

## Research Article

# On Some Inequalities of Uncertainty Principles Type in Quantum Calculus

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The aim of this paper is to generalize the  $q$ -Heisenberg uncertainty principles studied by Bettaibi et al. (2007), to state local uncertainty principles for the  $q$ -Fourier-cosine, the  $q$ -Fourier-sine, and the  $q$ -Bessel-Fourier transforms, then to provide an inequality of Heisenberg-Weyl-type for the  $q$ -Bessel-Fourier transform.

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## 1. Introduction

The uncertainty principle is a metatheorem in harmonic analysis that asserts, with the use of some inequalities, that a function and its Fourier transform cannot be sharply localized. We refer to the survey article by Folland and Sitaram [1] and the book of Havin and Jöricke [2] for various classical uncertainty principles of different nature which may be found in the literature.

In [3], the authors gave  $q$ -analogues of the Heisenberg uncertainty principle for the  $q$ -Fourier-cosine and the  $q$ -Fourier-sine transforms. One of the aims of this paper is to provide a generalization of their work next to state local uncertainty principles for various  $q$ -Fourier transforms.

This paper is organized as follows. In Section 2, we present some preliminaries results and notations that will be useful in the sequel. In Section 3, we prove a density theorem and a  $q$ -analogue of the Hausdorff-Young inequality. Then, we state a generalization of the  $q$ -Heisenberg uncertainty principle for the  $q$ -Fourier-cosine and the  $q$ -Fourier-sine transforms. In Section 4, we state local uncertainty principles for the  $q$ -Fourier-cosine,  $q$ -Fourier-sine, and

$q$ -Bessel-Fourier transforms. Then, we give a Heisenberg-Weyl-type inequality for some  $q$ -Bessel-Fourier transform.

## 2. Notations and preliminaries

Throughout this paper, we assume  $q \in ]0, 1[$ . We recall some usual notions and notations used in the  $q$ -theory (see [4, 5]). We refer to the book by Gasper and Rahman [4] for the definitions, notations, and properties of the  $q$ -shifted factorials and the  $q$ -hypergeometric functions.

We write  $\mathbb{R}_q = \{\pm q^n : n \in \mathbb{Z}\}$ ,  $\mathbb{R}_{q,+} = \{q^n : n \in \mathbb{Z}\}$ , and

$$[x]_q = \frac{1 - q^x}{1 - q}, \quad x \in \mathbb{C}, \quad [n]_q! = \frac{(q; q)_n}{(1 - q)^n}, \quad n \in \mathbb{N}. \quad (2.1)$$

The  $q$ -derivative of a function  $f$  is given by

$$(D_q f)(x) = \frac{f(x) - f(qx)}{(1 - q)x} \quad \text{if } x \neq 0, \quad (2.2)$$

$(D_q f)(0) = \lim_{k \rightarrow +\infty} (D_q f)(q^k)$ , provided that the limit exists.

The  $q$ -Jackson integrals from 0 to  $a$  and from 0 to  $\infty$ , of a function  $f$ , are (see [6])

$$\int_0^a f(x) d_q x = (1 - q)a \sum_{n=0}^{\infty} f(aq^n) q^n, \quad \int_0^{\infty} f(x) d_q x = (1 - q) \sum_{n=-\infty}^{\infty} f(q^n) q^n, \quad (2.3)$$

provided that the sums converge absolutely.

The  $q$ -Jackson integral in a generic interval  $[a, b]$  is given by (see [6])

$$\int_a^b f(x) d_q x = \int_0^b f(x) d_q x - \int_0^a f(x) d_q x. \quad (2.4)$$

The  $q$ -integration by parts rule is given, for suitable functions  $f$  and  $g$ , by

$$\int_a^b g(x) D_q f(x) d_q x = f(b)g(b) - f(a)g(a) - \int_a^b f(qx) D_q g(x) d_q x. \quad (2.5)$$

Jackson (see [6]) defined a  $q$ -analogue of the Gamma function by

$$\Gamma_q(x) = \frac{(q; q)_{\infty}}{(q^x; q)_{\infty}} (1 - q)^{1-x}, \quad x \neq 0, -1, -2, \dots \quad (2.6)$$

The third Jackson  $q$ -Bessel function (see [7, 8]) is

$$J_{\nu}(z; q^2) = \frac{z^{\nu}}{(1 - q^2)^{\nu} \Gamma_{q^2}(\nu + 1)} {}_1\phi_1(0; q^{2\nu+2}; q^2, q^2 z^2), \quad (2.7)$$

and the  $q$ -trigonometric functions ( $q$ -cosine and  $q$ -sine) are defined by (see [9])

$$\begin{aligned}\cos(x; q^2) &= \frac{\Gamma_{q^2}(1/2)}{q(1+q^{-1})^{1/2}} x^{1/2} J_{-1/2} \left( \frac{1-q}{q} x; q^2 \right) = \sum_{n=0}^{\infty} (-1)^n q^{n(n-1)} \frac{x^{2n}}{[2n]_q!}, \\ \sin(x; q^2) &= \frac{\Gamma_{q^2}(1/2)}{(1+q^{-1})^{1/2}} x^{1/2} J_{1/2} \left( \frac{1-q}{q} x; q^2 \right) = \sum_{n=0}^{\infty} (-1)^n q^{n(n-1)} \frac{x^{2n+1}}{[2n+1]_q!}.\end{aligned}\quad (2.8)$$

They verify

$$D_q \cos(x; q^2) = -\frac{1}{q} \sin(qx; q^2), \quad D_q \sin(x; q^2) = \cos(x; q^2). \quad (2.9)$$

We need the following spaces and norms.

