# A HADAMARD TYPE THEOREM FOR THE STATIC SPACE-TIME 

Walter Frattarolo


#### Abstract

Necessary and sufficient conditions on the static space-time are determined, in order that the exponential function $\exp _{p}$ be a diffeomorphism.


It is well known that the Hopf-Rinow theorem states that a connected and complete Riemannian manifold is geodesically connected.

We can express this concept saying that, if $M_{0}$ is such a manifold, then the exponential map

$$
\exp _{p}: T_{p} M_{0} \rightarrow M_{0}
$$

is onto, for every $p \in M_{0}$.
Hadamard theorem goes beyond this result and shows that, if $M_{0}$ is simply connected and complete and if its sectional curvature is $k \leq 0$, then the map $\exp _{p}$ is a diffeomorphism, for every $p \in M_{0}$; in particular:
i) There are not couples of conjugate points on $M_{0}$;
ii) $M_{0}$ is diffeomorphic to $\mathbb{R}^{n}\left(n=\operatorname{dim} M_{0}\right)$;
iii) For any couple of points $p, q \in M$, there exists a unique geodesic

$$
\gamma: \mathbb{R} \rightarrow M_{0}
$$

such that

$$
(\gamma(0)=p ; \quad \gamma(1)=q)
$$

The property (iii) above implies not only the existence of $\gamma$, that is to claim $\exp _{p}$ is onto, but also the uniqueness of it, that is: $\exp _{p}$ is one to one.

No result like Hadamard theorem is known about semi-Riemannian manifolds; on the contrary, there are counter-examples to the Hopf-Rinow theorem in the case of static Lorentz manifolds (e.g.: see the anti-De Sitter space, on [9], [13]).

[^0]However sufficient conditions which guarantee the validity of the Hopf-Rinow theorem for static space-times, have been established in [7].

The Lorentz metric $g=g(P)$ considered there applies to a manifold $M$ of the following form:

$$
M=M_{0} \times \mathbb{R}
$$

where $M_{0}$ is furnished with a Riemannian metric (index $=0$ ) denoted by: $\langle\cdot, \cdot\rangle_{M_{0}}$ and $g(P)$ is defined by:

$$
g(P)[(\zeta, \tau),(\zeta, \tau)]=\langle\zeta, \zeta\rangle-\beta\left(P_{0}\right) \tau^{2}
$$

for every $P=\left(P_{0}, t_{0}\right) \in M$ and $(\zeta, \tau) \in T_{p} M$; here $\beta$ is a positive real function on $M_{0}$.

In the following $g(P)$ will be denoted by $\langle\cdot, \cdot\rangle_{M}$, if the context is clear.
In this work we examine necessary and sufficient conditions on the metric $g(P)$ for the function $\exp _{p}$ to be a diffeomorphism, for any $P \in M$. Such conditions, on the contrary of what might be expected, are not a straightforward generalization of the Hadamard ones. Before showing our results, we recall a result of [7]. Some notation now.

Let $P=\left(x_{0}, t_{0}\right), Q=\left(x_{1}, t_{1}\right)$ be two events in $M$. We shall consider loop spaces $\Omega^{1}$ on $M_{0}$ and $M$, defined in the following way:
$\Omega^{1}=\Omega^{1}\left(M_{0}, x_{0}, x_{1}\right)=\left\{x:[0,1] \rightarrow M_{0}: \quad x\right.$ absolutely continuous and such that

$$
\left.\int_{0}^{1}\langle\dot{x}, \dot{x}\rangle_{M_{0}} d s<+\infty, \quad x(0)=x_{0}, x(1)=x_{1}\right\}
$$

$\Omega^{1}$ is a Riemannian manifold, modelled on Sobolev spaces, of curves in $M_{0}$.
The geodesic on $M$, joining $P$ and $Q$, are the critical points of the action functional $f$, defined by:

$$
f(\gamma)=\int_{0}^{1}\langle\dot{\gamma}(s), \dot{\gamma}(s)\rangle_{M} d s
$$

In [7] the following variational principle has been proved:
Let $M=M_{0} \times \mathbb{R}$ be a static space-time with the metric $\langle\cdot, \cdot\rangle_{M}$.
Let $P=\left(x_{0}, t_{0}\right), Q=\left(x_{1}, t_{1}\right)$ two points in $M$.
A curve $\gamma(s)=(x(s), t(s))$ is a critical point of $f$ (i.e.: a geodesic) if and only if:
.) $x=x(s)$ is a critical point of the functional $J$ on $\Omega^{1}$, defined by:

$$
J(x)=\int_{0}^{1}\langle\dot{x}, \dot{x}\rangle_{M_{0}} d s-\frac{\Delta^{2}}{\int_{0}^{1} \frac{d s}{\beta(x)}}
$$

where: $\Delta=t_{1}-t_{0}$;
..) $t=t(s)$ is the function such that:

$$
t(0)=0 \quad \text { and } \quad \dot{t}(s)=\Delta\left(\int_{0}^{1}(1 / \beta(x)) d s\right)^{-1}(1 / \beta(x(s)))=\Phi(x)
$$

Moreover if $\gamma(s)=(x(s), t(s))$ is a geodesic, then we have:

$$
f(\gamma)=J(x)
$$

This variational principle allows to use the Riemannian techniques for problems over the manifold $M$ and allows to apply the Lusternik and Schnirelman critical point theory too.

We want also to recall a result obtained in [7] which is of interest for what follows.

Using the same symbols as above, if $M_{0}$ is connected and complete and if $\beta \in C^{3}\left(M_{0} ; \mathbb{R}^{+}\right)$is a function limited from above and from below by positive constants, then the Lorentz manifold $(M, g)$ is geodesically connected.

