

## AN ELASTIC GREEN MATRIX FOR A SEMISTRIP

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*Abstract.* A Green matrix is constructed for one mixed two-dimensional problem of the elasticity theory. The results of work [1] are thereby defined more correctly.

*Key words:* system of equations of elasticity theory, Green matrix

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Assume that that we are seeking for a solution of the system of equations of the elasticity theory

$$\begin{aligned} (\lambda + 2\mu) \frac{\partial^2 u_1}{\partial x^2} + \mu \frac{\partial^2 u_1}{\partial y^2} + (\lambda + \mu) \frac{\partial^2 u_2}{\partial x \partial y} &= f_1(x, y), \\ (\lambda + \mu) \frac{\partial^2 u_1}{\partial x \partial y} + \mu \frac{\partial^2 u_2}{\partial x^2} + (\lambda + 2\mu) \frac{\partial^2 u_2}{\partial y^2} &= f_2(x, y) \end{aligned} \quad (1)$$

in a domain  $-\infty < x \leq 0$ ,  $0 \leq y \leq b$ , under the following boundary conditions, interesting from the standpoint of application

$$\begin{aligned} \mu \left( \frac{\partial u_1}{\partial y}(x, \alpha b) + \frac{\partial u_2}{\partial x}(x, \alpha b) \right) &= 0, \quad u_2(x, \alpha b) = 0, \\ &-\infty < x \leq 0, \quad \alpha = 0, 1, \\ u_1(0, y) = 0, \quad \mu \left( \frac{\partial u_1}{\partial y}(0, y) + \frac{\partial u_2}{\partial x}(0, y) \right) &= 0, \quad 0 \leq y \leq b, \end{aligned}$$

and that for  $x \rightarrow -\infty$  a solution is bounded.

Suppose that  $f_i(x, y) \in L_1((-\infty, 0], [0, b])$  and  $\int_{-\infty}^x \int_0^b f_i(\xi, \eta) d\xi d\eta = O((1-x)^{-1})$  as  $x \rightarrow -\infty$ ,  $i = 1, 2$ .

To obtain the Green matrix of the posed problem we will use the method and notation of the paper [1]. Let us define the vectors  $\mathbf{U}(x, y) = (u_1(x, y), u_2(x, y))$  and  $\mathbf{F}(x, y) = (f_1(x, y), f_2(x, y))$  (here and in what follows the vector transposition sign is omitted). Assume that

$$\begin{aligned} \mathbf{U}(x, y) &= \sum_{n=0}^{\infty} Q_n(y) \mathbf{U}_n(x), \quad \mathbf{F}(x, y) = \sum_{n=0}^{\infty} Q_n(y) \mathbf{F}_n(x), \\ Q_n(y) &= \text{diag}(\cos \nu y, \sin \nu y), \quad \mathbf{U}_n(x) = (u_{1n}(x), u_{2n}(x)), \\ \mathbf{F}_n(x) &= (f_{1n}(x), f_{2n}(x)), \quad \nu = \frac{\pi n}{b}. \end{aligned} \quad (3)$$

Owing to such a representation, conditions (2.1) are fulfilled. From (1)–(3) we conclude that to find a pair of functions  $u_{1n}$ ,  $u_{2n}$ ,  $n = 1, 2, \dots$ , it is required

to solve the system of ordinary differential equations

$$\begin{aligned} (\lambda + 2\mu)u''_{1n} - \mu\nu^2u_{1n} + (\lambda + \mu)\nu u'_{2n} &= f_{1n}, \\ -(\lambda + \mu)\nu u'_{1n} + \mu u''_{2n} - (\lambda + 2\mu)\nu^2u_{2n} &= f_{2n} \end{aligned} \tag{4}$$

under the boundary conditions by which

$$u_{1n}(0) = 0, \quad u'_{2n}(0) = 0, \tag{5}$$

must be fulfilled at the point  $x = 0$ , while for  $x \rightarrow -\infty$ ,  $u_{1n}$  and  $u_{2n}$  are bounded.

As to the case  $n = 0$ , here we have only one sought for function  $u_{10}(x)$  satisfying the first equation in (4) and the first equality in (5) and bounded for  $x \rightarrow -\infty$ .

