

EXPERIMENT ON NUMERICAL CONFORMAL MAPPING OF UNBOUNDED MULTIPLY CONNECTED DOMAIN IN FUNDAMENTAL SOLUTIONS METHOD

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ABSTRACT. We are concerned with the experiment on numerical conformal mappings. A potentially theoretical scheme in the fundamental solutions method, different from the conventional one, has been recently proposed for numerical conformal mappings of unbounded multiply connected domains. The scheme is based on the asymptotic theorem on extremal weighted polynomials. The scheme has the characteristic called “invariant and dual.” Applying the scheme for typical examples, we will show that the numerical results of high accuracy may be obtained.

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1. Introduction. The fundamental solutions method (or charge simulation method) has been applied to the problem in electrical engineering, numerical conformal mappings [2, 3, 4, 10, 12] and Dirichlet problems [10, 16, 17].

The principle of the method is the approximation of the solution by a linear combination of logarithmic potentials. Though the method requires only solving a system of simultaneous linear equations, it is possible to get a rather precise solution for boundary problems with respect to domains bounded by smooth curves.

Amano [2, 3] has recently proposed two kinds of schemes of approximations for the conformal mappings onto the domains with circular or radial cuts, respectively.

Kuhara [11, 12] has also established a construction method of the functions mapping multiply connected domains onto the rings with circular or radial slits, based upon the works of Bergman [5] and using the fundamental solutions method. The method is described from the two-dimensional electrostatic point of view.

A potentially theoretical scheme in the fundamental solutions method, different from the conventional one, has been recently proposed for numerical conformal mappings of unbounded multiply connected domains [9]. The scheme is based on the asymptotic theorem on extremal weighted polynomials [7, 8, 13, 14, 15]. The scheme has the characteristic called “invariant and dual” with respect to interior and exterior domains [9]. In this paper, applying the scheme for typical examples, we show that the numerical results of high accuracy may be obtained.

2. Scheme for numerical conformal mapping. In this section, we show the scheme computing the approximations for the conformal mappings of “unbounded” multiply connected domains, which has been recently proposed in [9].

Let D and D' denote unbounded multiply connected domains whose boundaries γ and γ' consist of Jordan curves γ_i and γ'_i ($i = 0(1)m$), respectively.

We assume that γ_0, γ'_0 enclose the origin. Let γ'_0 be a circle $\{w : |w| = r_0\}$. Let $f(z)$ map conformally D onto D' with the continuation to a bijection mapping from

$$D \cup \gamma \text{ to } D' \cup \gamma', \quad (2.1)$$

corresponding γ_i to γ'_i . $f(z)$ is uniquely determined under the condition $f(\infty) = \infty$, $f'(\infty) = 1$ [1].

We have proposed in [9] the following scheme of approximations of $f(z)$

$$f_n(z) = z \prod_{i=1}^n \left(1 - \frac{z_{n,i}}{z}\right)^{\alpha_i}, \quad \sum_{i=1}^n \alpha_i = 1, \quad (2.2)$$

where the charge points $\{z_{n,i}\}_{i=1}^n$ are appropriately chosen interior to γ .

When D' is $\{w : |w| > r_0\}$ with radial cuts $\sum_{j=1}^m \gamma'_j$, we propose the algorithm computing approximations of $f(z)$ as follows.

SCHEME 2.1. The approximation $f_n(z)$ of $f(z)$ may be obtained as follows:

(2a) $\{z_{n_j,i}^{(j)}\}_{i=1}^{n_j}$ and $\{\zeta_{n_j,i}^{(j)}\}_{i=1}^{n_j}$ with $n_0 = \dots = n_m = n$ are appropriately chosen interior to γ_j and on γ_j ($j = 0(1)m$), respectively.