(i)  $\mathcal{S}_{*q}(\mathbb{R}_q)$  is the space of even functions  $f$  on  $\mathbb{R}_q$  satisfying

$$\forall n, m \in \mathbb{N}, \quad P_{n,m,q}(f) = \sup_{x \in \mathbb{R}_q; 0 \leq k \leq n} |(1+x^2)^m D_q^k f(x)| < +\infty. \quad (2.10)$$

(ii)  $L_q^n(\mathbb{R}_{q,+}, x^{2\nu+1} d_q x)$ ,  $n \geq 1$ ,  $\nu \geq -1/2$ , is the set of all functions defined on  $\mathbb{R}_{q,+}$  such that

$$\|f\|_{n,\nu,q} = \left\{ \int_0^{\infty} |f(x)|^n x^{2\nu+1} d_q x \right\}^{1/n} < \infty. \quad (2.11)$$

(iii)  $L_q^n(\mathbb{R}_{q,+}) = L_q^n(\mathbb{R}_{q,+}, d_q x)$ ,  $n \geq 1$ , and  $\|\cdot\|_{n,q} = \|\cdot\|_{n,-1/2,q}$ .

(iv)  $L_q^\infty(\mathbb{R}_{q,+})$  is the set of all bounded functions on  $\mathbb{R}_{q,+}$ . We write  $\|f\|_{\infty,q} = \sup_{x \in \mathbb{R}_{q,+}} |f(x)|$ .

### 3. Generalization of the Heisenberg uncertainty principle

The  $q$ -Fourier-cosine and the  $q$ -Fourier-sine transforms are defined as (see [8, 9])

$$\mathcal{F}_q(f)(x) = c_q \int_0^{\infty} f(t) \cos(xt; q^2) d_q t, \quad {}_q\mathcal{F}(f)(x) = c_q \int_0^{\infty} f(t) \sin(xt; q^2) d_q t, \quad (3.1)$$

where

$$c_q = \frac{(1+q^{-1})^{1/2}}{\Gamma_{q^2}(1/2)}. \quad (3.2)$$

Letting  $q \uparrow 1$  subject to the condition  $(\text{Log}(1-q)/\text{Log}(q)) \in \mathbb{Z}$  gives, at least formally, the classical Fourier transforms (see [3, 10]). In the remainder of the present paper, we assume that this condition holds.

It was shown in [8, 9] that we have the following result.

**Proposition 3.1.** (1) For  $f \in L_q^1(\mathbb{R}_{q,+})$ , one has  $\mathcal{F}_q(f) \in L_q^\infty(\mathbb{R}_{q,+})$  and

$$\|\mathcal{F}_q(f)\|_{\infty,q} \leq \frac{(1+q^{-1})^{1/2}}{\Gamma_{q^2}(1/2)(q;q)_\infty^2} \|f\|_{1,q}. \quad (3.3)$$

(2)  $\mathcal{F}_q$  is an isomorphism of  $L_q^2(\mathbb{R}_{q,+})$  (resp.,  $S_{*,q}(\mathbb{R}_q)$ ) onto itself. Moreover, one has  $\mathcal{F}_q^{-1} = \mathcal{F}_q$  and the following Plancherel formula:

$$\|\mathcal{F}_q(f)\|_{2,q} = \|f\|_{2,q}, \quad f \in L_q^2(\mathbb{R}_{q,+}). \quad (3.4)$$

Similarly, it was shown in [3, 8] that the  $q$ -Fourier-sine transform verifies the following properties.

**Proposition 3.2.** (1) For  $f \in L_q^1(\mathbb{R}_{q,+})$ , one has  ${}_q\mathcal{F}(f) \in L_q^\infty(\mathbb{R}_{q,+})$  and

$$\|{}_q\mathcal{F}(f)\|_{\infty,q} \leq \frac{(1+q^{-1})^{1/2}}{\Gamma_{q^2}(1/2)(q;q)_\infty^2} \|f\|_{1,q}. \quad (3.5)$$

(2)  ${}_q\mathcal{F}$  is an isomorphism of  $L_q^2(\mathbb{R}_{q,+})$  onto itself; its inverse is given by  ${}_q\mathcal{F}^{-1} = (1/q^2) {}_q\mathcal{F}$ . One has the following Plancherel formula:

$$\|{}_q\mathcal{F}(f)\|_{2,q} = q \|f\|_{2,q}, \quad f \in L_q^2(\mathbb{R}_{q,+}). \quad (3.6)$$

Let us now state the following useful density result.

**Proposition 3.3.** For all  $n \geq 1$ ,  $S_{*,q}(\mathbb{R}_q)$  is dense in  $L_q^n(\mathbb{R}_{q,+})$ .

*Proof.* Let  $n \geq 1$  and  $f \in L_q^n(\mathbb{R}_{q,+})$ . For  $p \in \mathbb{N}$ , put  $f_p = f \cdot \chi_{[q^p, q^{-p}]}$ , where  $\chi_{[q^p, q^{-p}]}$  is the characteristic function of  $[q^p, q^{-p}]$ .

It is clear that for all  $p \in \mathbb{N}$ ,  $f_p \in S_{*,q}(\mathbb{R}_q)$  and  $|f - f_p|^n \leq |f|^n$ . So, the Lebesgue theorem implies that  $(f_p)_p$  converges to  $f$  in  $L_q^n(\mathbb{R}_{q,+})$ .  $\square$

*Remark 3.4.* Using the density of  $S_{*,q}(\mathbb{R}_q)$  in  $L_q^n(\mathbb{R}_{q,+})$  ( $n \geq 1$ ), one can see that the  $q$ -Fourier-cosine (resp.,  $q$ -Fourier-sine) transform has a unique continuous extension on  $L_q^n(\mathbb{R}_{q,+})$ , that will also be denoted as  $\mathcal{F}_q$  (resp.,  ${}_q\mathcal{F}$ ). We have the following  $q$ -analogue of the Hausdorff-Young inequality.

