

Linear stability of the incoherent solution and the transition formula for the Kuramoto-Daido model

By

Hayato CHIBA*

Abstract

The Kuramoto-Daido model, which describes synchronization phenomena, is a system of ordinary differential equations on N -torus defined as coupled harmonic oscillators, whose natural frequencies are drawn from some distribution function. In this paper, the continuous model for the Kuramoto-Daido model is introduced and the linear stability of its trivial solution (incoherent solution) is investigated. Kuramoto's transition point K_c , at which the incoherent solution changes the stability, is derived for an arbitrary distribution function for natural frequencies. It is proved that if the coupling strength K is smaller than K_c , the incoherent solution is asymptotically stable, while if K is larger than K_c , it is unstable.

§ 1. Introduction

Collective synchronization phenomena are observed in a variety of areas such as chemical reactions, engineering circuits and biological populations [16]. In order to investigate such a phenomenon, Kuramoto [9] proposed a system of ordinary differential equations

$$(1.1) \quad \frac{d\theta_i}{dt} = \omega_i + \frac{K}{N} \sum_{j=1}^N \sin(\theta_j - \theta_i), \quad i = 1, \dots, N,$$

where $\theta_i \in [0, 2\pi)$ denotes the phase of an i -th oscillator on a circle, $\omega_i \in \mathbf{R}$ denotes its natural frequency, $K > 0$ is the coupling strength, and where N is the number of oscillators. Eq.(1.1) is derived by means of the averaging method from coupled dynamical systems having limit cycles, and now it is called the *Kuramoto model*.

Received October 20, 2009.

2000 Mathematics Subject Classification(s):

Supported by a JSPS Fellowship for Young Scientists

*Faculty of Mathematics, Kyushu University, Fukuoka, 819-0395, Japan.

e-mail: chiba@math.kyushu-u.ac.jp

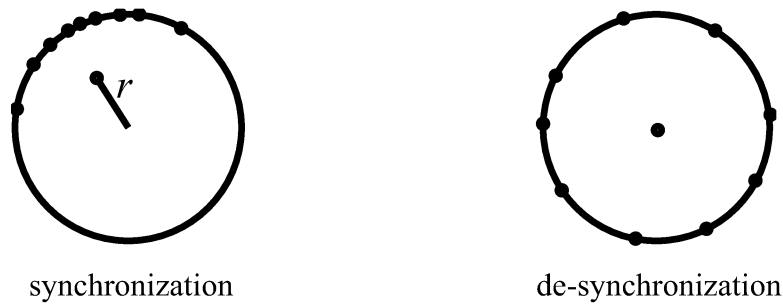


Figure 1. The order parameter for the Kuramoto model.

It is obvious that when $K = 0$, $\theta_i(t)$ and $\theta_j(t)$ rotate on a circle at different velocities unless ω_i is equal to ω_j , and it is true for sufficiently small $K > 0$. On the other hand, if K is sufficiently large, it is numerically observed that some of oscillators or all of them tend to rotate at the same velocity on average, which is called the *synchronization* [16, 18, 14]. If N is small, such a transition from de-synchronization to synchronization may be well revealed by means of the bifurcation theory [3, 11, 12]. However, if N is large, it is difficult to investigate the transition from the view point of the bifurcation theory and it is still far from understood.

In order to evaluate whether synchronization occurs or not, Kuramoto introduced the *order parameter* $r(t)e^{\sqrt{-1}\psi(t)}$ by

$$(1.2) \quad r(t)e^{\sqrt{-1}\psi(t)} := \frac{1}{N} \sum_{j=1}^N e^{\sqrt{-1}\theta_j(t)},$$

which gives the centroid of oscillators, where $r, \psi \in \mathbf{R}$. It seems that if synchronous state is formed, $r(t)$ takes a positive number, while if de-synchronization is stable, $r(t)$ is zero on time average (see Fig.1). Based on this observation and some formal calculation, Kuramoto conjectured a bifurcation diagram of $r(t)$ as follows:

Kuramoto's conjecture

Suppose that $N \rightarrow \infty$ and natural frequencies ω_i 's are distributed according to a probability density function $g(\omega)$. If $g(\omega)$ is an even and unimodal function, then the bifurcation diagram of $r(t)$ is given as Fig.2 (a); that is, if the coupling strength K is smaller than $K_c := 2/(\pi g(0))$, then $r(t) \equiv 0$ is asymptotically stable. On the other hand, if K is larger than K_c , there exists a positive constant r_c such that $r(t) = r_c$ is asymptotically stable. Near the transition point K_c , the scaling law of r_c is of $O((K - K_c)^{1/2})$.

Now the value $K_c = 2/(\pi g(0))$ is called the *Kuramoto's transition point*. See [10] and [18] for the Kuramoto's discussion.

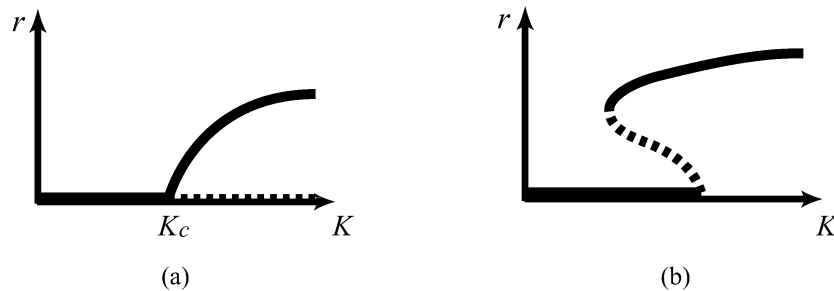


Figure 2. Typical bifurcation diagrams of the order parameter for the cases that (a) $g(\omega)$ is even and unimodal (b) $g(\omega)$ is even and bimodal. Solid lines denote stable solutions and dotted lines denote unstable solutions.

Significant papers of Strogatz et al. [19, 20, 15] partially confirmed the Kuramoto's conjecture. Though their arguments are not rigorous from a mathematical view point, almost all of them are justified as will be done in this paper. In [20], they introduced the continuous model for the Kuramoto model and investigated the linear stability of a trivial solution called the *incoherent solution*, which corresponds to the de-synchronous state $r \equiv 0$. They derived the Kuramoto's transition point $K_c = 2/(\pi g(0))$ and showed that if $K > K_c$, the incoherent solution is unstable in the linear level (i.e. nonlinear terms are neglected). When $K < K_c$, the linear operator T , which defines the linearized equation of the continuous model around the incoherent solution, has no eigenvalues. However, in [19], they found that an analytic continuation of the resolvent $(\lambda - T)^{-1}$ may have poles (*resonance poles*) on the left half plane, and they remarked a possibility that resonance poles induce exponential decay of the order parameter. In [15], the stability of the partially locked state, which corresponds to a solution with positive constant $r = r_c$, is investigated in the linear level.

Despite the active interest in the case that the distribution function $g(\omega)$ is even and unimodal, bifurcation diagrams of r for $g(\omega)$ other than the even and unimodal cases are not revealed well. Martens et al. [13] investigated the bifurcation diagram for a bimodal $g(\omega)$ which consists of two Lorentzian distributions. In particular, they found that stable partially locked states can coexist with stable incoherent solutions if K is slightly smaller than K_c (see Fig.2 (b)). Such a diagram seems to be common for any bimodal distributions.

A simple extension of the Kuramoto model defined to be

$$(1.3) \quad \frac{d\theta_i}{dt} = \omega_i + \frac{K}{N} \sum_{j=1}^N f(\theta_j - \theta_i), \quad i = 1, \dots, N,$$

is called the *Kuramoto-Daido model* [4, 5, 6, 7], where the 2π -periodic function f :

$\mathbf{R} \rightarrow \mathbf{R}$ is called the *coupling function*. Daido [7] investigated bifurcation diagrams of the order parameter for the Kuramoto-Daido model with even and unimodal $g(\omega)$ by a similar argument to Kuramoto's one. He found that if $f(\theta) \neq \sin \theta$, partially locked states may coexist with stable incoherent solutions even if $g(\omega)$ is even and unimodal.

