x \\
&-2_{s}^{*} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{s}^{*}} d x-4\left[\left\|u_{n}\right\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{u_{n}}^{t} u_{n}^{2} d x-\int_{\mathbb{R}^{3}} a(x)\left|u_{n}\right|^{p} d x\right. \\
&\left.-\int_{\mathbb{R}^{3}} \mu b(x)\left|u_{n}\right|^{q} d x-\int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{s}^{*}} d x\right] \\
&=-2\left\|u_{n}\right\|^{2}+(4-p) \int_{\mathbb{R}^{3}} a(x)\left|u_{n}\right|^{p} d x \\
&+(4-q) \int_{\mathbb{R}^{3}} \mu b(x)\left|u_{n}\right|^{q} d x+\left(4-2_{s}^{*}\right) \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{s}^{*}} d x \\
& \leq-2\left\|u_{n}\right\|^{2} \leq-2 \frac{S_{s}^{\frac{3}{2 s}}}{\left(2 C_{1 / 2}\right)^{\frac{3-2 s}{2 s}}} .
\end{aligned}
$$

From (3.1), we have $\lambda_{n} \rightarrow 0$. Thus $I^{\prime}\left(u_{n}\right) \rightarrow 0$. This means that $\left\{u_{n}\right\}$ is a $(P S)_{m}$ sequence for $I$, that is, $I\left(u_{n}\right) \rightarrow m$ and $I^{\prime}\left(u_{n}\right) \rightarrow 0$. By Lemma 2.2 if $m \in\left(0, m_{\infty}\right)$, then $u_{n} \rightarrow u$ in $H$ and thus $I(u)=m$ and $I^{\prime}(u)=0$. Hence, $m$ is attained by $u \in H \backslash\{0\}$. For this purpose, it is sufficient to prove $m<m_{\infty}$.

Similar argument as 2.3, we can obtain the equivalent characterization of the least energy $m$ :

$$
\begin{equation*}
m=\inf _{u \in H \backslash\{0\}} \max _{t \geq 0} I(t u) \tag{3.3}
\end{equation*}
$$

Let $R>0$ and $\gamma \in \mathbb{R}^{3}$ with $|\gamma|=1$. By (3.3), clearly, we have

$$
m \leq \sup _{t \geq 0} I\left(t u_{\infty}(x-R \gamma)\right)
$$

where $u_{\infty}$ is a positive ground state solution for limit problem 2.2). Since

$$
\begin{aligned}
& I\left(t u_{\infty}(x-R \gamma)\right) \\
& \leq \frac{t^{2}}{2}\left\|u_{\infty}(x-R \gamma)\right\|^{2}+C t^{4}\left\|u_{\infty}(x-R \gamma)\right\|^{4}-\frac{t^{2_{s}^{*}}}{2_{s}^{*}}\left\|u_{\infty}(x-R \gamma)\right\|_{2_{s}^{*}}^{2_{s}^{*}} \\
& =\frac{t^{2}}{2}\left\|u_{\infty}\right\|^{2}+C t^{4}\left\|u_{\infty}\right\|^{4}-\frac{t^{2}}{2_{s}^{*}}\left\|u_{\infty}\right\|_{2_{s}^{*}}^{2_{s}^{*}},
\end{aligned}
$$

there exist a small $t^{\prime}>0$ and a large $t^{\prime \prime}>0$ independent of $R$ and $\gamma$ such that

$$
\begin{equation*}
\sup _{t \in\left[0, t^{\prime}\right] \cup\left[t^{\prime \prime},+\infty\right)} I\left(t u_{\infty}(x-R \gamma)\right)<m_{\infty} \tag{3.4}
\end{equation*}
$$

On the other hand, by (A6), for any $u \in H$, we have

$$
\begin{aligned}
I(t u) \leq & I_{\infty}(t u)+\frac{t^{4}}{4} \int_{\mathbb{R}^{3}}(K(x)-1) \phi_{u}^{t} u^{2} d x-\frac{1}{p} t^{p} \int_{\mathbb{R}^{3}}(a(x)-1)|u|^{p} d x \\
& +\frac{t^{4}}{4} \int_{\mathbb{R}^{3}}\left(\phi_{u}^{t}-\hat{\phi}_{u}^{t}\right) u^{2} \mathrm{~d} x \\
= & I_{\infty}(t u)+\frac{t^{4}}{4} \int_{\mathbb{R}^{3}}(K(x)-1) \phi_{u}^{t} u^{2} d x+\frac{t^{4}}{4} \int_{\mathbb{R}^{3}}(K(x)-1) \hat{\phi}_{u}^{t} u^{2} \mathrm{~d} x \\
& -\frac{1}{p} t^{p} \int_{\mathbb{R}^{3}}(a(x)-1)|u|^{p} d x .
\end{aligned}
$$

Thus, choosing $u=u_{\infty}(x-R \gamma)$ in the inequality above and using $\left(K_{2}\right),\left(a_{1}\right)$, we obtain

$$
\begin{aligned}
& I\left(t u_{\infty}(x-R \gamma)\right) \\
& \leq I_{\infty}\left(t u_{\infty}\right)-\frac{t^{4}}{4} C_{0} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{k}} \int_{\mathbb{R}^{3}} \frac{K(y+R \gamma) u_{\infty}^{2}(y)}{|x-y|^{3-2 t}} d y\left|u_{\infty}(x)\right|^{2} d x d y \\
&-\frac{t^{4}}{4} C_{0} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{k}} \hat{\phi}_{u_{\infty}}^{t}(x)\left|u_{\infty}(x)\right|^{2} d x \\
&+\frac{1}{p} t^{p} C_{2} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{a}}\left|u_{\infty}(x)\right|^{p} d x \\
& \leq I_{\infty}\left(t u_{\infty}\right)-\frac{t^{4}}{4} C_{0} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{k}} \int_{\mathbb{R}^{3}} \frac{K(y+R \gamma) u_{\infty}^{2}(y)}{|x-y|^{3-2 t}}\left|u_{\infty}(x)\right|^{2} d x d y \\
&-\frac{t^{4}}{4} C_{0} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{k}} \hat{\phi}_{u_{\infty}}^{t}(x)\left|u_{\infty}(x)\right|^{2} d x \\
&+\frac{1}{p} t^{p} C_{2} C_{\sigma}^{p} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{a}} \frac{1}{(1+|x|)^{(3+2 s) p}} d x \\
& \leq I_{\infty}\left(t u_{\infty}\right)-\frac{t^{4}}{4} C_{0} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{k}} \hat{\phi}_{u_{\infty}}^{t}(x)\left|u_{\infty}(x)\right|^{2} d x \\
&+\frac{1}{p} t^{p} C_{2} C_{\sigma}^{p} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{a}} \frac{1}{(1+|x|)^{(3+2 s) p}} d x .
\end{aligned}
$$

Set $l(t)=I_{\infty}\left(t u_{\infty}\right), t \in(0, \infty)$. It is easy to verify that $\sup _{t \geq 0} l(t)=I_{\infty}\left(u_{\infty}\right)=$ $m_{\infty}$. Moreover, we have

$$
\begin{aligned}
\int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{k}} \hat{\phi}_{u_{\infty}}^{t}(x)\left|u_{\infty}(x)\right|^{2} d x & \geq \int_{|x| \leq 1} \frac{1}{(1+|x+R \gamma|)^{k}} \hat{\phi}_{u_{\infty}}^{t}(x)\left|u_{\infty}(x)\right|^{2} d x \\
& \geq C \int_{|x| \leq 1} \frac{1}{(2+R)^{k}} \hat{\phi}_{u_{\infty}}^{t}(x)\left|u_{\infty}(x)\right|^{2} d x \\
& \geq \frac{\widetilde{C}}{(2+R)^{k}}
\end{aligned}
$$

