## The Atiyah Extension of a Lie Algebra Deformation

## Ziv Ran<sup>\*</sup>

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**Abstract.** We construct an analogue of the 'Atiyah class', or 'Atiyah extension' in the context of deformation theory of sheaves of Lie algebras. Given a (suitable) sheaf  $\mathfrak{g}$  of Lie algebras on a space X, and a  $\mathfrak{g}$ -deformation  $\phi$  over a formally smooth base, our Atiyah extension is a Lie algebra extension of the algebra of base vector fields by the  $\phi$ -twist of  $\mathfrak{g}$ . It comes equipped with a representation of the  $\phi$ -twist on any (admissible)  $\mathfrak{g}$ -module. Mathematics Subject Index: Primary 32G99, secondary 17B55, 14D14. Key Words and Phrases: Sheaf of Lie algebras, deformation, Atiyah class, extension of Lie algebras.

The role of the Atiyah class (also known as Atiyah Chern class or Atiyah extension) in the study of vector bundles on manifolds is well known (see [3] for just one example and [6] for a discussion of that example from a viewpoint close to that of this note). The purpose of this note is to define an analogue (in fact, a generalization) of the Atiyah extension in the context of the deformation theory of Lie algebras. Given a sheaf of Lie algebras  $\mathfrak{g}$  on a space X, a  $\mathfrak{g}$ -deformation  $\phi$  (as reviewed below), say parametrized by a formally smooth local algebra  $S^{\wedge}$ , simultaneously determines deformations  $E^{\phi}$  for all  $\mathfrak{g}$ -modules E (including the adjoint module  $\mathfrak{g}$  itself, which yields a Lie algebra  $\mathfrak{g}^{\phi}$ ). Our Atiyah extension is a Lie algebra extension of the algebra of base vector fields by  $\mathfrak{g}^{\phi}$ , and comes equipped with a representation on  $E^{\phi}$  for any  $\mathfrak{g}$ -module E. This yields an interesting class of Lie algebras and representations. For X a smooth curve, some related notions were studied by Beilinson and Schechtman [2] in connection with Virasoro algebras. Even in the classical context of vector bundles, the realization of the Atiyah extension as a Lie algebra is apparently new.

To proceed with the basic definitions, let

$$f: X_B \to B$$

be a continuous mapping of Hausdorff spaces with fibres  $X_b = f^{-1}(b)$  and base B, which we assume endowed with a sheaf of local  $\mathbb{C}$ -algebras  $\mathcal{O}_B$ . A Lie

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pair  $(\mathfrak{g}_B, E_B)$  on  $X_B/S$  consists of a sheaf  $\mathfrak{g}_B$  of  $f^{-1}\mathcal{O}_B$ -Lie algebras (i.e. with  $f^{-1}\mathcal{O}_B$ - linear bracket), a sheaf  $E_B$  of  $f^{-1}\mathcal{O}_B$ -modules with  $f^{-1}\mathcal{O}_B$ linear  $\mathfrak{g}_B$ -action. This pair is said to be *admissible* if it admits compatible soft resolutions  $(\mathfrak{g}_B, E_B)$  such that  $\mathfrak{g}_B$  is a differential graded Lie algebra and  $E_B$  is a differential graded representation of  $\mathfrak{g}_B$ , and moreover,  $\Gamma(\mathfrak{g}_B), \Gamma(E_B)$ may be linearly topologized so that coboundaries (and cocycles) are closed, and the cohomology is of finite type as  $\mathcal{O}_B$ -module (and in particular vanishes in almost all degrees). Let's call such resolutions good. Note that if  $(\mathfrak{g}_B, E_B)$  is an admissible pair then for any  $b \in B$  the 'fibre'

$$(\mathfrak{g}_b, E_b) := (\mathfrak{g}_B, E_B) \otimes (\mathcal{O}_{B,b}/\mathfrak{m}_{B,b})$$

is an admissible pair on  $X_b$ .