All such studies by physicist are based on formal calculations and numerical simulations. The purpose of this paper is to justify and extend their results as mathematics for the Kuramoto-Daido model with any distribution function $g(\omega)$. The continuous model for the Kuramoto-Daido model is introduced and the linear stability of the incoherent solution is studied. In particular, the spectrum and the semigroup of a linear operator T , which is obtained by linearizing the continuous model around the incoherent solution, will be investigated in detail. At first, a formula for obtaining the transition point K_c for an arbitrary distribution $g(\omega)$ is derived. As a corollary, the Kuramoto's transition point $K_c = 2/(\pi g(0))$ is obtained if $g(\omega)$ is an even and unimodal function. If $K > K_c$, it is proved that the incoherent solution is unstable because the operator T has eigenvalues on the right half plane. It means that if the coupling strength K is large, the de-synchronous state is unstable and thus synchronization may occur. On the other hand, if $0 < K < K_c$, it will be shown that the spectrum of the operator T consists of the continuous spectrum and it lies on the imaginary axis. Thus the stability of the incoherent solution is nontrivial. Despite this fact, under appropriate assumptions for $g(\omega)$, the order parameter proves to decay exponentially because of existence of resonance poles on the left half plane as was expected by Strogatz et al. [19]. It suggests that in general, linear stability of a trivial solution of a linear equation on an infinite dimensional space is determined by not only the spectrum of the linear operator but also its resonance poles.

§ 2. Continuous model

In this section, we introduce a continuous model of the Kuramoto-Daido model and show a few properties of it.

Let us consider the Kuramoto-Daido model (1.3). We suppose that the coupling function f is a C^1 periodic function with the period 2π . It is expanded in a Fourier series as

$$(2.1) \quad f(\theta) = \sum_{l=-\infty}^{\infty} f_l e^{\sqrt{-1}l\theta}, \quad f_l \in \mathbf{C}.$$

We can suppose that $f_0 = 0$ without loss of generality because f_0 is renormalized into the constants ω_i . For the Kuramoto model ($f(\theta) = \sin \theta$), $f_{\pm 1} = \pm 1/(2\sqrt{-1})$ and $f_l = 0$ ($l \neq \pm 1$). Following Daido [7], we introduce the *generalized order parameters* \hat{Z}_k^0

by

$$(2.2) \quad \hat{Z}_k^0(t) = \frac{1}{N} \sum_{j=1}^N e^{\sqrt{-1}k\theta_j(t)}, \quad k = 0, \pm 1, \pm 2, \dots$$

In particular, \hat{Z}_1^0 is the order parameter defined in Section 1. By using them, Eq.(1.3) is rewritten as

$$(2.3) \quad \frac{d\theta_i}{dt} = \omega_i + K \sum_{l=-\infty}^{\infty} f_l \hat{Z}_l^0(t) e^{-\sqrt{-1}l\theta_i}.$$

Motivated by these equations, we introduce a continuous model of the Kuramoto-Daido model, which is an evolution equation of a probability density function $\rho_t = \rho_t(\theta, \omega)$ on $S^1 \times \mathbf{R}$ parameterized by $t \in \mathbf{R}$, as

$$(2.4) \quad \begin{cases} \frac{\partial \rho_t}{\partial t} + \frac{\partial}{\partial \theta} \left(\left(\omega + K \sum_{l=-\infty}^{\infty} f_l Z_l^0(t) e^{-\sqrt{-1}l\theta} \right) \rho_t \right) = 0, \\ Z_k^0(t) = \int_{\mathbf{R}} \int_0^{2\pi} e^{\sqrt{-1}k\theta} \rho_t(\theta, \omega) d\theta d\omega, \\ \rho_0(\theta, \omega) = h(\theta, \omega), \end{cases}$$

where $h(\theta, \omega)$ is an initial density function. The $Z_k^0(t)$ is a continuous version of $\hat{Z}_k^0(t)$, and we also call it the *generalized order parameter*. We can prove that Eq.(2.4) is proper in the sense that $\hat{Z}_k^0(t) \rightarrow Z_k^0(t)$ as $N \rightarrow \infty$ under some assumptions, although the proof is not given in this paper. If we regard

$$v := \omega + K \sum_{l=-\infty}^{\infty} f_l Z_l^0(t) e^{-\sqrt{-1}l\theta}$$

as a velocity field, Eq.(2.4) provides an equation of continuity $\partial \rho_t / \partial t + \partial(\rho_t v) / \partial \theta = 0$ known in fluid dynamics. It is easy to prove the law of conservation of mass:

$$(2.5) \quad \int_{\mathbf{R}} \int_0^{2\pi} \rho_t(\theta, \omega) d\theta = \int_{\mathbf{R}} \int_0^{2\pi} h(\theta, \omega) d\theta =: g(\omega).$$

A function g defined as above gives a probability density function for natural frequencies $\omega \in \mathbf{R}$ such that $\int_{\mathbf{R}} g(\omega) d\omega = 1$.

By using the characteristic curve method, Eq.(2.4) is formally integrated as follows: Consider the equation

$$(2.6) \quad \frac{dx}{dt} = \omega + K \sum_{l=-\infty}^{\infty} f_l Z_l^0(t) e^{-\sqrt{-1}lx}, \quad x \in S^1,$$

which defines a characteristic curve. Let $x = x(t, s; \theta, \omega)$ be a solution of Eq.(2.6) satisfying $x(s, s; \theta, \omega) = \theta$. Then, ρ_t is given as

$$(2.7) \quad \rho_t(\theta, \omega) = h(x(0, t; \theta, \omega), \omega) \exp \left[K \int_0^t \sum_{l=-\infty}^{\infty} \sqrt{-1} l f_l Z_l^0(s) e^{-\sqrt{-1} l x(s, t; \theta, \omega)} ds \right].$$

By using Eq.(2.7), it is easy to show the equality

$$(2.8) \quad \int_{\mathbf{R}} \int_0^{2\pi} a(\theta, \omega) \rho_t(\theta, \omega) d\theta d\omega = \int_{\mathbf{R}} \int_0^{2\pi} a(x(t, 0; \theta, \omega), \omega) h(\theta, \omega) d\theta d\omega,$$

for any continuous function $a(\theta, \omega)$. In particular, the generalized order parameters $Z_k^0(t)$ are rewritten as

$$(2.9) \quad Z_k^0(t) = \int_{\mathbf{R}} \int_0^{2\pi} e^{\sqrt{-1} k x(t, 0; \theta, \omega)} h(\theta, \omega) d\theta d\omega.$$

Substituting it into Eqs.(2.6), (2.7), we obtain

$$(2.10) \quad \frac{d}{dt} x(t, s; \theta, \omega) = \omega + K \int_{\mathbf{R}} \int_0^{2\pi} f(x(t, 0; \theta', \omega') - x(t, s; \theta, \omega)) h(\theta', \omega') d\theta' d\omega',$$

and

$$(2.11) \quad \rho_t(\theta, \omega) = h(x(0, t; \theta, \omega), \omega) \times \exp \left[K \int_0^t ds \cdot \int_{\mathbf{R}} \int_0^{2\pi} \frac{\partial f}{\partial \theta} \left(x(s, 0; \theta', \omega') - x(s, t; \theta, \omega) \right) h(\theta', \omega') d\theta' d\omega' \right],$$

respectively. Even if $h(\theta, \omega)$ is not differentiable, we consider Eq.(2.11) to be a weak solution of Eq.(2.4). It is easy in usual way to prove that the integro-ODE (2.10) has a unique solution for any $t > 0$, and this proves that the continuous model Eq.(2.4) has a unique weak solution (2.11) for an arbitrary initial data $h(\theta, \omega)$.

