

# Bounds for nonlinear eigenvalue problems \*

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## Abstract

We develop a technique for obtaining bounds on bifurcation curves of nonlinear boundary-value problems defined through nonlinear elliptic partial differential equations.

## 1 Introduction

Recently, we have obtained a variational characterization for the principal solution of a two point boundary-value problem in the line [1]. In particular, we proved the following result.

**Theorem 1.1** *Let the pair  $(\lambda, u)$  be the principal solution (i.e., with  $u(x) \geq 0$ ) of the two point boundary-value problem*

$$\frac{d^2u}{dx^2} + \lambda u = N(u) \quad (1.1)$$

*subject to  $u'(0) = u(1) = 0$ . Let  $u_m = u(0)$ , the sup-norm of the solution. Here  $N(u)$  is a general nonlinear term, which is continuous in  $(0, u_m)$ . Then,*

$$\lambda[u_m] = \max_{g \in D} \left( \int_0^{u_m} N(u)g(u) du + \frac{1}{2} \left( \int_0^{u_m} g'(u)^{1/3} du \right)^3 \right) / \int_0^{u_m} ug(u) du, \quad (1.2)$$

*where  $D = \{g \mid g \in C^1(0, u_m), g' > 0, g(0) = 0\}$ . Moreover, the maximum is attained at some  $\hat{g} \in D$ , which is unique up to a multiplicative constant.*

This theorem cannot be extended, as such, to higher dimensional boundary-value problems since the methods used in the proof depend heavily on the one dimensional character of (1.1). Nevertheless, at least for one particular three dimensional boundary-value problem (namely, the Thomas–Fermi equation) we were able to obtain a, suitably modified, variational characterization of the principal solution [2]. Thus, there are hopes that at least for some boundary-value problems defined through partial differential equations one can obtain a variational characterization of the principal solution. The purpose of this article

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is to illustrate how the methods used in [1, 2] can be extended to find bounds for the principal solution of boundary value problems defined through elliptic partial differential equations. Unfortunately, in general we fall short of obtaining a variational characterization. We will proceed through a well known example, since we believe the methods are well illustrated by it, and it is clear how to extend them to more general situations.

## 2 Nonlinear eigenvalue problem defined through a semilinear elliptic equation

Consider the boundary-value problem

$$-\Delta u + u^3 = \lambda u \quad \text{in } \Omega \subset \mathbb{R}^3, \quad (2.1)$$

with

$$u = 0 \quad \text{in } \partial\Omega. \quad (2.2)$$

Here  $\Omega$  is a bounded, smooth, domain in  $\mathbb{R}^3$ . Then we have,

**Theorem 2.1** *Let the pair  $(\lambda, u)$  be the principal solution (i.e., with  $u(x) \geq 0$ ) of the boundary-value problem (2.1), (2.2), then*

$$4\pi u(x) \leq \frac{2}{3\sqrt{3}} \lim_{\epsilon \rightarrow 0} \int_{\Omega \setminus B_\epsilon(x)} \frac{(\Delta g + \lambda g)^{3/2}}{g^{1/2}} dy, \quad (2.3)$$

where  $B_\epsilon(x) \subset \Omega$  is a ball of radius  $\epsilon$  centered at  $x$ . Here the function  $g$  satisfies,

$$g \in C^2(\Omega \setminus B_\epsilon(x)) \cap C^0(\overline{\Omega \setminus B_\epsilon(x)}), \quad (2.4)$$

$$g = 0 \quad \text{in } \partial\Omega, \quad (2.5)$$

$$g(y) \approx \frac{1}{|x-y|} \quad \text{in the neighborhood of } x, \quad (2.6)$$

(i.e.,  $g(y)$  behaves like the fundamental solution around  $x$ ),

$$g(y) > 0 \quad \text{and} \quad \Delta g + \lambda g > 0 \quad \text{in } \Omega \setminus B_\epsilon(x), \quad (2.7)$$

but is otherwise arbitrary.

