A CHARACTERIZATION OF THE ROGERS q-HERMITE POLYNOMIALS

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(Received October 14, 1993 and in revised form February 18, 1994)

ABSTRACT. In this paper we characterize the Rogers q-Hermite polynomials as the only orthogonal polynomial set which is also \mathcal{D}_q -Appell where \mathcal{D}_q is the Askey-Wilson finite difference operator.

KEY WORDS AND PHRASES. Orthogonal polynomials, generating functions, Askey-Wilson operator 1991 AMS SUBJECT CLASSIFICATION CODES. Primary 33D45, 33D05; Secondary 42A65

1. INTRODUCTION

Appell polynomials sets $\{P_n(x)\}$ are generated by the relation

$$A(t)e^{xt} = \sum_{n=0}^{\infty} P_n(x) t^n,$$
 (1.1)

where A(t) is a formal power series in t with A(0) = 1. This definition implies the equivalent property that

$$DP_n(x) = P_{n-1}(x), \qquad D = d/dx,$$
 (1.2)

Examples of such polynomial sets are

$$\left\{\frac{x^n}{n!}\right\}, \left\{\frac{B_n(x)}{n!}\right\}, \left\{\frac{H_n(x)}{2^n n!}\right\}$$
(1.3)

where $B_n(x)$ is the nth Bernoulli polynomial and $H_n(x)$ is the nth Hermite polynomials generated by

$$e^{2xt-t^2} = \sum_{n=0}^{\infty} H_n(x) \frac{t^n}{n!}.$$
 (1.4)

By an orthogonal polynomial set (OPS) we shall mean those polynomial sets which satisfy a three term recurrence relation of the form

$$P_{n+1}(x) = (A_n x + B_n) P_n(x) - C_n P_{n-1}(x), \qquad (n = 0, 1, 2, \cdots)$$
(1.5)

with $P_0(x) = 1$, $P_{-1}(x) = 0$, and $A_n A_{n-1} C_n > 0$.

By Favard's theorem [7] this is equivalent to the existence of a positive measure $d\alpha(x)$ such that

$$\int_{-\infty}^{\infty} P_n(x) P_m(x) \, d\alpha(x) = K_n \delta_{nm}. \tag{1.6}$$

As we see from the examples (1.3) some Appell polynomials are orthogonal and some are not. This prompted Angelesco [3] to prove that the only orthogonal polynomial sets which are also Appell is the Hermite polynomial set. This theorem was rediscovered by several authors later on (see, e.g., [10]).

There were several extensions and/or analogs of Appell polynomials that were introduced later. Some are based on changing the operator D in (1.2) into another differentiation-like operator or by replacing the generating relation (1.1) by a more general one. In most of these cases theorems like Angelesco's were given. For example Carlitz [6] proved that the Charlier polynomials are the only OPS which satisfy the difference relation

$$\Delta P_n(x) = P_{n-1}(x), \qquad (\Delta f(x) = f(x+1) - f(x)). \qquad (1.7)$$

See [1] for many other references.

A new and very interesting analog of Appell polynomials were introduced recently, as a biproduct of other considerations, by Ismail and Zhang [9]. In discussing the Askey-Wilson operator they defined a new q-analog of the exponential function e^{xt} . This we describe in the next section.

2. NOTATIONS AND DEFINITIONS

The Askey-Wilson operator is defined by

$$\mathcal{D}_q f(x) = \frac{\delta_q f(x)}{\delta_q x},\tag{2.1}$$

where $x = \cos \theta$ and

$$\delta_q g(e^{i\theta}) = g(q^{1/2} e^{i\theta}) - g(q^{-1/2} e^{i\theta}).$$
 (2.2)

We further assume that -1 < q < 1 and use the notation

$$(a;q)_0 = 1, \qquad (a;q)_n = (1-a)(1-qa)\cdots(1-aq^{n-1}), \quad (n=1,2,..)$$
 (2.3)

$$(a;q)_{\infty} = \prod_{k=0}^{\infty} (1-aq^{j}).$$
 (2.4)

There are two q-analogs of the exponential function e^x given by the infinite products

$$e_q(x) = \frac{1}{(x;q)_{\infty}} = \sum_{k=0}^{\infty} \frac{x^k}{(q;q)_k},$$
(2.5)

and

$$\frac{1}{e_q(x)} = (x;q)_{\infty} = \sum_{k=0}^{\infty} (-1)^k \frac{q^{\frac{1}{2}k(k-1)}}{(q;q)_k} x^k.$$
(2.6)

