

NONLINEAR DIFFERENTIAL EQUATIONS AND ALGEBRAIC SYSTEMS

LLOYD K. WILLIAMS

Department of Mathematics
Texas Southern University
Houston, Texas 77004
U.S.A.

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ABSTRACT. In this paper we obtain the general solution of scalar, first-order differential equations. The method is variation of parameters with asymptotic series and the theory of partial differential equations.

The result gives us a form like a differential quotient requiring only that a limit be taken. Like the familiar expression for the solution of linear, first order, ordinary equations, it is the same in all cases.

KEY WORDS AND PHRASES. *Riccati Equations, Abel Equations, Cauchy-Kowalewski Theorem, Cauchy-Kowalewski System, Universal Cauchy-Kowalewski System.*

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1. INTRODUCTION.

We present a unified treatment for the general scalar, first-order, ordinary differential equation

$$y' = G(x,y), G \in C^I.$$

Particular examples are linear equations, Riccati equations and Abel equations.

2. PRELIMINARIES.

We begin with the differential system

$$\left. \begin{array}{l} V_1' = f(V_1, V_2) = -V_1 V_2 \\ V_2' = h(V_1, V_2) = V_1 - V_2 \\ V \neq 0 \end{array} \right\} \quad (2.1)$$

with general solution $V_1 = V_1(x, c_1, c_2)$, $V_2 = V_2(x, c_1, c_2)$. Here c_1, c_2 are arbitrary constants.

Now let $x = x(t)$. Then we get

$$\left. \begin{array}{l} \dot{V}_1 = U_1 \dot{x} \quad \cdot = \frac{d}{dt} \\ V_2 = U_2 \dot{x} \\ U_1 = f(V_1, V_2), U_2 = h(V_1, V_2) \\ V_1 \neq 0 \end{array} \right\} \quad (2.2)$$

We are now ready to present the algebraic system referred to in the title.

3. THE CAUCHY-KOWALEWSKI SYSTEM.

Let $w_1 = w_1(t, \epsilon)$, $w_2 = w_2(t, \epsilon)$ be two functions of t and ϵ (at present unknown).

The functions V_1, V_2 have been given by (2.1). Finally two more unknown functions $K(w_1, w_2, t, \epsilon)$ and $L(w_1, w_2, t, \epsilon)$ will be defined by partial differential equations later. They will contain another variable, λ . It will be possible to substitute an arbitrary $G(w_1, t)$ for λ to solve specific equations.

DEFINITION. The system of algebraic equations

$$\left. \begin{array}{l} \text{(a) } w_1 - K(w_1, w_2, t, \epsilon) V_1 = 0 \\ \text{(b) } w_2^2 - L(w_1, w_2, t, \epsilon) - V_2 = 0 \\ \text{(c) } x = w_1 + tw_2. \\ \qquad w_1 \neq 0 \end{array} \right\} \quad (3.1)$$

is called the Cauchy-Kowalewski system, for a specific $G(w_1, t)$. Using λ we will get a universal system.

Under suitable conditions on the functions K and L , we can solve it for $w_1 = w_1(t, \epsilon)$ and $w_2 = w_2(t, \epsilon)$. We proceed by defining these functions as solutions of appropriate partial differential equations. We will derive these functions $L(w_1, w_2, t, \epsilon, \lambda)$ and $K(w_1, w_2, t, \epsilon, \lambda)$ and regard them as fixed like universal constants.

4. THE FIRST FUNCTION K IN THE CAUCHY-KOWALEWSKI SYSTEM.

We differentiate 3(a-b) with respect to t to get expressions for \dot{w}_1, \dot{w}_2 . Denoting the expression for \dot{w}_1 by R we get

$$\dot{w}_1 = R \quad (4.1)$$

To simplify notation, let $K = \alpha$ in (4.1) and get

$$\dot{w}_1 = R = \frac{A_1 L_2 + A_2 L_3 + A_3}{-A_2 L_1 + A_4 L_2 + A_5} \quad (4.1a)$$

Some of the $A_i, i = 1, \dots, 5$ are given explicitly later. These are not partial derivatives. By contrast,

$$L_1 = \frac{\partial L}{\partial w_1} \quad \text{etc.}$$

Now let $z = L - w_2^2$ and note that from (2.1), 3(a-b) we have $f = \frac{w_1}{\alpha} z$, $h = \frac{w_1}{\alpha} + z$ in the new notation.