**Theorem 3.5.** Let  $n \in ]1, 2]$  (resp.,  $n = 1$ ) and  $m = n/(n-1)$  (resp.,  $m = \infty$ ) be the dual exponent of  $n$ . For all  $f$  in  $L_q^n(\mathbb{R}_{q,+})$ , the functions  $\mathcal{F}_q(f)$  and  ${}_q\mathcal{F}(f)$  belong to  $L_q^m(\mathbb{R}_{q,+})$ , and one has

$$\|\mathcal{F}_q(f)\|_{m,q} \leq C_1 \|f\|_{n,q}, \quad \|{}_q\mathcal{F}(f)\|_{m,q} \leq C_2 \|f\|_{n,q}, \quad (3.7)$$

where

$$C_1 = \left( \frac{(1+q^{-1})^{1/2}}{\Gamma_{q^2}(1/2)(q; q)_\infty^2} \right)^{1-2((n-1)/n)}, \quad C_2 = \left( \frac{(1+q^{-1})^{1/2}}{\Gamma_{q^2}(1/2)(q; q)_\infty^2} \right)^{1-2((n-1)/n)} q^{2((n-1)/n)}. \quad (3.8)$$

*Proof.* The result is a direct consequence of [11, Theorem 1.3.4, page 35], and Propositions 3.1 and 3.2, by taking  $S_{*,q}(\mathbb{R}_q)$  as a set of simple functions.  $\square$

The following lemma gives relations between the two Fourier  $q$ -trigonometric transforms.

**Lemma 3.6.** (1) For  $f \in L_q^2(\mathbb{R}_{q,+})$  such that  $D_q f \in L_q^2(\mathbb{R}_{q,+})$ , one has

$${}_q \mathcal{F}(D_q f)(\lambda) = -\frac{\lambda}{q} \mathcal{F}_q(f)\left(\frac{\lambda}{q}\right), \quad \lambda \in \mathbb{R}_{q,+}. \quad (3.9)$$

(2) Additionally, if  $\lim_{n \rightarrow +\infty} f(q^n) = 0$ , then

$$\mathcal{F}_q(D_q f)(\lambda) = \frac{\lambda}{q^2} {}_q \mathcal{F}(f)(\lambda), \quad \lambda \in \mathbb{R}_{q,+}. \quad (3.10)$$

*Proof.* The same steps as in the proof of [3, Lemma 2]; the  $q$ -integration by parts rule and the fact that

$$\int_0^\infty f(t) d_q t = \lim_{n \rightarrow +\infty} \int_0^{q^{-n}} f(t) d_q t \quad (3.11)$$

give the result.  $\square$

In [3], the authors proved the following  $q$ -analogues of the Heisenberg uncertainty principle.

**Theorem 3.7.** Let  $f$  be in  $L_q^2(\mathbb{R}_{q,+})$  such that  $D_q f$  is in  $L_q^2(\mathbb{R}_{q,+})$ . Then,

$$\|t f\|_{2,q} \|\lambda \mathcal{F}_q(f)\|_{2,q} \geq \frac{q}{q^{3/2} + 1} \|f\|_{2,q}^2. \quad (3.12)$$

In addition, if  $\lim_{n \rightarrow +\infty} f(q^n) = 0$ , one has

$$\|t f\|_{2,q} \|\lambda {}_q \mathcal{F}(f)\|_{2,q} \geq \frac{q}{q^{-3/2} + 1} \|f\|_{2,q}^2. \quad (3.13)$$

Now, we are in a position to generalize Theorem 3.7. One obvious way to generalize it is to replace the  $L_q^2$  norms by  $L_q^n$  norms. This is the purpose of the following result.

**Theorem 3.8.** For  $1 \leq n \leq 2$  and  $f \in L_q^2(\mathbb{R}_{q,+})$ , one has

$$\|f\|_{2,q}^2 \leq C'_1 \|x f\|_{n,q} \|\lambda \mathcal{F}_q(f)\|_{n,q}, \quad (3.14)$$

$$\|f\|_{2,q}^2 \leq C'_2 \|x f\|_{n,q} \|\lambda {}_q \mathcal{F}(f)\|_{n,q}, \quad (3.15)$$

where

$$C'_1 = q^{-1+1/n}(1 + q^{-(n+1)/n})C_2, \quad C'_2 = q^{-1}(1 + q^{-(n+1)/n})C_1, \quad (3.16)$$

with  $C_1$  and  $C_2$  being given by (3.8).

*Proof.* The case  $n = 2$  has been dealt with in Theorem 3.7. Now, assume  $1 \leq n < 2$  and let  $m$  be the dual exponent of  $n$ . Let  $f \in S_{*,q}(\mathbb{R}_q)$  such that  $\lim_{t \rightarrow 0} f(t) = 0$ . From the relation

$$D_q(f\bar{f})(t) = D_q f(t)\bar{f}(t) + f(qt)D_q \bar{f}(t), \quad (3.17)$$

the  $q$ -integration by parts rule, and the Hölder inequality, we have, since  $t|f(t)|^2$  tends to 0 as  $t$  tends to  $\infty$  in  $\mathbb{R}_{q,+}$ ,

$$\begin{aligned} \frac{1}{q} \int_0^\infty |f(t)|^2 d_q t &= \left| \int_0^\infty t D_q(f\bar{f})(t) d_q t \right| \\ &\leq \int_0^\infty |t D_q f(t)\bar{f}(t)| d_q t + \int_0^\infty |t f(qt) D_q \bar{f}(t)| d_q t \\ &\leq \left( \int_0^\infty |t\bar{f}(t)|^n d_q t \right)^{1/n} \left( \int_0^\infty |D_q f(t)|^m d_q t \right)^{1/m} \\ &\quad + \left( \int_0^\infty |t f(qt)|^n d_q t \right)^{1/n} \left( \int_0^\infty |D_q \bar{f}(t)|^m d_q t \right)^{1/m}. \end{aligned} \quad (3.18)$$

However, the change of variable  $u = qt$  gives

$$\left( \int_0^\infty |t f(qt)|^n d_q t \right)^{1/n} = q^{-(n+1)/n} \left( \int_0^\infty |t f(t)|^n d_q t \right)^{1/n}. \quad (3.19)$$

So,

$$\frac{1}{q} \int_0^\infty |f(t)|^2 d_q t \leq (1 + q^{-(n+1)/n}) \|t f\|_{n,q} \|D_q(f)\|_{m,q}. \quad (3.20)$$