Throughout this paper, we suppose that the initial data $h(\theta, \omega)$ is of the form $h(\theta, \omega) = \hat{h}(\theta)g(\omega)$. This assumption corresponds to the assumption for the Kuramoto-Daido model (1.3) that initial values $\{\theta_j(0)\}_{j=1}^N$ and natural frequencies $\{\omega_j\}_{j=1}^N$ are independently distributed. This is a physically natural assumption used in many literatures. In this case, $\rho_t(\theta, \omega)$ is written as $\rho_t(\theta, \omega) = \hat{\rho}_t(\theta, \omega)g(\omega)$, where

$$(2.12) \quad \hat{\rho}_t(\theta, \omega) = \hat{h}(x(0, t; \theta, \omega)) \times \exp \left[K \int_0^t ds \cdot \int_{\mathbf{R}} \int_0^{2\pi} \frac{\partial f}{\partial \theta} \left(x(s, 0; \theta', \omega') - x(s, t; \theta, \omega) \right) \hat{h}(\theta') g(\omega') d\theta' d\omega' \right],$$

and $\hat{\rho}_t(\theta, \omega)$ satisfies the same equation as Eq.(2.4).

§ 3. Linear stability of the incoherent solution

A trivial solution of the continuous model (2.4), which is independent of θ and t , is given by $\rho_t(\theta, \omega) = g(\omega)/(2\pi)$, or equivalently $\hat{\rho}_t(\theta, \omega) = 1/(2\pi)$. It is called the *incoherent solution*, which corresponds to the de-synchronized state. Note that in this case $r = 0$. In this section, we investigate the stability of the incoherent solution and the order parameter.

Let

$$(3.1) \quad Z_j(t, \omega) := \int_0^{2\pi} e^{\sqrt{-1}j\theta} \hat{\rho}_t(\theta, \omega) d\theta$$

be the Fourier coefficients of $\hat{\rho}_t(\theta, \omega)$. Then, $Z_0(t, \omega) = 1$ and Z_j , $j = \pm 1, \pm 2, \dots$ satisfy the differential equations

$$\begin{aligned} \frac{dZ_j}{dt} &= \sqrt{-1}j\omega Z_j + \sqrt{-1}jK \sum_{-\infty}^{\infty} f_l Z_l^0(t) Z_{j-l} \\ &= \sqrt{-1}j\omega Z_j + \sqrt{-1}jK f_j Z_j^0(t) + \sqrt{-1}jK \sum_{l \neq j} f_l Z_l^0(t) Z_{j-l}. \end{aligned}$$

The incoherent solution corresponds to the zero solution $Z_j \equiv 0$ for $j = \pm 1, \pm 2, \dots$. Since $|Z_j(t, \omega)| \leq 1$, $Z_j(t, \omega)$ is in the Hilbert space $L^2(\mathbf{R}, g(\omega)d\omega)$ for every t :

$$\|Z_j\|_{L^2(\mathbf{R}, g(\omega)d\omega)}^2 = \int_{\mathbf{R}} |Z_j(t, \omega)|^2 g(\omega) d\omega \leq 1.$$

Thus we linearize the above equation as an evolution equation on $L^2(\mathbf{R}, g(\omega)d\omega)$

$$(3.2) \quad \frac{dZ_j}{dt} = (j\sqrt{-1}\mathcal{M} + j\sqrt{-1}K f_j \mathcal{P}) Z_j, \quad j = \pm 1, \pm 2, \dots,$$

where $\mathcal{M} : q(\omega) \mapsto \omega q(\omega)$ is the multiplication operator on $L^2(\mathbf{R}, g(\omega)d\omega)$ and \mathcal{P} is the projection on $L^2(\mathbf{R}, g(\omega)d\omega)$ defined to be

$$(3.3) \quad \mathcal{P}q(\omega) = \int_{\mathbf{R}} q(\omega) g(\omega) d\omega.$$

If we put $P_0(\omega) \equiv 1$, \mathcal{P} is also written as $\mathcal{P}q(\omega) = (q, P_0)$, where (\cdot, \cdot) is the inner product on $L^2(\mathbf{R}, g(\omega)d\omega)$:

$$(3.4) \quad (q_1, q_2) := \int_{\mathbf{R}} q_1(\omega) \overline{q_2(\omega)} g(\omega) d\omega.$$

Note that the order parameter is given as $Z_1^0 = (Z_1, P_0)$. To determine the stability of the incoherent solution and the order parameter, we have to investigate the spectrum and the semigroup of the operator $T_j := j\sqrt{-1}\mathcal{M} + j\sqrt{-1}K f_j \mathcal{P}$.

§ 3.1. Analysis of the operator $\sqrt{-1}\mathcal{M}$

If $f_j = 0$, $T_j = j\sqrt{-1}\mathcal{M}$. It is known that the multiplication operator \mathcal{M} on $L^2(\mathbf{R}, g(\omega)d\omega)$ is self-adjoint and its spectrum is given by $\sigma(\mathcal{M}) = \text{supp}(g) \subset \mathbf{R}$, where $\text{supp}(g)$ is a support of the density function g . Thus the spectrum of $j\sqrt{-1}\mathcal{M}$ is

$$(3.5) \quad \sigma(j\sqrt{-1}\mathcal{M}) = j\sqrt{-1} \cdot \text{supp}(g) = \{j\sqrt{-1}\lambda \mid \lambda \in \text{supp}(g)\} \subset \sqrt{-1}\mathbf{R}.$$

The semi-group $e^{j\sqrt{-1}\mathcal{M}t}$ generated by $j\sqrt{-1}\mathcal{M}$ is given as $e^{j\sqrt{-1}\mathcal{M}t}q(\omega) = e^{j\sqrt{-1}\omega t}q(\omega)$. In particular, we obtain

$$(3.6) \quad (e^{j\sqrt{-1}\mathcal{M}t}q_1, q_2) = \int_{\mathbf{R}} e^{j\sqrt{-1}\omega t} q_1(\omega) \overline{q_2(\omega)} g(\omega) d\omega$$

for any $q_1, q_2 \in L^2(\mathbf{R}, g(\omega)d\omega)$. This is the Fourier transform of the function $q_1(\omega) \overline{q_2(\omega)} g(\omega)$. Thus if $q_1(\omega) \overline{q_2(\omega)} g(\omega)$ is real analytic on \mathbf{R} and has an analytic continuation to a neighborhood of the real axis, then $(e^{j\sqrt{-1}\mathcal{M}t}q_1, q_2)$ decays exponentially as $t \rightarrow \infty$, while if $q_1(\omega) \overline{q_2(\omega)} g(\omega)$ is C^r , then it decays as $O(1/t^r)$ (see Vilenkin [21]).

These facts are summarized as follows:

Proposition 3.1. *Suppose that $f_j = 0$ and Eq.(3.2) is reduced to $dZ_j/dt = j\sqrt{-1}\mathcal{M}Z_j$. A solution of this equation with an initial value $q(\omega) \in L^2(\mathbf{R}, g(\omega)d\omega)$ is given by $Z_j(t) = e^{j\sqrt{-1}\mathcal{M}t}q(\omega) = e^{j\sqrt{-1}\omega t}q(\omega)$. In particular the linearized order parameter $Z_1^0(t) = (e^{j\sqrt{-1}\mathcal{M}t}q, P_0)$ decays exponentially as $t \rightarrow \infty$ if $g(\omega)$ and $q(\omega)$ have analytic continuations to a neighborhood of the real axis.*

The resolvent $(\lambda - j\sqrt{-1}\mathcal{M})^{-1}$ of the operator $j\sqrt{-1}\mathcal{M}$ is calculated as

$$(3.7) \quad ((\lambda - j\sqrt{-1}\mathcal{M})^{-1}q_1, q_2) = \int_{\mathbf{R}} \frac{1}{\lambda - j\sqrt{-1}\omega} q_1(\omega) \overline{q_2(\omega)} g(\omega) d\omega.$$

We define the function $D(\lambda)$ to be

$$(3.8) \quad D(\lambda) = ((\lambda - j\sqrt{-1}\mathcal{M})^{-1}P_0, P_0) = \int_{\mathbf{R}} \frac{1}{\lambda - j\sqrt{-1}\omega} g(\omega) d\omega$$

(recall that $P_0(\omega) \equiv 1$). It is holomorphic in $\mathbf{C} \setminus \sigma(j\sqrt{-1}\mathcal{M})$ and will play an important role in the later calculation.