The following equation is of fundamental importance. We arbitrarily set

$$A_2 = 2w_1w_2\alpha_1 - w_1t\alpha_1 + w_1\alpha_3 + tfa^2 = \varepsilon \quad (4.2)$$

where $K = \alpha = \alpha(L, w_1, w_2, t, \varepsilon)$ and $\alpha_1 = \frac{\partial \alpha}{\partial L}$, etc., for real $\varepsilon > 0$.

By the Cauchy-Kowalewski theorem [See e.g. (2.1)] let $\alpha_0 = \alpha_0(L, w_1, w_2, t, \varepsilon)$ be an analytic solution of (4.2). Further, we will write

$$\bar{A}_i = A_i(\alpha_0), \quad i = 1, 2, 3, 4, 5.$$

Let $\alpha_0 = \sum_{n=0}^{\infty} c_n \varepsilon^n$ where $c_n = c_n(L, w_1, w_2, t)$ are analytic. Before imposing conditions on c_0 we give the following definitions.

$$\text{DEFINITION.} \quad \lim_{\varepsilon \rightarrow 0} L \left[\left(\frac{w_1}{\alpha} + z \right) (w_1\alpha_{04} - w_1w_2\alpha_{02} + w_2\alpha_0) \right] = S_1(L, w_1, w_2, t).$$

Two more of the \bar{A}_i will now be given explicitly.

$$\bar{A}_1 = \left(\frac{w_1}{\alpha} + z \right) (w_1\alpha_{04} - w_1w_2\alpha_{02} + w_2\alpha_0)$$

$$\bar{A}_4 = w_1\alpha_{02} + \alpha_0^2 f - \alpha_0 - w_1h\alpha_{01}$$

$$\text{DEFINITION.} \quad \lim_{\varepsilon \rightarrow 0} L \bar{A}_1 = \Delta.$$

$$\text{DEFINITION.} \quad \lim_{\varepsilon \rightarrow 0} L (\bar{A}_1 - G(w_1, t)\bar{A}_4) = S_2(L, w_1, w_2, t).$$

The conditions on c_0 can be stated now as follows:

$$(1) \quad c_0 \neq 0, \quad (2) \quad S_1(L, w_1, w_2, t) \neq 0, \quad (3) \quad \Delta \neq 0.$$

Substituting $\alpha_0 = \sum_{n=0}^{\infty} c_n \varepsilon^n$ in (4.2) we get

$$2w_1w_2c_{01} - w_1t\left(\frac{w_1}{c_0} + z\right)c_{01} + w_1c_{03} + tw_1zc_0 = 0 \quad (4.3)$$

of which some solutions are given

$$H[\beta(c_0, z, w_1), w_2 + \frac{1}{t} P(c_0, w_1, \beta(c_0, z, w_1))] = \text{constant} \tag{4.3a}$$

where

- (1) H is arbitrary
- (2) β satisfies the partial differential equation

$$c_0 z \beta_1 + \left(\frac{w_1}{c_0} + z\right) \beta_2 = 0$$

$$(w_1 \beta_3 + c_0 \beta_1 \neq 0)$$

(3) P is defined as follows: first solve $\beta(c_0, w_1, z) = a$ for $z = Q(c_0, w_1, a)$. Then set

$$P = \int \frac{d c_0}{c_0 Q(c_0, w_1, a)} .$$

THEOREM 1. The function H can be chosen analytic in (4.3a) so that conditions (2.1), (2.2), (3.1) hold for c_0 .