In the sequel we shall write simply:

$$
J(x)=\Phi_{1}(x)-\Delta^{2} \Phi_{2}(x)
$$

where:

$$
\begin{aligned}
\Phi_{1}(x) & =\int_{0}^{1}\langle\dot{x}, \dot{x}\rangle_{M_{0}} d s \\
\Phi_{2}(x) & =\frac{1}{\int_{0}^{1} \frac{d s}{\beta(x)}}
\end{aligned}
$$

Our main result is the following.
Theorem (1). Let $M, M_{0}, \mathbf{R}, g$, be defined as above. Suppose that:
i) $M_{0}$ is simple connected and complete;
ii) $\beta \in C^{3}\left(M_{0} ; \mathbb{R}^{+}\right)$is a function limited from above and from below by positive constants;
iii) the sectional curvature on $M_{0}$ is $k \leq 0$.

Under such assumptions, if the function $\exp _{p}: T_{p} M \rightarrow M$ is a diffeomorphism for any point $P \in M$, then the Riemannian Hessian $H_{R}^{\beta}\left(P_{0}\right)$ is a negative semidefinite quadratic form, for any point $P_{0} \in M_{0}$.

Viceversa, if the Riemannian Hessian $H_{R}^{\beta}\left(P_{0}\right)$ is a negative definite quadratic form for any point $P_{0} \in M_{0}$, then the function $\exp _{p}$ is a diffeomorphism for any point $P \in M$.

Now we get to prove Theorem (1) and some complementary results.
Proposition (2). Let : $\beta \in C^{2}\left(M_{0} ; \mathbb{R}^{+}\right) ; P_{0} \in M_{0} ; R$ be the curvature tensor on $M$. Let $\tau_{1}$ be the unitary vector field along $\mathbb{R}$, tangent to $\mathbb{R}$. The following results are valid:
i) $\left\langle R_{v, \tau_{1}} v, \tau_{1}\right\rangle_{\left(P_{0}, t_{0}\right)}=\frac{1}{2} H_{R}^{\beta}\left(P_{0}\right)[v, v]-\left(\langle v, \operatorname{grad} \sqrt{\beta}\rangle_{P_{0}}\right)^{2}, \forall t_{0} \in \mathbb{R}, \forall v \in T_{P_{0}} M_{0}$.
ii) $H_{R}^{\beta}\left(P_{0}\right)$ is a negative semidefinite form iff:

$$
\left\langle R_{v, \tau_{1}} v, \tau_{1}\right\rangle_{\left(P_{0}, t_{0}\right)} \leq-\left(\langle\operatorname{grad} \sqrt{\beta}, v\rangle_{P_{0}}\right)^{2},
$$

$$
\forall t_{0} \in R, \forall v \in T_{P_{O}} M_{0} .
$$

Proof: (ii) is a straightforward consequence of (i). In order to prove (i), we recall that:

$$
\begin{aligned}
\left\langle R_{v, \tau_{1}}, v, \tau_{1}\right\rangle_{\left(P_{0}, t_{0}\right)} & =\frac{-1}{\sqrt{\beta\left(P_{0}\right)}} H_{R}^{\sqrt{\beta}}\left(P_{0}\right)[v, v] \cdot\left\langle\tau_{1}, \tau_{1}\right\rangle_{M} \\
& =\left(\frac{-1}{\sqrt{\beta\left(P_{0}\right)}} H_{R}^{\sqrt{\beta}}\left(P_{0}\right)[v, v]\right) \cdot\left(-\beta\left(P_{0}\right)\right) \\
& =\sqrt{\beta\left(P_{0}\right)} H_{R}^{\sqrt{\beta}}\left(P_{0}\right)[v, v], \quad \forall v \in T_{P_{0}} M_{0} .
\end{aligned}
$$

Now we observe that:

$$
\begin{align*}
& H_{R}^{\sqrt{\beta}}\left(P_{0}\right)[v, v]=\left\langle D_{v}(\operatorname{grad} \sqrt{\beta}), v\right\rangle_{P_{0}}=  \tag{2}\\
& =\frac{1}{2}\left\langle D_{v}\left(\frac{1}{\sqrt{\beta}} \operatorname{grad} \beta\right), v\right\rangle_{P_{0}} \\
& =\frac{1}{2}\left\langle\left(\left\langle v, \operatorname{grad} \frac{1}{\sqrt{\beta(P)}}\right\rangle_{M_{0}}\right) \operatorname{grad} \beta+\frac{1}{\sqrt{\beta(P)}} D_{v}(\operatorname{grad} \beta), v\right\rangle_{P_{0}} \\
& =\left(-\frac{1}{4}\right) \frac{1}{\beta\left(P_{0}\right) \sqrt{\beta\left(P_{0}\right)}}\left(\langle v, \operatorname{grad} \beta\rangle_{P_{0}}\right)^{2}+\frac{1}{2 \sqrt{\beta\left(P_{0}\right)}}\left\langle D_{v}(\operatorname{grad} \beta), v\right\rangle_{P_{0}} \\
& =\left(-\frac{1}{4}\right) \frac{1}{\beta\left(P_{0}\right) \sqrt{\beta\left(P_{0}\right)}}\left(\langle v, \operatorname{grad} \beta\rangle_{P_{0}}\right)^{2}+\frac{1}{2 \sqrt{\beta\left(P_{0}\right)}} H_{R}^{\beta}\left(P_{0}\right)[v, v],
\end{align*}
$$

$$
\forall v \in T_{P_{0}} M_{0} .
$$

(i) follows from (1) and (2).