By Lemma 2.4, we have

$$
\int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{a}} \frac{1}{(1+|x|)^{(3+2 s) p}} d x \leq \frac{C}{(1+R)^{\alpha}}
$$

where $\alpha=\min \{a,(3+2 s) p\}$. Thus

$$
\sup _{t^{\prime} \leq t \leq t^{\prime \prime}} I\left(t u_{\infty}(x-R \gamma)\right) \leq m_{\infty}-\frac{\left(t^{\prime}\right)^{4}}{4} C_{0} \widetilde{C} \frac{1}{(2+R)^{k}}+\frac{1}{p}\left(t^{\prime \prime}\right)^{p} C_{2} C_{\sigma}^{p} C \frac{1}{(1+R)^{\alpha}}
$$

By $0<k<\alpha$, there exists $\hat{R}>0$ large such that for $R>\hat{R}$,

$$
\sup _{t^{\prime} \leq t \leq t^{\prime \prime}} I\left(t u_{\infty}(x-R \gamma)\right)<m_{\infty}, \quad \forall|\gamma|=1
$$

Thus, combing with (3.4), for $R>\hat{R}$, we have

$$
\sup _{t \geq 0} I\left(t u_{\infty}(x-R \gamma)\right)<m_{\infty}, \quad \forall|\gamma|=1
$$

which yields $m<m_{\infty}$. The remaining of the proof of Theorem 1.1 is to show that the solution $u \in H$ is positive.

Proof of Theorem 1.2. The argument is similar to the on in Theorem 1.1, we only need to prove for $R>0$ large, $\sup _{t \geq 0} I\left(t u_{\infty}(x-R \gamma)\right)<m_{\infty}$ uniformly in $\gamma$. Clearly, there exist $0<t^{\prime}<t^{\prime \prime}$ independent of $R$ and $\gamma$ such that

$$
\sup _{t \in\left[0, t^{\prime}\right] \cup\left[t^{\prime \prime},+\infty\right)} I\left(t u_{\infty}(x-R \gamma)\right)<m_{\infty}
$$

On the other hand, by (A3), (A5), (A7), and we have for any $\sigma>0$, there exist $C_{\sigma}>0$ such that

$$
\begin{aligned}
& \sup _{t \in\left[t^{\prime}, t^{\prime \prime}\right]} I\left(t u_{\infty}(x-R \gamma)\right) \\
& \leq \sup _{t \geq 0} I_{\infty}\left(t u_{\infty}\right)+\frac{C_{1}\left(t^{\prime \prime}\right)^{4}}{4} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{d}} \int_{\mathbb{R}^{3}} \frac{K(y+R \gamma) u_{\infty}^{2}(y)}{|x-y|^{3-2 t}} d y\left|u_{\infty}(x)\right|^{2} d x \\
& \quad+\frac{C_{1}\left(t^{\prime \prime}\right)^{4}}{4} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{d}} \hat{\phi}_{u_{\infty}}^{t}(x)\left|u_{\infty}(x)\right|^{2} d x \\
& \quad+\frac{1}{p}\left(t^{\prime \prime}\right)^{p} C_{2} C_{\sigma}^{P} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{a}} \frac{1}{(1+|x|)^{(3+2 s) p}} d x \\
& \quad-\mu \frac{C_{3}\left(t^{\prime}\right)^{q}}{q} \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{b}}\left|u_{\infty}(x)\right|^{q} d x .
\end{aligned}
$$

By calculations, we have

$$
\begin{align*}
& \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{b}}\left|u_{\infty}(x)\right|^{q} d x \\
& \geq \int_{|x| \leq 1} \frac{1}{(1+|x+R \gamma|)^{b}}\left|u_{\infty}(x)\right|^{q} d x  \tag{3.5}\\
& \geq \int_{|x| \leq 1} \frac{1}{|2+R|^{b}}\left|u_{\infty}(x)\right|^{q} d x \geq C \frac{1}{(2+R)^{b}}
\end{align*}
$$

From Lemma 2.4 we obtain

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{a}} \frac{1}{(1+|x|)^{(3+2 s) p}} d x \leq \frac{\alpha}{(1+R)^{\alpha}} \tag{3.6}
\end{equation*}
$$

where $\alpha=\min \{a,(3+2 s) p\}$.

By Hölder's inequality, (A3) and (3.6), we have

$$
\begin{align*}
& \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{d}} \int_{\mathbb{R}^{3}} \frac{K(y+R \gamma) u_{\infty}^{2}(y)}{|x-y|^{3-2 t}} d y\left|u_{\infty}(x)\right|^{2} d x \\
& \leq\left(1+C_{1}\right) \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{d}} \int_{\mathbb{R}^{3}} \frac{u_{\infty}^{2}(y)}{|x-y|^{3-2 t}} d y\left|u_{\infty}(x)\right|^{2} d x \\
& =\left(1+C_{1}\right) \int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{d}} \hat{\phi}_{u_{\infty}}^{t}(x)\left|u_{\infty}(x)\right|^{2} d x  \tag{3.7}\\
& \leq C\left\|\hat{\phi}_{u_{\infty}}^{t}(x)\right\|_{2_{s}^{*}}\left[\int_{\mathbb{R}^{3}}\left(\frac{1}{(1+|x+R \gamma|)^{d}} \frac{1}{(1+|x|)^{6+4 s}}\right)^{\frac{6}{3+2 s}} d x\right]^{\frac{3+2 s}{6}} \\
& \leq C\left[\frac{1}{(1+R)^{m}}\right]^{\frac{3+2 s}{6}} \leq C \frac{1}{(1+R)^{\beta}},
\end{align*}
$$

where $m=\min \left\{\frac{6 d}{3+2 s}, 12\right\}$. Similar as the above argument, we obtain

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} \frac{1}{(1+|x+R \gamma|)^{d}} \hat{\phi}_{u_{\infty}}^{t}(x)\left|u_{\infty}(x)\right|^{2} d x \leq C \frac{1}{(1+R)^{\beta}}, \tag{3.8}
\end{equation*}
$$

where $\beta=\min \{d, 6+4 s\}$. By (3.5)-(3.8), we have

$$
\sup _{t \in\left[t^{\prime}, t^{\prime \prime}\right]} I\left(t u_{\infty}(x-R \gamma)\right) \leq m_{\infty}-\bar{C}_{1} \frac{1}{(2+R)^{b}}+\bar{C}_{2} \frac{1}{(1+R)^{\alpha}}+\bar{C}_{3} \frac{1}{(1+R)^{\beta}},
$$

where $\bar{C}_{1}, \bar{C}_{2}, \bar{C}_{3}$ are positive constants. Since $b<\min \{\alpha, \beta\}$, we obtain that there exists $R_{0}>0$ such that for $R>R_{0}, \sup _{t \geq 0} I\left(t u_{\infty}(x-R \gamma)\right)<m_{\infty}$ uniformly in $\gamma$. The proof is complete.

