

SECOND-ORDER DIFFERENTIAL EQUATIONS WITH ASYMPTOTICALLY SMALL DISSIPATION AND PIECEWISE FLAT POTENTIALS

ALEXANDRE CABOT, HANS ENGLER, SÉBASTIEN GADAT

Pour Alban, né le 27 mars 2008

ABSTRACT. We investigate the asymptotic properties as $t \rightarrow \infty$ of the differential equation

$$\ddot{x}(t) + a(t)\dot{x}(t) + \nabla G(x(t)) = 0, \quad t \geq 0$$

where $x(\cdot)$ is \mathbb{R} -valued, the map $a : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ is non increasing, and $G : \mathbb{R} \rightarrow \mathbb{R}$ is a potential with locally Lipschitz continuous derivative. We identify conditions on the function $a(\cdot)$ that guarantee or exclude the convergence of solutions of this problem to points in $\operatorname{argmin} G$, in the case where G is convex and $\operatorname{argmin} G$ is an interval. The condition

$$\int_0^\infty e^{-\int_0^t a(s) ds} dt < \infty$$

is known to be necessary for convergence of trajectories. We give a slightly stronger condition that is sufficient.

1. INTRODUCTION

In this note, we study the differential equation

$$\ddot{x}(t) + a(t)\dot{x}(t) + \nabla G(x(t)) = 0, \quad t \geq 0 \tag{1.1}$$

where $x(\cdot)$ is \mathbb{R} -valued, the map $G : \mathbb{R} \rightarrow \mathbb{R}$ is at least of class C^1 , and $a : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ is a non increasing function. In a previous paper [3], we studied this differential equation in a finite- or infinite-dimensional Hilbert space \mathcal{H} . We are interested in the case where $a(t) \rightarrow 0$ as $t \rightarrow \infty$. Broadly speaking, convergence of solutions can be expected if $a(t) \rightarrow 0$ sufficiently slowly. One of the questions left open in that paper was whether solutions converge to a limit if the property

$$\int_0^\infty e^{-\int_0^t a(s) ds} dt = \infty \tag{1.2}$$

does *not* hold and if $\operatorname{argmin} G$ consists of more than just one point. In this note, we give a positive answer to this question, in the one dimensional case.

2000 *Mathematics Subject Classification.* 34G20, 34A12, 34D05.

Key words and phrases. Differential equation; dissipative dynamical system; vanishing damping; asymptotic behavior.

©2009 Texas State University - San Marcos.

Published April 15, 2009.

2. PRELIMINARY FACTS

Throughout this paper, we will denote by $G : \mathbb{R} \rightarrow \mathbb{R}$ a \mathcal{C}^1 function for which the derivative G' is Lipschitz continuous, uniformly on bounded sets. The function $a : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ will always be assumed to be continuous and non-increasing. We also define the energy

$$\mathcal{E}(t) = G(x(t)) + \frac{1}{2}|\dot{x}(t)|^2.$$

Here are some basic results for solutions of (1.1) from [3].

For any $(x_0, x_1) \in \mathbb{R}^2$, the problem (1.1) has a unique solution $x(\cdot) \in \mathcal{C}^2([0, T], \mathbb{R})$ satisfying $x(0) = x_0, \dot{x}(0) = x_1$ on some maximal time interval $[0, T) \subset [0, \infty)$. For every $t \in [0, T)$, the energy identity holds

$$\frac{d}{dt}\mathcal{E}(t) = -a(t)|\dot{x}(t)|^2.$$

If in addition G is bounded from below, then

$$\int_0^T a(t)|\dot{x}(t)|^2 dt < \infty, \tag{2.1}$$

and the solution exists for all $T > 0$. If also $G(\xi) \rightarrow \infty$ as $|\xi| \rightarrow \infty$ (i.e. if G is *coercive*), then all solutions to (1.1) remain bounded together with their first and second derivatives for all $t > 0$. The bound depends only on the initial data. If a solution x to (1.1) converges toward some $\bar{x} \in \mathbb{R}$, then $\lim_{t \rightarrow \infty} \dot{x}(t) = \lim_{t \rightarrow \infty} \ddot{x}(t) = 0$ and $G'(\bar{x}) = 0$. If $\int_0^\infty a(s) ds < \infty$ and if $\inf G > -\infty$, then solutions $x(\cdot)$ of (1.1) for which $(x(0), \dot{x}(0)) \notin \operatorname{argmin} G \times \{0\}$ cannot converge to a point in $\operatorname{argmin} G$.