PROOF. Let $\gamma = w_2 + \frac{1}{t} P$ and then (4.3a) becomes $H(\beta, \gamma) = \text{constant}$. The partial derivatives of c_0 are computed from (4.3a) and from them we see that $H_\gamma \neq 0$ implies that $\frac{\partial c_0}{\partial t} \neq 0$, so condition (2.1) holds. Further, $\Delta = L \bar{A}_1 = 0$ implies $(\frac{P}{t} + w_2) H_\gamma = 0$. So $H_\gamma \neq 0$ implies $\Delta \neq 0$. Thus (2.1), (2.2) hold if merely $H_\gamma \neq 0$. Now $S_1 = 0$ implies that $tw_2(w_1 \beta_3 + c_0 \beta_1) H_\beta + H_\gamma = 0$. Since $w_1 \beta_3 + c_0 \beta_1 \neq 0$, we can choose H so that $S_1 \neq 0$. This completes the proof.

Summarizing the results of this section, $K = \alpha = \alpha_0$ can be defined as the solution of (4.2) where H is analytic, $c_0 \neq 0$, $S_1 \neq 0$, and $\Delta \neq 0$. To solve (3.1) however, we must define L.

5. SOLUTION OF THE CAUCHY-KOWALEWSKI SYSTEM.

To solve the system (3.1), we must now define the function $L(w_1, w_2, t, \epsilon)$.

Setting $\dot{w} = G$, $\alpha = \alpha_0$ and $\bar{A}_2 = \epsilon$, (4.2) in (4.1a) suggests defining L by

$$\epsilon GL_1 + (\bar{A}_1 - G\bar{A}_4)L_2 + \epsilon L_3 = G\bar{A}_5 - \bar{A}_3.$$

$L_1 = \frac{\partial L}{\partial w_1}$, etc. This does not seem to be feasible. Instead, letting ϵ tend to zero leads to

$$L_2 = \frac{\partial L}{\partial w_2} = \frac{G\bar{A}_5 - \bar{A}_3}{\bar{A}_1 - G\bar{A}_4} \quad (5.1)$$

This will be used to define L .

Let λ be a new variable and consider

$$L_2 = \frac{\lambda\bar{A}_5 - \bar{A}_3}{\bar{A}_1 - \lambda\bar{A}_4} \quad (5.2)$$

Note that the right side of (5.2) is analytic where $w_1 \neq 0$ and $\bar{A}_1 - \lambda\bar{A}_4 \neq 0$. So let $L = \bar{L}(w_1, w_2, t, \epsilon, \lambda) = P_1(w_2) + P_2(w_1, w_2, t, \epsilon, \lambda)$ be an analytic solution on (5.2) and assume that none of the expressions Δ , S_1 , c_0 vanish when $L \equiv P_1(w_2)$.

Now since the value of $\frac{\partial}{\partial w_2}(\bar{L}(w_1, w_2, t, \epsilon, \lambda))$ for $\lambda = G(w_1, t)$ is the same as $\frac{\partial}{\partial w_2}(\bar{L}(w_1, w_2, t, \epsilon, G(w_1, t)))$ we see that $\bar{\bar{L}}(w_1, w_2, t, \epsilon) \equiv \bar{L}(w_1, w_2, t, \epsilon, G(w_1, t))$ is a solution of (5.1) for any G . Moreover $\bar{\bar{L}} \in C^I$ since G is continuous and \bar{L} is analytic. Let $K_G = \alpha_0(\bar{\bar{L}}, w_1, w_2, t)$ and $L = \bar{\bar{L}}$.

We now prove the solvability near suitable points of the Cauchy-Kowalewski system. The variable λ gives our functions the universal character referred to previously.

LEMMA I. Let (a, b, c) be such that $S_1(P_1(b)a, b, c) \neq 0$. Then, for small t , the Jacobian of (3.1) is nonzero at (a, b, c, ϵ) .