## 4. Proof of Theorem 1.3

Let $h(t)=I\left(t u_{\infty}(x-R \gamma)\right), t \in(0, \infty), \gamma \in \mathbb{R}^{3}$ with $|\gamma|=1$. Form the proof of Theorem 1.2, we know there exists $R_{0}>0$ such that for $R>R_{0}$, there exists $\varepsilon(R)>0$ satisfying

$$
\sup _{t \geq 0} h(t) \leq m_{\infty}-\varepsilon(R)<m_{\infty} \quad \text { uniformly in } \gamma
$$

For any fixing $R$ and $\gamma$, it is easy to check that $h(t)$ attains its maximum at a unique point $t=t_{\infty}$. Hence, we define a mapping $F_{R}: S^{2}=\left\{\gamma \in \mathbb{R}^{3}:|\gamma|=1\right\} \rightarrow M$ by

$$
F_{R}(\gamma)=t_{\infty} u_{\infty}(x-R \gamma)
$$

Immediately we have the following Lemma.
Lemma 4.1. There exists $R_{0}>0$ such that for $R>R_{0}$, there exists $\varepsilon(R)>0$ satisfying $F_{R}\left(S^{2}\right) \subset\left\{u \in M: I(u) \leq m_{\infty}-\varepsilon(R)\right\}$ uniformly in $\gamma \in S^{2}$.

For $u \in H$, we define a map $\Phi: H \rightarrow H$ by

$$
\Phi(u)(x):=\frac{1}{\left|B_{1}(x)\right|} \int_{B_{1}(x)}|u(y)| d y, \quad \forall x \in \mathbb{R}^{3}
$$

where $\left|B_{1}(x)\right|$ is the Lebesgue measure of $B_{1}(x)$. Let

$$
\hat{u}(x)=\left[\Phi(u)(x)-\frac{1}{2} \max _{x \in \mathbb{R}^{3}} \Phi(u)(x)\right]^{+}
$$

and $\beta: H \backslash\{0\} \rightarrow \mathbb{R}^{3}$ given by

$$
\beta(u)=\frac{1}{\|\hat{u}\|_{1}} \int_{\mathbb{R}^{3}} x \hat{u}(x) d x
$$

Obviously, $\beta(u)$ is well defined for all $u \in H \backslash\{0\}$ and $\beta(u)$ has a compact support in $\mathbb{R}^{3}$. Moreover, $\beta(u)$ is continuous in $H \backslash\{0\}$ and satisfies the following properties.
Lemma 4.2. (i) For any $t \neq 0$ and $u \in H \backslash\{0\}, \beta(t u)=\beta(u)$.
(ii) For any $z \in \mathbb{R}^{3}$ and $u \in H \backslash\{0\}, \beta(u(x-z))=\beta(u)+z$.

Define a functional $J: H \rightarrow \mathbb{R}$ given as follows
$J(u)=\frac{1}{2}\|u\|^{2}+\frac{1}{4} \int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x-\frac{1}{p} \int_{\mathbb{R}^{3}} a(x)|u|^{p} d x-\frac{1}{2_{s}^{*}} \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x, \quad u \in H$.
Lemma 4.3. $m_{0}:=\inf _{M_{0}} J(u)=m_{\infty}$ is not attained, where

$$
M_{0}=\left\{u \in H \backslash\{0\}:\left[J^{\prime}(u), u\right]=0\right\}
$$

Proof. First, we show that for any $u \in M_{0}$, there exists a unique $0<\tau \leq 1$ such that $\tau u \in M_{\infty}$. Indeed, by $u \in M_{0}$ and $\tau u \in M_{\infty}$, we have

$$
\begin{equation*}
\|u\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x=\int_{\mathbb{R}^{3}} a(x)|u|^{p} d x+\int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x \tag{4.1}
\end{equation*}
$$

and then

$$
\begin{align*}
\tau^{p} \int_{\mathbb{R}^{3}} a(x)|u|^{p} d x+\tau^{2_{s}^{*}} \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x & \leq \tau^{p} \int_{\mathbb{R}^{3}}|u|^{p} d x+\tau^{2_{s}^{*}} \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x  \tag{4.2}\\
& =\tau^{2}\|u\|^{2}+\tau^{4} \int_{\mathbb{R}^{3}} \hat{\phi}_{u}^{t} u^{2} d x
\end{align*}
$$

From (A3) and $K(x) \geq 1$ for any $x \in \mathbb{R}^{3}$, it follows that

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} \hat{\phi}_{u}^{t} u^{2} d x \leq \int_{\mathbb{R}^{3}} K(x) \hat{\phi}_{u}^{t} u^{2} d x \leq \int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x \tag{4.3}
\end{equation*}
$$

If $\tau>1$, by (4.1), 4.2 and (4.3), we deduce that

$$
\begin{aligned}
\tau^{4}\left(\|u\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x\right) & \geq \tau^{4}\left(\|u\|^{2}+\int_{\mathbb{R}^{3}} \hat{\phi}_{u}^{t} u^{2} d x\right) \\
& \geq \tau^{p}\left(\int_{\mathbb{R}^{3}} a(x)|u|^{p} d x+\int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x\right) \\
& =\tau^{p}\left(\|u\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x\right)
\end{aligned}
$$

which yields $\tau \leq 1$, this achieves a contradiction. Hence $\tau \leq 1$ and the claim is true.

For $u \in M_{0}$, using (4.3), we have

$$
\begin{aligned}
J(u) & =J(u)-\frac{1}{p}\left[J^{\prime}(u), u\right] \\
& =\left(\frac{1}{2}-\frac{1}{p}\right)\|u\|^{2}+\left(\frac{1}{4}-\frac{1}{p}\right) \int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x+\left(\frac{1}{p}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x \\
& \geq\left(\frac{1}{2}-\frac{1}{p}\right)\|\tau u\|^{2}+\left(\frac{1}{4}-\frac{1}{p}\right) \tau^{4} \int_{\mathbb{R}^{3}} \hat{\phi}_{u}^{t} u^{2} d x+\left(\frac{1}{p}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}}|\tau u|^{2_{s}^{*}} d x \\
& =I_{\infty}(\tau u)-\frac{1}{p}\left[I_{\infty}^{\prime}(\tau u), \tau u\right]
\end{aligned}
$$

$$
=I_{\infty}(\tau u) \geq m_{\infty}
$$

which implies that $m_{0} \geq m_{\infty}$.
Next we prove $m_{0} \leq m_{\infty}$. Let $w_{n}=u_{\infty}\left(.-z_{n}\right)$, where $z_{n} \in \mathbb{R}^{3}$ with $\left|z_{n}\right| \rightarrow \infty$. We claim that for $w_{n} \in M_{\infty}$, there exists $\tau_{n} \geq 1$ such that $\tau_{n} w_{n} \in M_{0}$. In fact, from $w_{n} \in M_{\infty}$ and $\tau_{n} w_{n} \in M_{0}$, it holds

$$
\left\|w_{n}\right\|^{2}+\int_{\mathbb{R}^{3}} \hat{\phi}_{w_{n}}^{t} w_{n}^{2} d x=\int_{\mathbb{R}^{3}}\left|w_{n}\right|^{p} d x+\int_{\mathbb{R}^{3}}\left|w_{n}\right|^{2_{s}^{*}} d x
$$

and then

$$
\begin{aligned}
& \tau_{n}^{p} \int_{\mathbb{R}^{3}}\left|w_{n}\right|^{p} d x+\tau_{n}^{2_{s}^{*}} \int_{\mathbb{R}^{3}}\left|w_{n}\right|^{2_{s}^{*}} d x \\
& \geq \tau_{n}^{p} \int_{\mathbb{R}^{3}} a(x)\left|w_{n}\right|^{p} d x+\tau_{n}^{2_{s}^{*}} \int_{\mathbb{R}^{3}}\left|w_{n}\right|^{2_{s}^{*}} d x \\
& =\tau_{n}^{2}\left\|w_{n}\right\|^{2}+\tau_{n}^{4} \int_{\mathbb{R}^{3}} K(x) \phi_{w_{n}}^{t} w_{n}^{2} d x
\end{aligned}
$$