On the other hand, we have  $D_q(f) = \mathcal{F}_q[\mathcal{F}_q(D_q(f))] = q^{-2} \mathcal{F}_q[\mathcal{F}_q(D_q(f))]$  since  $D_q(f)$  is in  $L_q^2(\mathbb{R}_{q,+})$ . Then, by using Lemma 3.6 and the  $q$ -analogue of the Hausdorff-Young inequality, we obtain

$$\begin{aligned} \|D_q(f)\|_{m,q} &\leq C_1 \|\mathcal{F}_q(D_q(f))\|_{n,q} = \frac{C_1}{q^2} \|\lambda_q \mathcal{F}(f)\|_{n,q}, \\ \|D_q(f)\|_{m,q} &\leq q^{-2} C_2 \|\mathcal{F}_q(D_q(f))\|_{n,q} = q^{-2} C_2 \left\| \frac{\lambda}{q} \mathcal{F}_q(f) \left( \frac{\lambda}{q} \right) \right\|_{n,q} = q^{-2+1/n} C_2 \|\lambda \mathcal{F}_q(f)\|_{n,q}. \end{aligned} \quad (3.21)$$

Thus,

$$\begin{aligned} \|f\|_{2,q}^2 &\leq q^{-1} (1 + q^{-(n+1)/n}) C_1 \|t f\|_{n,q} \|\lambda_q \mathcal{F}(f)\|_{n,q}, \\ \|f\|_{2,q}^2 &\leq q^{-1+1/n} (1 + q^{-(n+1)/n}) C_2 \|t f\|_{n,q} \|\lambda \mathcal{F}_q(f)\|_{n,q}. \end{aligned} \quad (3.22)$$

Now, let  $f \in L_q^2(\mathbb{R}_{q,+})$ ; it is easy to see that for all  $p \in \mathbb{N}$ ,  $f_p = f \chi_{[q^p, q^{p+1}]} \in S_{*,q}(\mathbb{R}_q)$ ,  $\lim_{t \rightarrow 0} f_p(t) = 0$ , and  $(f_p)_p$  converges to  $f$  in  $L_q^2(\mathbb{R}_{q,+})$ . Moreover, if the right-hand side of (3.14) (resp., (3.15)) is finite, then the functions  $t f$  and  $\lambda \mathcal{F}_q(f)$  (resp.,  $\lambda_q \mathcal{F}(f)$ ) are in  $L_q^n(\mathbb{R}_{q,+})$ , and they are limits in  $L_q^n(\mathbb{R}_{q,+})$  (as  $p$  tends to  $\infty$ ) of  $t f_p$  and  $\lambda \mathcal{F}_q(f_p)$  (resp.,  $\lambda_q \mathcal{F}(f_p)$ ), respectively. Finally, the substitution of  $f_p$  in (3.22) and a passage to the limit when  $p$  tends to  $\infty$  complete the proof.  $\square$

#### 4. Local uncertainty principles

In the literature, the first classical local inequalities were obtained by Faris (see [12]) in 1978, and they were generalized by Price (see [13, 14]) in 1983 and 1987. In this section, we will generalize Price's results by giving their  $q$ -analogues.

##### 4.1. Local uncertainty principles for the $q$ -Fourier trigonometric transforms

**Theorem 4.1.** *If  $0 < a < 1/2$ , there is a constant  $K = K(a, q)$  such that for all bounded subset  $E$  of  $\mathbb{R}_{q,+}$  and all  $f \in L_q^2(\mathbb{R}_{q,+})$ , one has*

$$\int_E |\mathcal{F}_q(f)(\lambda)|^2 d_q \lambda \leq K |E|^{2a} \|x^a f\|_{2,q}^2. \quad (4.1)$$

Here,  $|E| = \int_0^\infty \chi_E(x) d_q x$  and  $K = ((\tilde{c}_q / \sqrt{[1-2a]_q})((1-2a)/2a))^{4a} (1/(1-2a)^2)$ , where  $\tilde{c}_q = (1+q^{-1})^{1/2} / \Gamma_{q^2}(1/2)(q; q)_\infty^2$ .

*Proof.* For  $r > 0$ , let  $\chi_r = \chi_{[0,r]}$  be the characteristic function of  $[0, r]$  and  $\tilde{\chi}_r = 1 - \chi_r$ . Then, for  $r > 0$ , we have, since  $f \cdot \chi_r \in L_q^1(\mathbb{R}_{q,+})$ ,

$$\begin{aligned} \left( \int_E |\mathcal{F}_q(f)(\lambda)|^2 d_q \lambda \right)^{1/2} &= \|\mathcal{F}_q(f) \chi_E\|_{2,q} \leq \|\mathcal{F}_q(f \cdot \chi_r) \chi_E\|_{2,q} + \|\mathcal{F}_q(f \cdot \tilde{\chi}_r) \chi_E\|_{2,q} \\ &\leq |E|^{1/2} \|\mathcal{F}_q(f \cdot \chi_r)\|_{\infty,q} + \|\mathcal{F}_q(f \cdot \tilde{\chi}_r)\|_{2,q}, \end{aligned} \quad (4.2)$$

and by the use of the Hölder inequality, we obtain

$$\begin{aligned} \|\mathcal{F}_q(f \cdot \chi_r)\|_{\infty,q} &\leq \tilde{c}_q \|f \cdot \chi_r\|_{1,q} \\ &= \tilde{c}_q \|x^{-a} \chi_r \cdot x^a f\|_{1,q} \leq \tilde{c}_q \|x^{-a} \chi_r\|_{2,q} \|x^a f\|_{2,q} \leq \frac{\tilde{c}_q}{\sqrt{[1-2a]_q}} r^{1/2-a} \|x^a f\|_{2,q}. \end{aligned} \quad (4.3)$$

On the other hand, since  $f \in L_q^2(\mathbb{R}_{q,+})$ , we have  $f \cdot \tilde{\chi}_r \in L_q^2(\mathbb{R}_{q,+})$ , and by the Plancherel formula, we get

$$\|\mathcal{F}_q(f \cdot \tilde{\chi}_r)\|_{2,q} = \|f \cdot \tilde{\chi}_r\|_{2,q} = \|x^{-a} \tilde{\chi}_r \cdot x^a f\|_{2,q} \leq \|x^{-a} \tilde{\chi}_r\|_{\infty,q} \|x^a f\|_{2,q} \leq r^{-a} \|x^a f\|_{2,q}. \quad (4.4)$$