(2b) When $\alpha_i^{(j)}$ ($i = 0(1)n_j, j = 0(1)m$) are the solutions of a system of $(m+1)(n+1)$ simultaneous linear equations using Dirichlet-Neumann and charge conditions [11, 12]:

$$\begin{aligned} \alpha_0^{(0)} + \log |\zeta_{n_0,k}^{(0)}| + \sum_{j=0}^m \sum_{i=1}^{n_j} \alpha_i^{(j)} \log \left| 1 - \frac{z_{n_j,i}^{(j)}}{\zeta_{n_0,k}^{(0)}} \right| &= 0 \quad (k = 1(1)n_0), \\ \alpha_0^{(l)} + \arg(\zeta_{n_l,k}^{(l)}) + \sum_{j=0}^m \sum_{i=1}^{n_j} \alpha_i^{(j)} \arg \left(1 - \frac{z_{n_j,i}^{(j)}}{\zeta_{n_l,k}^{(l)}} \right) &= 0 \quad (k = 1(1)n_l, l = 1(1)m), \\ \sum_{i=1}^{n_0} \alpha_i^{(0)} &= 1, \quad \sum_{i=1}^{n_j} \alpha_i^{(j)} = 0 \quad (j = 1(1)m), \end{aligned} \quad (2.3)$$

the charges at $\{z_{n_j,i}^{(j)}\}_{i=1}^{n_j}$ are given by $\{\alpha_i^{(j)}\}_{i=1}^{n_j}$ ($j = 1(1)m$), respectively.

(2c) The approximation $f_n(z)$ is represented by

$$f_n(z) = z \prod_{j=0}^m \prod_{i=1}^{n_j} \left(1 - \frac{z_{n_j,i}^{(j)}}{z}\right)^{\alpha_i^{(j)}}, \quad \sum_{i=1}^{n_0} \alpha_i^{(0)} = 1, \quad \sum_{i=1}^{n_j} \alpha_i^{(j)} = 0 \quad (j = 1(1)m). \quad (2.4)$$

Note that the approximations

$$\alpha_0^{(0)} \simeq -\log r_0, \quad \alpha_0^{(j)} \simeq -\theta_j \quad (j = 1(1)m) \quad (2.5)$$

hold, where θ_j is the argument of γ'_j .

The invariant scheme of approximations has been first shown for the numerical Dirichlet problem by Murota [16, 17]. It is physically natural and mathematically reasonable.

The solutions of a system of simultaneous linear equations in Scheme 2.1 are also invariant in the sense that the transformation $z \rightarrow az$ ($a > 0$) implies

$$\begin{aligned} \alpha_0^{(0)} &\rightarrow \alpha_0^{(0)} + \log a, \\ \alpha_0^{(j)} &\rightarrow \alpha_0^{(j)} \quad (j = 1(1)m), \\ \alpha_i^{(j)} &\rightarrow \alpha_i^{(j)} \quad (i = 1(1)n, j = 0(1)m). \end{aligned} \tag{2.6}$$

3. A numerical example. We consider a function

$$w = f(z) = z + \frac{1}{z-6} \left(z = f^{-1}(w) = \frac{(7-5\sqrt{(w-8)/(w-4)})}{(1-\sqrt{(w-8)/(w-4)})} \right) \tag{3.1}$$

mapping D onto D' , where D' is $\{w : |w| > 1\}$ with a radial cut

$$\{w : 4 \leq \text{Re}(w) \leq 8, \text{Im}(z) = 0\}. \tag{3.2}$$

This corresponds to Dirichlet-Neumann problem and easy to check the accuracy of the approximation. We apply Scheme 2.1 to compute the approximations of $f(z)$.

The charge points interior to γ_1 and the collocation points on γ_1 are so chosen that

$$l \exp\left(\frac{2\pi j(i-1)}{n}\right) + 6, \quad j = \sqrt{-1}, \quad (i = 1(1)n) \tag{3.3}$$

with $l = 0.75, 0.5, 0.25$ and $l = 1$, respectively.

On the other hand, the charge points interior to γ_0 and the collocation points on γ_0 are the images of the points

$$l \exp\left(\frac{2\pi j(i-1)}{n}\right), \quad j = \sqrt{-1}, \quad (i = 1(1)n) \tag{3.4}$$

being distributed on $\{w : |w| = 0.75, 0.5, 0.25, 1\}$ under the inverse function $f^{-1}(w)$ of $f(z)$, respectively.

For $n = 13$, we solve a system of simultaneous linear equations (2.3) and obtain the approximations $f_n(z)$.