If $\tau_{n}<1$, then

$$
\begin{aligned}
\tau_{n}^{p}\left(\int_{\mathbb{R}^{3}}\left|w_{n}\right|^{p} d x+\int_{\mathbb{R}^{3}}\left|w_{n}\right|^{2_{s}^{*}} d x\right) & \geq \tau_{n}^{4}\left(\left\|w_{n}\right\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{w_{n}}^{t} w_{n}^{2} d x\right) \\
& \geq \tau_{n}^{4}\left(\left\|w_{n}\right\|^{2}+\int_{\mathbb{R}^{3}} \hat{\phi}_{w_{n}}^{t} w_{n}^{2} d x\right) \\
& =\tau_{n}^{4}\left(\int_{\mathbb{R}^{3}}\left|w_{n}\right|^{p} d x+\int_{\mathbb{R}^{3}}\left|w_{n}\right|^{2_{s}^{*}} d x\right)
\end{aligned}
$$

which leads to a contradiction with $\tau_{n}<1$. Hence $\tau_{n} \geq 1$ and the claim holds.
By the definition of $m_{0}$ and $\tau_{n} u_{n} \in M_{0}$, we have

$$
\begin{aligned}
m_{0} \leq J\left(\tau_{n} w_{n}\right)= & \frac{1}{2} \tau_{n}^{2}\left\|u_{\infty}\right\|^{2}+\frac{1}{4} \tau_{n}^{4} \int_{\mathbb{R}^{3}} K(x) \phi_{w_{n}}^{t} w_{n}^{2} d x \\
& -\frac{1}{p} \tau_{n}^{p} \int_{\mathbb{R}^{3}} a(x)\left|w_{n}\right|^{p} d x-\frac{1}{2_{s}^{*}} \tau_{n}^{2_{s}^{*}} \int_{\mathbb{R}^{3}}\left|u_{\infty}(x)\right|^{2_{s}^{*}} d x
\end{aligned}
$$

By Lebesgue dominated convergence Theorem, we deduce that

$$
\begin{align*}
& \lim _{n \rightarrow \infty} \int_{\mathbb{R}^{3}} K(x) \phi_{w_{n}}^{t} w_{n}^{2} d x \\
& =\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{3}} K\left(x+z_{n}\right) \int_{\mathbb{R}^{3}} \frac{K\left(y+z_{n}\right) u_{\infty}^{2}(y)}{|x-y|^{3-2 t}} d y\left|u_{\infty}(x)\right|^{2} d x d y \\
& =\int_{\mathbb{R}^{3}} \hat{\phi}_{u_{\infty}}^{t}(x)\left|u_{\infty}(x)\right|^{2} d x  \tag{4.4}\\
& \lim _{n \rightarrow \infty} \int_{\mathbb{R}^{3}} a(x)\left|w_{n}\right|^{p} d x \\
& =\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{3}} a\left(x+z_{n}\right)\left|u_{\infty}(x)\right|^{p} d x \\
& \\
& =\int_{\mathbb{R}^{3}}\left|u_{\infty}(x)\right|^{p} d x
\end{align*}
$$

If $\tau_{n} \rightarrow 1$, we obtain $m_{0} \leq \lim _{n \rightarrow \infty} J\left(t_{n} w_{n}\right)=I_{\infty}\left(u_{\infty}\right)=m_{\infty}$, form which we see that $m_{0}=m_{\infty}$. Thus we only need to prove $\tau_{n} \rightarrow 1$. By $\tau_{n} w_{n} \in M_{0}$, with $\tau_{n} \geq 1$, we have

$$
\tau_{n}^{4}\left(\left\|w_{n}\right\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{w_{n}}^{t} w_{n}^{2} d x\right)
$$

$$
\begin{aligned}
& \geq \tau_{n}^{2}\left\|w_{n}\right\|^{2}+\tau_{n}^{4} \int_{\mathbb{R}^{3}} K(x) \phi_{w_{n}}^{t} w_{n}^{2} d x \\
& =\tau_{n}^{p} \int_{\mathbb{R}^{3}} a(x)\left|w_{n}\right|^{p} d x+\tau_{n}^{2_{s}^{*}} \int_{\mathbb{R}^{3}}\left|w_{n}\right|^{2_{s}^{*}} d x \\
& \geq \tau_{n}^{p}\left(\int_{\mathbb{R}^{3}} a(x)\left|w_{n}\right|^{p} d x+\int_{\mathbb{R}^{3}}\left|w_{n}\right|^{2_{s}^{*}} d x\right) .
\end{aligned}
$$

Thus, by (4.4), we deduce that

$$
\begin{aligned}
1 \leq \tau_{n}^{p-4} & \leq \frac{\left\|w_{n}\right\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{w_{n}}^{t} w_{n}^{2} d x}{\int_{\mathbb{R}^{3}} a(x)\left|w_{n}\right|^{p} d x+\int_{\mathbb{R}^{3}}\left|w_{n}\right|^{2_{s}^{*}} d x} \\
& =\frac{\left\|u_{\infty}\right\|^{2}+\int_{\mathbb{R}^{3}} \hat{\phi}_{u_{\infty}}^{t}(x)\left|u_{\infty}(x)\right|^{2} d x+o(1)}{\int_{\mathbb{R}^{3}}\left|u_{\infty}(x)\right|^{p} d x+o(1)+\int_{\mathbb{R}^{3}}\left|u_{\infty}(x)\right|^{2_{s}^{*}} d x}
\end{aligned}
$$

which yields $\tau_{n} \rightarrow 1$ by using $u_{\infty} \in M_{\infty}$.
Next we prove $m_{0}$ is not attained. Assume by contradiction that there exists $u_{0} \in M_{0}$ such that $m_{0}=J\left(u_{0}\right)$. We claim $J^{\prime}\left(u_{0}\right)=0$. Set $\widetilde{G}(u)=\left[J^{\prime}(u), u\right]$, By the Lagrange multipliers Theorem, we obtain $\lambda \in \mathbb{R}$ such that $J^{\prime}\left(u_{0}\right)-\lambda \widetilde{G}^{\prime}\left(u_{0}\right) \rightarrow 0$, similar to the of proof of Theorem 1.1. we have $J^{\prime}\left(u_{0}\right)=0$. Note that if $u_{0}$ is singchanging, by Remark 5.6 in [34], we see that $J\left(u_{0}\right) \geq 2 m_{0}$, a contradiction. Thus we may assume that $u_{0} \geq 0$ in $H$ and $u_{0} \not \equiv 0$, we claim $u_{0}>0$, by the definition of $\phi_{u_{0}}^{t}(x)$, there exists $C>0$ such that