§ 3.2. Analysis of the operator $T_j = j\sqrt{-1}\mathcal{M} + j\sqrt{-1}K f_j \mathcal{P}$

In what follows, we suppose that $f_j \neq 0$. The domain $D(T_j)$ of T_j is given by $D(\mathcal{M}) \cap D(\mathcal{P}) = D(\mathcal{M})$. Since \mathcal{M} is self-adjoint and since \mathcal{P} is bounded, T_j is a closed

operator [8]. Let $\varrho(T_j)$ be the resolvent set of T_j and $\sigma(T_j) = \mathbf{C} \setminus \varrho(T_j)$ the spectrum. Since T_j is closed, there is no residual spectrum. Let $\sigma_p(T_j)$ and $\sigma_c(T_j)$ be the point spectrum (the set of eigenvalues) and the continuous spectrum of T_j , respectively.

Proposition 3.2. (i) Eigenvalues λ of T_j are given as roots of

$$(3.9) \quad D(\lambda) = \frac{1}{j\sqrt{-1}Kf_j}, \quad \lambda \in \mathbf{C} \setminus \sigma(j\sqrt{-1}\mathcal{M}).$$

(ii) The continuous spectrum of T_j is given by

$$(3.10) \quad \sigma_c(T_j) = \sigma(j\sqrt{-1}\mathcal{M}) = j\sqrt{-1} \cdot \text{supp}(g).$$

Proof. (i) Suppose that $\lambda \in \sigma_p(T_j) \setminus \sigma(j\sqrt{-1}\mathcal{M})$. Then, there exists $x \in L^2(\mathbf{R}, g(\omega)d\omega)$ such that

$$\lambda x = (j\sqrt{-1}\mathcal{M} + j\sqrt{-1}Kf_j\mathcal{P})x, \quad x \neq 0.$$

Since $\lambda \notin \sigma(j\sqrt{-1}\mathcal{M})$, $(\lambda - j\sqrt{-1}\mathcal{M})^{-1}$ is defined and the above is rewritten as

$$\begin{aligned} x &= (\lambda - j\sqrt{-1}\mathcal{M})^{-1}j\sqrt{-1}Kf_j\mathcal{P}x \\ &= j\sqrt{-1}Kf_j(x, P_0)(\lambda - j\sqrt{-1}\mathcal{M})^{-1}P_0(\omega). \end{aligned}$$

By taking the inner product with $P_0(\omega)$, we obtain

$$(3.11) \quad 1 = j\sqrt{-1}Kf_j((\lambda - j\sqrt{-1}\mathcal{M})^{-1}P_0, P_0) = j\sqrt{-1}Kf_jD(\lambda).$$

This proves that roots of Eq.(3.9) is in $\sigma_p(T_j)$. The corresponding eigenvector is given by $x = (\lambda - j\sqrt{-1}\mathcal{M})^{-1}P_0(\omega) = 1/(\lambda - j\sqrt{-1}\omega)$. If $\lambda \in \sqrt{-1}\mathbf{R}$, $x \notin L^2(\mathbf{R}, g(\omega)d\omega)$. Thus there are no eigenvalues on the imaginary axis.

(ii) This follows from the fact that the essential spectrum is stable under the bounded perturbation and that there are no eigenvalues on $\sigma(j\sqrt{-1}\mathcal{M})$, see [8]. \square

§ 3.3. Eigenvalues of the operator T_j and the transition point formula

Our next task is to calculate roots of Eq.(3.9) to obtain eigenvalues of $T_j = j\sqrt{-1}\mathcal{M} + j\sqrt{-1}Kf_j\mathcal{P}$. By putting $\lambda = x + \sqrt{-1}y$, $x, y \in \mathbf{R}$, Eq.(3.9) is rewritten as

$$(3.12) \quad \begin{cases} \int_{\mathbf{R}} \frac{x}{x^2 + (j\omega - y)^2} g(\omega)d\omega = -\frac{\text{Im}(f_j)}{jK|f_j|^2}, \\ \int_{\mathbf{R}} \frac{j\omega - y}{x^2 + (j\omega - y)^2} g(\omega)d\omega = -\frac{\text{Re}(f_j)}{jK|f_j|^2}. \end{cases}$$

In what follows, we suppose that $\text{Im}(f_j) < 0$. The case $\text{Im}(f_j) \geq 0$ will be treated in Sec.3.5. The next lemma is easily obtained.

Lemma 3.3.

- (i) When $\text{Im}(f_j) < 0$, λ satisfies $\text{Re}(\lambda) > 0$ for any $K > 0$.
- (ii) If $K > 0$ is sufficiently large, there exists at least one eigenvalue λ near infinity.
- (iii) If $K > 0$ is sufficiently small, there are no eigenvalues.

Proof. Part (i) of the lemma immediately follows from the first equation of Eq.(3.12). To prove part (ii) of the lemma, note that if $|\lambda|$ is large, Eq.(3.9) is rewritten as

$$\frac{1}{\lambda} + O\left(\frac{1}{\lambda^2}\right) = \frac{1}{j\sqrt{-1}Kf_j}.$$

Thus the Rouché's theorem proves that Eq.(3.9) has a root $\lambda \sim j\sqrt{-1}Kf_j$ if $K > 0$ is sufficiently large. To prove part (iii) of the lemma, we see that the left hand side of the first equation of Eq.(3.12) is bounded for any $x, y \in \mathbf{R}$. To do so, let $G(\omega)$ be the primitive function of $g(\omega)$ and fix $\delta > 0$ small. The left hand side of the first equation of Eq.(3.12) is calculated as

$$\begin{aligned} & \int_{\mathbf{R}} \frac{xg(\omega)d\omega}{x^2 + (j\omega - y)^2} \\ &= \int_{y/j+\delta}^{\infty} \frac{xg(\omega)d\omega}{x^2 + (j\omega - y)^2} + \int_{-\infty}^{y/j-\delta} \frac{xg(\omega)d\omega}{x^2 + (j\omega - y)^2} + \int_{y/j-\delta}^{y/j+\delta} \frac{xg(\omega)d\omega}{x^2 + (j\omega - y)^2} \\ &= \int_{y/j+\delta}^{\infty} \frac{xg(\omega)d\omega}{x^2 + (j\omega - y)^2} + \int_{-\infty}^{y/j-\delta} \frac{xg(\omega)d\omega}{x^2 + (j\omega - y)^2} \\ & \quad + \frac{x}{x^2 + j^2\delta^2} (G(y/j + \delta) - G(y/j - \delta)) + \int_{y/j-\delta}^{y/j+\delta} \frac{2jx(j\omega - y)}{(x^2 + (j\omega - y)^2)^2} G(\omega)d\omega. \end{aligned}$$

The first three terms in the right hand side above are bounded for any $x, y \in \mathbf{R}$. Since G is continuous, there exists a number ξ such that the last term is rewritten as

$$\int_{y/j-\delta}^{y/j+\delta} \frac{2jx(j\omega - y)}{(x^2 + (j\omega - y)^2)^2} G(\omega)d\omega = 2j\delta \cdot \frac{2x\xi}{(x^2 + \xi^2)^2} G(y/j + \xi/j).$$

This is bounded for any $x, y \in \mathbf{R}$. Now we have proved that the left hand side of the first equation of Eq.(3.12) is bounded for any $x > 0$, although the right hand side tends to infinity as $K \rightarrow +0$. Thus Eq.(3.9) has no roots if K is small. \square

Lemma 3.3 shows that if $K > 0$ is sufficiently large, the trivial solution $Z_j = 0$ of the system $dZ_j/dt = T_j Z_j$ is unstable because of the eigenvalues with positive real

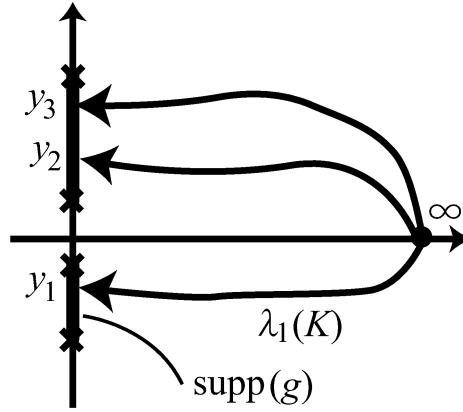


Figure 3. A schematic view of behavior of roots λ of Eq.(3.9) when K decreases.