So,

$$\left( \int_E |\mathcal{F}_q(f)(\lambda)|^2 d_q \lambda \right)^{1/2} \leq \left( \frac{\tilde{c}_q}{\sqrt{[1-2a]_q}} |E|^{1/2} r^{1/2-a} + r^{-a} \right) \|x^a f\|_{2,q}. \quad (4.5)$$

The desired result is obtained by minimizing the right-hand side of the previous inequality over  $r > 0$ .  $\square$

**Corollary 4.2.** For  $0 < a < 1/2$  and  $b > 0$ , there is a constant  $K_{a,b}$  such that for all  $f \in L_q^2(\mathbb{R}_{q,+})$ , one has

$$\|f\|_{2,q}^{(a+b)} \leq K_{a,b} \|x^a f\|_{2,q}^b \|\lambda^b \mathcal{F}_q(f)\|_{2,q}^a. \quad (4.6)$$

*Proof.* For  $r > 0$ , put  $E_r = [0, r[ \cap \mathbb{R}_{q,+}$  and  $\tilde{E}_r = [r, +\infty[ \cap \mathbb{R}_{q,+}$ . It is easy to see that  $E_r$  is a bounded subset of  $\mathbb{R}_{q,+}$  and  $|E_r| \leq r$ .

Then, from the Plancherel formula and Theorem 4.1, we have

$$\begin{aligned} \|f\|_{2,q}^2 &= \|\mathcal{F}_q(f)\|_{2,q}^2 \\ &= \int_{E_r} |\mathcal{F}_q(f)|^2(\lambda) d_q \lambda + \int_{\tilde{E}_r} |\mathcal{F}_q(f)|^2(\lambda) d_q \lambda \\ &\leq K r^{2a} \|x^a f\|_{2,q}^2 + r^{-2b} \|\lambda^b \mathcal{F}_q(f)\|_{2,q}^2. \end{aligned} \quad (4.7)$$

Choosing  $r > 0$  so as to minimize the right-hand side of the inequality, we obtain  $\|f\|_{2,q}^2 \leq (K_{a,b} \|x^a f\|_{2,q}^b \|\lambda^b \mathcal{F}_q(f)\|_{2,q}^a)^{2/(a+b)}$ , with  $K_{a,b} = ((a/b)^{b/(a+b)} + (b/a)^{a/(a+b)})^{(a+b)/2} K^{b/2}$ , and  $K$  is the constant given in Theorem 4.1.  $\square$

In the same way, one can prove the following local uncertainty principle for the  $q$ -Fourier-sine transform.

**Theorem 4.3.** If  $0 < a < 1/2$ , there is a constant  $K' = K'(a, q)$  such that for all bounded subset  $E$  of  $\mathbb{R}_{q,+}$  and all  $f \in L_q^2(\mathbb{R}_{q,+})$ , one has

$$\int_E |\mathcal{F}_q(f)(\lambda)|^2 d_q \lambda \leq K' |E|^{2a} \|x^a f\|_{2,q}^2, \quad (4.8)$$

where  $K' = ((\tilde{c}_q / \sqrt{[1 - 2a]_q}) ((1 - 2a) / 2qa))^{4a} [1 + 2qa / (1 - 2a)]^2$ .

**Corollary 4.4.** For  $0 < a < 1/2$  and  $b > 0$ , there is a constant  $K'_{a,b}$  such that for all  $f \in L_q^2(\mathbb{R}_{q,+})$ , one has

$$\|f\|_{2,q}^{(a+b)} \leq K'_{a,b} \|x^a f\|_{2,q}^b \|\lambda^b \mathcal{F}_q(f)\|_{2,q}^a, \quad (4.9)$$

with  $K'_{a,b} = ((a/b)^{b/(a+b)} + (b/a)^{a/(a+b)})^{(a+b)/2} (K')^{b/2} q^{-(a+b)}$ .

*Proof.* The same steps of Corollary 4.2 give the result.  $\square$

**Theorem 4.5.** If  $a > 1/2$ , there is a constant  $K_1 = K_1(a, q)$  such that for all bounded subset  $E$  of  $\mathbb{R}_{q,+}$  and  $f \in L_q^2(\mathbb{R}_{q,+})$ , one has

$$\int_E |\mathcal{F}_q(f)(\lambda)|^2 d_q \lambda \leq K_1 |E| \|f\|_{2,q}^{(2-1/a)} \|x^a f\|_{2,q}^{1/a}, \quad (4.10)$$

$$\int_E |\mathcal{F}_q(f)(\lambda)|^2 d_q \lambda \leq K_1 |E| \|f\|_{2,q}^{(2-1/a)} \|x^a f\|_{2,q}^{1/a}. \quad (4.11)$$

The proof of this result needs the following lemmas.

**Lemma 4.6.** *Suppose  $a > 1/2$ , then for all  $f \in L_q^2(\mathbb{R}_{q,+})$ , such that  $x^a f \in L_q^2(\mathbb{R}_{q,+})$ ,*

$$\|f\|_{1,q}^2 \leq K_2 [\|f\|_{2,q}^2 + \|x^a f\|_{2,q}^2], \quad (4.12)$$

where  $K_2 = K_2(a, q) = (1 - q)((q^{2a}, q^{2a}, -q, -q^{2a-1}; q^{2a})_\infty / (q, q^{2a-1}, -q^{2a}, -1; q^{2a})_\infty)$ .