Now let S be an augmented  $\mathcal{O}_B$ -algebra of finite type as  $\mathcal{O}_B$ -module, with maximal ideal  $\mathfrak{m}_S$  (below we shall also consider the case where S is an inverse limit of such algebras, hence is complete noetherian rather than finite type). By a relative  $\mathfrak{g}_B$ -deformation of  $E_B$ , parametrized by S we mean a sheaf  $E_B^{\phi}$  of S-modules on  $X_B$ , together with a maximal atlas of the following data

- An open covering  $(U_{\alpha})$  of  $X_B$ .
- $\mathcal{S}$ -isomorphisms  $\Phi_{\alpha}: E^{\phi}|_{U_{\alpha}} \xrightarrow{\sim} E|_{U_{\alpha}} \otimes_{\mathcal{O}_B} \mathcal{S}.$
- For each  $\alpha, \beta$ , a lifting of  $\Phi_{\beta} \circ \Phi_{\alpha}^{-1} \in Aut(E|_{U_{\alpha} \cap U_{\beta}} \otimes_{\mathcal{O}_{B}} \mathcal{S})$  to an element  $\Psi_{\alpha,\beta} \in \exp(\mathfrak{g}_{B} \otimes \mathfrak{m}_{\mathcal{S}}(U_{\alpha} \cap U_{\beta}))$ . If  $\mathfrak{g}_{B}$  acts faithfully on  $E_{B}$  then the  $\Psi_{\alpha,\beta}$  are uniquely determined by the  $\Phi_{\alpha}$  and form a cocycle; in general we require additionally that the  $\Psi_{\alpha,\beta}$  form a cocycle.

Note that if  $(\mathfrak{g}_B, E_B)$  is admissible then, as in the absolute case, for any relative deformation  $\phi$  there is a good resolution  $(E^{\cdot}, \partial)$  of E and a resolution of  $E^{\phi}$  of the form

(1) 
$$E^0 \otimes_{\mathcal{O}_B} \mathcal{S} \xrightarrow{\partial + \phi} E^1 \otimes_{\mathcal{O}_B} \mathcal{S} \cdots$$

with

$$\phi \in \Gamma(\mathfrak{g}^1_B) \otimes \mathfrak{m}_{\mathcal{S}}.$$

We call such a resolution a good resolution of  $E^{\phi}$ . Let  $(\mathfrak{g}_B, E_B)$  be an admissible pair on  $X_B/B$ ,  $\mathcal{S}$  a finite-length  $\mathcal{O}_B$ -algebra, and  $E^{\phi}$  an admissible  $\mathfrak{g}_B$ -deformation parametrized by  $\mathcal{S}$ . There is a corresponding deformation  $\mathfrak{g}^{\phi}$ , and clearly  $\mathfrak{g}^{\phi}$  is a Lie algebra acting on  $E^{\phi}$ . We ignore momentarily the status of  $\mathcal{E}^{\phi}$  as a deformation and just view it as a  $\mathfrak{g}^{\phi}$ -module over

$$X_{\mathcal{S}} = X_B \times_B \operatorname{Spec}(\mathcal{S}).$$

Let  $\operatorname{Spec}(\mathcal{S}_1)$  be the first infinitesimal neighborhood of the diagonal in

$$\operatorname{Spec}(\mathcal{S}) \times_{\mathcal{O}_B} \operatorname{Spec}(\mathcal{S}) = \operatorname{Spec}_{\mathcal{O}_B}(\mathcal{S} \otimes_{\mathcal{O}_B} \mathcal{S})$$

with projections

$$p, q: \operatorname{Spec}(\mathcal{S}_1) \to \operatorname{Spec}(\mathcal{S}).$$

Then  $p_*q^*E^{\phi}$  may be viewed as a first-order  $\mathfrak{g}^{\phi}$ -deformation of  $E^{\phi}$  and we let

(2) 
$$AC(\phi) \in H^1(\mathfrak{g}^{\phi} \otimes \mathfrak{m}_{\mathcal{S}_1}) = H^1(\mathfrak{g}^{\phi} \otimes_{\mathcal{S}} \Omega_{\mathcal{S}/B})$$

be the associated (first-order) Kodaira-Spencer class. We call  $AC(\phi)$  the Atiyah class of the deformation  $\phi$ .