Lemma (3). Assume:
i) $M_{0}$ is simply connected and complete;
ii) $\beta \in C^{3}\left(M_{0} ; \mathbb{R}^{+}\right)$is a function limited from above and from below by positive constants;
iii) the sectional curvature on $M$ is $K \leq 0$.

If the Riemannian Hessian $H_{R}^{\beta}\left(P_{0}\right)$ is a negative definite quadratic form for any point $P_{0} \in M_{0}$, then the quadratic form $J^{\prime \prime}(x)$ is positive definite for any curve $x=x(s)(s \in I)$ which is a critical point of the functional $J$ and for any couple of extreme points $x(0), x(1) \in M_{0}$.

Proof: Let $x$ be a critical point of $J$. We shall show:

1) $J^{\prime \prime}(x)$ is a positive definite form if $\Phi_{2}^{\prime \prime}(x)$ is negative definite;
2) $\Phi_{2}^{\prime \prime}(x)$ is negative definite if the integral functional:

$$
I(x)[v, v]=\int_{0}^{1} H_{R}^{\beta}(x)[v, v] d s
$$

whose domain is the set of vector fields with compact support, tangent to $M_{0}$ along $x$, is a negative definite form.

If $v=v(s)$ is a vector field having compact support and tangent to $M_{0}$ along $x$, we have:

$$
\begin{gathered}
.) \Phi_{1}(x)=\int_{0}^{1}\langle\dot{x}, \dot{x}\rangle_{x(s)} d s \\
. .) \quad \Phi_{1}^{\prime \prime}(x)[v, v]=2 \int_{0}^{1}\left\langle D_{s} v, D_{s} v\right\rangle_{x(s)} d s-2 \int_{0}^{1}\left\langle R_{\dot{x}, v} \dot{x}, v\right\rangle_{x(s)} d s
\end{gathered}
$$

Besides that, we have:

$$
\begin{aligned}
& .) \quad \Phi_{2}(x)=\frac{1}{\int_{0}^{1} \frac{d s}{\beta(x(s))}}, \\
& . .) \quad \Phi_{2}^{\prime}(x)[v]=\frac{1}{\left(\int_{0}^{1} \frac{d s}{\beta(x(s))}\right)^{2}} \int_{0}^{1} \frac{\langle\operatorname{grad} \beta, v\rangle_{x(s)}}{\beta^{2}(x(s))} d s \\
& \cdots) \\
& +\frac{1}{\left(\int_{0}^{1} \frac{d s}{\beta(x(s))}\right)^{2}} \int_{0}^{1}(-2) \frac{\left(\langle\operatorname{grad} \beta, v\rangle_{x(s)}\right)^{2}}{\beta^{3}(x(s))} d s+\frac{2}{\left(\int_{0}^{1} \frac{d s}{\beta(x(s))}\right)^{3}}\left(\int_{0}^{1} \frac{\langle\operatorname{grad} \beta, v\rangle_{x(s)}}{\beta^{2}(x(s))} d s\right)^{2} \\
& \left(\int_{0}^{1} \frac{d s}{\beta(x(s))}\right)^{2}
\end{aligned} \int_{0}^{1} \frac{H_{R}^{\beta}(x)[v, v]}{\beta^{2}(x(s))} d s . .
$$

Now (1) comes straight, because:

$$
J^{\prime \prime}(x)=\Phi_{1}^{\prime \prime}(x)-\Delta \Phi_{2}^{\prime \prime}(x)
$$

Also (2) is soon found, as a consequence of the following

$$
\begin{aligned}
& \frac{2}{\left(\int_{0}^{1} \frac{d s}{\beta(x(s))}\right)^{3}}\left(\int_{0}^{1} \frac{\langle\operatorname{grad} \beta, v\rangle_{x(s)}}{\beta^{2}(x(s))} d s\right)^{2}-\frac{2}{\left(\int_{0}^{1} \frac{d s}{\beta(x(s))}\right)^{2}} \int_{0}^{1} \frac{\left(\langle\operatorname{grad} \beta, v\rangle_{x(s)}\right)^{2}}{\beta^{3}(x(s))} d s= \\
& =\frac{2}{\left(\int_{0}^{1} \frac{d s}{\beta(x(s))}\right)^{2}}\left\{\frac { 1 } { \int _ { 0 } ^ { 1 } \frac { d s } { \beta ( x ( s ) ) } } \left(\int_{0}^{1} \frac{\left.\langle\operatorname{grad} \beta, v\rangle_{x(s)}^{\beta^{1 / 2}(x(s)) \beta^{3 / 2}(x(s))}\right)^{2}+}{} \begin{array}{l}
\left.\quad-\int_{0}^{1} \frac{\left(\langle\operatorname{grad} \beta, v\rangle_{x(s)}\right)^{2}}{\beta^{3}(x(s))} d s\right\} \leq \\
\leq \frac{2}{\left(\int_{0}^{1} \frac{d s}{\beta(x(s))}\right)^{2}}\left\{\left(\int_{0}^{1} \frac{d s}{\beta(x(s))} \int_{0}^{1} \frac{\left(\langle\operatorname{grad} \beta, v\rangle_{x(s)}\right)^{2}}{\beta^{3}(x(s))} d s\right) \cdot \frac{1}{\int_{0}^{1} \frac{d s}{\beta(x(s))}}+\right. \\
\left.\quad-\int_{0}^{1} \frac{\left(\langle\operatorname{grad} \beta, v\rangle_{x(s)}\right)^{2}}{\beta^{3}(x(s))} d s\right\}=0 .
\end{array}\right.\right.
\end{aligned}
$$

Lemma (4). Assume:
i) $M_{0}$ is simply connected and complete;
ii) $\beta \in C^{3}\left(M_{0} ; \mathbb{R}^{+}\right)$is a function limited from above and from below by positive constants;
iii) the sectional curvature on $M_{0}$ is $k \leq 0$.