*Proof.* From [15, Example 1], and the Hölder inequality, we have

$$\|f\|_{1,q}^2 = \left[ \int_0^{+\infty} (1 + x^{2a})^{1/2} |f(x)| (1 + x^{2a})^{-1/2} d_q x \right]^2 \leq K_2 [\|f\|_{2,q}^2 + \|x^a f\|_{2,q}^2], \quad (4.13)$$

where  $K_2 = \int_0^{+\infty} (1 + x^{2a})^{-1} d_q x = (1 - q)((q^{2a}, q^{2a}, -q, -q^{2a-1}; q^{2a})_\infty / (q, q^{2a-1}, -q^{2a}, -1; q^{2a})_\infty)$ .  $\square$

**Lemma 4.7.** *Suppose  $a > 1/2$ , then for all  $f \in L_q^2(\mathbb{R}_{q,+})$ , such that  $x^a f \in L_q^2(\mathbb{R}_{q,+})$ , one has*

$$\|f\|_{1,q} \leq K_3 \|f\|_{2,q}^{(1-1/2a)} \|x^a f\|_{2,q}^{1/2a}, \quad (4.14)$$

where  $K_3 = K_3(a, q) = [2aK_2(2aq - q)^{1/2a-1}]^{1/2}$ .

*Proof.* For  $s \in \mathbb{R}_{q,+}$ , define the function  $f_s$  by  $f_s(x) = f(sx)$ ,  $x \in \mathbb{R}_{q,+}$ .

We have  $\|f_s\|_{1,q} = s^{-1} \|f\|_{1,q}$ ,  $\|x^a f_s\|_{2,q}^2 = s^{-2a-1} \|x^a f\|_{2,q}^2$ .

Replacement of  $f$  by  $f_s$  in Lemma 4.6 gives

$$\|f\|_{1,q}^2 \leq K_2 [s \|f\|_{2,q}^2 + s^{-2a+1} \|x^a f\|_{2,q}^2]. \quad (4.15)$$

Now, for all  $r > 0$ , put  $\alpha(r) = \text{Log}(r)/\text{Log}(q) - E(\text{Log}(r)/\text{Log}(q))$ . We have  $s = (r/q^{\alpha(r)}) \in \mathbb{R}_{q,+}$  and  $r \leq s < r/q$ . Then, for all  $r > 0$ ,

$$\|f\|_{1,q}^2 \leq K_2 \left[ \frac{r}{q} \|f\|_{2,q}^2 + r^{-2a+1} \|x^a f\|_{2,q}^2 \right]. \quad (4.16)$$

The right-hand side of this inequality is minimized by choosing

$$r = (2a - 1)^{1/2a} q^{1/2a} \|f\|_{2,q}^{-1/a} \|x^a f\|_{2,q}^{1/a}. \quad (4.17)$$

When this is done, we obtain the result.  $\square$

*Proof of Theorem 4.5.* Since the proofs of the two statements are similar, it is sufficient to prove (4.11).

Let  $E$  be a bounded subset of  $\mathbb{R}_{q,+}$ . When the right-hand side of the inequality (4.11) is finite, Lemma 4.6 implies that  $f \in L_q^1(\mathbb{R}_{q,+})$ ; so  $\mathcal{F}_q(f)$  is defined and bounded on  $\mathbb{R}_{q,+}$ . Using

Proposition 3.1, Lemma 4.7, and the fact that

$$\int_E |\mathcal{F}_q(f)(\lambda)|^2 d_q \lambda \leq |E| \|\mathcal{F}_q(f)\|_{\infty, q}^2, \quad (4.18)$$

we obtain the result with  $K_1 = ((1 + q^{-1})/\Gamma_q^2(1/2)(q; q)_\infty^4)K_3^2$ .  $\square$

*Remark 4.8.* By the same technique as in the proof of Corollary 4.2, we can show that Theorem 4.5 leads to inequalities (4.6) and (4.9) with some different constants.

#### 4.2. Local uncertainty principles for the $q$ -Bessel-Fourier transform

The  $q$ -Bessel-Fourier transform is defined (see [16]) for  $f \in L_q^1(\mathbb{R}_{q,+}, x^{2\nu+1} d_q x)$  by

$$\mathcal{F}_{\nu, q}(f)(\lambda) = c_{\nu, q} \int_0^\infty f(x) j_\nu(\lambda x; q^2) x^{2\nu+1} d_q x, \quad (4.19)$$

where

$$j_\nu(z; q^2) = (1 - q^2)^\nu \Gamma_{q^2}(\nu + 1) ((1 - q)q^{-1}z)^{-\nu} J_\nu((1 - q)q^{-1}z; q^2) \quad (4.20)$$

is the normalized third Jackson  $q$ -Bessel function, and

$$c_{\nu, q} = \frac{(1 + q^{-1})^{-\nu}}{\Gamma_{q^2}(\nu + 1)}. \quad (4.21)$$

It was shown in [10] that for  $\nu \geq -1/2$ , we have the following result.

**Theorem 4.9.** (1) For  $f \in L_q^1(\mathbb{R}_{q,+}, x^{2\nu+1} d_q x)$ , one has  $\mathcal{F}_{\nu, q}(f) \in L_q^\infty(\mathbb{R}_{q,+})$  and

$$\|\mathcal{F}_{\nu, q}(f)\|_{\infty, q} \leq \frac{c_{\nu, q}}{(q; q^2)_\infty^2} \|f\|_{1, \nu, q}. \quad (4.22)$$

(2)  $\mathcal{F}_{\nu, q}$  is an isomorphism of  $L_q^2(\mathbb{R}_{q,+}, x^{2\nu+1} d_q x)$  onto itself,  $\mathcal{F}_{\nu, q}^{-1} = q^{4\nu+2} \mathcal{F}_{\nu, q}$ , and one has the following Plancherel formula:

$$\forall f \in L_q^2(\mathbb{R}_{q,+}, x^{2\nu+1} d_q x), \quad \|\mathcal{F}_{\nu, q} f\|_{2, \nu, q} = q^{2\nu+1} \|f\|_{2, \nu, q}. \quad (4.23)$$

The following result states a local uncertainty principle for the  $q$ -Bessel-Fourier transform.