*Proof:* Pick any function  $g$  satisfying (2.4), (2.5), (2.6) and (2.7). If we multiply (2.1) by  $g$ , and integrate over  $\Omega \setminus B_\epsilon(x)$  we obtain,

$$-\int_{\Omega \setminus B_\epsilon(x)} g \Delta u dy + \int_{\Omega \setminus B_\epsilon(x)} g u^3 dy = \lambda \int_{\Omega \setminus B_\epsilon(x)} g u dy. \quad (2.8)$$

Using Green's formula and the boundary conditions (2.2) and (2.5) satisfied by  $u$  and  $g$  respectively, we have

$$\int_{\Omega \setminus B_\epsilon(x)} (g \Delta u - u \Delta g) dy = - \int_{\partial B_\epsilon(x)} (g \nabla u - u \nabla g) \cdot \hat{n} dS, \quad (2.9)$$

where  $\hat{n}$  is the exterior normal to the surface of the ball  $B_\epsilon(x)$  and  $dS$  its surface element. Since  $u \in C^2(\Omega)$ , it follows from (2.6) that the limit of the right side of (2.9) as  $\epsilon$  goes to zero is given by  $-4\pi u(x)$ . Thus,

$$\lim_{\epsilon \rightarrow 0} \int_{\Omega \setminus B_\epsilon(x)} (g\Delta u - u\Delta g) dy = -4\pi u(x). \quad (2.10)$$

Hence, using (2.8), (2.9), and (2.10) we have

$$4\pi u(x) = \lim_{\epsilon \rightarrow 0} \int_{\Omega \setminus B_\epsilon(x)} ((\Delta g + \lambda g)u - gu^3) dy. \quad (2.11)$$

Let us denote  $h \equiv \Delta g + \lambda g$ , which is positive by assumption (2.7). For fixed  $y$ , consider the integrand of (2.11) as a function of  $u$ . Maximizing the integrand with respect to  $u$  we get

$$hu - u^3 g \leq \frac{2}{3\sqrt{3}} \frac{h^{3/2}}{g^{1/2}} \quad (2.12)$$

and (2.7) follows from here.

**Remark.** For general facts about bifurcation problems defined through ordinary differential equations see [5]. For general bifurcation problems see [6].

As an application, consider  $\Omega$  to be the unit ball in  $\mathbb{R}^3$ . We will use Theorem 2.1 to find an estimate for the principal branch  $(\lambda, u_m)$  (with  $u(x) \geq 0$ ) of the nonlinear eigenvalue problem (2.1), (2.2), in this case. Here  $u_m$  denotes the sup-norm of the solution, which occurs at zero (the center of the ball). In fact, for any balanced, smooth domain  $\Omega$ , the sup-norm of the positive solution of (2.1), (2.2), is attained at the center of balance, in this case the origin of the ball.

For our purpose, take  $g(x) = \cos(\pi r/2)/r$ , with  $r = |x|$ . Clearly, this  $g$  satisfies the hypothesis (2.4), (2.5), and (2.6). An elementary computation shows that

$$\Delta g = -\frac{\pi^2}{4}g, \quad \text{for } r \neq 0. \quad (2.13)$$

Thus, the corresponding  $h$  will be nonnegative as long as  $\lambda \geq \pi^2/4$ . It turns out that this is not a restriction, since the nonlinear eigenvalue problem (2.1), (2.2) has a nontrivial positive solution if and only if  $\lambda$  is larger than the first Dirichlet eigenvalue of  $\Omega$ , which for the case of the unit ball is  $\pi^2$ . Setting  $x = 0$ , and using the function  $g$  picked above, we conclude from (2.3)

$$u_m \equiv u(0) \leq \frac{4\sqrt{3}}{9\pi} \left(1 - \frac{2}{\pi}\right) \left(\lambda - \frac{\pi^2}{4}\right)^{3/2}, \quad (2.14)$$

which gives the following lower bound for the *nonlinear eigenvalue*  $\lambda$ ,

$$\lambda \geq \frac{\pi^2}{4} + \left[ \frac{3\sqrt{3}}{4} \pi \frac{u_m}{(1 - 2/\pi)} \right]^{2/3} = \frac{\pi^2}{4} + cu_m^{2/3}, \quad (2.15)$$

with  $c \approx 5.015$ . A better multiplicative constant can be obtained in (2.15) using a different trial function  $g$ . In fact take

$$g(x) = \frac{1}{r}(1 - r), \quad (2.16)$$

where  $r = |x|$ , as before. This function  $g$  satisfies (2.4), (2.5), and (2.6). Moreover,  $\Delta g = 0$  for  $r > 0$ , and  $g > 0$  for  $0 < r < 1$ . Hence, from (2.3) we get,  $u_m \equiv u(0) \leq \lambda^{3/2}/(9\sqrt{3})$ , i.e.,

$$\lambda \geq 3^{5/3}u_m^{2/3} \approx 6.240u_m^{2/3}, \quad (2.17)$$

which is better than (2.15) for large values of  $u_m$ .