The accuracy of the errors is estimated by

$$|f_n(z) - f(z)| \tag{3.5}$$

at the points on γ_1

$$\exp\left(\frac{2\pi j(i-1)}{n} + \frac{\pi j}{n}\right) + 6, \quad j = \sqrt{-1}, \quad (i = 1(1)n) \tag{3.6}$$

and the images of the points on γ_0

$$\exp\left(\frac{2\pi j(i-1)}{n} + \frac{\pi j}{n}\right), \quad j = \sqrt{-1}, \quad (i = 1(1)n). \tag{3.7}$$

under the function $f^{-1}(w)$.

By the maximum principle for the analytic functions, it is sufficient that the errors are estimated only on the boundary.

The numerical results are presented for the following cases (under a minor modification of the scheme in order to keep the continuity of the argument).

(3a) The charge distribution with $l = 0.75$.

We show the charges in (2b) interior to y_1 and y_0 , respectively as follows:

TABLE 3.1. The charge distribution with $n = 13$, $l = 0.75$ on y_1 .

$-6.095578422303619D - 002$	$5.357980215700247D - 003$	$-5.370095750188362D - 002$
$2.111998739887658D - 002$	$-2.743036534771264D - 002$	$5.611121564384160D - 002$
$2.902003170269616D - 002$	$2.902003170269579D - 002$	$5.611121564384106D - 002$
$-2.743036534771176D - 002$	$2.111998739887570D - 002$	$-5.370095750188297D - 002$
$5.357980215700064D - 003$		

TABLE 3.2. The charge distribution with $n = 13$, $l = 0.75$ on y_0 .

$7.693733748274377D - 002$	$7.693054597318545D - 002$	$7.692230818952048D - 002$
$7.691974075923927D - 002$	$7.691943391695984D - 002$	$7.691958631030785D - 002$
$7.691971610941527D - 002$	$7.691971610941528D - 002$	$7.691958631030788D - 002$
$7.691943391695967D - 002$	$7.691974075923903D - 002$	$7.692230818952112D - 002$
$7.693054597318509D - 002$		

Note that

$$\frac{1}{13} = 7.692307692307693D - 002. \quad (3.8)$$

The errors on y_1 and y_0 are as follows:

TABLE 3.3. The errors with $n = 13$, $l = 0.75$ on y_1 .

$7.906455437100468D - 002$	$7.272028918873658D - 002$	$6.874725258323458D - 002$
$5.361607892270346D - 002$	$5.209733883293302D - 002$	$4.984156926710034D - 002$
$4.894182325162078D - 002$	$4.984156926709897D - 002$	$5.209733883293163D - 002$
$5.361607892270290D - 002$	$6.874725258323204D - 002$	$7.272028918873554D - 002$
$7.906455437100488D - 002$		

TABLE 3.4. The errors with $n = 13$, $l = 0.75$ on y_0 .

$1.917221469173348D - 003$	$1.917212886788780D - 003$	$1.917182701862004D - 003$
$1.916825188166824D - 003$	$1.916448045831497D - 003$	$1.916203808715182D - 003$
$1.916121835051632D - 003$	$1.916203808714936D - 003$	$1.916448045830375D - 003$
$1.916825188167161D - 003$	$1.917182701862035D - 003$	$1.917212886788246D - 003$
$1.917221469172982D - 003$		

Furthermore,

$$\alpha_0^{(1)} - (-\theta_1) = \alpha_0^{(1)} - (-0) = 1.741206340272162D - 003, \tag{3.9}$$

$$\alpha_0^{(0)} - (-\log r_0) = \alpha_0^{(0)} - (-\log 1) = -1.040834085586084D - 017, \tag{3.10}$$

which shows that (2.5) holds with high accuracy.

(3b) The charge distribution with $l = 0.5$.

We show the charges in (2b) interior to γ_1 and γ_0 , respectively as follows:

TABLE 3.5. The charge distribution with $n = 13, l = 0.5$ on γ_1 .

5.676888435560761D - 002	-1.336464863115225D - 001	6.408202622703688D - 002
-1.158867927979212D - 001	9.337154978330532D - 002	-4.215724143899066D - 002
1.058525023602893D - 001	1.058525023602879D - 001	-4.215724143899442D - 002
9.337154978331164D - 002	-1.158867927979256D - 001	6.408202622703865D - 002
-1.336464863115230D - 001		

TABLE 3.6. The charge distribution with $n = 13, l = 0.5$ on γ_0 .