We shall also use the function

$$\Psi_n(x) = i^n (iq^{(1-n)/2} e^{i\theta}; q)_n (iq^{(1-n)/2} e^{-i\theta}; q)_n, \qquad (2.7)$$

so that

$$\Psi_{2n}(x) = \prod_{k=0}^{n-1} \left[4x^2 + (1 - q^{2n-1-2k})(1 - q^{1-2n+2k}) \right]$$

$$\Psi_{2n+1}(x) = 2x \prod_{k=0}^{n-1} \left[4x^2 - (1 - q^{2n-2k})(1 - q^{-2n+2k}) \right]$$

$$4x^2 \Psi_n(x) = \Psi_{n+2}(x) + (1 - q^{n+1})(1 - q^{-n-1})\Psi_n(x)$$
(2.8)

Thus

$$\mathcal{D}_{q}\Psi_{n}(x) = 2q^{(1-n)/2}\frac{1-q^{n}}{1-q} \Psi_{n-1}(x).$$
(2.9)

and

$$\mathcal{D}_q\left[x \ \Psi_n(x)\right] = \frac{q^{(1+n)/2} - q^{-(n+1)/2}}{q^{1/2} - q^{-1/2}} \ 2x \ \Psi_{n-1}(x). \tag{2.10}$$

Iterating (2.9) we get

$$\mathcal{D}_{q}^{k}\Psi_{n}(x) = 2^{k}q^{\frac{1}{4}k(k+1)-\frac{1}{2}nk}\frac{(q;q)_{n}}{(q;q)_{n-k}(1-q)^{k}}\Psi_{n-k}(x).$$
(2.11)

The Ismail-Zhang q-analog of the exponential function [9] is

$$\mathcal{E}(x) = \sum_{n=0}^{\infty} \frac{q^{n(n-1)/4}(1-q)^n}{2^n(q;q)_n} \Psi_n(x) t^n.$$
(2.12)

It follows from (2.12) and (2.9) that

$$\mathcal{D}_q \mathcal{E}(x) = t \ \mathcal{E}(x). \tag{2.13}$$

This suggested to Ismail and Zhang to define the \mathcal{D}_q -Appell polynomials as those, in analogy with (1.1), defined by

$$A(t)\mathcal{E}(x) = \sum_{n=0}^{\infty} P_n(x) t^n, \qquad (2.14)$$

so that

$$\mathcal{D}_q P_n(x) = P_{n-1}(x).$$
 (2.15)

An example of such a set is the Rogers q-Hermite polynomials, $\{H_n(x|q)\}$, (see [2, 4, 8]).

$$\prod_{n=0}^{\infty} \left(1 - 2xtq^n + t^2 q^{2n} \right)^{-1} = \sum_{n=0}^{\infty} H_n(x|q) \frac{t^n}{(q;q)_n}.$$
(2.16)

They satisfy the three term recurrence relation

$$H_{n+1}(x|q) = 2xH_n(x|q) - (1-q^n)H_{n-1}(x|q), \qquad n = 0, 1, 2, 3, \dots$$
(2.17)

with $H_0(x|q) = 1$, $H_{-1}(x|q) = 0$.

3. THE MAIN RESULT

We now state our main result:

Theorem 1. The orthogonal polynomial sets which are also \mathcal{D}_q -Appell, i.e., satisfy (2.15) or (2.14) is the set of the Rogers q-Hermite polynomials.

Proof Let $\{Q_n(x)\}$ be a polynomial set which is both orthogonal and \mathcal{D}_q -Appell. That is $\{Q_n(x)\}$ satisfy (2.14) and (1.5).

We next note that (2.16) implies that

$$h_n(x|q) = \frac{(1-q)^n q^{n(n-1)/4}}{2^n(q;q)_n} H_n(x|q)$$
(3.1)

satisfy

$$\mathcal{D}_q h_n(x|q) = h_{n-1}(x|q), \tag{3.2}$$

so that $\{h_n(x|q)\}$ is a \mathcal{D}_q -Appell polynomial set and at the same time is an OPS satisfying the three term recurrence relation

$$(1-q^{n+1})h_{n+1}(x|q) = (1-q)q^{n/2}xh_n(x|q) - \frac{1}{4}(1-q)^2q^{n-1/2}h_{n-1}(x|q)$$
(3.3)

It also follows from (2.14) that any two polynomial sets $\{R_n(x)\}$ and $\{S_n(x)\}$, in that class are related by $R_n(x) = \sum_{k=0}^n c_{n-k} S_k(x)$. Thus the solution to our problem may be expressed as

$$Q_n(x) = \sum_{k=0}^n a_{n-k} h_k(x|q).$$
 (3.4)

for some sequence of real constants $\{a_n\}$. We may assume without loss of generality that $a_0 = 1$. The three term recurrence relation satisfied by $\{Q_n(x)\}$ is

$$(1-q^{n+1})Q_{n+1}(x) = \left((1-q)q^{n/2}x + \beta_n\right)Q_n(x) - \gamma_n Q_{n-1}(x), \tag{3.5}$$

with $Q_0(x) = 1$, $Q_{-1}(x) = 0$. Thus $Q_1(x) = x + \beta_0 = a_1 + h_1(x|q)$, from which it follows that $a_1 = \beta_0$.