parts. Our purpose in this subsection is to determine the bifurcation point $K_c^{(j)}$, which is the minimum value of K such that if $K < K_c^{(j)}$, the operator T_j has no eigenvalues on the right half plane. To calculate eigenvalues $\lambda = \lambda(K)$ explicitly is difficult in general. However, note that since zeros of a holomorphic function do not vanish because of the argument principle, $\lambda(K)$ disappears if and only if it is absorbed into the continuous spectrum $\sigma(j\sqrt{-1}\mathcal{M})$, on which $D(\lambda)$ is not holomorphic. This fact suggests that to determine $K_c^{(j)}$, it is sufficient to investigate Eq.(3.9) or Eq.(3.12) near the imaginary axis. Since we are interested in $\lambda(K)$ absorbed into $\sigma(j\sqrt{-1}\mathcal{M}) \subset \sqrt{-1}\mathbf{R}$, take the limit $x \rightarrow +0$ in Eq.(3.12):

$$(3.13) \quad \begin{cases} \lim_{x \rightarrow +0} \int_{\mathbf{R}} \frac{x}{x^2 + (j\omega - y)^2} g(\omega) d\omega = -\frac{\text{Im}(f_j)}{jK|f_j|^2}, \\ \lim_{x \rightarrow +0} \int_{\mathbf{R}} \frac{j\omega - y}{x^2 + (j\omega - y)^2} g(\omega) d\omega = -\frac{\text{Re}(f_j)}{jK|f_j|^2}. \end{cases}$$

These equations determine K_n and y_n such that one of the eigenvalues $\lambda_n(K)$ converges to $\sqrt{-1}y_n$ as $K \rightarrow K_n + 0$ (see Fig.3). To calculate them, we need the next lemma.

Lemma 3.4. (i) Suppose that $\lambda_n(K) \rightarrow \sqrt{-1}y_n$ as $K \rightarrow K_n$. Then, $g(\omega)$ is continuous at $\omega = y_n$.

(ii) If $g(\omega)$ is continuous at $\omega = y$, then

$$(3.14) \quad \lim_{x \rightarrow +0} \int_{\mathbf{R}} \frac{x}{x^2 + (j\omega - y)^2} g(\omega) d\omega = \pi g(y/j)/j.$$

Proof. To prove (i), suppose that $g(\omega)$ is discontinuous at $\omega = 0$ without loss of generality.

STEP 1: At first, we suppose that $g(\omega)$ is piecewise continuous. Put $g(+0) = h_+$, $g(-0) = h_-$ and $h_+ \neq h_-$. In this case, for any $\varepsilon > 0$, there exists $\delta > 0$ such

that if $-\delta < \omega < 0$, then $|g(\omega) - h_-| < \varepsilon$ and if $0 < \omega < \delta$, then $|g(\omega) - h_+| < \varepsilon$. For Eq.(3.9), we suppose $|\lambda| = |x + \sqrt{-1}y| < \delta$ and $y > 0$. The case $y < 0$ is treated in a similar manner. We calculate $D(\lambda)$ as

$$\begin{aligned}
D(\lambda) &= \int_{\delta}^{\infty} \frac{g(\omega)}{\lambda - j\sqrt{-1}\omega} d\omega + \int_{-\infty}^{-\delta} \frac{g(\omega)}{\lambda - j\sqrt{-1}\omega} d\omega \\
&\quad + \int_0^{\delta} \frac{1}{\lambda - j\sqrt{-1}\omega} (g(\omega) - h_+ + h_+) d\omega + \int_{-\delta}^0 \frac{1}{\lambda - j\sqrt{-1}\omega} (g(\omega) - h_- + h_-) d\omega \\
&= \int_{\delta}^{\infty} \frac{g(\omega)}{\lambda - j\sqrt{-1}\omega} d\omega + \int_{\delta}^{\infty} \frac{g(-\omega)}{\lambda + j\sqrt{-1}\omega} d\omega \\
&\quad + \int_0^{\delta} \frac{1}{\lambda - j\sqrt{-1}\omega} (g(\omega) - h_+) d\omega + \int_0^{\delta} \frac{1}{\lambda + j\sqrt{-1}\omega} (g(-\omega) - h_-) d\omega \\
(3.15) \quad &+ \int_0^{\delta} \frac{1}{\lambda - j\sqrt{-1}\omega} (h_+ - h_-) d\omega + h_- \int_0^{\delta} \frac{d\omega}{\lambda - j\sqrt{-1}\omega} + h_- \int_0^{\delta} \frac{d\omega}{\lambda + j\sqrt{-1}\omega}.
\end{aligned}$$

Since $|\lambda| < \delta$, there exists a positive number M , which is independent of λ , such that

$$\int_{\delta}^{\infty} \frac{g(\pm\omega)}{\lambda \mp j\sqrt{-1}\omega} d\omega < M.$$

Thus $|D(\lambda)|$ is estimated as

$$\begin{aligned}
|D(\lambda)| &\geq |h_+ - h_-| \int_0^{\delta} \frac{d\omega}{|\lambda - j\sqrt{-1}\omega|} - \varepsilon \int_0^{\delta} \frac{d\omega}{|\lambda - j\sqrt{-1}\omega|} - \varepsilon \int_0^{\delta} \frac{d\omega}{|\lambda + j\sqrt{-1}\omega|} \\
&\quad - \left| \frac{h_-}{j} \log \left(\frac{\sqrt{-1}\lambda + j\delta}{\sqrt{-1}\lambda - j\delta} \right) \right| - 2M.
\end{aligned}$$

Since $y > 0$, $|\lambda - j\sqrt{-1}\omega| < |\lambda + j\sqrt{-1}\omega|$. This shows that

$$(3.16) \quad |D(\lambda)| \geq (|h_+ - h_-| - 2\varepsilon) \int_0^{\delta} \frac{d\omega}{|\lambda - j\sqrt{-1}\omega|} - \left| \frac{h_-}{j} \log \left(\frac{\sqrt{-1}\lambda + j\delta}{\sqrt{-1}\lambda - j\delta} \right) \right| - 2M.$$

The right hand side tends to infinity as $\lambda \rightarrow 0$ if $2\varepsilon < |h_+ - h_-|$. This proves that Eq.(3.9) has no roots at $\lambda = 0$ for positive K .

STEP 2: In general, since $g(\omega)$ is a non-negative measurable function, there exists a monotonic increasing sequence $\{g_n(\omega)\}_{n=1}^{\infty}$ of non-negative simple functions such that $g_n(\omega) \rightarrow g(\omega)$ for each ω . In particular if $g(\omega)$ is discontinuous at $\omega = 0$, we can choose $\{g_n(\omega)\}_{n=1}^{\infty}$ so that $g_n(\omega)$ is discontinuous at $\omega = 0$ for any $n \in \mathbf{N}$. Then, the proof is done in the same way as STEP 1 by approximating $g(\omega)$ by $g_n(\omega)$.

(ii) The formula Eq.(3.14) is proved in Ahlfors [1]. \square

Let (y, K) be one of the solutions of Eq.(3.13). Since $g(\omega)$ is continuous at $\omega = y$, substituting it into the first equation of Eq.(3.13) yields

$$(3.17) \quad \pi g(y/j) = -\frac{\text{Im}(f_j)}{K|f_j|^2}.$$

Substituting $K = -\text{Im}(f_j)/(\pi|f_j|^2g(y/j))$ obtained from the above into the second equation of Eq.(3.13) results in

$$(3.18) \quad \lim_{x \rightarrow 0} \int_{\mathbf{R}} \frac{j\omega - y}{x^2 + (j\omega - y)^2} g(\omega) d\omega = \frac{\pi \text{Re}(f_j)}{j \text{Im}(f_j)} g(y/j).$$

This equation for y determines imaginary parts to which $\lambda(K)$ converges as $\text{Re}(\lambda(K)) \rightarrow +0$. Let y_1, y_2, \dots be roots of Eq.(3.18). Then,

$$(3.19) \quad K_n = \frac{-\text{Im}(f_j)}{\pi|f_j|^2g(y_n/j)}, \quad n = 1, 2, \dots$$

give the values such that $\text{Re}(\lambda(K)) \rightarrow 0$ as $K \rightarrow K_n + 0$. Now we obtain the next theorem.