For the remainder of this note we shall assume that $\operatorname{argmin} G \neq \emptyset$. Without loss of generality, we may assume that $\min_{\mathbb{R}} G = 0$ and $G(0) = 0$. If for some $\rho \in \mathbb{R}_+$ and $z \in \operatorname{argmin} G$

$$\forall x \in \mathbb{R}, \quad G(x) - G(z) \leq \rho G'(x)(x - z)$$

then it is possible to show that any solution x to the differential equation (1.1) satisfies

$$\int_0^\infty a(t) \mathcal{E}(t) dt < \infty.$$

Since $t \mapsto \mathcal{E}(t)$ is decreasing, this estimate implies that $\mathcal{E}(t) \rightarrow \min G = 0$ as $t \rightarrow \infty$, provided that $\int_0^\infty a(t) dt = \infty$. If now $\operatorname{argmin} G = \{\bar{x}\}$ is a singleton, then trajectories must converge to \bar{x} under fairly weak additional conditions. The reader is referred to [3] for details.

3. CONVEX POTENTIALS WITH NON-UNIQUE MINIMA

In this section, we investigate the convergence of the trajectories of (1.1) when $\operatorname{argmin} G$ is *not* a singleton. While the previous discussion shows that $\int_0^\infty a(s) ds = \infty$ is a necessary condition for trajectories to converge to a point in $\operatorname{argmin} G$, this condition is clearly not sufficient, as the particular case $G \equiv 0$ shows. In this case, the solution is given by

$$x(t) = x(0) + \dot{x}(0) \int_0^t e^{-\int_0^s a(u) du} ds$$

and the solution x converges if and only if (1.2) does not hold. Therefore it is natural to ask whether for a general potential G , the trajectory x is convergent if

this condition does not hold. The potential G is assumed to have all the properties listed in the previous section. A general result of non-convergence of the trajectories under the condition (1.2) is shown in [3]. There, we assume that G is coercive, $\inf_{\mathbb{R}} G = 0$, $\operatorname{argmin} G = [\alpha, \beta]$ for some $\alpha < \beta$, and that G is non-increasing on $(-\infty, \alpha]$ and non-decreasing on $[\beta, \infty)$. It is also assumed that a satisfies condition (1.2). Then either a solution satisfies $(x(0), \dot{x}(0)) \in [\alpha, \beta] \times \{0\}$, or else the ω -limit set $\omega(x_0, \dot{x}_0)$ contains $[\alpha, \beta]$ and hence the trajectory x does not converge.

We now ask if the converse assertion is true: do the trajectories x of (1.1) converge if (1.2) does not hold? We give a positive answer when the map a satisfies the following stronger condition

$$\int_0^\infty e^{-\theta \int_0^s a(u) du} ds < \infty, \quad (3.1)$$

for some $\theta \in (0, 1)$.

Theorem 3.1. *Let $G : \mathbb{R} \rightarrow \mathbb{R}$ be a convex function of class \mathcal{C}^1 such that G' is Lipschitz continuous on the bounded sets of \mathbb{R} . Assume that $\operatorname{argmin} G = [\alpha, \beta]$ with $\alpha < \beta$ and that there exists $\delta > 0$ such that*

$$\forall \xi \in (-\infty, \alpha], \quad G'(\xi) \leq 2\delta(\xi - \alpha) \quad \text{and} \quad \forall \xi \in [\beta, \infty), \quad G'(\xi) \geq 2\delta(\xi - \beta).$$

Let $a : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ be a differentiable non increasing map such that $\lim_{t \rightarrow \infty} a(t) = 0$ and such that condition (3.1) holds for some positive $\theta < 1$. Then, for any solution x to the differential equation (1.1), $\lim_{t \rightarrow \infty} x(t)$ exists.

Proof. We may assume without loss of generality that $\alpha = 0, \beta = 1$. The conditions on G imply that it is coercive, hence $\lim_{t \rightarrow \infty} \mathcal{E}(t) = 0$ and $|x(t)| \leq M$ for some $M > 0$, for all $t \in \mathbb{R}_+$.

Define the set $\mathcal{T} = \{t \geq 0 \mid \dot{x}(t) = 0\}$. We shall show that either $\mathcal{T} = [0, \infty)$ or \mathcal{T} is a finite set. Assume first that \mathcal{T} has an accumulation point t^* . Then $\dot{x}(t^*) = 0$ and $\ddot{x}(t^*) = 0$ by Rolle's Theorem. Since then $\dot{x}(t^*) = \ddot{x}(t^*) = G'(x(t^*)) = 0$, $x(\cdot)$ must be constant by forward and backward uniqueness, $\mathcal{T} = [0, \infty)$, and clearly the limit exists. Therefore we may now assume that \mathcal{T} is discrete. If \mathcal{T} is a finite set, then \dot{x} does not change sign for sufficiently large t , and the trajectory x has a limit.