PROOF. If the Jacobian of (3.1) = 0, then

$$-\bar{A}_2 L_1 + \bar{A}_4 L_2 + \bar{A}_5 = 0 \tag{5.3}$$

The subsidiary equations of (5.3) are:

$$\frac{dw_1}{-\bar{A}_2} = \frac{dw_2}{\bar{A}_4} = \frac{dL}{-\bar{A}_5}, \quad \text{so that} \quad \frac{dL}{dw_2} = \frac{-\bar{A}_5}{\bar{A}_4}$$

But from (5.1),
$$\frac{dL}{dw_2} = \frac{\bar{G}\bar{A}_5 - \bar{A}_3}{\bar{A}_1 - \bar{G}\bar{A}_4} .$$

Thus $\bar{A}_1 \bar{A}_5 - \bar{A}_3 \bar{A}_4 = 0$.

But $\bar{A}_1 \bar{A}_5 - \bar{A}_3 \bar{A}_4 = (w_1 \alpha_{o4} - w_1 w_2 \alpha_{o2} + w_2 \alpha_o) (\frac{w_1}{\alpha_o} + z) \epsilon$. So

$$\lim_{\epsilon \rightarrow 0} (w_1 \alpha_{o4} - w_1 w_2 \alpha_{o2} + w_2 \alpha_o) (\frac{w_1}{\alpha_o} + z) = 0. \quad \text{However}$$

$$\lim_{\epsilon \rightarrow 0} (w_1 \alpha_{o4} - w_1 w_2 \alpha_{o2} + w_2 \alpha_o) (\frac{w_1}{\alpha_o} + z) = S_1(P_1(w_2), w_1, w_2, t) \neq 0 \quad \text{and the}$$

proof is complete.

We next consider continuity in order to apply the implicit function theorem to (3.1). We first observe that $\lim_{\epsilon \rightarrow 0} \bar{A}_1 \neq 0$. If $\lim_{\epsilon \rightarrow 0} \bar{A}_4 = 0$, then

$$\lim_{\epsilon \rightarrow 0} (\bar{A}_1 - \bar{G}\bar{A}_4) \neq 0.$$

Now consider the case where $\lim_{\epsilon \rightarrow 0} \bar{A}_4 \neq 0$, but $\lim_{\epsilon \rightarrow 0} (\bar{A}_1 - \bar{G}\bar{A}_4) = 0$.

LEMMA II. There is at most one function G such that $\lim_{\epsilon \rightarrow 0} (\bar{A}_1 - \bar{G}\bar{A}_4) = 0$.

PROOF. $\bar{L}(w_1, w_2, t, \epsilon) = \bar{L}(w_1, w_2, t, G(w_1, t)) = P_1(w_2) + \epsilon P_2(w_1, w_2, t, \epsilon, G(w_1, t))$.

So it and its partials with respect to w_1, w_2, t do not contain G as $\epsilon \rightarrow 0$. Since

$$\alpha_o = \sum_{n=0}^{\infty} c_n(L, w_1, w_2, t) = c_o(L, w_1, w_2, t) + c_1(L, w_1, w_2, t) + c_2(L, w_1, w_2, t)\epsilon^2 + \dots,$$

the same holds for it.

Thus $\lim_{\epsilon \rightarrow 0} \bar{A}_1$ and $\lim_{\epsilon \rightarrow 0} \bar{A}_4$ are independent of G.

So $G = \frac{\lim_{\epsilon \rightarrow 0} \bar{A}_1}{\lim_{\epsilon \rightarrow 0} \bar{A}_4}$. This completes the proof.

In the sequel, we ignore this possible exception and assume that

$\lim_{\epsilon \rightarrow 0} (\bar{A}_1 - G\bar{A}_4) \neq 0$ for any G.

LEMMA III. If (a,b,c) is such that $S_2(P_1(b), a, b, c) \neq 0$, there is an $\epsilon > 0$ such that the left sides of (3.1) are C^I at (a,b,c,ε).