Theorem 3.5. *Suppose that $\text{Im}(f_j) < 0$. Let y_1, y_2, \dots be roots of Eq.(3.18). Put*

$$(3.20) \quad K_c^{(j)} := \inf_n K_n = \frac{-\text{Im}(f_j)}{\pi|f_j|^2 \sup_n g(y_n/j)}.$$

If $0 < K < K_c^{(j)}$, the operator T_j has no eigenvalues, while if K is slightly larger than $K_c^{(j)}$, T_j has eigenvalues on the right half plane.

Note that $\inf_n K_n$ is positive because of Lemma.3.3 (iii). As a corollary, we obtain the transition point (bifurcation point to the partially locked state) conjectured by Kuramoto [10]:

Corollary 3.6 (Kuramoto's transition point). *Suppose that the probability density function $g(\omega)$ is even and $\max g(\omega) = g(0)$. If $\text{Re}(f_1) = 0$ and $\text{Im}(f_1) = -1/2$ (it corresponds to $f(\theta) = \sin \theta$ in Eq.(1.3)), then $K_c^{(1)}$ defined as above is given by*

$$(3.21) \quad K_c^{(1)} = \frac{2}{\pi g(0)}.$$

When $K > K_c^{(1)}$, the solution $Z_1 = 0$ of Eq.(3.2) is unstable.

§ 3.4. Semi-group generated by the operator T_1 ($\text{Im}(f_1) < 0$)

Since we are interested in the dynamics of the order parameter $Z_1^0(t) = (Z_1, P_0)$, in what follows, we consider only $j = 1$ while cases $j = 2, 3, \dots$ are investigated in the same way. Theorem 3.5 shows that $K_c^{(1)}$ is the least bifurcation point and the

trivial solution $Z_1(t, \omega) = 0$ of Eq.(3.2) is unstable if K is slightly larger than $K_c^{(1)}$. If $0 < K < K_c^{(1)}$, the spectrum of T_1 is on the imaginary axis: $\sigma(T_1) = \sigma(\sqrt{-1}\mathcal{M})$, and thus the dynamics of Z_1 is nontrivial. In this subsection, we investigate the dynamics of Z_1 and the order parameter for $0 < K < K_c^{(1)}$. We will see that the order parameter may decay exponentially even if the spectrum lies on the imaginary axis because of existence of resonance poles.

Since $\sqrt{-1}\mathcal{M}$ has the semi-group $e^{\sqrt{-1}\mathcal{M}t}$ and since \mathcal{P} is bounded, the operator $T_1 = \sqrt{-1}\mathcal{M} + \sqrt{-1}Kf_1\mathcal{P}$ also generates the semi-group (Kato [8]), say e^{T_1t} . A solution of Eq.(3.2) with an initial value $q(\omega) \in L^2(\mathbf{R}, g(\omega)d\omega)$ is given by $e^{T_1t}q(\omega)$. The e^{T_1t} is calculated by using the Laplace inversion formula

$$(3.22) \quad e^{T_1t} = \lim_{y \rightarrow \infty} \frac{1}{2\pi\sqrt{-1}} \int_{x-\sqrt{-1}y}^{x+\sqrt{-1}y} e^{\lambda t} (\lambda - T_1)^{-1} d\lambda,$$

where $x > 0$ is chosen so that the contour is to the right of the spectrum of T_1 (Yosida [22]). At first, let us calculate the resolvent $(\lambda - T_1)^{-1}$.

Lemma 3.7. *For any $q(\omega) \in L^2(\mathbf{R}, g(\omega)d\omega)$, the equality*

$$(3.23) \quad F_0(\lambda) := ((\lambda - T_1)^{-1}q, P_0) = \frac{((\lambda - \sqrt{-1}\mathcal{M})^{-1}q, P_0)}{1 - \sqrt{-1}Kf_1D(\lambda)}$$

holds.

Proof. Put

$$R(\lambda)q := (\lambda - T_1)^{-1}q = (\lambda - \sqrt{-1}\mathcal{M} - \sqrt{-1}Kf_1\mathcal{P})^{-1}q,$$

which yields

$$\begin{aligned} (\lambda - \sqrt{-1}\mathcal{M})R(\lambda)q &= q + \sqrt{-1}Kf_1\mathcal{P}R(\lambda)q \\ &= q + \sqrt{-1}Kf_1(R(\lambda)q, P_0)P_0. \end{aligned}$$

This is rearranged as

$$R(\lambda)q = (\lambda - \sqrt{-1}\mathcal{M})^{-1}q + \sqrt{-1}Kf_1(R(\lambda)q, P_0)(\lambda - \sqrt{-1}\mathcal{M})^{-1}P_0.$$

By taking the inner product with P_0 , we obtain

$$(R(\lambda)q, P_0) = ((\lambda - \sqrt{-1}\mathcal{M})^{-1}q, P_0) + \sqrt{-1}Kf_1(R(\lambda)q, P_0)D(\lambda).$$

This proves Eq.(3.23). □

Let $Z_1^0(t) = (Z_1, P_0)$ be the order parameter with the initial condition $Z_1(0, \omega) = q(\omega)$. Eqs.(3.22) and (3.23) show that $Z_1^0(t)$ is given by

$$(3.24) \quad Z_1^0(t) = (e^{T_1 t} q, P_0) = \lim_{y \rightarrow \infty} \frac{1}{2\pi\sqrt{-1}} \int_{x-\sqrt{-1}y}^{x+\sqrt{-1}y} e^{\lambda t} \frac{((\lambda - \sqrt{-1}\mathcal{M})^{-1}q, P_0)}{1 - \sqrt{-1}K f_1 D(\lambda)} d\lambda.$$

One of the effective way to calculate the integral above is to use the residue theorem. Recall that the resolvent $(\lambda - T_1)^{-1}$ is holomorphic on $\mathbf{C} \setminus \sigma(T_1)$. Since we assume that $0 < K < K_c^{(1)}$, T_1 has no eigenvalues and the continuous spectrum lies on the imaginary axis : $\sigma(T_1) = \sigma(\sqrt{-1}\mathcal{M}) = \sqrt{-1} \cdot \text{supp}(g)$. Thus the integrand $e^{\lambda t} F_0(\lambda)$ in Eq.(3.24) is holomorphic on the right half plane and may not be holomorphic on $\sigma(T_1)$. However, under assumptions below, we can show that $F_0(\lambda)$ has an analytic continuation $F_1(\lambda)$ through the line $\sigma(T_1)$ from right to left. Then, $F_1(\lambda)$ may have poles on the left half plane (the second Riemann sheet of the resolvent), which are called the *resonance poles* [17]. The resonance pole μ affects the integral in Eq.(3.24) through the residue theorem (see Fig.4). In this manner, the order parameter $Z_1^0(t)$ can decay with the exponential rate $\text{Re}(\mu)$. Such an exponential decay caused by resonance poles is well known in the theory of Schrödinger operators [17], and for the Kuramoto model, it is investigated numerically by Strogatz *et al.* [19] and Balmforth *et al.* [2].

At first, we construct an analytic continuation of the function $F_0(\lambda)$.