To solve problem (4),(5), we use the well-known method [2]. First we consider the homogeneous system corresponding to (4). The system of its fundamental solutions forms the rectangular matrix  $\Phi_n = (\phi^n_{ij})$ ,  $i = 1, 2$ ,  $j = 1, 2, 3, 4$ , whose elements are

$$\begin{aligned} \phi^n_{11} &= -\phi^n_{21} = e^{\nu x}, \quad \phi^n_{12} = \phi^n_{22} = e^{-\nu x}, \\ \phi^n_{i3} &= [(-1)^i(\lambda + \mu)\nu x + (i - 1)(\lambda + 3\mu)] e^{\nu x}, \\ \phi^n_{i4} &= [(\lambda + \mu)\nu x - (i - 1)(\lambda + 3\mu)] e^{-\nu x}, \quad i = 1, 2. \end{aligned}$$

Using fundamental solutions, we come to the conclusion that the general solution of system (4) has the form

$$\mathbf{U}_n(x) = \int_{-\infty}^x S_n(x, \xi) \mathbf{F}_n(\xi) d\xi + \Phi_n(x) D_n, \quad n = 1, 2, \dots, \tag{6}$$

where  $S_n(x, \xi) = (s^n_{ij})$  is the second order matrix,  $i, j = 1, 2$ , whose elements are defined by the relations

$$\begin{aligned} s^n_{ii}(x, \xi) &= \frac{m}{2} \left( \frac{1}{\nu} (\lambda + 3\mu) \operatorname{sh}\nu(x - \xi) \right. \\ &\quad \left. + (-1)^i(\lambda + \mu)(x - \xi) \operatorname{ch}\nu(x - \xi) \right), \quad i = 1, 2, \\ -s^n_{12}(x, \xi) &= s^n_{21}(x, \xi) = \frac{m}{2} (\lambda + \mu)(x - \xi) \operatorname{sh}\nu(x - \xi), \\ m &= (\mu(\lambda + 2\mu))^{-1}, \end{aligned} \tag{7}$$

and  $D_n = (d_i)$  is the column-matrix of arbitrary constants,  $i = 1, 2, 3, 4$ .

Taking into account (6), let us choose elements  $D_n$  so that (5) and the condition of the boundedness of  $u_{1n}$  and  $u_{2n}$  for  $x \rightarrow -\infty$  be fulfilled. As a result, we obtain

$$\begin{aligned} d_1 &= \frac{m}{2} \int_{-\infty}^0 \left[ \left( \frac{1}{\nu} (\lambda + 3\mu) \operatorname{sh}\nu\xi - (\lambda + \mu)\xi \operatorname{ch}\nu\xi \right) f_1(\xi) \right. \\ &\quad \left. + (\lambda + \mu)\xi \operatorname{sh}\nu\xi f_2(\xi) \right] d\xi, \\ d_2 = d_4 &= 0, \quad d_3 = \frac{m}{2\nu} \int_{-\infty}^0 (\operatorname{sh}\nu\xi f_1(\xi) - \operatorname{ch}\nu\xi f_2(\xi)) d\xi. \end{aligned} \tag{8}$$

Formulas (6)–(8) imply

$$\mathbf{U}_n(x) = \int_{-\infty}^0 g_n(x, \xi) \mathbf{F}_n(\xi) d\xi, \quad (9)$$

where  $g_n(x, \xi) = (g_{ij}^n(x, \xi))$  is the second order matrix,  $i, j = 1, 2$ , moreover,

$$g_{ij}^n(x, \xi) = \begin{cases} \gamma_{ij}^n(x, \xi) + s_{ij}^n(x, \xi) & \text{for } x \geq \xi, \\ \gamma_{ij}^n(x, \xi) & \text{for } x \leq \xi, \end{cases} \quad (10)$$

and

$$\begin{aligned} \gamma_{11}^n(x, \xi) &= m(p(x - \xi) - q(x - \xi) - p(x + \xi) + q(x + \xi)), \\ \gamma_{12}^n(x, \xi) &= m(p(x - \xi) + p(x + \xi)), \\ \gamma_{21}^n(x, \xi) &= m(p(x + \xi) - p(x - \xi)), \\ \gamma_{22}^n(x, \xi) &= -m(p(x - \xi) + q(x - \xi) + p(x + \xi) + q(x + \xi)), \end{aligned} \quad (11)$$

with the notation

$$p(u) = \frac{1}{4}(\lambda + \mu)ue^{\nu u}, \quad q(u) = \frac{1}{4\nu}(\lambda + 3\mu)e^{\nu u}.$$

Formula (9) holds for  $n = 0$  too, where  $g_0(x, \xi) = (g_{ij}^0(x, \xi))$ ,  $i, j = 1, 2$ , and in the cases  $x \geq \xi$  and  $x \leq \xi$   $g_{11}^0$  is equal respectively to  $m\mu x$  and  $m\mu\xi$ , while for other elements of the matrix  $g_0(x, \xi)$  we have  $g_{12}^0 = g_{21}^0 = g_{22}^0 = 0$  in both cases.