It remains to consider the case $\mathcal{T} = \{t_n \mid n \in \mathbb{N}\}$, where the t_n are increasing and tend to ∞ . We want to show that this is impossible. Observe that at each t_n , \dot{x} must change its sign and $G'(x(t_n)) \neq 0$, since otherwise also $\ddot{x}(t_n) = 0$ and we would again have a stationary solution. Without loss of generality, we can assume that $\dot{x}(0) < 0$, $x(0) < 0$ and therefore $x(t_0) < 0$. Since $G'(x(t_0)) < 0$, equation (1.1) shows that $\ddot{x}(t_0) > 0$, hence the map \dot{x} is positive on (t_0, t_1) , $x(t_1) > 1$, \dot{x} is negative on (t_1, t_2) , and so on.

The argument so far shows that $G'(x(t))$ vanishes on a union of infinitely many disjoint closed intervals,

$$\{t \mid 0 \leq x(t) \leq 1\} = \cup_{k \geq 0} [u_{2k}, u_{2k+1}]$$

where $0 < t_0 < u_0$ and $u_{2k-1} < t_k < u_{2k}$ for $k = 1, 2, \dots$. Let us observe that, for every $k \in \mathbb{N}$,

$$1 = |x(u_{2k+1}) - x(u_{2k})| = \int_{u_{2k}}^{u_{2k+1}} |\dot{x}(t)| dt \leq |u_{2k+1} - u_{2k}| \max_{t \geq u_{2k}} |\dot{x}(t)|.$$

Since $\lim_{t \rightarrow \infty} \dot{x}(t) = 0$, we deduce that $\lim_{k \rightarrow \infty} |u_{2k+1} - u_{2k}| = \infty$.

We next observe that for $u_{2k} \leq t \leq u_{2k+1}$ the function $v = \dot{x}$ satisfies $\dot{v}(t) + a(t)v(t) = 0$ and hence

$$\forall t \in [u_{2k}, u_{2k+1}], \quad \dot{x}(t) = \dot{x}(u_{2k})e^{-\int_{u_{2k}}^t a(\tau)d\tau}. \quad (3.2)$$

Claim 3.2. *There is a constant γ such that $u_{2k+2} - u_{2k+1} \leq \gamma$ for all $k \in \mathbb{N}$.*

To show this claim, fix $k \in \mathbb{N}$ and assume that $t \in [u_{2k+1}, u_{2k+2}]$. Assume for now that k is odd and thus $x(t) \leq 0$. Define the quantity $A(t) = \exp\left(\frac{1}{2}\int_0^t a(s)ds\right)$ and set $y(t) = A(t)x(t)$. Then y is the solution of the differential equation

$$\ddot{y}(t) + A(t)G'\left(\frac{y(t)}{A(t)}\right) - \left(\frac{a^2(t)}{4} + \frac{\dot{a}(t)}{2}\right)y(t) = 0, \quad (3.3)$$

and satisfies $y(u_{2k+1}) = y(u_{2k+2}) = 0$ and $\dot{y}(u_{2k+1}) = A(u_{2k+1})\dot{x}(u_{2k+1}) < 0$. Since the map a converges to 0, we can choose k large enough so that $a(t) < 2\sqrt{\delta}$ for every $t \in [u_{2k+1}, u_{2k+2}]$. On the other hand, the assumption on G' shows that, for every $t \in [u_{2k+1}, u_{2k+2}]$,

$$A(t)G'\left(\frac{y(t)}{A(t)}\right) \leq 2\delta y(t).$$

Recalling finally that $\dot{a}(t) \leq 0$ for every $t \geq 0$, we deduce from (3.3) that

$$\forall t \in [u_{2k+1}, u_{2k+2}], \quad \ddot{y}(t) + \delta y(t) \geq 0.$$

The unique solution z of the differential equation $\ddot{z}(t) + \delta z(t) = 0$ with the same initial conditions as y has the first zero larger than u_{2k+1} at $u_{2k+1} + \frac{\pi}{\sqrt{\delta}}$. By a standard comparison argument, we deduce that y vanishes before z does, hence

$$u_{2k+2} \leq u_{2k+1} + \gamma, \quad \gamma = \frac{\pi}{\sqrt{\delta}}.$$

The same argument applies if k is even. This proves the claim.

Claim 3.3. *There is a $k_0 \in \mathbb{N}$ such that for $k \geq k_0$*

$$|\dot{x}(u_{2k+2})| \leq |\dot{x}(u_{2k})|e^{-\theta \int_{u_{2k}}^{u_{2k+2}} a(s)ds}.$$

where θ is as in (3.1).