Lemma 3.8. *Suppose that the probability density function $g(\omega)$ and an initial condition $q(\omega)$ are real analytic on \mathbf{R} . If $g(\omega)$ and $q(\omega)$ have meromorphic continuations $g^*(\lambda)$ and $q^*(\lambda)$ to the upper half plane, respectively, then the function $F_0(\lambda)$ defined on the right half plane has the meromorphic continuation $F_1(\lambda)$ to the left half plane, which is given by*

$$(3.25) \quad F_1(\lambda) = \frac{((\lambda - \sqrt{-1}\mathcal{M})^{-1}q, P_0) + 2\pi q^*(-\sqrt{-1}\lambda)g^*(-\sqrt{-1}\lambda)}{1 - \sqrt{-1}K f_1 D(\lambda) - 2\pi\sqrt{-1}K f_1 g^*(-\sqrt{-1}\lambda)}.$$

Proof. By the formula (3.14), we obtain

$$(3.26) \quad \lim_{\text{Re}(\lambda) \rightarrow +0} ((\lambda - \sqrt{-1}\mathcal{M})^{-1}q, P_0) - \lim_{\text{Re}(\lambda) \rightarrow -0} ((\lambda - \sqrt{-1}\mathcal{M})^{-1}q, P_0) = 2\pi q(\text{Im}(\lambda)) \cdot g(\text{Im}(\lambda)).$$

Thus the meromorphic continuation of $((\lambda - \sqrt{-1}\mathcal{M})^{-1}q, P_0)$ from right to left is given by

$$(3.27) \quad \begin{cases} ((\lambda - \sqrt{-1}\mathcal{M})^{-1}q, P_0) & (\text{Re}(\lambda) > 0), \\ ((\lambda - \sqrt{-1}\mathcal{M})^{-1}q, P_0) + 2\pi q^*(-\sqrt{-1}\lambda)g^*(-\sqrt{-1}\lambda) & (\text{Re}(\lambda) < 0). \end{cases}$$

This proves Eq.(3.25). □

Poles of $F_1(\lambda)$ (resonance poles) on the left half plane are given as roots of the equation

$$(3.28) \quad D(\lambda) + 2\pi g^*(-\sqrt{-1}\lambda) = \frac{1}{\sqrt{-1}Kf_1}, \quad \operatorname{Re}(\lambda) < 0$$

and poles of the function $q^*(-\sqrt{-1}\lambda)$. In the next theorem, we suppose for simplicity that $q^*(-\sqrt{-1}\lambda)$ has no poles. Now we calculate the order parameter $Z_1^0(t)$.

Theorem 3.9. *For Eq.(3.2) with $j = 1$, suppose that*

(i) $\operatorname{Im}(f_1) < 0$ and $0 < K < K_c^{(1)}$.

(ii) *the probability density function $g(\omega)$ is real analytic on \mathbf{R} and has meromorphic continuation $g^*(\lambda)$ to the upper half plane.*

(iii) *an initial condition $q(\omega)$ is real analytic on \mathbf{R} and has analytic continuation $q^*(\lambda)$ to the upper half plane.*

(iv) *there exists a positive number δ such that $|F_1(\lambda)| \rightarrow 0$ as $|\lambda| \rightarrow \infty$ in the angular domains*

$$(3.29) \quad |\arg(\lambda)| \leq \delta, \quad |\arg(\lambda) - \pi| \leq \delta.$$

(v) *there exist positive constants D and β such that*

$$(3.30) \quad |F_1(\lambda)| \leq De^{\beta|\lambda|}$$

in the angular domain $\pi/2 + \delta \leq \arg(\lambda) \leq 3\pi/2 - \delta$.

Then, there exist resonance poles of T_1 on the left half plane. Let $\alpha_1, \alpha_2, \dots$ be resonance poles such that $|\alpha_1| \leq |\alpha_2| \leq \dots$. Then, there exists a positive constant t_0 such that the order parameter is given by

$$(3.31) \quad Z_1^0(t) = (e^{T_1 t} q, P_0) = \sum_{n=1}^{\infty} p_n(t) e^{\alpha_n t}, \quad t > t_0$$

where $p_n(t)$ is a polynomial in t . In particular, $Z_1^0(t)$ decays exponentially as $t \rightarrow \infty$.

Proof. At first, we prove the existence of resonance poles. Resonance poles are roots of Eq.(3.28), which is the analytic continuation of the equation (3.9) for $j = 1$. Thus one of the resonance poles is obtained as a continuation of an eigenvalue $\lambda(K)$. Recall that $\lambda(K)$ converges into the imaginary axis as $K \rightarrow K_c^{(1)} + 0$. To prove that there exists a resonance pole on the left half plane when $K < K_c^{(1)}$, we have to show that $\lambda(K)$ does not stay on the imaginary axis for $K < K_c^{(1)}$. Differentiating Eq.(3.9) with respect to K , we obtain

$$(3.32) \quad \lambda'(K) \int_{\mathbf{R}} \frac{1}{(\lambda - \sqrt{-1}\omega)^2} g(\omega) d\omega = \frac{1}{\sqrt{-1}K^2 f_1},$$

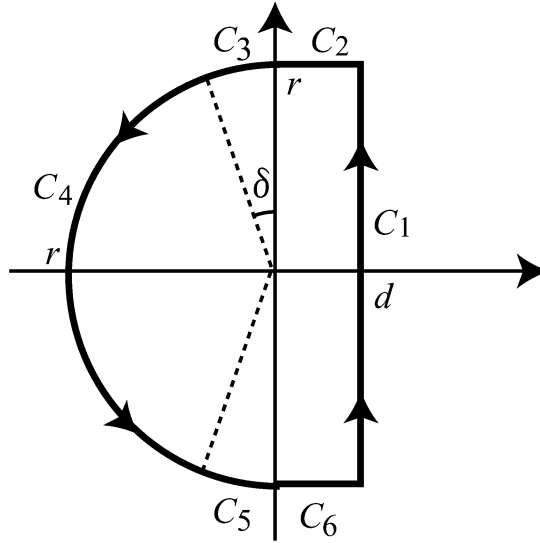


Figure 4. The contour for the Laplace inversion formula.

which proves that $\lambda'(K) \neq 0$. Further, roots y of Eq.(3.18), which determines eigenvalues on the imaginary axis, are isolated because both side of Eq.(3.18) are analytic with respect to y . This means that $\lambda(K)$ can not move along the imaginary axis. This proves that an eigenvalue $\lambda(K)$ gets across the imaginary axis from right to left as K decreases from $K_c^{(1)}$, which gives a root of Eq.(3.28). Note that there may exist resonance poles which are not continuations of eigenvalues (see Example 3.11).

Next, let us prove Eq.(3.31). Let $d > 0$ be a small number and r sufficiently large number. Take paths C_1 to C_6 as are shown in Fig.4:

$$\begin{aligned}
 C_1 &= \{d + \sqrt{-1}y \mid -r \leq y \leq r\}, \\
 C_2 &= \{x + \sqrt{-1}r \mid 0 \leq x \leq d\}, \\
 C_3 &= \{re^{\sqrt{-1}\theta} \mid \pi/2 \leq \theta \leq \pi/2 + \delta\}, \\
 C_4 &= \{re^{\sqrt{-1}\theta} \mid \pi/2 + \delta \leq \theta \leq 3\pi/2 - \delta\}, \\
 C_5 &= \{re^{\sqrt{-1}\theta} \mid \pi/2 + \delta \leq \theta \leq 3\pi/2 - \delta\}, \\
 C_6 &= \{x + \sqrt{-1}r \mid 0 \leq x \leq d\},
 \end{aligned}$$

and C_5 and C_6 are defined in a similar way to C_3 and C_2 , respectively. We put $C = \sum_{j=1}^6 C_j$.