$$
\begin{aligned}
\phi_{u_{0}}^{t}(x) & =\int_{|x-y| \geq 1} \frac{K(y) u_{0}^{2}(y)}{|x-y|^{3-2 t}} d y+\int_{|x-y|<1} \frac{K(y) u_{0}^{2}(y)}{|x-y|^{3-2 t}} d y \\
& \leq C\left\|u_{0}\right\|_{2}^{2}+C \int_{|x-y|<1} \frac{1}{|x-y|^{3-2 t}} d y<+\infty
\end{aligned}
$$

and $|g| \leq C\left(\left|u_{0}\right|+\left|u_{0}\right|^{q-1}\right)$, where $g(x)=a(x)\left|u_{0}(x)\right|^{p-2} u_{0}(x)+\left|u_{0}(x)\right|^{2_{s}^{*}-2} u_{0}(x)-$ $u_{0}(x)-K(x) \phi_{u_{0}}^{t}(x) u_{0}(x)$. Then it follows from [22, Proposition 3.4] that there exists $\sigma \in(0,1)$ such that $u_{0} \in C^{0, \sigma}$. Let $\omega$ satisfy $-\Delta \omega=-u_{0}-K(x) \phi_{u_{0}}^{t} u+$ $a(x)\left|u_{0}\right|^{p-2} u_{0}+\left|u_{0}\right|^{2_{s}^{*}-2} u_{0} \in C^{0, \sigma}$. By the Hölder regularity theory for the Laplacian, we have $\omega \in C^{2, \sigma}$. It follows from $2 s+\sigma>1$ that $(-\Delta)^{1-s} \omega \in C^{1,2 s+\sigma-1}$. Then, since $(-\Delta)^{s}\left(u_{0}-(-\Delta)^{1-s} \omega\right)=0$, the function $u-(-\Delta)^{1-s} \omega$ is harmonic and we obtain $u_{0}$ has the same regularity as $(-\Delta)^{1-s} \omega$. That is, $u_{0} \in C^{1,2 s+\sigma-1}$. The regularity obtained above implies that

$$
(-\Delta)^{s} u_{0}=-\int_{\mathbb{R}^{3}} \frac{u_{0}(x+y)+u_{0}(x-y)-2 u_{0}(x)}{|y|^{3+2 s}} d y
$$

Assume that there exists $x_{0} \in \mathbb{R}^{3}$ such that $u_{0}\left(x_{0}\right)=0$, then by $u_{0} \not \equiv 0$ and $u_{0} \geq 0$,

$$
(-\Delta)^{s} u\left(x_{0}\right)=-\int_{\mathbb{R}^{3}} \frac{u\left(x_{0}+y\right)+u\left(x_{0}-y\right)}{|y|^{3+2 s}} d y<0
$$

However, noting that $-\Delta u_{0}=-u_{0}-K(x) \phi_{u_{0}}^{t} u_{0}+a(x)\left|u_{0}\right|^{p-2} u_{0}+\left|u_{0}\right|^{2_{s}^{*}-2} u_{0}$ we obtain $-\Delta u_{0}\left(x_{0}\right)=0$, which is a contradiction. Therefore, $u_{0}>0$.

From the above proof, we see that for $u_{0} \in M_{0}$, there exists a unique $\tau_{0} \leq 1$ such that $\tau_{0} u_{0} \in M_{\infty}$. Thus,

$$
\begin{aligned}
m_{\infty} & \leq I_{\infty}\left(\tau_{0} u_{0}\right) \\
& =I_{\infty}\left(\tau_{0} u_{0}\right)-\frac{1}{p}\left[I_{\infty}^{\prime}\left(\tau_{0} u_{0}\right), \tau_{0} u_{0}\right]
\end{aligned}
$$

$$
\begin{aligned}
& =\left(\frac{1}{2}-\frac{1}{p}\right)\left\|\tau_{0} u_{0}\right\|^{2}+\left(\frac{1}{4}-\frac{1}{p}\right) \tau_{0}^{4} \int_{\mathbb{R}^{3}} \hat{\phi}_{u_{0}}^{t} u_{0}^{2} d x+\left(\frac{1}{p}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}}\left|\tau_{0} u_{0}\right|^{2_{s}^{*}} d x \\
& \leq\left(\frac{1}{2}-\frac{1}{p}\right)\left\|u_{0}\right\|^{2}+\left(\frac{1}{4}-\frac{1}{p}\right) \int_{\mathbb{R}^{3}} K(x) \phi_{u_{0}}^{t} u_{0}^{2} d x+\left(\frac{1}{p}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}}\left|u_{0}\right|^{2_{s}^{*}} d x \\
& =J\left(u_{0}\right)-\frac{1}{p}\left[J^{\prime}\left(u_{0}\right), u_{0}\right]=J\left(u_{0}\right)=m_{0} .
\end{aligned}
$$

From $m_{0}=m_{\infty}$, it follows that

$$
\begin{aligned}
& \left(\frac{1}{2}-\frac{1}{p}\right)\left\|\tau_{0} u_{0}\right\|^{2}+\left(\frac{1}{4}-\frac{1}{p}\right) \tau_{0}^{4} \int_{\mathbb{R}^{3}} \hat{\phi}_{u_{0}}^{t} u_{0}^{2} d x+\left(\frac{1}{p}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}}\left|\tau_{0} u_{0}\right|^{2_{s}^{*}} d x \\
& =\left(\frac{1}{2}-\frac{1}{p}\right)\left\|u_{0}\right\|^{2}+\left(\frac{1}{4}-\frac{1}{p}\right) \int_{\mathbb{R}^{3}} K(x) \phi_{u_{0}}^{t} u_{0}^{2} d x+\left(\frac{1}{p}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}}\left|u_{0}\right|^{2_{s}^{*}} d x
\end{aligned}
$$

that is

$$
\begin{aligned}
& \tau_{0}^{2}\left\|u_{0}\right\|^{2}+\tau_{0}^{4} \int_{\mathbb{R}^{3}} \hat{\phi}_{u_{0}}^{t} u_{0}^{2} d x+\tau_{0}^{2_{s}^{*}} \int_{\mathbb{R}^{3}}\left|u_{0}\right|^{2_{s}^{*}} d x \\
& =\left\|u_{0}\right\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{u_{0}}^{t} u_{0}^{2} d x+\int_{\mathbb{R}^{3}}\left|u_{0}\right|^{2_{s}^{*}} d x
\end{aligned}
$$

Thus

$$
\begin{aligned}
& \left(1-\tau_{0}^{2}\right)\left\|u_{0}\right\|^{2}+\int_{\mathbb{R}^{3}}(K(x)-1) \phi_{u_{0}}^{t} u_{0}^{2} d x+\int_{\mathbb{R}^{3}}\left(\phi_{u_{0}}^{t}-\hat{\phi}_{u_{0}}^{t}\right) u_{0}^{2} d x \\
& +\left(1-\tau_{0}^{4}\right) \int_{\mathbb{R}^{3}} \hat{\phi}_{u_{0}}^{t} u_{0}^{2} d x+\left(1-\tau_{0}^{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}}\left|u_{0}\right|^{2_{s}^{*}} d x \\
& =\left(1-\tau_{0}^{2}\right)\left\|u_{0}\right\|^{2}+\int_{\mathbb{R}^{3}}(K(x)-1)\left(\phi_{u_{0}}^{t}+\hat{\phi}_{u_{0}}^{t}\right) u_{0}^{2} d x \\
& \quad+\left(1-\tau_{0}^{4}\right) \int_{\mathbb{R}^{3}} \hat{\phi}_{u_{0}}^{t} u_{0}^{2} d x+\left(1-\tau_{0}^{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}}\left|u_{0}\right|^{2_{s}^{*}} d x=0
\end{aligned}
$$

by $\tau_{0} \leq 1$, so

$$
\int_{\mathbb{R}^{3}}(K(x)-1)\left(\phi_{u_{0}}^{t}+\hat{\phi}_{u_{0}}^{t}\right) u_{0}^{2} d x=0
$$

this contradicts $u_{0}$ being positive, $K(x) \geq 1$ and meas $\left\{x \in \mathbb{R}^{3}: K(x)>1\right\}>0$.
Lemma 4.4. There exists $\rho_{0}>0$ such that for $u \in M_{0}$ satisfying $J(u) \leq m_{\infty}+\rho_{0}$, it holds $|\beta(u)|>0$.