**Theorem 4.10.** For  $\nu \geq -1/2$  and  $0 < a < \nu + 1$ , there is a constant  $K_{a, \nu} = K(a, \nu, q)$  such that for all  $f \in L_q^2(\mathbb{R}_{q,+}, x^{2\nu+1} d_q x)$  and all bounded subset  $E$  of  $\mathbb{R}_{q,+}$ , one has

$$\int_E |\mathcal{F}_{\nu, q}(f)(\lambda)|^2 \lambda^{2\nu+1} d_q \lambda \leq K_{a, \nu} |E|_\nu^{a/(v+1)} \|x^a f\|_{2, \nu, q}^2. \quad (4.24)$$

Here,  $|E|_\nu = \int_0^\infty \chi_E(x) x^{2\nu+1} d_q x$ ,  $\tilde{c}_{\nu, q} = c_{\nu, q}/(q; q^2)_\infty^2$ , and

$$K_{a, \nu} = \left( \frac{\tilde{c}_{\nu, q}}{\sqrt{[2\nu + 2 - 2a]_q}} \right)^{2a/(v+1)} \left[ \left( \frac{aq^{2\nu+1}}{\nu + 1 - a} \right)^{1-a/(v+1)} + q^{2\nu+1} \left( \frac{aq^{2\nu+1}}{\nu + 1 - a} \right)^{-a/(v+1)} \right]^2. \quad (4.25)$$

*Proof.* Let  $\nu \geq -1/2$ ,  $0 < a < \nu + 1$ ,  $f \in L_q^2(\mathbb{R}_{q,+}, x^{2\nu+1} d_q x)$ , and let  $E$  be a bounded subset of  $\mathbb{R}_{q,+}$ . For  $r > 0$ , we have, since  $f \cdot \chi_r \in L_q^1(\mathbb{R}_{q,+}, x^{2\nu+1} d_q x)$ ,

$$\begin{aligned} \left( \int_E |\mathcal{F}_{\nu,q}(f)(\lambda)|^2 \lambda^{2\nu+1} d_q \lambda \right)^{1/2} &= \|\mathcal{F}_{\nu,q}(f) \chi_E\|_{2,\nu,q} \\ &\leq \|\mathcal{F}_{\nu,q}(f \cdot \chi_r) \chi_E\|_{2,\nu,q} + \|\mathcal{F}_{\nu,q}(f \cdot \tilde{\chi}_r) \chi_E\|_{2,\nu,q} \\ &\leq |E|_q^{1/2} \|\mathcal{F}_{\nu,q}(f \cdot \chi_r)\|_{\infty,q} + \|\mathcal{F}_{\nu,q}(f \cdot \tilde{\chi}_r)\|_{2,\nu,q}. \end{aligned} \quad (4.26)$$

However, by the use of the Hölder inequality, we obtain

$$\begin{aligned} \|\mathcal{F}_{\nu,q}(f \cdot \chi_r)\|_{\infty,q} &\leq \tilde{c}_{\nu,q} \|f \cdot \chi_r\|_{1,q} \\ &= \tilde{c}_q \|x^{-a} \chi_r \cdot x^a f\|_{1,\nu,q} \\ &\leq \tilde{c}_{\nu,q} \|x^{-a} \chi_r\|_{2,\nu,q} \|x^a f\|_{2,\nu,q}. \end{aligned} \quad (4.27)$$

Now, if  $k$  is the integer such that  $q^k \leq r < q^{k+1}$ , we get, since  $a < \nu + 1$ ,

$$\|x^{-a} \chi_r\|_{2,\nu,q}^2 = \int_0^\infty x^{-2a} \chi_r(x) x^{2\nu+1} d_q x = \int_0^{q^k} x^{2\nu+1-2a} d_q x = \frac{q^{2k(\nu+1-a)}}{[2\nu+2-2a]_q} \leq \frac{r^{2(\nu+1-a)}}{[2\nu+2-2a]_q}. \quad (4.28)$$

Then,

$$\|\mathcal{F}_{\nu,q}(f \cdot \chi_r)\|_{\infty,q} \leq \frac{\tilde{c}_{\nu,q}}{\sqrt{[2\nu+2-2a]_q}} r^{(\nu+1-a)} \|x^a f\|_{2,\nu,q}. \quad (4.29)$$

On the other hand, since  $f \in L_q^2(\mathbb{R}_{q,+}, x^{2\nu+1} d_q x)$ , we have  $f \cdot \tilde{\chi}_r \in L_q^2(\mathbb{R}_{q,+}, x^{2\nu+1} d_q x)$ , and by the Plancherel formula (4.23), we obtain

$$\begin{aligned} \|\mathcal{F}_{\nu,q}(f \cdot \tilde{\chi}_r)\|_{2,\nu,q} &= q^{2\nu+1} \|f \cdot \tilde{\chi}_r\|_{2,\nu,q} = q^{2\nu+1} \|x^{-a} \tilde{\chi}_r \cdot x^a f\|_{2,\nu,q} \\ &\leq q^{2\nu+1} \|x^{-a} \tilde{\chi}_r\|_{\infty,q} \|x^a f\|_{2,\nu,q} \leq q^{2\nu+1} r^{-a} \|x^a f\|_{2,\nu,q}. \end{aligned} \quad (4.30)$$

So,

$$\left( \int_E |\mathcal{F}_{\nu,q}(f)(\lambda)|^2 \lambda^{2\nu+1} d_q \lambda \right)^{1/2} \leq \left( \frac{\tilde{c}_{\nu,q}}{\sqrt{[2\nu+2-2a]_q}} |E|_q^{1/2} r^{(\nu+1-a)} + q^{2\nu+1} r^{-a} \right) \|x^a f\|_{2,\nu,q}. \quad (4.31)$$

By minimization of the right-hand side of the previous inequality over  $r > 0$  and by easy computation, we obtain the desired result.  $\square$

**Theorem 4.11.** For  $\nu \geq -1/2$  and  $a > \nu + 1$ , there exists a constant  $K'_{a,\nu}$  such that for all bounded subset  $E$  of  $\mathbb{R}_{q,+}$  and all  $f$  in  $L^2_q(\mathbb{R}_{q,+}, x^{2\nu+1}d_qx)$ , one has

$$\int_E |\mathcal{F}_{\nu,q}(f)(\lambda)|^2 \lambda^{2\nu+1} d_q\lambda \leq K'_{a,\nu} |E| \|f\|_{2,\nu,q}^{2(1-(\nu+1)/a)} \|x^a f\|_{2,\nu,q}^{2((\nu+1)/a)}. \tag{4.32}$$