To prove this, pick k_0 so large that for all $k \geq k_0$,

$$(1 - \theta)(u_{2k+2} - u_{2k}) \geq \gamma\theta.$$

This is possible since $u_{2k+2} - u_{2k} \rightarrow \infty$ as $k \rightarrow \infty$. Since a is non-increasing, this implies that

$$\begin{aligned} \theta \int_{u_{2k+1}}^{u_{2k+2}} a(\tau)d\tau &\leq \gamma\theta a(u_{2k+1}) \\ &\leq (1 - \theta)(u_{2k+1} - u_{2k})a(u_{2k+1}) \\ &\leq (1 - \theta) \int_{u_{2k}}^{u_{2k+1}} a(\tau)d\tau \end{aligned}$$

and hence

$$\theta \int_{u_{2k}}^{u_{2k+2}} a(\tau)d\tau \leq \int_{u_{2k}}^{u_{2k+1}} a(\tau)d\tau.$$

Then for $k \geq k_0$,

$$\begin{aligned} |\dot{x}(u_{2k+2})| &\leq |\dot{x}(u_{2k+1})| = |\dot{x}(u_{2k})| e^{-\int_{u_{2k}}^{u_{2k+1}} a(s) ds} \\ &\leq |\dot{x}(u_{2k})| e^{-\theta \int_{u_{2k}}^{u_{2k+2}} a(s) ds} \end{aligned}$$

proving the claim.

Claim 3.4. *If the set \mathcal{T} is unbounded, there must exist a constant C , depending on \mathcal{T} and on $x(0), \dot{x}(0)$ such that for all $t \geq 0$*

$$|\dot{x}(t)| \leq C e^{-\theta \int_0^t a(s) ds}. \tag{3.4}$$

By making sure that C is sufficiently large, we only have to prove the estimate for $t \geq u_{2k_0}$. First assume that $u_{2k} \leq t \leq u_{2k+1}$ for some k . Then from (3.2)

$$|\dot{x}(t)| \leq |\dot{x}(u_{2k})| e^{-\int_{u_{2k}}^t a(s) ds} \leq |\dot{x}(u_{2k})| e^{-\theta \int_{u_{2k}}^t a(s) ds}.$$

Using induction, we deduce from Claim 3.3 that

$$|\dot{x}(t)| \leq |\dot{x}(u_{2k_0})| e^{-\theta \int_{u_{2k_0}}^t a(s) ds} = C_1 e^{-\theta \int_0^t a(s) ds}$$

with $C_1 = |\dot{x}(u_{2k_0})| e^{\theta \int_0^{u_{2k_0}} a(s) ds}$. Next consider the case where $u_{2k+1} < t \leq u_{2k+2}$ for some k . Then

$$|\dot{x}(t)| \leq |\dot{x}(u_{2k+1})| \leq C_1 e^{-\theta \int_0^{u_{2k+1}} a(s) ds} \leq C_1 e^{\theta \int_{u_{2k+1}}^{u_{2k+2}} a(\tau) d\tau} e^{-\theta \int_0^t a(s) ds}.$$

Due to Claim 3.2, $e^{\theta \int_{u_{2k+1}}^{u_{2k+2}} a(\tau) d\tau} \leq C_2$ for all k , for some constant C_2 . Estimate (3.4) now follows for $t \geq u_{2k_0}$ with $C = C_1 C_2$. By enlarging C further, the estimate follows for all $t \geq 0$.

Let us now conclude the proof of the theorem. From assumption (3.1) and estimate (3.4), we derive that $\dot{x} \in L^1(0, \infty)$. Hence $\lim_{t \rightarrow \infty} x(t)$ exists, contradicting the initial assumption. Therefore $\lim_{t \rightarrow \infty} x(t)$ exists after all, and the theorem has been proved. \square

Remark 3.5. *Note that the map $t \mapsto \frac{c}{t+1}$ with $c > 1$ satisfies condition (3.1) for every $\theta \in (\frac{1}{c}, 1)$. In fact, if merely $a(t) \geq \frac{c}{t+1}$ for t large enough for some $c > 1$, then condition (3.1) is satisfied. Consider next the family of maps $a : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ defined by*

$$a(t) = \frac{1}{t+1} + \frac{d}{(t+1) \ln(t+2)},$$

for some $d > 0$. It is immediate to check that condition (1.2) holds if and only if $d \in (0, 1]$. Thus non-stationary trajectories of (1.1) do not converge when $d \in (0, 1]$. But condition (3.1) is never satisfied, for any $\theta \in (0, 1)$ and $d > 0$, and the convergence of trajectories remains an open question. Thus there remains a “logarithmic” gap between the criteria for existence and non-existence of limits.