Proof. Assume by the contrary that there exists $\left\{u_{n}\right\} \subset M_{0}$ such that $J\left(u_{n}\right) \rightarrow$ $m_{\infty}=m_{0}$ and $|\beta(u)|=0$. Similar to the proof Theorem 1.1, we can derive by the Lagrange multipliers Theorem that $J^{\prime}\left(u_{n}\right) \rightarrow 0$. We omit the proof here. Similar to the proof Lemma 2.3 , we obtain $u_{n} \rightharpoonup u$ weakly in $H, J^{\prime}(u)=0$, and

$$
\begin{equation*}
m_{\infty}-J(u)=I_{\infty}\left(v_{n}\right)+o(1) \text { and } I_{\infty}^{\prime}\left(v_{n}\right)=o(1), \tag{4.5}
\end{equation*}
$$

where $v_{n}=u_{n}-u$.
For the sequence $\left\{v_{n}\right\}$, two cases may occur: $\left\|v_{n}\right\|_{r} \rightarrow 0$ for any $r \in\left(2,2_{s}^{*}\right)$, or there $y_{n} \in \mathbb{R}^{3}$ with $\left|y_{n}\right| \rightarrow \infty$ such that $v_{n}\left(.+y_{n}\right) \rightharpoonup v \neq 0$ weakly in $H$. By virtue of $J^{\prime}(u)=0$, we can deduce that $J(u) \geq 0$. From Lemma 2.3. we see that $m_{\infty}<\frac{s}{3} S_{s}^{\frac{3}{2 s}}$. Thus $m_{\infty}-J(u)<\frac{s}{3} S_{s}^{\frac{3}{2 s}}$.

If $\left\|v_{n}\right\|_{r} \rightarrow 0$ for any $r \in\left(2,2_{s}^{*}\right)$, by 4.6), we have $m_{\infty}-J(u)=\frac{1}{2}\left\|v_{n}\right\|^{2}-$ $\frac{1}{2_{s}^{*}}\left\|v_{n}\right\|_{2_{s}^{*}}^{2_{s}^{*}}+o(1)$ and $\left\|v_{n}\right\|^{2}-\left\|v_{n}\right\|_{2_{s}^{*}}^{2_{2}^{*}}=o(1)$. Up to a subsequence, we may assume
that $\left\|v_{n}\right\|^{2} \rightarrow l$ and then $\left\|v_{n}\right\|_{2_{s}^{*}}^{2_{s}^{*}} \rightarrow l$. If $l>0$, by the definition of $S_{s}$, we obtain $l \geq S_{s}^{\frac{3}{2 s}}$. So $m_{\infty}-J(u)=\frac{1}{2}\left\|v_{n}\right\|^{2}-\frac{1}{2_{s}^{*}}\left\|v_{n}\right\|_{2_{s}^{*}}^{2_{2}^{*}}=\frac{s}{3} l \geq \frac{s}{3} S_{s}^{\frac{3}{2 s}}$, a contradiction with $m_{\infty}-J(u)<\frac{s}{3} S_{s}^{\frac{3}{2 s}}$. Thus, $l=0$ and then $u_{n} \rightarrow u$ in $H$, we obtain $m_{0}=J(u)$, a contradiction with $m_{0}$ is not attained. Therefore, $v_{n}\left(.+y_{n}\right) \rightharpoonup v \neq 0$ weakly in $H$. Similar to the proof Lemma 2.3 , we can deduce that

$$
\begin{gathered}
m_{\infty}-J(u)=I_{\infty}\left(v_{n}\left(.+y_{n}\right)\right)+o(1) \\
I_{\infty}^{\prime}\left(v_{n}\left(.+y_{n}\right)\right)=o(1)
\end{gathered}
$$

Hence, $I_{\infty}^{\prime}(v)=0$ and by using Fatou's Lemma, we have

$$
\begin{aligned}
m_{\infty}-J(u)= & I_{\infty}\left(v_{n}\left(.+y_{n}\right)\right)-\frac{1}{4}\left[I_{\infty}^{\prime}\left(v_{n}\left(.+y_{n}\right)\right), v_{n}\left(.+y_{n}\right)\right]+o(1) \\
= & \frac{1}{4} \|\left(\left.v_{n}\left(.+y_{n}\right) \|^{2}+\left(\frac{1}{4}-\frac{1}{p}\right) \int_{\mathbb{R}^{3}} \right\rvert\,\left(\left.v_{n}\left(.+y_{n}\right)\right|^{p} d x\right.\right. \\
& \left.+\left(\frac{1}{4}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}} \right\rvert\,\left(\left.v_{n}\left(.+y_{n}\right)\right|^{2_{s}^{*}} d x+o(1)\right. \\
\geq & \frac{1}{4}\|v\|^{2}+\left(\frac{1}{4}-\frac{1}{p}\right) \int_{\mathbb{R}^{3}}|v|^{p} d x+\left(\frac{1}{4}-\frac{1}{2_{s}^{*}}\right) \int_{\mathbb{R}^{3}}|v|^{2_{s}^{*}} d x \\
= & \left.I_{\infty}(v)-\frac{1}{4}\left(I_{\infty}^{\prime}(v), v\right)\right)=I_{\infty}(v) \geq m_{\infty}
\end{aligned}
$$

Combining with $J(u) \geq 0$, we obtain $J(u)=0$ and then $v_{n}\left(.+y_{n}\right)=u_{n}\left(.+y_{n}\right) \rightarrow v$ in $H$. By Lemma 4.2, we have

$$
\beta(v(x))+o(1)=\beta\left(u_{n}\left(x+y_{n}\right)\right)=\beta\left(u_{n}\right)-y_{n}=-y_{n}
$$

Which yields $|\beta(v(x))|=\infty$, this leads to a contradiction.
Lemma 4.5. There exists $\mu_{0}>0$ small such that for $\mu \in\left(0, \mu_{0}\right)$, we have $|\beta(u)|>0$ for $u \in\left\{u \in M: I(u)<m_{\infty}\right\}$.
Proof. Let $u \in M$ be such that $I(u)<m_{\infty}$, then we have

$$
\begin{equation*}
m_{\infty}>I(u)=I(u)-\frac{1}{4}\left[I^{\prime}(u), u\right] \geq \frac{1}{4}\|u\|^{2} \tag{4.6}
\end{equation*}
$$