A cochain representative for  $AC(\phi)$  may be constructed as follows. Let

 $\phi \in \Gamma(\mathfrak{g}^1) \otimes \mathfrak{m}_{\mathcal{S}}$ 

be a Kodaira-Spencer cochain corresponding to  $E^{\phi}$ , satisfying the integrability condition

$$\partial \phi = -\frac{1}{2}[\phi, \phi].$$

Let

$$d_{\mathcal{S}}: \Gamma(\mathfrak{g}^1) \otimes \mathfrak{m}_{\mathcal{S}} \to \Gamma(\mathfrak{g}^1) \otimes \Omega_{\mathcal{S}/B}$$

be the map induced by exterior derivative on  $\mathfrak{m}_{\mathcal{S}}$ . Set

(3) 
$$\psi = d_{\mathcal{S}}(\phi).$$

Then

$$AC(\phi) = [\psi].$$

Note that differentiating the integrability condition for  $\phi$  yields

$$\partial \psi = -[\phi, \psi].$$

Since  $(\mathfrak{g}, \partial + \mathrm{ad}(\phi))$  is a resolution of  $\mathfrak{g}^{\phi}$ , this means that  $\psi$  is a cocycle for  $\mathfrak{g}^{\phi}$ .

**Example 1.** Let  $X_B/B$  be a family of complex manifolds and let E be a vector bundle on  $X_B$  with a  $\mathfrak{g}$ -structure. To recall what that means, let  $G(E) = ISO(\mathbb{C}^r, E)$ ,  $r = \operatorname{rk}(E)$  be the associated principal bundle, i.e. the open subset of the geometric vector bundle  $\mathfrak{hom}(\mathbb{C}^r, E)$  consisting of fibrewise isomorphisms, with the obvious action of  $GL_r$ . Let  $\mathfrak{D}(E)$  be the sheaf of  $GL_r$ -invariant vector fields on G(E), which may also be identified as the sheaf of relative derivations of  $(E, \mathcal{O}_X)$ , consisting of pairs  $(v, a), v \in T_{X_B/B}, a \in Hom_{\mathbb{C}}(E, E)$  such that

$$a(fe) = fa(e) + v(f)e, \ \forall f \in \mathcal{O}_X, e \in E.$$

Note that  $\mathfrak{D}(E)$  is an extension of Lie algebras

(4) 
$$0 \to \mathfrak{gl}(E) \to \mathfrak{D}(E) \to T_{X_B/B} \to 0$$

Then a  $\mathfrak{g}$ -structure on E is a Lie subalgebra  $\hat{\mathfrak{g}} \subseteq \mathfrak{D}(E)$  which fits in an exact sequence

Note that in this case a maximal integral submanifold  $\hat{G}$  of  $\hat{\mathfrak{g}}$  yields a principal subbundle of G(E) with structure group  $G = \exp(\mathfrak{g})$  and conversely such a principal subbundle with Lie algebra  $\mathfrak{g}$  yields a  $\mathfrak{g}$ -structure. Now let  $P^m$  be the m-th neighborhood of the diagonal in  $X_B \times_B X_B$ , viewed as an  $\mathcal{O}_X$ -algebra via the 1st projection  $p_1$ . Then clearly a  $\mathfrak{g}$ -structure on E yields a structure of  $\mathfrak{g}$ -deformation on  $P^m(E) = P^m \otimes p_2^*(E)$ , parametrized by  $P^m$ , for any m, and as above this admits a good (Dolbeault) resolution. We denote this deformation by  $P^m(E,\mathfrak{g})$ .

In particular, taking

$$\mathcal{S} = P_X^1 = \mathcal{O}_{X \times_B X} / \mathcal{I}_{\Delta_X}^2 = \mathcal{O}_X \oplus \Omega_{X_B/B}$$

(the standard 1st order deformation of  $\mathcal{O}_X$ ), we get a first-order relative  $\mathfrak{g}$ -deformation  $P^1(E,\mathfrak{g})$  parametrized by  $\mathcal{S}$ . Note that in this case  $\Omega_{\mathcal{S}/B} = \Omega_{X_B/B}$ and its  $\mathcal{S}$ -module structure factors through  $\mathcal{O}_X$ . Thus the Atiyah-Chern class

$$AC(P^1(E,\mathfrak{g})) \in H^1(\mathfrak{g} \otimes \Omega_{X_B/B})$$

and it is easy to see that it coincides with the usual Atiyah-Chern class of the  $\mathfrak{g}$ -structure E which may be defined, e.g. differential-geometrically in terms of a  $\mathfrak{g}$ -connection (and which reduces to the usual Atiyah-Chern class if  $\mathfrak{g} = \mathfrak{gl}(E)$ , cf. [1]). Indeed our good resolution in this case takes the form