7.690431354619905D - 002	7.691784500736430D - 002	7.692657263421879D - 002
7.692635000860265D - 002	7.692605635027436D - 002	7.692558594969530D - 002
7.692543327674563D - 002	7.692543327674435D - 002	7.692558594969630D - 002
7.692605635027493D - 002	7.692635000860015D - 002	7.692657263422217D - 002
7.691784500736203D - 002		

Note that

$$\frac{1}{13} = 7.692307692307693D - 002. \tag{3.11}$$

The errors on γ_1 and γ_0 are as follows:

TABLE 3.7. The errors with $n = 13, l = 0.5$ on γ_1 .

1.686595817949471D - 002	1.594668522932475D - 002	1.426456576379659D - 002
1.245942289990706D - 002	1.124520330394002D - 002	1.730254947067371D - 002
1.872540245176557D - 002	1.730254947067469D - 002	1.124520330394140D - 002
1.245942289990679D - 002	1.426456576379672D - 002	1.594668522932589D - 002
1.686595817949507D - 002		

TABLE 3.8. The errors with $n = 13, l = 0.5$ on γ_0 .

6.070869763412672D - 005	6.733732155656440D - 005	6.897867302145204D - 005
6.621157003193045D - 005	6.260984607098040D - 005	6.000919726996125D - 005
5.909032538686798D - 005	6.000919726934812D - 005	6.260984607020485D - 005
6.621157003170747D - 005	6.897867302240548D - 005	6.733732155617585D - 005
6.070869763390371D - 005		

Furthermore,

$$\alpha_0^{(1)} - (-\theta_1) = 7.787235018993428D - 005, \quad (3.12)$$

$$\alpha_0^{(0)} - (-\log r_0) = 3.426078865054194D - 017, \quad (3.13)$$

which shows that (2.5) holds with high accuracy.

(3c) The charge distribution with $l = 0.25$.

We show the charges in (2b) interior to γ_1 and γ_0 , respectively as follows:

TABLE 3.9. The charge distribution with $n = 13$, $l = 0.25$ on γ_1 .

$-4.896753338030436D - 002$	$-7.545409516631382D - 002$	$-4.502745232832742D - 002$
$-6.625722438826202D - 002$	$-2.034627278488312D - 002$	$-2.970889592876331D - 003$
$2.345397009507822D - 001$	$2.345397009509353D - 001$	$-2.970889593038040D - 003$
$-2.034627278472483D - 002$	$-6.625722438839048D - 002$	$-4.502745232824451D - 002$
$-7.545409516635260D - 002$		

TABLE 3.10. The charge distribution with $n = 13$, $l = 0.25$ on γ_0 .

$7.692306130528526D - 002$	$7.692308599892384D - 002$	$7.692307475080235D - 002$
$7.692307800922398D - 002$	$7.692307646531414D - 002$	$7.692307723249565D - 002$
$7.692307689059898D - 002$	$7.692307689069573D - 002$	$7.692307723235956D - 002$
$7.692307646552272D - 002$	$7.692307800894760D - 002$	$7.692307475107583D - 002$
$7.692308599875447D - 002$		

Note that

$$\frac{1}{13} = 7.692307692307693D - 002. \quad (3.14)$$

The errors on γ_1 and γ_0 are as follows:

TABLE 3.11. The errors with $n = 13$, $l = 0.25$ on γ_1 .

$2.620143221187881D - 005$	$2.460270478748571D - 005$	$2.165285311487771D - 005$
$1.774249777583574D - 005$	$1.308661337754165D - 005$	$7.354889113223171D - 006$
$2.287964484537497D - 009$	$7.354889114253348D - 006$	$1.308661337918933D - 005$
$1.774249777787493D - 005$	$2.165285311546687D - 005$	$2.460270478816393D - 005$
$2.620143221202827D - 005$		

TABLE 3.12. The errors with $n = 13$, $l = 0.25$ on γ_0 .