If the quadratic form $J^{\prime \prime}(x)$ is positive definite, for any curve $x=x(s)(s \in I)$ which is a critical point of the functional $J$, and for any couple of extreme points $x(0), x(1) \in M_{0}$, then the Riemannian Hessian $H_{R}^{\beta}\left(P_{0}\right)$ is a negative semidefinite quadratic form, for any point $P_{0} \in M_{0}$.

Proof: Let $P_{0} \in M_{0}$. By [7], there exists a geodesic $\gamma(s)=(x(s), t(s))$, $s \in I$, contained in the manifold $M$ and starting from $P_{0}$. Such geodesic depends on Cauchy data for a system of second order differential equations, then it depends on the point $P_{0}$, on the initial instant $t(0)=t_{0}$ and on the initial velocity $\dot{\gamma}(0)=(\dot{x}(0), \dot{t}(0))$. Now, as we are dealing with a static metric, it results (see
introduction):

$$
\begin{gathered}
t(s)=t(0)+\frac{\int_{0}^{s} \frac{d \rho}{\beta(x(\rho))}}{\frac{d \rho}{\beta(x(\rho))}} \Delta, \quad\left\{\begin{array}{l}
s \in I \\
\Delta=t(1)-t(0),
\end{array}\right. \\
\dot{t}(s)=\frac{1}{\int_{0}^{1} \frac{d \rho}{\beta(x)}} \cdot \frac{\Delta}{\beta(x(s))},
\end{gathered}
$$

then $\dot{t}(s)$ depends on $x(s)$ and on $\Delta$; therefore $\gamma(s)$ depends on the vector $\mu_{1}=\dot{x}(0)$, on $\Delta$ and on the point $\left(P_{0}, t_{0}\right) \in M$.

Fixing $t(0)=t_{0}$ and $\Delta$, we fix the temporal extrema of the geodesic arc $\gamma(s)$, $s \in I$, whereas fixing $P$ and $\mu_{1}$, we fix its spatial extrema; the vector $\mu(0)=\mu_{0}$, such that:

$$
\mu_{1}=\dot{x}(0)=\|\dot{x}(0)\|_{T_{P_{0}}} M_{0} \mu_{0}
$$

determines the direction of $\gamma$.
Stated that, let $\mu_{1}=\dot{x}(0)$ and $\Delta$ be fixed and $\gamma=\gamma_{\mu_{1} ; \Delta}$ be the corresponding geodesic, with extreme points:

$$
\begin{aligned}
& P=\gamma(0)=(x(0), t(0)), \quad x(0)=P_{0}, \\
& Q=\gamma(1)=((x(1), t(1)) .
\end{aligned}
$$

The curve $x=x(s), s \in I$ is a critical point of $J$. Let $\nu_{0} \in T_{P_{0}} M_{0}$ be a unitary vector and let:

$$
\left\{\nu=\nu(s) \mid \forall s \in I: \nu(s) \in T_{x(s)} M_{0} ;\|\nu(s)\|_{T_{x(s)} M_{0}}=1\right\}
$$

be the unitary vector field, obtained by parallel translation of $\nu_{0}$ along $x$. We consider now a sequence of functions:

$$
\begin{align*}
& s: I \rightarrow \mathbb{R} \quad k \in \mathbb{N}  \tag{1}\\
& s_{k}(\sigma)=\frac{1}{k(k+1)} \sigma+\frac{1}{k+1}
\end{align*}
$$

and set:

$$
\begin{aligned}
x_{k}(\sigma) & =x\left(s_{k}(\sigma)\right), \\
t_{k}(\sigma) & =t\left(s_{k}(\sigma)\right) ;
\end{aligned}
$$

from (1) it follows:

$$
\begin{aligned}
& \dot{x}_{k}(\sigma)=\left.\frac{1}{k(k+1)} \dot{x}(s)\right|_{s=s_{k}(\sigma)}, \\
& \dot{t}_{k}(\sigma)=\left.\frac{1}{k(k+1)} \dot{t}(s)\right|_{s=s_{k}(\sigma)} ;
\end{aligned}
$$

besides that, we have:

$$
\begin{aligned}
\Delta_{k} & =\int_{0}^{1} \dot{t}_{k}(\sigma) d \sigma=\int_{1 / k+1}^{1 / k}\left(\frac{\dot{t}(s)}{k(k+1)}\right) k(k+1) d s \\
& =\int_{1 / k+1}^{1 / k} \dot{t}(s) d s=\int_{1 / k+1}^{1 / k} \frac{c}{\beta(x)} d s \\
& =\frac{\int_{1 / k+1}^{1 / k} \frac{1}{\beta(x)} d s}{\int_{0}^{1} \frac{1}{\beta(x)} d s} \Delta
\end{aligned}
$$

here, we have used: $\beta(x) \dot{t}=c$ (see the introduction).
The curve $\gamma_{k}=\left(x_{k}, t_{k}\right)$ are monotone linear reparametrizations of pieces of $\gamma=(x, t)$, therefore they are all geodesics and the sequence $\left\{x_{k}\right\}$ is made of critical points of $J$.