$$E^0 \otimes (\mathcal{O}_X \oplus \Omega_{X_B/B}) \to E^1 \otimes (\mathcal{O}_X \otimes \Omega_{X_B/B}).$$

with differential

$$\begin{pmatrix} \bar{\partial} & \phi \\ 0 & \bar{\partial} \end{pmatrix}$$

and note that in this case  $\phi = \psi$  since  $\mathfrak{m}_{\mathcal{S}} = \Omega_{\mathcal{S}}$ . Assuming E is endowed with a  $\bar{\partial}$ - connection, the parallel lift of a section e of E to  $E^0 \otimes (\mathcal{O}_X \oplus \Omega_{X_B/B})$  is given by  $(e, \nabla e)$  and consequently we have

$$\phi(e) = [\bar{\partial}, \nabla](e).$$

Thus

(5) 
$$\psi = [\bar{\partial}, \nabla]$$

In other words, for any section v of  $T_{X_B/B}$ , holomorphic or not, we have

$$\psi \neg v = [\bar{\partial}, \nabla_v]$$

**Example 2.** Consider an ordinary first-order deformation  $\phi$  of a complex manifold X, corresponding to an algebra S of exponent 1. Suppose this deformation comes from a geometric family

$$\pi: \mathcal{X} \to Y$$

with  $\mathcal{X}, Y$  smooth,  $S = \mathcal{O}_{Y,0}/\mathfrak{m}_{Y,0}^2$ . Then it is easy to see that  $AC(\phi)$  corresponds to the extension

$$0 \to T_X \to \mathfrak{D}_\pi \to T_0 Y \otimes \mathbb{C}_X \to 0$$

where  $\mathfrak{D}_{\pi}$  is the subsheaf of  $T_{\mathcal{X}} \otimes \mathcal{O}_X$  consisting of 'vector fields locally constant in the normal direction', i.e. those derivations  $\mathcal{O}_{\mathcal{X}} \to \mathcal{O}_X$  that preserve the subsheaf  $\pi^{-1}\mathcal{O}_Y \subset \mathcal{O}_X$ . Ran

The last example suggests an interpretation of the Atiyah class as an extension also in the general case. To state this, let  $\phi$  be a relative deformation parametrized by S as above, and set

$$I = \operatorname{Ann}(\Omega_{\mathcal{S}/B}) \subset \mathcal{S}, \mathcal{S}' = \mathcal{S}/I, \phi' = \phi \otimes_{\mathcal{S}} \mathcal{S}'$$

and let  $\Omega^{vv}_{\mathcal{S}/B}$  denote the double dual as  $\mathcal{S}'$ -module. Note that

$$\Omega^{vv}_{\mathcal{S}/B} = \mathrm{Der}_{\mathcal{O}_B}(\mathcal{S}, \mathcal{S}')^v$$

( dual as left  $\mathcal{S}'$ -module).

We will also consider the analogous situation over a formally smooth, complete noetherian augmented local  $\mathcal{O}_B$ -algebra  $\mathcal{S}^{\wedge}$  (which is thus locally a power series algebra over  $\mathcal{O}_B$ ), where of course dual means as (left)  $\mathcal{S}^{\wedge}$ -module.

**Theorem**. (a) Let S be an augmented  $\mathcal{O}_B$ -algebra of finite type as  $\mathcal{O}_B$ -module, and let  $\phi$  be a relative  $\mathfrak{g}_B$ -deformation parametrized by S. Then the image of  $AC(\phi)$  in  $H^1(\mathfrak{g}_B^{\phi'} \otimes \Omega_{S/B}^{vv})$  corresponds to an extension of S' modules