By virtue of (3) and (9)

$$\mathbf{U}(x, y) = \int_{-\infty}^0 \int_0^b G(x, y, \xi, \eta) \mathbf{F}(\xi, \eta) d\xi d\eta,$$

where the sought Green matrix  $G(x, y, \xi, \eta)$  is defined by the relation

$$G(x, y, \xi, \eta) = \frac{1}{b} \sum_{n=0}^{\infty} \varepsilon_n Q_n(y) g_n(x, \xi) Q_n(\eta). \quad (12)$$

Here  $\varepsilon_n$  is equal to 1 for  $n = 0$ , and to 2 in other cases. To obtain an explicit form of the elements  $G_{ij}(x, y, \xi, \eta)$  of the matrix  $G(x, y, \xi, \eta)$ ,  $i, j = 1, 2$ , we use (7), (10)–(12), the well-known formulas [3]

$$\begin{aligned} \sum_{n=1}^{\infty} t^n \cos n\theta &= (1 - t \cos \theta)(1 - 2t \cos \theta + t^2)^{-1}, \\ \sum_{n=1}^{\infty} t^n n^{-1} \cos n\theta &= -\frac{1}{2} \ln(1 - 2t \cos \theta + t^2), \\ t^2 &< 1, \quad 0 < \theta < 2\pi, \end{aligned}$$

and the notation

$$z = x + iy, \quad \zeta = \xi + i\eta, \quad \omega(u) = e^{\frac{\pi u}{b}}, \quad P(u) = \operatorname{Re}(1 - \omega(u)), \quad S(u) = \operatorname{Im}\omega(u), \\ E(u) = |1 - \omega(u)|^2, \quad Q(u) = P(u)E^{-1}(u), \quad T(u) = S(u)E^{-1}(u).$$

When  $x \leq \xi$ , for the diagonal elements we have

$$G_{ii}(x, y, \xi, \eta) = \frac{2-i}{b} m\mu \operatorname{Re}\zeta + \frac{m}{8\pi} (\lambda + 3\mu) \ln \left[ E(z - \zeta)E(z - \bar{\zeta}) \left( E(z + \zeta) \times \right. \right. \\ \left. \left. \times E(z + \bar{\zeta}) \right)^{(-1)^i} \right] - \frac{m}{4b} (\lambda + \mu) \sum_{k=-1,+1} k^i \operatorname{Re}(z + k\zeta) (Q(z + k\zeta) + Q(z + k\bar{\zeta})), \\ i = 1, 2,$$

while the nondiagonal ones are defined by the formula

$$G_{ij}(x, y, \xi, \eta) = \frac{m}{4b} (\lambda + \mu) \sum_{k=-1,+1} k^i \operatorname{Re}(z + k\zeta) (T(z + k\zeta) - T(z + k\bar{\zeta})), \\ i, j = 1, 2, \quad i \neq j.$$

For  $x \geq \xi$  these relations readily imply formulas for  $G_{ij}(x, y, \xi, \eta)$  if we use the well-known property of symmetry: in the diagonal and nondiagonal elements we should make the replacement  $x \leftrightarrow \xi$  and, in addition to this, the nondiagonal elements should be interchanged. We obtain

$$G_{ii}(x, y, \xi, \eta) = \frac{2-i}{b} m\mu \operatorname{Re}z + \frac{m}{8\pi} (\lambda + 3\mu) \ln \left[ E(-z + \zeta)E(-z + \bar{\zeta}) \times \right. \\ \left. \times \left( E(z + \zeta)E(z + \bar{\zeta}) \right)^{(-1)^i} \right] - \frac{m}{4b} (\lambda + \mu) \sum_{k=-1,+1} k^{i+1} \operatorname{Re}(z + k\zeta) (Q(kz + \zeta) + \\ + Q(kz + \bar{\zeta})), \quad i = 1, 2, \\ G_{ij}(x, y, \xi, \eta) = \frac{m}{4b} (\lambda + \mu) \sum_{k=-1,+1} k^{i+1} \operatorname{Re}(z + k\zeta) (T(kz + \zeta) - T(kz + \bar{\zeta})), \\ i, j = 1, 2, \quad i \neq j.$$

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