We conclude with some remarks on convergence results in dimension $n > 1$. It is possible to extend the non-convergence result given at the beginning of this section to the case where the differential equation is given in a Hilbert space \mathcal{H} , see [3]. However, it is not clear how to prove that $\lim_{t \rightarrow \infty} x(t)$ exists, in a general Hilbert space \mathcal{H} and for the case where G is convex and $\text{argmin} G$ is not a singleton. Since in this case $|\dot{x}(t)| \leq \sqrt{2\mathcal{E}(t)}$, it appears natural to derive convergence results from suitable estimates for $\mathcal{E}(t)$. In [3], we give conditions that imply $\mathcal{E}(t) \leq Da(t)$ for all t , for some constant $D > 0$. However, since we must also assume that

$\int_0^\infty a(s)ds = \infty$, these estimates are not strong enough to guarantee the convergence of trajectories.

One could try to extend the proof of Theorem 3.1. Set $a_1(t) = a(t) \cdot \chi_S(x(t))$, where χ_S is the characteristic function of $S = \operatorname{argmin} G$, then $\frac{d}{dt}\mathcal{E}(t) \leq -2a_1(t)\mathcal{E}(t)$, and hence $\mathcal{E}(t) \leq \mathcal{E}(0)e^{-2\int_0^t a_1(s)ds}$. If the function $t \mapsto e^{-\int_0^t a_1(s)ds}$ can be shown to be in $L^1(0, \infty)$, it would follow that $|\dot{x}|$ is integrable, implying the convergence of trajectories. This works in the one-dimensional case since the behavior of trajectories is quite simple. However, if $\dim \mathcal{H} > 1$, it is difficult to satisfy this property, since trajectories corresponding to (1.1) can be expected to behave like trajectories of a billiard problem in $S = \operatorname{argmin} G$ for large times.

When the map a is constant and positive, it is established in [1, 2] that the trajectories of (1.1) are weakly convergent if the potential $G : \mathcal{H} \rightarrow \mathbb{R}$ is convex and $\operatorname{argmin} G \neq \emptyset$, in an arbitrary Hilbert space \mathcal{H} . The key ingredient of the proof is the Opial lemma [4], which allows the authors of these papers to prove convergence even if $|\dot{x}(\cdot)|$ is only in $L^2(0, \infty)$ and not in $L^1(0, \infty)$. However, if e.g. $a(t) = \frac{c}{t+1}$, then Opial's lemma requires that we show $\int_0^\infty (t+1)|\dot{x}(t)|^2 dt < \infty$, while (2.1) implies only $\int_0^\infty \frac{1}{t+1}|\dot{x}(t)|^2 dt < \infty$. Hence there remains a gap if arguments similar to those in [1] or [2] are to be used. It is unclear how this gap can be closed.

REFERENCES

- [1] F. Alvarez, On the minimizing property of a second order dissipative system in Hilbert spaces, *SIAM J. on Control and Optimization*, 38 (2000), n° 4, 1102-1119.
- [2] H. Attouch, X. Goudou, P. Redont, The heavy ball with friction method: I the continuous dynamical system, *Communications in Contemporary Mathematics*, 2 (2000), n° 1, 1-34.
- [3] A. Cabot, H. Engler, S. Gadat, On the long time behavior of second order differential equations with asymptotically small dissipation, *Trans. of the Amer. Math. Soc.*, in press. <http://arxiv.org/abs/0710.1107>
- [4] Z. Opial, Weak convergence of the sequence of successive approximations for nonexpansive mappings, *Bull. of the American Math. Society*, 73, (1967), 591-597.

ALEXANDRE CABOT

DÉPARTEMENT DE MATHÉMATIQUES, UNIVERSITÉ MONTPELLIER II, CC 051, PLACE EUGÈNE BATAILLON, 34095 MONTPELLIER CEDEX 5, FRANCE

E-mail address: acabot@math.univ-montp2.fr

HANS ENGLER

DEPARTMENT OF MATHEMATICS, GEORGETOWN UNIVERSITY, BOX 571233, WASHINGTON, DC 20057, USA

E-mail address: engler@georgetown.edu

SÉBASTIEN GADAT

INSTITUT DE MATHÉMATIQUES DE TOULOUSE, UNIVERSITÉ PAUL SABATIER, 118, ROUTE DE NARBONNE 31062 TOULOUSE CEDEX 9, FRANCE

E-mail address: Sebastien.Gadat@math.ups-tlse.fr