PROOF. Based on analytic properties of V_1, V_2, \bar{L}, K_G and the nonvanishing of S_2 , we will not give details.

Choosing constant values for w_1, w_2 in (3.1), we can get $c_1(\epsilon), c_2(\epsilon)$ so that left sides vanishes and apply the implicit function theorem to (3.1). Then we solve for $w_1(t, \epsilon)$ and $w_2(t, \epsilon)$. Here c_1, c_2 come from equation (2.1) of section 2.

6. THE PRINCIPAL DIFFERENTIAL EQUATION.

We now consider the differential equation

$$\frac{dy}{dx} = y' = g(x, y) \tag{6.1}$$

DEFINITION. $W_1(t) = \lim_{\epsilon \rightarrow 0} w_1(t, \epsilon)$.

It will be shown that $W_1(t)$ satisfies (6.1). Of course we change y, x to W_1, t respectively.

We begin this process with

THEOREM II. Let $S_1 \neq 0$ at $(\bar{w}_1, \bar{w}_2, \bar{t})$. Then $\frac{d}{dt} w_1(t, \epsilon) \rightarrow G(\bar{w}_1, \bar{t})$ as $\epsilon \rightarrow 0$.

PROOF. $\bar{L} = P_1(w_2) + \epsilon P_2(w_1, w_2, t, G(w_1, t))$ so that $\frac{\partial \bar{L}}{\partial w_1} \rightarrow 0$ as $\epsilon \rightarrow 0$ and also $\frac{\partial \bar{L}}{\partial t} \rightarrow 0$ as $\epsilon \rightarrow 0$. Thus $\bar{L}_1, \bar{L}_3 \rightarrow 0$ as $\epsilon \rightarrow 0$.

Now (5.1) $\bar{A}_1 L_2 + \bar{A}_3 = G(\bar{A}_4 L_2 + \bar{A}_5)$.

$$\text{Also } \bar{A}_4 L_2 + \bar{A}_5 = \frac{\bar{A}_1 \bar{A}_5 - \bar{A}_3 \bar{A}_4}{\bar{A}_1 - G\bar{A}_4} = \frac{S_1 \epsilon}{\bar{A}_1 - G\bar{A}_4}.$$

$$\text{Thus } R = \frac{\bar{A}_1 L_2 + \bar{A}_2 L_3 + \bar{A}_3}{-\bar{A}_2 L_1 + \bar{A}_4 L_2 + \bar{A}_5} = \frac{\epsilon L_3 + \bar{A}_1 L_2 + \bar{A}_3}{-\epsilon L_1 + \bar{A}_4 L_2 + \bar{A}_5}.$$

$$\text{So } \dot{w}_1 = \frac{\epsilon L_3 + G(\bar{A}_4 L_2 + \bar{A}_5)}{-\epsilon L_1 + (\bar{A}_4 L_2 + \bar{A}_5)} = \frac{(\bar{A}_1 - G\bar{A}_4)L_3 + GS_1}{-(\bar{A}_1 - G\bar{A}_4)L_1 + S_1}.$$

Therefore $\dot{w}_1 \rightarrow \frac{GS_1}{S_1}$ as $\epsilon \rightarrow 0$ and $S_1 \neq 0$. This completes the proof.

By the last theorem, $L_{\epsilon \rightarrow 0} \frac{d}{dt} w_1(t, \epsilon) = L_{\epsilon \rightarrow 0} G(w_1(t, \epsilon), t) = G(L_{\epsilon \rightarrow 0} w_1(t, \epsilon), t) =$

$G(W_1(t), t)$.