The bound embodied in (2.3) is a local upper bound on the principal solution of the boundary-value problem (2.1), (2.2). To end with this example, we will consider a function  $g(y)$  depending parametrically on  $x$  to produce an  $x$  dependent bound in (2.3). What we will use as a trial  $g$  will be the fundamental solution of the Laplacian in a ball of radius 1. For the ball of radius  $R$ , the fundamental solution can be constructed using the method of images. It is given explicitly by

$$G_x(y) = \frac{1}{|y-x|} - \frac{R}{|x|} \frac{1}{|y - R^2x/|x|^2|} \quad (2.18)$$

(see, e.g., [3], pp. 19–20). Clearly this particular function satisfies all the hypothesis. Moreover,  $\Delta G_x(y) = 0$  for  $y \neq x$ . Thus, using  $g(y) = G_x(y)$ , with  $R = 1$  in (2.3), together with Newton's theorem (i.e.,  $\int d\Omega_y (1/|y-x|) = 4\pi/\max(|x|, |y|)$ , where the integral is over the sphere of radius  $|y|$  and  $d\Omega_y$  is the invariant measure on the sphere), we get at once

$$u(x) \leq \frac{\lambda^{3/2}}{3\sqrt{3}}(1 - |x|^2). \quad (2.19)$$

The bound (2.19), with a better constant, can be obtained using comparison theorems. In fact, using the comparison function  $1 - |x|^2$ , and the maximum principle (see e.g., [7, 4]), one can show that

$$u(x) \leq \frac{\lambda^{3/2}}{9\sqrt{3}}(1 - |x|^2). \quad (2.20)$$

Notice that at  $x = 0$  the bound (2.20) is precisely the bound (2.17) obtained above.

If instead of (2.1), (2.2), one considers a more general boundary-value problem

$$-\Delta u + f(u) = \lambda u \quad \text{in } \Omega \subset \mathbb{R}^3, \quad (2.21)$$

with

$$u = 0 \quad \text{in } \partial\Omega, \quad (2.22)$$

where  $\Omega$ , as before, is a bounded, smooth, domain in  $\mathbb{R}^3$ , and  $f$  is a positive, continuous, increasing function, with  $f(u)/u \rightarrow \infty$  as  $u \rightarrow \infty$ , then theorem 2.1 holds with (2.3) replaced by

$$4\pi u(x) \leq \lim_{\epsilon \rightarrow 0} \int_{\Omega \setminus B_\epsilon(x)} f_L\left(\frac{h}{g}\right) g \, dy. \quad (2.23)$$

Here,  $f_L$  denotes the Legendre transform of the function  $f$ , and  $h \equiv \Delta g + \lambda g$  as before.

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## References

- [1] R.D. Benguria and M.C. Depassier. A variational method for nonlinear eigenvalue problems. *Contemporary Mathematics*, 217:1–17, 1998.
- [2] R.D. Benguria and J.M. Yáñez. Variational principle for the chemical potential in the Thomas–Fermi model. *J.Phys. A: Math. Gen.*, 31:585–593, 1998.
- [3] D. Gilbarg and N.S. Trudinger. Partial Differential Equations. Second Edition. Grundlehren der mathematischen Wissenschaften 224. Springer–Verlag, Berlin, 1998.
- [4] M.H. Protter and H.F. Weinberger. Maximum principles in differential equations. Prentice–Hall, Englewood Cliffs, NJ., 1967.
- [5] P. H. Rabinowitz. Nonlinear Sturm–Liouville Problems for Second Order Ordinary Differential Equations. *Comm. Pure Applied Math.*, XXIII:939–961, 1970.
- [6] P. H. Rabinowitz (Ed.) Applications of Bifurcation Theory. Academic Press, New York, 1977.
- [7] D.H. Sattinger. Topics in stability and bifurcation theory. Lecture notes in mathematics, vol. 309. Springer–Verlag, Berlin, 1973.

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