Using the conditions (A4), (A6) and $u \in M$, we obtain for any $\varepsilon>0$, there exists $C_{\varepsilon}>0$ such that

$$
\begin{align*}
\|u\|^{2} & \leq \int_{\mathbb{R}^{3}} a(x)|u|^{p} d x+\mu \int_{\mathbb{R}^{3}} b(x)|u|^{q} d x+\int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x \\
& \leq(1+\mu)\left[\varepsilon \int_{\mathbb{R}^{3}}|u|^{2} d x+C_{\varepsilon} \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x\right] . \tag{4.7}
\end{align*}
$$

Choose $\varepsilon \in(0,1 / 4)$, we have $\frac{1}{2}\|u\|^{2} \leq(1+\mu) C_{\varepsilon} \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x$ for $\mu \in(0,1)$. In fact, if $\varepsilon \in\left(0, \frac{1}{4}\right)$ and $\mu \in(0,1)$, we obtain $0<(1+\mu) \varepsilon<1 / 2$, and then $\frac{1}{2}-(1+\mu) \varepsilon>0$. Thus, it holds

$$
\frac{1}{2} \int_{\mathbb{R}^{3}}|\nabla u|^{2} d x+\left(\frac{1}{2}-(1+\mu) \varepsilon\right) \int_{\mathbb{R}^{3}}|u|^{2} d x \geq 0
$$

that is

$$
\frac{1}{2}\|u\|^{2} \leq\|u\|^{2}-(1+\mu) \varepsilon \int_{\mathbb{R}^{3}}|u|^{2} d x
$$

by (4.7), we have

$$
\|u\|^{2}-(1+\mu) \varepsilon \int_{\mathbb{R}^{3}}|u|^{2} d x \leq(1+\mu) C_{\varepsilon} \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x
$$

so

$$
\frac{1}{2}\|u\|^{2} \leq(1+\mu) C_{\varepsilon} \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x
$$

Thus, by the definition of $S_{s}$, there exists $L_{0}>0$ independent of $\mu \in(0,1)$ such that

$$
\begin{equation*}
\int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x \geq \frac{L_{0}}{(1+\mu)^{\frac{3}{2 s}}} \tag{4.8}
\end{equation*}
$$

Similar to the argument of Lemma 4.3, we can deduce that for any $u \in M$, there exists a unique $\tau(u) \geq 1$ such that $\tau(u) u \in M_{0}$. Then

$$
\begin{aligned}
& \tau^{4}(u)\left(\|u\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x\right) \\
& \geq \tau^{2}(u)\|u\|^{2}+\tau^{4}(u) \int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x \\
& =\tau^{p}(u) \int_{\mathbb{R}^{3}} a(x)|u|^{p} d x+\tau^{2_{s}^{*}}(u) \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x \\
& \geq \tau^{2_{s}^{*}}(u) \int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x
\end{aligned}
$$

which implies that

$$
\tau^{2_{s}^{*}-4}(u) \leq \frac{\|u\|^{2}+\int_{\mathbb{R}^{3}} K(x) \phi_{u}^{t} u^{2} d x}{\int_{\mathbb{R}^{3}}|u|^{2_{s}^{*}} d x}
$$

Together with 4.6) and 4.8, we derive there exists $\bar{C}>0$ independent of $\mu \in(0,1)$ such that

$$
\begin{equation*}
1 \leq \tau^{2_{s}^{*}-4}(u) \leq \bar{C}(1+\mu)^{\frac{3}{2 s}} \tag{4.9}
\end{equation*}
$$

Note that for $u \in M$ with $I(u)<m_{\infty}$, thus

$$
m_{\infty}>I(u)=\sup _{t \geq 0} I(t u) \geq I(t(u) u)=J(t(u) u)-\mu \frac{t^{q}(u)}{q} \int_{\mathbb{R}^{3}} b(x)|u|^{q} d x
$$

By (4.6) and 4.9), there exists a small $\mu_{0} \in(0,1)$ such that $\mu \in\left(0, \mu_{0}\right)$,

$$
J(t(u) u)<m_{\infty}+\mu \frac{t^{q}(u)}{q} \int_{\mathbb{R}^{3}} b(x)|u|^{q} d x \leq m_{\infty}+\rho_{0}
$$

Form Lemma 4.4, we have $|\beta(t(u) u)|>0$. Hence, Lemma 4.2 implies $|\beta(u)|>0$.
Lemma 4.6. For $\mu \in\left(0, \mu_{0}\right)$, define $G:\left\{u \in M: I(u)<m_{\infty}\right\} \rightarrow S^{2}$ by $G(u)=$ $\frac{\beta(u)}{|\beta(u)|}$. Then for $R>R_{0}$ and $\mu \in\left(0, \mu_{0}\right)$, the map

$$
G \circ F_{R}: S^{2} \rightarrow S^{2} ; y \rightarrow G \circ\left(F_{R}(y)\right)
$$

is homotopic to the identity.
Proof. Similar to the argument of [1, 37, Proposition 2.9], we define the map $\zeta(\theta, y)$ : $[0,1] \times S^{2} \rightarrow S^{2}$ by

$$
\zeta(\theta, y)= \begin{cases}G\left((1-2 \theta) F_{R}(y)+2 \theta u_{\infty}(x-R y)\right), & \theta \in[0,1 / 2) \\ G\left(u_{\infty}\left(x-\frac{R}{2(1-\theta)} y\right),\right. & \theta \in[1 / 2,1) \\ y, & \theta=1\end{cases}
$$

By the definition of $G$ and Lemma 2.7, tt is not difficult to check that $\zeta(\theta, y) \in$ $C\left([0,1] \times S^{2}, S^{2}\right), \zeta(0, y)=G \circ\left(F_{R}(y)\right)$ for $y \in S^{2}$ and $\zeta(1, y)=y$ for $y \in S^{2}$. The proof is complete.

Proof of Theorem 1.3. Form Lemma 2.7. Lemma 4.1 and Lemma 4.6, we have that for $R>R_{0}$ and $\mu \in\left(0, \mu_{0}\right)$, it holds

$$
\operatorname{cat}\left(\left\{u \in M: I(u) \leq m_{\infty}-\varepsilon(R)\right\}\right) \geq 2
$$

Then by Lemma 2.3 and Lemma 2.6 , we see that $I$ admits at least two nontrivial critical point in $\left\{u \in M: I(u)<m_{\infty}\right\}$.

Acknowledgments. This work is supported by NSFC grant 11501403 and fund program for the Scientific Activities of Selected Returned Overseas Professionals in Shanxi Province (2018), and by the Natural Science Foundation of Shanxi Province (No. 201901D111085).