Now we use an arbitrarily fixed function $\varphi \in C_{0}^{\infty}\left(I, \mathbb{R}^{+}\right)$and build up the following vector fields with compact support:

$$
\begin{cases}v_{k}(s)=\left.v_{k}(x(s))\right|_{s=s_{k}(\sigma)}=\varphi\left(\sigma_{k}(s)\right) \nu(s), & \frac{1}{k+1} \leq s \leq \frac{1}{k}, \\ v_{k}(s)=0, & s \in I \backslash\left[\frac{1}{k+1}, \frac{1}{k}\right]\end{cases}
$$

here, $\sigma_{k}$ is the inverse function of $s_{k}$.
We can also write: $v_{k}(\sigma)=\varphi(\sigma) \nu\left(s_{k}(\sigma)\right)$; then observe that:

$$
\begin{aligned}
\int_{0}^{1}\left\langle D_{\sigma} v_{k}, D_{\sigma} v_{k}\right\rangle_{x_{k}(\sigma)} d \sigma= & \int_{0}^{1}\left\langle D_{\dot{x}_{k}} v_{k}, D_{\dot{x}_{k}} v_{k}\right\rangle_{x_{k}(\sigma)} d \sigma= \\
= & \int_{0}^{1}\left\langleD _ { \dot { x } _ { k } } \left(\varphi(\sigma) \nu\left(s_{k}(\sigma)\right), D_{\dot{x}_{k}}\left(\varphi(\sigma) \nu\left(s_{k}(\sigma)\right)\right\rangle_{x_{k}(\sigma)} d \sigma\right.\right. \\
= & \int_{0}^{1}\left\langle\dot{\varphi}(\sigma) \nu\left(s_{k}(\sigma)\right)+\varphi(\sigma) D_{\dot{x}_{k}} \nu\left(s_{k}(\sigma)\right), \dot{\varphi}(\sigma) \nu\left(s_{k}(\sigma)\right)\right. \\
& \left.+\varphi(\sigma) D_{\dot{x}_{k}} \nu\left(s_{k}(\sigma)\right)\right\rangle_{x_{k}(\sigma)} d \sigma \\
= & \int_{0}^{1}\left(\frac{d \varphi}{d \sigma}\right)^{2} d \sigma=\frac{1}{k(k+1)} \int_{1 / k+1}^{1 / k}\left(\frac{d \varphi}{d s}\right)^{2} d s
\end{aligned}
$$

indeed:

$$
\begin{aligned}
\left\langle\dot{x}_{k}, \varphi\right\rangle_{x_{k}(\sigma)} & =\left\langle x_{k}^{*}\left(\frac{d}{d \sigma}\right), \varphi\right\rangle_{x_{k}(\sigma)} \\
& =\left\langle\frac{d}{d \sigma}, \varphi \circ x_{k}\right\rangle_{x_{k}(\sigma)}=\frac{d}{d \sigma} \varphi\left(x_{k}(\sigma)\right)=\frac{d \varphi(\sigma)}{d \sigma} .
\end{aligned}
$$

Therefore, we have (see Lemma (3)):

$$
\begin{aligned}
& \text { (3) } J^{\prime \prime}\left(x_{k}\right)\left[v_{k}, v_{k}\right]=\Phi_{1}^{\prime \prime}\left(x_{k}\right)\left[v_{k}, v_{k}\right]-\Delta_{k}^{2} \Phi_{2}^{\prime \prime}\left(x_{k}\right)\left[v_{k}, v_{k}\right]= \\
& =2 \int_{0}^{1}\left\langle D_{\sigma} v_{k}, D_{\sigma} v_{k}\right\rangle_{x_{k}(\sigma)} d \sigma-2 \int_{0}^{1}\left\langle R_{\dot{x}_{k}, v_{k}} \dot{x}_{k}, v_{k}\right\rangle_{x_{k}(\sigma)} d \sigma+ \\
& -\Delta_{k}^{2}\left\{\frac{2}{\left(\int_{0}^{1} \frac{d \sigma}{\beta\left(x_{k}\right)}\right)^{3}}\left(\int_{0}^{1} \frac{\left\langle\operatorname{grad} \beta, v_{k}\right\rangle_{x_{k}(\sigma)}}{\beta^{2}\left(x_{k}\right)} d \sigma\right)^{2}+\right. \\
& -\frac{2}{\left(\int_{0}^{1} \frac{d \sigma}{\beta\left(x_{k}\right)}\right)^{2}} \int_{0}^{1} \frac{\left(\left\langle\operatorname{grad} \beta, v_{k}\right\rangle_{x_{k}(\sigma)}\right)^{2}}{\beta^{3}\left(x_{k}\right)} d \sigma+ \\
& \left.+\frac{1}{\left(\int_{0}^{1} \frac{d \sigma}{\beta\left(x_{k}\right)}\right)^{2}} \int_{0}^{1} \frac{1}{\beta^{2}\left(x_{k}\right)} H_{R}^{\beta}\left(x_{k}\right)\left[v_{k}, v_{k}\right] d \sigma\right\}>0
\end{aligned}
$$

whence:

$$
\begin{align*}
& 2 \int_{0}^{1}\left(\frac{d \varphi}{d \sigma}\right)^{2} d \sigma-2 \int_{0}^{1}\left\langle R_{\dot{x}_{k}, v_{k}}, \dot{x}_{k}, v_{k}\right\rangle_{x_{k}(\sigma)} d \sigma+  \tag{4}\\
&-\left(\frac{\Delta_{k}}{\int_{0}^{1} \frac{d \sigma}{\beta\left(x_{k}\right)}}\right)^{2} \int_{0}^{1} \frac{1}{\beta^{2}\left(x_{k}\right)} H_{R}^{\beta}\left(x_{k}\right)\left[v_{k}, v_{k}\right] d \sigma> \\
&> 2\left(\frac{\Delta_{k}}{\int_{0}^{1} \frac{d \sigma}{\beta\left(x_{k}\right)}}\right)^{2}\left\{\frac{1}{\int_{0}^{1} \frac{d \sigma}{\beta\left(x_{k}\right)}}\left(\int_{0}^{1} \frac{\left\langle\operatorname{grad} \beta, v_{k}\right\rangle_{x_{k}(\sigma)}}{\beta^{2}\left(x_{k}\right)} d \sigma\right)^{2}+\right. \\
&\left.\quad-\int_{0}^{1} \frac{\left(\left\langle\operatorname{grad} \beta, v_{k}\right\rangle_{x_{k}(\sigma)}\right)^{2}}{\beta^{3}\left(x_{k}\right)} d \sigma\right\}
\end{align*}
$$