*Proof.* Since  $a > \nu + 1$ , the same steps as in the proof of Theorem 4.5 and the relation (4.22) give the result with

$$K'_{a,\nu} = \frac{(q^{2a}, q^{2a}, -q^{2\nu+2}, -q^{2(a-\nu-1)}; q^{2a})_\infty c'_{\nu,q'}}{(q^{2\nu+2}, q^{2(a-\nu-1)}, -q^{2a}, -1; q^{2a})_\infty} c'_{\nu,q'}$$

$$c'_{\nu,q} = (1 - q) \left( \frac{c_{\nu,q}}{(q; q^2)_\infty} \right)^2 \left( \frac{a}{\nu + 1} - 1 \right)^{(\nu+1)/a} \left( \frac{a}{a - \nu - 1} \right) q^{-2(\nu+1)((a-\nu-1)/a)}. \tag{4.33}$$

□

**Corollary 4.12.** For  $\nu \geq -1/2$  and  $a, b > 0$ , there is a constant  $K_{a,b,\nu} = K(a, b, \nu, q)$  such that for all  $f \in L^2_q(\mathbb{R}_{q,+}, x^{2\nu+1}d_qx)$ , one has

$$\|f\|_{2,\nu,q}^{(a+b)} \leq K_{a,b,\nu} \|x^a f\|_{2,\nu,q}^b \|\lambda^b \mathcal{F}_{\nu,q}(f)\|_{2,\nu,q}^a, \tag{4.34}$$

with

$$K_{a,b,\nu} = \begin{cases} \left[ \left( \frac{b}{a} \right)^{a/(a+b)} + \left( \frac{a}{b} \right)^{b/(a+b)} \right]^{(a+b)/2} (K_{a,\nu})^{b/2} \frac{q^{-(2\nu+1)(a+b)}}{([2\nu + 2]_q)^{ab/2(\nu+1)}} & \text{if } a < \nu + 1, \\ \left( \frac{K'_{a,\nu}}{[2\nu + 2]_q} \right)^{ab/(2\nu+2)} \left( q^{-(4\nu+2)} \left[ \left( \frac{b}{\nu + 1} \right)^{(\nu+1)/(\nu+b+1)} + \left( \frac{b}{\nu + 1} \right)^{-b/(\nu+b+1)} \right] \right)^{a(\nu+b+1)/2(\nu+1)} & \text{if } a > \nu + 1, \end{cases} \tag{4.35}$$

where  $K_{a,\nu}$  (resp.,  $K'_{a,\nu}$ ) is the constant given in Theorem 4.10 (resp., Theorem 4.11).

*Proof.* For  $r > 0$ , we put  $E_r = [0, r[ \cap \mathbb{R}_{q,+}$  and  $\tilde{E}_r = [r, +\infty[ \cap \mathbb{R}_{q,+}$ .

We have  $E_r$  is a bounded subset of  $\mathbb{R}_{q,+}$  and  $|E_r|_\nu \leq r^{2\nu+2}/[2\nu + 2]_q$ . Then, the Plancherel formula (4.23) and Theorems 4.10 and 4.11 lead to

$$q^{4\nu+2} \|f\|_{2,\nu,q}^2 = \|\mathcal{F}_{\nu,q}(f)\|_{2,\nu,q}^2 = \int_{E_r} |\mathcal{F}_{\nu,q}(f)|^2(\lambda) \lambda^{2\nu+1} d_q\lambda + \int_{\tilde{E}_r} |\mathcal{F}_{\nu,q}(f)|^2(\lambda) \lambda^{2\nu+1} d_q\lambda$$

$$\leq \begin{cases} K_{a,\nu} |E_r|_\nu^{a/(\nu+1)} \|x^a f\|_{2,\nu,q}^2 + r^{-2b} \|\lambda^b \mathcal{F}_{\nu,q}(f)\|_{2,\nu,q}^2 & \text{if } a < \nu + 1, \\ K'_{a,\nu} |E_r| \|f\|_{2,\nu,q}^{2(a-\nu-1)/a} \|x^a f\|_{2,\nu,q}^{2(\nu+1)/a} + r^{-2b} \|\lambda^b \mathcal{F}_{\nu,q}(f)\|_{2,\nu,q}^2 & \text{if } a > \nu + 1, \end{cases}$$

$$\leq \begin{cases} \frac{K_{a,\nu}}{[2\nu + 2]_q^{a/(\nu+1)}} r^{2a} \|x^a f\|_{2,\nu,q}^2 + r^{-2b} \|\lambda^b \mathcal{F}_{\nu,q}(f)\|_{2,\nu,q}^2 & \text{if } a < \nu + 1, \\ K'_{a,\nu} \frac{r^{2\nu+2}}{[2\nu + 2]_q} \|f\|_{2,\nu,q}^{2(a-\nu-1)/a} \|x^a f\|_{2,\nu,q}^{2(\nu+1)/a} + r^{-2b} \|\lambda^b \mathcal{F}_{\nu,q}(f)\|_{2,\nu,q}^2 & \text{if } a > \nu + 1. \end{cases} \tag{4.36}$$

The desired result follows by minimizing the right expressions over  $r > 0$ . □

Remark that when  $a = b = 1$ , we obtain a Heisenberg-Weyl-type inequality for the  $q$ -Bessel-Fourier transform.

**Corollary 4.13.** For  $\nu \geq -1/2$ ,  $\nu \neq 0$ , one has for all  $f \in L_q^2(\mathbb{R}_{q,+}, x^{2\nu+1}d_qx)$ ,

$$\|f\|_{2,\nu,q}^2 \leq K_{1,1,\nu} \|xf\|_{2,\nu,q} \|\lambda \mathcal{F}_{\nu,q}(f)\|_{2,\nu,q}. \quad (4.37)$$

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