Putting (3.4) in (3.5) and using (3.3) to replace $xh_k(x|q)$ in terms of $h_{k+1}(x|q)$ and $h_{k-1}(x|q)$ we get, on equating coefficients of $h_k(x|q)$,

$$(1-q^{(n-k+1)/2})(1+q^{(n+1+k)/2})a_{n+1-k}-\beta_n a_{n-k}+\left[\gamma_n-\frac{1}{4}(1-q)^2q^{(n+k)/2}\right]a_{n-k-1}=0,\quad(3.6)$$

valid for all n and k = 0, 1, 2, ..., n + 1 provided we interpret $a_{-1} = a_{-2} = 0$. It is easy to see that this system of equations is equivalent to the solution of our problem.

Putting k = n in (3.6) we get

$$\beta_n = (1 - q^{\frac{1}{2}})(1 + q^{n + \frac{1}{2}})a_1. \tag{3.7}$$

Hence if $\beta_0 = 0$ then $\beta_n = 0$ for all n. In fact if $\beta_m = 0$ for any n = m then $\beta_n = 0$ for all n.

Now we treat these two cases seperately.

Case I. $(\beta_0 = 0)$.

The system (3.6) can now be written as

$$(1-q^{(k+1)/2})(1+q^{n+\frac{1}{2}(1-k)})a_{k+1}+\left[\gamma_n-\frac{1}{4}(1-q)^2q^{n-\frac{1}{2}k}\right]a_{k-1}=0.$$
(3.8)

Since $a_1 = 0$ then it follows from (3.8) that $a_{2k+1} = 0$ for all k. In particular we get

$$\gamma_n = \frac{1}{4}(1-q)^2 q^{n-\frac{1}{2}} - a_2(1-q)(1+q^n), \qquad (3.9)$$

so that if $a_2 = 0$ then

$$Q_n(x) = h_n(x|q). \tag{3.10}$$

Now we show that $a_2 \neq 0$ leads to contradiction. To do this replace k by 2k - 1. We get

$$(1-q^k)(1+q^{n-k+1})a_{2k} + \left[\frac{1}{4}(1-q)^2q^{n-\frac{1}{2}}(1-q^{1-k}) - a_2(1-q)(1+q^n)\right]a_{2k-2} = 0.$$
(3.11)

Keep k fixed and let $n \to \infty$. We get $(1-q^k)a_{2k} = (1-q)a_2a_{2k-2}$. Thus

$$a_{2k} = \frac{(1-q)^k}{(q;q)_k} a_2^k. \tag{3.12}$$

Putting this value in (3.11) we get $q^{1-k} = 1$. This is a contradiction and Case I is finished. Case II $(\beta_0 \neq 0)$. We start with (3.6) we get, assuming $a_1 \neq 0$,

$$\gamma_n = \frac{1}{4} (1-q)^2 q^{n-\frac{1}{2}} + (1-q^{\frac{1}{2}})(1+q^{n+\frac{1}{2}})a_1^2 - (1-q)(1+q^n)a_2.$$
(3.13)

Putting this value of γ_n and the value of β_n in (3.7) in (3.6), and finally equating coefficients of q^n and the terms independent of n we get the pair of equation systems

$$\left(1-q^{(k+1)/2}\right)a_{k+1}-\left(1-q^{\frac{1}{2}}\right)a_{1}a_{k}+\left\{\left(1-q^{\frac{1}{2}}\right)a_{1}^{2}-\left(1-q\right)a_{2}\right\}a_{k-1}=0$$
(3.14)

and

$$(1 - q^{(k+1)/2})a_{k+1} - (1 - q^{\frac{1}{2}})q^{k/2}a_1a_k +$$

$$\left\{\frac{1}{4}(1 - q)^2 q^{-\frac{1}{2}}(q^{(k-1)/2} - 1) + q^{k/2}(1 - q^{\frac{1}{2}})a_1^2 - (1 - q)q^{(k-1)/2}a_2\right\}a_{k-1} = 0$$
(3.15)

Eliminating a_{k+1} in these equations we get

$$(1-q^{\frac{1}{2}})(1-q^{k/2})a_1a_k + \left\{(1-q)a_2(1-q^{(k-1)/2}) - (1-q^{\frac{1}{2}})(1-q^{k/2})a_1^2 - \frac{1}{4}(1-q)^2q^{-\frac{1}{2}}(1-q^{(k-1)/2})\right\}a_{k-1} = 0.$$
(3.16)

This equation is of the form $(1 - q^{k/2})a_1a_k = c(1 - bq^{k/2})a_{k-1}$ so that the general solution of (3.16) is

$$a_{k} = c^{k} \frac{(bq^{\frac{1}{2}}; q^{\frac{1}{2}})_{k}}{(q^{\frac{1}{2}}; q^{\frac{1}{2}})_{k}}$$
(3.17)

Putting this in (3.14) we get that b = 0. On the other hand (3.15) gives that $c^2 = \frac{1}{4}(1-q)^2q^{-\frac{1}{2}}$. Finally putting those values of a_k in (3.13) we get that $\gamma_n = 0$ which is a contradiction.