and using (2):
(5) $\quad 2 \int_{0}^{1}\left(\frac{d \varphi}{d \sigma}\right)^{2} d \sigma-2 \int_{0}^{1}\left\langle R_{\dot{x}_{k}, v_{k}}, \dot{x}_{k}, v_{k}\right\rangle_{x_{k}(\sigma)} d \sigma+$

$$
-\left(\frac{1}{k(k+1)}\right)^{2}\left(\frac{\Delta}{\int_{0}^{1} \frac{d s}{\beta(x)}}\right)^{2} \inf _{\sigma \in I}\left(\frac{1}{\beta^{2}\left(x_{k}\right)} H_{R}^{\beta}\left(x_{k}\right)[\nu, \nu]\right) \int_{0}^{1} \varphi^{2}(\sigma) d \sigma>
$$

$>2\left(\frac{1}{k(k+1)}\right)^{2}\left(\frac{\Delta}{\int_{0}^{1} \frac{d s}{\beta(x)}}\right)^{2}\left\{\frac{k(k+1)}{\int_{1 / k+1}^{1 / k} \frac{d s}{\beta(x)}}\left(\int_{1 / k+1}^{1 / k} \frac{\left\langle\operatorname{grad} \beta, v_{k}\right\rangle_{x(s)}}{\beta^{2}(x)} d s\right)^{2}+\right.$

$$
\left.-k(k+1) \int_{1 / k+1}^{1 / k} \frac{\left\langle\operatorname{grad} \beta, v_{k}\right\rangle_{x(s)}}{\beta^{3}(x)} d s\right\}
$$

Afterwards, we are going to apply the mean theorem for continuous functions to the second part of inequality (5); so we obtain:

$$
\begin{align*}
& 2 \int_{0}^{1}\left(\frac{d \phi}{d \sigma}\right)^{2} d \sigma-2 \int_{0}^{1}\left\langle R_{\dot{x}_{k}, v_{k}} \dot{x}_{k}, v_{k}\right\rangle_{x_{k}(\sigma)} d \sigma+  \tag{6}\\
& -\left(\frac{\Delta}{k(k+1) \int_{0}^{1} \frac{d s}{\beta(x)}}\right)^{2} \inf _{\frac{1}{k+1} \leq s \leq \frac{1}{k}}\left(\frac{1}{\beta^{2}(x)} H_{R}^{\beta}(x)[\nu, \nu]\right) \int_{0}^{1} \varphi^{2}(\sigma) d \sigma> \\
& >2\left(\frac{\Delta}{k(k+1) \int_{0}^{1} \frac{d s}{\beta(x)}}\right)^{2}\left\{\beta\left(x\left(s_{1}^{(k)}\right)\right) \frac{\left(\left\langle\operatorname{grad} \beta, v_{k}\right\rangle_{x\left(s_{2}^{(k)}\right)}\right)^{2}}{\beta^{4}\left(x\left(s_{2}^{(k)}\right)\right)}-\frac{\left(\left\langle\operatorname{grad} \beta, v_{k}\right\rangle_{x\left(s_{3}^{(k)}\right)}\right)^{2}}{\beta^{3}\left(x\left(s_{3}^{(k)}\right)\right)}\right\} ;
\end{align*}
$$

here, the $\left.s_{i} \in\right] \frac{1}{k+1}, \frac{1}{k}[$ are suitable points.
At last observe that, by the completeness of $M$, the geodesic $\gamma=\gamma(s)$ can be extended from $s \in I$ up to $s \in \mathbb{R}$; after making that, let's divide the inequality (6) by $\Delta^{2}$ and go to the limit for $(\Delta \rightarrow+\infty)$, then let's multiplicate the resulting inequality by $(k(k+1))^{2}$ and go again to the limit for $(k \rightarrow+\infty)$. We shall obtain the thesis of the lemma.

Proposition (5). We suppose that

1) $\gamma=(x, t)$ is a geodesic contained in $M$;
2) $f(\gamma)=\int_{0}^{1}\langle\dot{\gamma}, \dot{\gamma}\rangle_{M} d s$;
3) $\Phi(x)=t$ is the function obtained in [7] ( see the introduction);
let us denote
4) $v$ a vector field having compact support and tangent to $M_{0}$ along $x$;
5) $(w, \tau)$ a vector field having compact support and tangent to $M$ along $\gamma$.