But also it is true [2: P.461] that

$$L_{\epsilon \rightarrow 0} \frac{d}{dt} w_1(t, \epsilon) = \frac{d}{dt} (L_{\epsilon \rightarrow 0} w_1(t, \epsilon)) = W_1'(t).$$

$$\text{So } W_1'(t) = G(W_1(t), t) \tag{6.2}$$

7. PARTICULAR AND GENERAL SOLUTIONS OF $y' = G(x, y)$.

7(a) PARTICULAR SOLUTIONS. Let $J(w_1, t) \in C^I$,

$$L^*(w_1, w_2, t) = \bar{L}(w_1, w_2, t, \epsilon, J(w_1, t)) \text{ and } \alpha^*(w_1, w_2, t) = \alpha_0(L^*, w_1, w_2, t).$$

Let Q be the set of points in (w_1, w_2, t) -space where

- (1) $w_1 \neq 0$
- (2) $c_0 \neq 0$
- (3) $S_1 \neq 0$
- (4) $S_2 \neq 0$.

Let \bar{Q} be the projection of Q on the (w_1, t) plane.

The Universal Cauchy-Kowalewski System

DEFINITION. $\bar{\alpha}(w_1, w_2, t, \epsilon, \lambda) = \alpha_0(\bar{L}, w_1, w_2, t)$.

DEFINITION. $F_1 \equiv w_1 - \bar{\alpha}V_1(w_1 + tw_2, c_1, c_2)$.

DEFINITION. $F_2 \equiv w_2^2 - \bar{L}(w_1, w_2, t, \epsilon, \lambda) - V_2(w_1 + tw_2, c_1, c_2)$.

DEFINITION. $F_3 \equiv \bar{A}_1 - J(w_1, t)\bar{A}_4$ with λ replaced by $J(w_1, t)$.

DEFINITION. The system
$$\begin{cases} F_1 = 0 \\ F_2 = 0 \\ F_3 \neq 0 \end{cases}$$
 is also called the Universal

Cauchy-Kowalewski System.

We refer to it in the following

THEOREM III. Let $P \in \bar{Q}$. There is a region in which the solution through P of $\dot{w}_1 = J(w_1, t)$ is determined as follows:

- (1) In F_1, F_2 replace λ by $J(w_1, t)$ and c_1, c_2 by suitable functions of ϵ .
- (2) Equate the results in (2.1) to zero.
- (3) Solve the resulting system for $w_1(t, \epsilon)$ and $w_2(t, \epsilon)$.
- (4) Take the limit of $w_1(t, \epsilon)$ as $\epsilon \rightarrow 0$.

PROOF. Let $P = (a, t_0)$, $P \in \bar{Q}$. Since $c_0(P_1(b), a, b, t_0) \neq 0$, there is an ϵ such that $\alpha^*(a, b, t_0, \epsilon) \neq 0$. Let (\bar{v}_1, \bar{v}_2) be a solution of (2.1) such that

$$\left\{ \begin{array}{l} \bar{v}_1(a + t_0 b) = \frac{a}{\alpha^*(a, b, t_0, \epsilon)} \\ \bar{v}_2(a + t_0 b) = b^2 - L^*(a, b, t_0, \epsilon) \end{array} \right\}$$

Solve the system:

$$\left\{ \begin{array}{l} (1) \quad V_1(a + t_0 b, c_1, c_2) - \frac{a}{\alpha^*(a, b, t_0, \epsilon)} = 0 \\ (2) \quad V_2(a + t_0 b, c_1, c_2) - b^2 + L^*(a, b, t_0, \epsilon) = 0 \end{array} \right\}$$

to get suitable $c_1 = c_1(\epsilon)$, $c_2 = c_2(\epsilon)$.

Since $S_1 \neq 0$ our system has nonzero Jacobian. We solve for $w_1(t, \epsilon)$ and get the result.

7(b) GENERAL SOLUTIONS. Alternatively, eliminating w_2 from the Universal Cauchy-Kowalewski System we get

$$X(w_1, t, \epsilon, \lambda, c_1, c_2) = 0 \quad (7.1)$$

where c_1, c_2 are constants.

The general solution of a specific equation is obtained as follows:

- (1) Replace λ by $G(w_1, t)$ in (7.1).
- (2) Take the limit as $\epsilon \rightarrow 0$ of the result.

X is derived from L and K and is like the familiar differential quotient in generality.

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