$4.903447136123618D - 010$	$1.019052540368634D - 009$	$9.883685502913276D - 010$
$7.553265530481330D - 010$	$4.968949958792792D - 010$	$2.547095062921054D - 010$
$8.554679187256638D - 011$	$2.547094246284032D - 010$	$4.968949414229016D - 010$
$7.553265708226896D - 010$	$9.883683875455573D - 010$	$1.019053538929026D - 009$
$4.903451511212511D - 010$		

Furthermore,

$$\alpha_0^{(1)} - (-\theta_1) = 2.378106242066085D - 009, \tag{3.15}$$

$$\alpha_0^{(0)} - (-\log r_0) = 8.673617379884036D - 019, \tag{3.16}$$

which shows that (2.5) holds with high accuracy.

The numerical example shows:

(3d) The data present the distribution of the charges and errors of the approximations, which are symmetric with respect to the real axis and with high accuracy.

(3e) When the charges and collocation points are distributed uniformly (see the definition in [6]) as (3.4), the example shows that the approximations

$$\alpha_i \simeq \frac{1}{n} \quad (i = 1, 2, \dots, n), \quad \alpha_0^{(0)} \simeq -\log r_0, \quad f(z) \simeq f_n(z) \tag{3.17}$$

would hold with high accuracy.

We have examined the case of odd n . The case of $n = 12$ is shown in the following (3f).

(3f) The charge distribution with $l = 0.25$.

We show the charges in (2b) interior to γ_1 and γ_0 , respectively as follows:

TABLE 3.13. The charge distribution with $n = 12, l = 0.25$ on γ_1 .

-1.942606947010962D - 001	5.972416806033705D - 002	-1.896074309191883D - 001
7.405835693044212D - 002	-1.569787980380907D - 001	2.128037039665000D - 001
1.942606947011143D - 001	2.128037039664623D - 001	-1.569787980380575D - 001
7.405835693042352D - 002	-1.896074309191738D - 001	5.972416806032727D - 002

TABLE 3.14. The charge distribution with $n = 12, l = 0.25$ on γ_0 .

8.333347749250593D - 002	8.333322686829960D - 002	8.333337977798844D - 002
8.333331344289655D - 002	8.333334554423848D - 002	8.333332530161311D - 002
8.333334063736710D - 002	8.333332530170494D - 002	8.333334554409336D - 002
8.333331344306075D - 002	8.333337977783503D - 002	8.333322686839680D - 002

Note that

$$\frac{1}{12} = 8.333333333333333D - 002. \tag{3.18}$$

The errors on γ_1 and γ_0 are as follows:

TABLE 3.15. The errors with $n = 12, l = 0.25$ on γ_1 .

4.889879193915697D - 004	4.566691529354238D - 004	4.003418633466522D - 004
3.341897580605150D - 004	2.746913382174855D - 004	2.377726507971854D - 004
2.377726507949406D - 004	2.746913382176728D - 004	3.341897580660071D - 004
4.003418633444508D - 004	4.566691529379670D - 004	4.889879193896632D - 004

Furthermore,

$$\alpha_0^{(1)} - (-\theta_1) = 8.050910065412068D - 009, \quad (3.19)$$

$$\alpha_0^{(0)} - (-\log r_0) = 1.001802807376606D - 016, \quad (3.20)$$

which shows that (2.5) holds with high accuracy.

TABLE 3.16. The errors with $n = 12$, $l = 0.25$ on y_0 .

2.472695830800364D - 009	2.547021968349248D - 009	1.880133758085106D - 009
1.954951566801633D - 009	1.938261615460128D - 009	1.892427103310218D - 009
1.892426551711589D - 009	1.938261234766441D - 009	1.954951344522423D - 009
1.880133600747165D - 009	2.547022633553721D - 009	2.472696052933940D - 009

We have shown the data exactly for the convenience of the readers in order to follow the numerical experiment in Fortran 90 with double precision.

The numerical calculation has been performed in *MsDevf90* (PC9821-NEC).

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4. Concluding remark. The numerical scheme of approximations in this paper has been applied for a doubly connected domains applying the charge simulation method.

The numerical experiments for other doubly connected domains and $m(\neq 2)$ -multiply connected ones will be furthermore studied in a future paper.

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