Then:

$$
J^{\prime \prime}(x)[v, w]=f^{\prime \prime}(\gamma)\left[\left(v, \Phi^{\prime}(x)[v]\right),(w, \tau)\right]
$$

Proof: On smooth curves $y=y(s)$ contained in $M_{0}$, it is identically (see [7]):

$$
\begin{equation*}
\left.f_{t}^{\prime}(y, t)\right|_{t=\Phi(y)}=0 \tag{1}
\end{equation*}
$$

therefore we have identically:

$$
\frac{d}{d y} f_{t}^{\prime}(y, \Phi(y))=0
$$

that is, for any vector field $w=w(s)$ having compact support and tangent to $M_{0}$ along $y$, for any vertical vector field $\tau$ along $\Phi(y)$, having compact support, we have:

$$
\begin{equation*}
f_{t x}^{\prime \prime}(y, \Phi(y))[w, \tau]+f_{t t}^{\prime \prime}(y, \Phi(y))\left[\Phi^{\prime}(y)[w], \tau\right]=0 \tag{2}
\end{equation*}
$$

Now recall that:

$$
\begin{aligned}
f^{\prime \prime}(\gamma)[(v, & \left.\left.\tau_{1}\right),\left(w, \tau_{2}\right)\right]= \\
= & \int_{0}^{1}\left\{2\left\langle D_{s} v, D_{s} w\right\rangle_{M_{0}}-2\left\langle R_{\dot{x}, v} \dot{x}, w\right\rangle_{M_{0}}-H_{R}^{\beta}(x)[v, w] \dot{t}^{2}\right\} d s \\
& -\int_{0}^{1}\left\{2\langle\operatorname{grad} \beta, v\rangle_{M_{0}} \dot{t} \dot{\tau}_{2}+2\langle\operatorname{grad} \beta, w\rangle_{M_{0}} \dot{t} \dot{\tau}_{1}\right\} d s-\int_{0}^{1} 2 \beta(x) \dot{\tau}_{1} \dot{\tau}_{2} d s \\
= & f_{x x}^{\prime \prime}(\gamma)[v, w]+f_{x t}^{\prime \prime}(\gamma)\left[v, \tau_{2}\right]+f_{x t}^{\prime \prime}(\gamma)\left[w, \tau_{1}\right]+f_{t t}^{\prime \prime}(\gamma)\left[\tau_{1}, \tau_{2}\right] \\
= & \left(f_{x x}^{\prime \prime}(\gamma)[v, w]+f_{x t}^{\prime \prime}(\gamma)\left[w, \tau_{1}\right]\right)+\left(f_{x t}^{\prime \prime}(\gamma)\left[v, \tau_{2}\right]+f_{t t}^{\prime \prime}(\gamma)\left[\tau_{1}, \tau_{2}\right]\right)
\end{aligned}
$$

so that, (2) implies:

$$
\begin{align*}
& f^{\prime \prime}(x, \Phi(x))\left[\left(v, \Phi^{\prime}(x)[v]\right),\left(w, \tau_{2}\right)\right]=  \tag{3}\\
& \quad=f_{x x}^{\prime \prime}(x, \Phi(x))[v, w]+f_{x, t}^{\prime \prime}(x, \Phi(x))\left[w, \Phi^{\prime}(x)[v]\right]
\end{align*}
$$

Recall also that:

$$
\begin{aligned}
& J(x)= f(x, \Phi(x)) \\
&\cdot \cdot) \\
& J^{\prime}(x)[v]= f_{x}^{\prime}(x, \Phi(x))[v]+f_{t}^{\prime}(x, \Phi(x))\left[\Phi^{\prime}(x)[v]\right] \\
&\cdots) J^{\prime \prime}(x)[v, w]= f_{x x}^{\prime \prime}(x, \Phi(x))[v, w]+f_{x t}^{\prime \prime}(x, \Phi(x))\left[v, \Phi^{\prime}(x)[w]\right]+ \\
&+f_{x t}^{\prime \prime}(x, \Phi(x))\left[w, \Phi^{\prime}(x)[v]\right]+f_{t t}^{\prime \prime}(x, \Phi(x))\left[\Phi^{\prime}(x)[v], \Phi^{\prime}(x)[w]\right] ;
\end{aligned}
$$

here, $(\cdots)$ can be obtained easily, using a variation of $x$ corresponding to the directions of $v$ and $w$ and observing that $\beta(x) \dot{t}=$ constant (see [7]). At this point, the thesis springs out from (2), (3), ( $\cdot$ ).

Corollary (6). The following a) and b) are equivalent.
a) $J^{\prime \prime}(x)[v, w]=0$, for any vector field $w$;
b) $f^{\prime \prime}(\gamma)\left[\left(v, \Phi^{\prime}(x)\right)[v],(w, \tau)\right]=0$, for any vector field $(w, \tau)$.

So that, the null spaces of $J^{\prime \prime}(x)$ and of $f^{\prime \prime}(\gamma)$ have the same dimension.
Proof: The equivalence $a) \leftrightarrow b$ ) is straightforward.

We shall show the second statement. Let $\left(v, \tau_{1}\right)$ belong to the null space of $f^{\prime \prime}(\gamma)$ :

$$
\begin{aligned}
f^{\prime \prime}(\gamma)\left[\left(v, \tau_{1}\right),\left(w, \tau_{2}\right)\right]= & \left(f_{x x}^{\prime \prime}(\gamma)[v, w]+f_{x t}^{\prime \prime}(\gamma)\left[w, \tau_{1}\right]\right)+ \\
& +\left(f_{x t}^{\prime \prime}(\gamma)\left[v, \tau_{2}\right]+f_{t t}^{\prime \prime}(\gamma)\left[\tau_{1}, \tau_{2}\right]\right)=0, \quad \forall\left(w, \tau_{2}\right) .
\end{aligned}
$$

So, it is identically:

$$
\begin{cases}f_{x x}^{\prime \prime}(\gamma)[v, w]+f_{t t}^{\prime \prime}(\gamma)\left[w, \tau_{1}\right]=0, & \forall w  \tag{1}\\ f_{x t}^{\prime \prime}(\gamma)\left[v, \tau_{2}\right]+f_{t t}^{\prime \prime}(\gamma)\left[\tau_{1}, \tau_{2}\right]=0, & \forall \tau_{2}\end{cases}
$$