(6) 
$$0 \to \mathfrak{g}_B^{\phi'} \to \mathfrak{D}(\phi) \to f^{-1} \mathrm{Der}_{\mathcal{O}_B}(\mathcal{S}, \mathcal{S}') \to 0$$

and for any admissible  $\mathfrak{g}_B$ -module  $E_B$  there is a natural action pairing

$$\mathfrak{D}(\phi) \times E_B^{\phi} \to E_B^{\phi'}.$$

(b) If  $\phi^{\wedge}$  is a relative  $\mathfrak{g}_{B}$ -deformation parametrized by a formally smooth noetherian  $\mathcal{O}_{B}$ -algebra  $\mathcal{S}^{\wedge}$ , then the image of  $AC(\phi^{\wedge})$  in  $H^{1}(\mathfrak{g}_{B}^{\phi^{\wedge}} \otimes \Omega_{\mathcal{S}^{\wedge}/B}^{vv})$  corresponds to an extention of  $\mathcal{S}^{\wedge}$ -Lie algebras

(7) 
$$0 \to \mathfrak{g}_B^{\phi^{\wedge}} \to \mathfrak{D}(\phi^{\wedge}) \xrightarrow{\nu} f^{-1}T_{\mathcal{S}^{\wedge}} \to 0$$

where  $T_{S^{\wedge}} = \text{Der}_{\mathcal{O}_B}(S^{\wedge}, S^{\wedge})$ , and for any admissible  $\mathfrak{g}_B$ -module  $E_B$ ,  $\mathfrak{D}(\phi^{\wedge})$  acts on  $E_B^{\phi^{\wedge}}$  satisfying the rule

(8) 
$$d(f.v) = f.d(v) + \nu(d)(f).v, \forall d \in \mathfrak{D}(\phi^{\wedge}), f \in \mathcal{S}^{\wedge}, v \in E_B^{\phi^{\wedge}}$$

**Proof.** For brevity we shall work out the formal case, the artinian case being similar. For convenience, we will drop the *B* substript. We let  $(\mathfrak{g}, E)$  be a pair (differential graded Lie algebra, differential graded module) forming a soft resolution of  $(\mathfrak{g}, E)$ ; also let  $(C \cdot, \partial)$  be a soft resolution of  $f^{-1}\mathcal{O}_B$ , and note that  $\mathfrak{g}$  is a *C*-module. Then clearly  $\mathfrak{D}(\phi^{\wedge})$ , i.e. the extension corresponding to  $AC(\phi^{\wedge})$  is resolved by the complex

$$\mathfrak{D}^{\cdot}(\phi^{\wedge}) = \mathfrak{g}^{\cdot} \otimes \mathcal{S}^{\wedge} \oplus C^{\cdot} \otimes T_{\mathcal{S}^{\wedge}}$$

with differential given by the matrix

$$\begin{pmatrix} \partial + \phi^{\wedge} & \psi^{\wedge} \\ 0 & \partial \end{pmatrix}$$

Ran

where  $\psi^{\wedge} = d_{\mathcal{S}^{\wedge}}(\phi^{\wedge})$  as in (3), which defines in an obvious way a map  $C^i \otimes T_{\mathcal{S}^{\wedge}} \to \mathfrak{g}^{i+1} \otimes \mathcal{S}^{\wedge}$ .

Now we claim that  $\mathfrak{D}^{\cdot}(\phi^{\wedge})$  is a differential graded Lie algebra: indeed since  $\mathfrak{g} \otimes \mathcal{S}^{\wedge}$  and  $T_{\mathcal{S}^{\wedge}} \otimes C^{\cdot}$  with the induced differentials are clearly differential graded Lie algebra's (in the latter case, the bracket is induced by that of  $T_{\mathcal{S}^{\wedge}}$ ), and  $T_{\mathcal{S}^{\wedge}} \otimes C^{\cdot}$  acts on  $\mathfrak{g} \otimes \mathcal{S}^{\wedge}$  via the action of  $T_{\mathcal{S}^{\wedge}}$  on  $\mathcal{S}^{\wedge}$  and the  $C^{\cdot}$ -module structure of  $\mathfrak{g}^{\cdot}$ , it suffices to show that  $\psi^{\wedge}$  is a derivation, which comes from the following calculation:

$$\psi^{\wedge}([v_1, v_2]) = [v_1, v_2](\phi^{\wedge}) = v_1(v_2(\phi^{\wedge})) - v_2(v_1(\phi^{\wedge}))$$
$$= v_1(\psi^{\wedge}(v_2)) - v_2(\psi^{\wedge}(v_1)).$$