## References

[1] S. Adachi, K. Tanaka; Four positive solutions for the semilinear elliptic equation: $-\Delta u+u=$ $a(x) u^{p}+f(x)$ in $\mathbb{R}^{N}$, Calc. Var. Partial Differential Equations, 11 (2000), 63-95.
[2] A. Azzollini, P. d'Avenia, A. Pomponio; On the Schrödinger-Maxwell equations under the effect of a general nonlinear term, Ann. Inst. H. Poincaré Anal. Non Linéaire, 27 (2010), 779-791.
[3] A. Azzollini, A. Pomponio; Ground state solutions for the nonlinear Schrödinger-Maxwell equations, J. Math. Anal. Appl., 345 (2008), 90-108.
[4] V. Benci, D. Fortunato; An eigenvalue problem for the Schrödinger-Maxwell equations, Nonlinear Anal., 11 (1998), 283-293.
[5] G. Cerami, G. Vaira; Positive solutions for some non-autonomous Schrödinger-Poisson systems, J. Differential Equations., 248 (2010), 521-543.
[6] G. Cerami, D. Passaseo; The effect of concentrating potentials in some singularly perturbed problems, Calc. Var. Partial Differential Equations., 17 (2003), 257-281.
[7] L. Caffarelli, L. Silvestre; An extension problem related to the fractional Laplacian, Comm. Partial Differential Equations, 32 (2007), 1245-1260.
[8] G. M. Coclite; A multiplicity result for the nonlinear Schrödinger-Maxwell equations, Commun. Appl. Anal., 7 (2003), 417-423.
[9] S. Y. A. Chang, M. del Mar Gonzjäalez; Fractional Laplacian in conformal geometry, Adv. Math., 226 (2011), 1410-1432.
[10] T. D'Aprile, D. Mugnai; Solitary waves for nonlinear Klein-Gordon-Maxwell and SchrodingerMaxwell equations, Proc. Roy. Soc. Edinburgh Sect. A., 134 (2004), 893-906.
[11] T. D'Aprile, D. Mugnai; Non-existence results for the coupled Klein-Gordon-Maxwell equations, Adv. Nonlinear Stud., 4 (2004), 307-322.
[12] J. Davila, M. del Pino, S. Dipierro, E. Valdinoci; Concentration phenomena for the nonlocal Schrödinger equation with Dirichlet datum, Anal. PDE., 8 (2015), 1165-1235.
[13] S. Dipierro, M. Medina, E. Valdinoci; Fractional elliptic problems with critical growth in the whole of $\mathbb{R}^{N}$, Lecture Notes. Scuola Normale Superiore di Pisa (New Series), 15. Edizioni della Normale, Pisa., (2017).
[14] M. M. Fall, F. Mahmoudi, E. Valdinoci; Ground states and concentration phenomena for the fractional Schrödinger equation, Nonlinearity, 28 (2015), 1937-1961.
[15] A. R. Giammetta; Fractional Schrödinger-Poisson-Slater system in one dimension, e-print., arXiv:1405.2796v1.
[16] L. Huang, E. M. Rocha, J. Chen; Two positive solutions of a class of Schrödinger-Poisson system with indefinite nonlinearity, J. Differential Equations, 255 (2013), 2463-2483.
[17] Y. Jiang, H. Zhou; Schrödinger-Poisson system with steep potential well, J. Differential Equations, 251 (2011), 582-608.
[18] T. Jin, Y. Li, J. Xiong; On a fractional Nirenberg problem, part I: blow up analysis and compactness of solutions, J. Eur. Math. Soc., 16 (2014), 1111-1171.
[19] H. Kikuchi; On the existence of a solution for elliptic system related to the MaxwellSchrödinger equations, Nonlinear Anal., 67 (2007), 1445-1456.
[20] H. Kikuchi; Fractional quantum mechanics and Lévy path integrals, Physics Letters A., 268 (2000), 298-305.
[21] N. Laskin; Fractional quantum mechanics and Lévy path integrals, Physical Review., 66 (2002), 56108-0.
[22] Z. S. Liu, J. J. Zhang; Multiplicity and concentration of solutions for the following critical fractional Schrödinger-Poisson systems with critical growth, ESAIM: Control, Optim. Calc. Var., 23 (2017), 1515-1542.
[23] R. Metzler, J. Klafter; The random walks guide to anomalous diffusion: a fractional dynamics approach, Phys. Rep., 339 (2000), 1-77.
[24] D. Ruiz; The Schrödinger-Poisson equation under the effect of a nonlinear local term, $J$. Funct. Anal., 237 (2006), 655-674.
[25] S. Secchi; Ground state solutions for nonlinear fractional Schrödinger equations in $\mathbb{R}^{n}, J$. Math. Phys., 54 (2013), 031501.
[26] L. Silvestre; Regularity of the obstacle problem for a fractional power of the Laplace operator, Comm. Pure Appl. Math., 60 (2007), 67-112.
[27] J. Sun, T. F. Wu, Z. Feng; Two Positive Solutions to Non-autonomous Schrödinger-Poisson Systems, Nonlinearity, 32 (2019), 4002-4032.
[28] Y. Su, H. Chen, S. Liu, X. Fang; Fractional Schrödinger-Poisson systems with weighted Hardy potential and critical exponent, Electron. J. Differential Equations, 2020 no. 1 (2020), 1-17.
[29] J. Sun, T. F. Wu, Z. Feng; Multiplicity of positive solutions for a nonlinear SchrödingerPoisson system, J. Differential Equations, 1 (2016), 586-627.
[30] K. M. Teng; Multiple solutions for a class of fractional Schrödinger equation in $\mathbb{R}^{n}$, Nonlinear Anal. Real World Appl., 21 (2015), 76-86.
[31] K. M. Teng, X. M. He; Ground state solutions for fractional Schrödinger equations with critical Sobolev exponent, Commu. Pure Appl. Anal., 15 (2016), 991-1008.
[32] K. M. Teng; Existence of ground state solutions for the nonlinear fractional SchrödingerPoisson system with critical Sobolev exponent, J. Differential Equations, 261 (2016), 30613106.
[33] K. M. Teng; Ground state solutions for the nonlinear fractional Schrödinger-Poisson system, Applicable Anal., 98 (2019), 1959-1996.
[34] K. M. Teng, C. X. Ye; Ground state and sign-changing solutions for fractional SchrödingerPoisson system with critical growth, Complex Variables and Elliptic Equations, 65 (2019), 1-34.
[35] Z. Wang, H. Zhou, Positive solution for a nonlinear stationary Schrödinger-Poisson system in $\mathbb{R}^{3}$, Discrete Contin. Dyn. Syst., 18 (2007), 809-816.
[36] M. Willem, Minimax Theorems, Birkhäuser, Boston, 1996.
[37] J. Zhang, Ground state and multiple solutions for SchrödingeršCPoisson equations with critical nonlinearity, J. Math. Anal. Appl., 440 (2016), 466-482.
[38] J. J. Zhang, J. M. DO Ó, M. Squassina, Fractional Schrödinger-Poisson systems with a general subcritical or critical nonlinearity, Adv. Nonlinear Stud., 16 (2016), 15-30.
[39] L. Zhao, H. Liu, F. Zhao, Existence and concentration of solutions for the Schrödinger-Poisson equations with steep well potential, J. Differential Equations., 255 (2013), 1-23.
[40] L. Zhao, F. Zhao, On the existence of solutions for the Schrödinger-Poisson equations, J. Math. Anal. Appl., 346 (2008), 155-169.

Lintao Liu
Department of Mathematics, Taiyuan University of Technology, Taiyuan, Shanxi 030024, China

Email address: 956484600@qq.com
Kaimin Teng
Department of Mathematics, Taiyuan University of Technology, Taiyuan, Shanxi 030024, China

Email address: tengkaimin2013@163.com


[^0]:    2010 Mathematics Subject Classification. 35R11, 35B38.
    Key words and phrases. Fractional Schrödinger-Poisson system; variational methods;
    critical growth; Lusternik-Schnirelman category.
    (C) 2021 Texas State University.

    Submitted April 10, 2020. Published February 1, 2021.