Let $\alpha_1, \alpha_2, \dots, \alpha_{n(C)}$ be resonance poles inside the closed curve C , where we assume that there are no resonance poles on the curve C by deforming it slightly if necessary. Let $R_1(t), R_2(t), \dots, R_{n(C)}(t)$ be corresponding residues of $e^{\lambda t} F_1(\lambda)$, respectively. Note that if α_j is a pole of $F_1(\lambda)$ of order m_j , $R_j(t)$ is of the form $R_j(t) = p_j(t)e^{\alpha_j t}$ with a

polynomial $p_j(t)$ of degree $m_j - 1$. By the residue theorem, we have

$$2\pi\sqrt{-1} \sum_{j=1}^{n(C)} R_j(t) = \int_{C_6+C_1+C_2} e^{\lambda t} F_0(\lambda) d\lambda + \int_{C_3+C_4+C_5} e^{\lambda t} F_1(\lambda) d\lambda.$$

The integral $\int_{C_1} e^{\lambda t} F_0(\lambda) d\lambda / (2\pi\sqrt{-1})$ converges to $Z_1^0(t)$ as $r \rightarrow \infty$. It is easy to show that the integrals along C_2, C_3, C_5, C_6 tend to zero as $r \rightarrow \infty$ because of the assumption (iv). We have to estimate the integral along C_4 as

$$\begin{aligned} \left| \int_{C_4} e^{\lambda t} F_1(\lambda) d\lambda \right| &\leq \int_{\pi/2+\delta}^{3\pi/2-\delta} r e^{rt \cos \theta} |F_1(re^{\sqrt{-1}\theta})| d\theta \\ &\leq \max_{\pi/2+\delta \leq \theta \leq 3\pi/2-\delta} |F_1(re^{\sqrt{-1}\theta})| \int_{\pi/2+\delta}^{3\pi/2-\delta} r e^{rt \cos \theta} d\theta \\ &\leq D e^{\beta r} \int_{\delta}^{\pi/2} 2r e^{-rt \sin \phi} d\phi \\ &\leq D e^{\beta r} \int_{\delta}^{\pi/2} 2r_2 e^{-2rt\phi/\pi} d\phi \\ (3.33) \quad &\leq D e^{\beta r} \cdot \frac{\pi}{t} \left(e^{-2rt\delta/\pi} - e^{-rt} \right). \end{aligned}$$

Thus if $t > t_0 := \max\{\beta, \pi\beta/(2\delta)\}$, this integral tends to zero as $r \rightarrow \infty$. \square

Example 3.10. If $g(\omega)$ is a rational function, the assumptions are satisfied when $q^*(\lambda)$ is bounded on the upper half plane. In this case, the number of resonance poles is finite and thus Eq.(3.31) becomes finite sum. For example if $g(\omega) = 1/(\pi(1 + \omega^2))$ is the Lorentzian distribution, a resonance pole is given by $\lambda = \sqrt{-1}Kf_1 - 1$ (a root of Eq.(3.28)). Therefore $Z_1^0(t)$ decays with the exponential rates $\text{Re}(\sqrt{-1}Kf_1 - 1)$.

Example 3.11. If $g(\omega)$ is the Gaussian distribution, the assumptions are satisfied when $q^*(\lambda)$ is of exponential type; that is, there exist positive constants C and β such that $|q^*(\lambda)| \leq C e^{\beta|\lambda|}$. Since the analytic continuation $g^*(\lambda)$ has an essential singularity at infinity, there exist infinitely many resonance poles and they accumulate at infinity.

§ 3.5. Semi-group generated by the operator T_1 ($\text{Im}(f_1) \geq 0$)

In Sec.3.1 and Sec.3.4, we investigate the semi-group generated by the operator $T_1 = \sqrt{-1}\mathcal{M} + \sqrt{-1}Kf_1\mathcal{P}$ for the cases $f_1 = 0$ and $\text{Im}(f_1) < 0$, respectively. In this subsection, we consider the case $\text{Im}(f_1) \geq 0$.

Theorem 3.12. Suppose that the assumptions (ii) to (v) of Thm.3.9 hold. If $\text{Im}(f_1) \geq 0$, for an arbitrarily fixed $K > 0$, the order parameter $Z_1^0(t) = (e^{T_1 t} q, P_0)$ decays

exponentially as $t \rightarrow \infty$.

We show an idea of the proof. If $\text{Im}(f_1) = 0$, $T_1/\sqrt{-1} = \mathcal{M} + K\text{Re}(f_1)\mathcal{P}$ is self-adjoint and a rank one perturbation of the multiplication \mathcal{M} . By Theorem X-4.3 in [8], $T_1/\sqrt{-1}$ and \mathcal{M} are unitarily equivalent. Since $(e^{\sqrt{-1}\mathcal{M}t}q, P_0)$ decays exponentially (see Sec.3.1), we can prove that so is $(e^{T_1t}q, P_0)$.

If $\text{Im}(f_1) > 0$, change the parameter as $K \mapsto -K$. Then, the problem is reduced to the case $K < 0$ and $\text{Im}(f_1) < 0$, and Thm.3.12 is proved in a similar manner to the proof of Thm.3.9.

References

- [1] L. V. Ahlfors, *Complex analysis. An introduction to the theory of analytic functions of one complex variable*, McGraw-Hill Book Co., New York, 1978.
- [2] N. J. Balmforth, R. Sassi, A shocking display of synchrony, *Phys. D*, **143** (2000), 21–55.
- [3] H. Chiba, D. Pazó, Stability of an $[N/2]$ -dimensional invariant torus in the Kuramoto model at small coupling, *Phys. D*, **238** (2009), 1068–1081.
- [4] H. Daido, Order Function and Macroscopic Mutual Entrainment in Uniformly Coupled Limit-Cycle Oscillators, *Prog. Theor. Phys.*, **88** (1992), 1213–1218.
- [5] H. Daido, Critical Conditions of Macroscopic Mutual Entrainment in Uniformly Coupled Limit-Cycle Oscillators, *Prog. Theor. Phys.*, **89** (1993), 929–934.
- [6] H. Daido, Generic scaling at the onset of macroscopic mutual entrainment in limit-cycle oscillators with uniform all-to-all coupling, *Phys. Rev. Lett.*, **73** (1994), 760.
- [7] H. Daido, Onset of cooperative entrainment in limit-cycle oscillators with uniform all-to-all interactions: bifurcation of the order function, *Phys. D*, **91** (1996), 24–66.
- [8] T. Kato, *Perturbation theory for linear operators*, Springer-Verlag, Berlin, 1995.
- [9] Y. Kuramoto, Self-entrainment of a population of coupled non-linear oscillators, *International Symposium on Mathematical Problems in Theoretical Physics, Lecture Notes in Phys.*, **39**. Springer, Berlin, 1975.
- [10] Y. Kuramoto, *Chemical oscillations, waves, and turbulence*, Springer Series in Synergetics, 19. Springer-Verlag, Berlin, 1984.
- [11] Y. Maistrenko, O. Popovych, O. Burylko, P. A. Tass, Mechanism of desynchronization in the finite-dimensional Kuramoto model, *Phys. Rev. Lett.*, **93** (2004) 084102.
- [12] Y. L. Maistrenko, O. V. Popovych, P. A. Tass, Chaotic attractor in the Kuramoto model, *Int. J. of Bif. and Chaos*, **15** (2005) 3457–3466.
- [13] E. A. Martens, E. Barreto, S. H. Strogatz, E. Ott, P. So, T. M. Antonsen, Exact results for the Kuramoto model with a bimodal frequency distribution, *Phys. Rev. E*, **79** (2009) 026204.
- [14] R. E. Mirollo, S. H. Strogatz, The spectrum of the locked state for the Kuramoto model of coupled oscillators, *Phys. D*, **205** (2005), 249–266.
- [15] R. Mirollo, S. H. Strogatz, The spectrum of the partially locked state for the Kuramoto model, *J. Nonlinear Sci.*, **17** (2007), 309–347.
- [16] A. Pikovsky, M. Rosenblum, J. Kurths, *Synchronization: A Universal Concept in Nonlinear Sciences*, Cambridge University Press, Cambridge, 2001.

- [17] M. Reed, B. Simon, *Methods of modern mathematical physics IV. Analysis of operators*, Academic Press, New York-London, 1978.
- [18] S. H. Strogatz, From Kuramoto to Crawford: exploring the onset of synchronization in populations of coupled oscillators, *Phys. D*, **143** (2000), 1–20.
- [19] S. H. Strogatz, R. E. Mirollo, P. C. Matthews, Coupled nonlinear oscillators below the synchronization threshold: relaxation by generalized Landau damping, *Phys. Rev. Lett.*, **68** (1992), 2730–2733.
- [20] S. H. Strogatz, R. E. Mirollo, Stability of incoherence in a population of coupled oscillators, *J. Statist. Phys.*, **63** (1991), 613–635.
- [21] N. J. Vilenkin, *Special functions and the theory of group representations*, American Mathematical Society, 1968.
- [22] K. Yosida, *Functional analysis*, Springer-Verlag, Berlin, 1995.