Comparing the second equation of the above system (1) with the equation (2) of the proposition (5), we obtain:

$$
\begin{array}{lll} 
& f_{x t}^{\prime \prime}(\gamma)\left[v, \tau_{2}\right]+f_{t t}^{\prime \prime}(\gamma)\left[\tau_{1}, \tau_{2}\right]=f_{x t}^{\prime \prime}(\gamma)\left[v, \tau_{2}\right]+f_{t t}^{\prime \prime}(\gamma)\left[\Phi^{\prime}(x)[v], \tau_{2}\right], \quad \forall \tau_{2} \\
\Rightarrow & f_{t t}^{\prime \prime}(\gamma)\left[\tau_{1}, \tau_{2}\right]=f_{t t}^{\prime \prime}(\gamma)\left[\Phi^{\prime}(x)[v], \tau_{2}\right], \quad \forall \tau_{2} \\
\Rightarrow & f_{t t}^{\prime \prime}(\gamma)\left[\tau_{1}-\Phi^{\prime}(x)[v], \tau_{2}\right]=0, \quad \forall \tau_{2} \\
\Rightarrow & \int_{0}^{1} \beta(x) \frac{d}{d s}\left(\tau_{1}-\Phi^{\prime}(x)[v]\right) \dot{\tau}_{2} d s=0, \quad \forall \tau_{2} .
\end{array}
$$

If $\left(\tau_{2}=\tau_{1}-\Phi^{\prime}(x)[v]\right)$, then:

$$
\int_{0}^{1}\left\|\frac{d}{d s}\left(\tau_{1}-\Phi^{\prime}(x)[v]\right)\right\|_{T_{\gamma(s)} M}^{2} d s=0
$$

so,

$$
\tau_{1}=\Phi^{\prime}(x)[v]+\text { constant } ;
$$

but, as $\tau_{1}(0)=0=\left.\Phi^{\prime}(x)[v]\right|_{s=0}$, then:

$$
\tau_{1}=\Phi^{\prime}(x)[v]
$$

Remark (7). A vector field $(v, \tau)$ belongs to the null space of $f^{\prime \prime}(\gamma)$ iff it is:

$$
(v, \tau)=\left(v, \Phi^{\prime}(x)[v]\right)
$$

and $v$ belongs to the null space of $J^{\prime \prime}(x)$.
Proof of Theorem (1): We already know, by [7], that the function $\exp _{p}$ is onto, for any point $P \in M$. In order to prove it is one to one, it will suffice to show that all the critical points of the functional $J$ are minima; afterwards, the unicity of the geodesic joining two given points on $M$ will follow from well known results of the critical points theory (see e.g.: [8], Theorem ( $6,5,3$ ), page 354).

Viceversa, if $\exp _{p}$ is a diffeomorphism for any point $P \in M$, then no couple of conjugate points with respect to the action functional exists, on any geodesic $\gamma$ contained in $M$; then corollary (6) and remark (7) imply the lack of couples of conjugate points, with respect to the functional $J$, on the critical curves of $J$.

This fact shows that $\exp _{p}$ is a diffeomorphism for any $P \in M$, iff the form $J^{\prime \prime}$ is positive definite on any critical curve of $J$, independently of the extreme points of that, so Lemmas (3) and (4) complete the proof.

ACKNOWLEDGEMENTS - I would like to thank V. Benci, D. Fortunato and F. Giannoni for many enlightening discussions and the "Istituto di Matematiche Applicate" of the faculty of engineering of the University of Pisa for its warm hospitality.

## REFERENCES

[1] Alber, S.I. - The topology of functional manifolds and the calculus of variations in the large, Russian Math. Surveys, 25 (1970), 51-117.
[2] Avez, A. - Essais de qéométrie Riemannienne hyperbolique: Applications à la relativité générale, Inst. Fourier, 105 (1963), 1.
[3] Benci, V. - Periodic solutions of Lagrangian systems on compact manifolds, J. Diff. Eq., 63 (1986), 135-161.
[4] Benci, V. and Fortunato, D. - Existence of geodesics for the Lorentz metric of a stationary gravitational field, Ann. Inst. H. Poincaré, Anal. non linéaire, 1990.
[5] Benci, V. and Fortunato, D. - Periodic trajectories for the Lorentz metric of a static gravitational field, Progress in Nonlin. Diff. Eq., vol. 4, Birkhäuser, 1990.
[6] Benci, V. and Fortunato, D. - On the existence of infinitely many geodesics on space-time manifolds, Adv. Math., to appear.
[7] Benci, V., Fortunato, D. and Giannoni, F. - On the existence of multiple geodesics in static space-times, Ann. Inst. H. Poincaré, Anal. non linéaire, 8 (1991).
[8] Berger, M.S. - Non-linearity and functional analysis, Academic Press, 1977.
[9] Hawking, S.W. and Ellis, G.F. - The large scale structure of space-time, Cambridge Univ. Press, 1973.
[10] Milnor, J. - Morse theory, Ann. Math. Studies, 51, Princeton Univ. Press, 1963.
[11] Nash, J. - The embedding problem for Riemannian manifolds, Ann. Math., 63 (1956), 20-63.
[12] O'Nelle, B. - Semi-Riemannian geometry with applications to relativity, Academic Press, London, 1983.
[13] Penrose, R. - Techniques of differential topology in relativity, Conf. Board Math. Sc. 7, SIAM Philadelphia, 1972.
[14] Schwartz, J.T. - Non linear functional analysis, Gordon and Breach, N.Y., 1969.

[^1]
[^0]:    Received: February 7, 1992; Revised: May 22, 1992.

[^1]:    Walter Frattarolo,
    Facolta di Ingegneria, Istituto di Matematiche Applicate "U. Dini", Via Bonano, 25 B, I 56126 Pisa - ITALIA