Now since  $\mathfrak{D}^{\cdot}(\phi^{\wedge})$  is a differential graded Lie algebra, the fact that it acts on  $E^{\phi^{\wedge}}$  follows from the fact that the differential of  $\mathfrak{D}^{\cdot}(\phi^{\wedge})$  is just commutator with the differential on the resolution of  $E^{\phi^{\wedge}}$ , i.e.  $\partial + \phi^{\wedge}$ . To check the latter, it is firstly clear on the  $\mathfrak{g}^{\cdot} \otimes \mathcal{S}^{\wedge}$  summand; for the other summand, take  $v \in T_{\mathcal{S}^{\wedge}} \otimes C^{\cdot}$ . Then

$$[v, \partial + \phi^{\wedge}] = [v, \partial] + [v, \phi^{\wedge}] = \partial(v) + \psi^{\wedge}(v)$$

This shows that the obvious term-by-term pairing induces a pairing of complexes

$$\mathfrak{D}^{\cdot}(\phi^{\wedge}) \times (E^{\cdot}, \partial + \phi^{\wedge}) \to (E^{\cdot}, \partial + \phi^{\wedge}),$$

whence a pairing  $\mathfrak{D}(\phi^{\wedge}) \times E^{\phi^{\wedge}} \to E^{\phi^{\wedge}}$ ; that this is in fact a Lie action is clear from the fact that the corresponding assertion holds term-by-term. This completes the proof.

**Example 2 bis.** If  $X_B/B$  is a family of complex manifolds and  $\mathfrak{g}_B = T_{X/B}$  is the vertical tangent algebra acting on  $\mathcal{O}_X$ ,  $\phi^{\wedge}$  is a deformation of complex structure corresponding geometrically to  $\mathcal{X} \to \operatorname{Spec}(\mathcal{S})$ ,  $\mathfrak{g}^{\phi^{\wedge}}$  is just the vertical tangent algebra  $T_{\mathcal{X}/\mathcal{S}}$ , i.e. the derivations of  $\mathcal{O}_{\mathcal{X}}$  killing  $\mathcal{S}$  (cf. [6]), and  $\mathfrak{D}(\phi^{\wedge})$  is the algebra of derivations of  $\mathcal{O}_{\mathcal{X}}$  leaving  $\mathcal{S}$  invariant.

As a specific example, we consider a so-called *Schiffer variation*, where X is a compact Riemann surface,  $S = \mathbb{C}[[s]]$ , and for some point  $p \in X$ , we pick a local coordinate z centered at p and set

$$\phi = s \frac{\partial}{\partial z} d\bar{z}$$

(the corresponding 1st order deformation corresponds to the bicanonical image of p under the identification  $H^1(T_X)^* = H^0(K_X^{\otimes 2})$ ). This yields a formal deformation  $\mathcal{X}/\mathcal{S}$ , where a holomorphic function near p locally has the form  $A(z - s\bar{z})$  where A is a power series with coefficients in  $\mathcal{S}$ . A local holomorphic generator of the relative tangent algebra  $T_{\mathcal{X}/\mathcal{S}}$  is

$$v = \frac{1}{1 - s\bar{s}} \Big( \frac{\partial}{\partial z} + \bar{s} \frac{\partial}{\partial \bar{z}} \Big).$$

502

Ran

We seek a lift of  $\frac{\partial}{\partial s}$  that is holomorphic and horizontal, i.e. kills both  $z - s\bar{z}$  and its conjugate. By a direct computation, such a horizontal lift is given by

$$w = \frac{\partial}{\partial s} + \bar{z}\frac{\partial}{\partial z} + s\bar{s}v.$$

Then locally, the Lie algebra  $\mathfrak{D}(\phi)$  is determined as  $\mathcal{O}_{\mathcal{X}} v \oplus \mathcal{S} w$  with [v, w] = 0, v annihilating  $\mathcal{S}$  and w acting on  $\mathcal{O}_{\mathcal{X}}$  in the natural way. Globally, one could either glue these local data with a Čech twist corresponding to  $\phi$ , or take the standard differential graded Lie algebra Dolbeault resolution of  $\mathfrak{D}(\phi)$  and twist the differential by adding  $\phi$ .

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Ziv Ran University of California, Riverside CA 92521, USA ziv.ran@ucr.edu

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