This completes the proof of the theorem.

4. GENERATING FUNCTION

We obtain, for the q-Hermite polynomials, a generating function of the form (2.14). More specifically we prove

Theorem 2. Let $H_n(x|q)$ be the nth Rogers q-Hermite polynomial. Then we have

$$\sum_{n=0}^{\infty} \frac{q^{n(n-1)/4}}{(q;q)_n} H_n(x|q) t^n = (t^2 q^{-\frac{1}{2}};q^2)_{\infty} \mathcal{E}(x).$$
(4.1)

Proof. Let $A(t) = 1 + a_1t + a_2t^2 + a_3t^3 + \cdots$ and

$$A(t)\mathcal{E}(x) = \sum_{n=0}^{\infty} h_n(x|q)t^n.$$
(4.2)

Then we get

$$h_n(x|q) = \sum_{k=0}^n a_{n-k} c_k \Psi_k(x).$$
 (4.3)

where

$$c_{k} = \frac{(1-q)^{k}}{2^{k}(q;q)_{k}} q^{k(k-1)/4}.$$
(4.4)

To calculate the coefficients $\{a_n\}$ we first iterate (3.3) we get

$$4x^{2}h_{n}(x|q) = \frac{4}{(1-q)^{2}}(1-q^{n+1})(1-q^{n+2})q^{-n-\frac{1}{2}}h_{n+2}(x|q)$$

$$+ (2-q^{n}-q^{n+1})h_{n}(x|q) + \frac{(1-q)^{2}}{4}q^{n-\frac{3}{2}}h_{n-2}(x|q).$$
(4.5)

Putting (4.3) in (4.5), using (2.6) and then equating coefficients of $\Psi_k(x)$ we get after some simplification

$$\frac{4}{(1-q)^2}q^{-n-\frac{1}{2}}(1-q^{n-k+2})(1-q^{n+k+1})a_{n+2-k}+$$

$$q^{-k-1}\left\{1+q^{2k+2}-q^{n+k+1}-q^{n+k+2}\right\}a_{n-k}+$$

$$\frac{(1-q)^2}{4}q^{n-\frac{3}{2}}a_{n-2-k}=0 \qquad (k=0,1,...,n+2).$$
(4.6)

By direct calculation of a_1 , a_2 , a_3 we see easily that $a_1 = a_3 = 0$. Thus (4.6) shows that $a_{2k+1} = 0$ for all k.

Furthermore we can easily verify that

$$a_{2j} = (-1)^j \frac{(1-q)^{2j}}{2^{2j}(q^2;q^2)_j} q^{j(j-\frac{3}{2})} \qquad (j=0,1,2,3,\dots$$
(4.7)

Hence

$$A(t) = \sum_{j=0}^{\infty} (-1)^{j} \frac{q^{j(j-1)}}{(q^{2};q^{2})_{j}} \left(\frac{(1-q)^{2}t^{2}}{4}q^{-\frac{1}{2}}\right)^{j}$$

$$= \left(\frac{(1-q)^{2}}{4}t^{2}q^{-\frac{1}{2}};q^{2}\right)_{\infty}.$$
(4.8)

After some rescaling we get the theorem.

As a corollary of (4.1) we state the pair of inverse relations

$$\Psi_n(x) = \sum_k \frac{(q;q)_n q^{k(k-n)}}{(q^2;q^2)_k (q;q)_{n-2k}} H_{n-2k}(x|q), \qquad (4.9)$$

$$H_n(x|q) = \sum_k (-1)^k \frac{(q;q)_n q^{k(2k-n-1)}}{(q^2;q^2)_k(q;q)_{n-2k}} \Psi_{n-2k}(x).$$
(4.10)

These follows from the identities (2.5) and (2.6)

Formula (4.10) and (2.11) give

$$H_n(x|q) = \frac{1}{e_{q^2}(\frac{(1-q)^2}{4}q^{-\frac{1}{2}}\mathcal{D}_q^2)}\Psi_n(x).$$
(4.11)

This is a q-analog of the formula

$$e^{-D^2}x^n = H_n(x)$$

for the regular Hermite polynomials (1.4).

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Email: waleed@euler.math.ualberta.ca Research was partially supported by NSERC (Canada) grant A2975