

# BEHAVIOR OF RADIALY SYMMETRIC SOLUTIONS OF A SYSTEM RELATED TO CHEMOTAXIS

TOSHITAKA NAGAI 永井 敏隆

Department of Mathematics, Kyushu Institute of Technology, Tobata, Kitakyushu 804  
JAPAN

TAKASI SENBA 仙葉 隆

Department of Applied Mathematics, Miyazaki University, Kibana, Miyazaki 889-21  
JAPAN

## 1. Introduction

We consider time-global existence and blow-up of solutions of the following system related to chemotaxis

$$\begin{cases} b_t = \nabla \cdot (\nabla b - \chi b \nabla \phi(s)) & \text{in } \Omega \times (0, \infty), \\ 0 = \Delta s - s + b & \text{in } \Omega \times (0, \infty), \end{cases} \quad (1.1)$$

under the conditions

$$\begin{cases} \frac{\partial b}{\partial n} = \frac{\partial s}{\partial n} = 0 & \text{on } \partial\Omega \times (0, \infty), \\ b(\cdot, 0) = b_0 & \text{in } \Omega, \end{cases} \quad (1.2)$$

where  $\chi$  is a positive constant and  $\phi$  is a smooth function on  $(0, \infty)$  with  $\phi' > 0$ . The system is a simplified Keller-Segel model. Keller-Segel model was introduced by Keller and Segel [11] to describe the initiation of chemotactic aggregation of cellular slime molds. On Keller-Segel model and simplified Keller-Segel models, time-local existence of the solutions has been studied by [19] and blow-up of the solutions has been studied by [4, 10, 9, 14, 18].

The domain  $\Omega$  and the non-trivial initial function  $b_0$  are only confined to the following case:

(A1)  $\Omega$  is the open ball of radius  $L$  with center at the origin in  $\mathbf{R}^N$ .

(A2)  $b_0$  is smooth and nonnegative on  $\bar{\Omega}$ , and is radially symmetric when  $N \geq 2$ .

Under these assumptions, there exists a unique solution  $(b(x, t), s(x, t))$  to (1.1) and (1.2) defined maximal interval of existence  $[0, T_{max})$ , which is radially symmetric in  $x$  when  $N \geq 2$ , smooth in  $\bar{\Omega} \times (0, T_{max})$  and  $b(x, t) > 0$ ,  $s(x, t) > 0$  for  $(x, t) \in \Omega \times (0, T_{max})$ . If  $T_{max} < \infty$ ,

$$\limsup_{t \rightarrow T_{max}} (\|b(\cdot, t)\|_{L^\infty} + \|s(\cdot, t)\|_{L^\infty}) = \infty,$$

by which we mean that  $(b(x, t), s(x, t))$  blows up in finite time.

**Theorem 1** *Let  $N = 1$  and  $\phi$  be smooth on  $(0, \infty)$ . Then the solution  $(b, s)$  to (1.1), (1.2) is globally bounded, that is,  $T_{max} = \infty$  and  $(b, s)$  satisfies*

$$\sup_{t \geq 0} (\|b(\cdot, t)\|_{L^\infty} + \|s(\cdot, t)\|_{L^\infty}) < \infty.$$

We put

$$M_a(t) = \int_{\Omega} b(x, t) |x|^a dx \quad \text{for } 0 \leq t < T_{max},$$

where  $a$  is a positive constant. That is called the moment of order  $a$ , of  $b(\cdot, t)$ .

**Theorem 2** *Assume  $\phi(s) = s^p$  ( $p > 0$ ), (A1) and (A2).*

(1)  $N = 2$  :

(a) *If  $0 < p < 1$ , then the solution is globally bounded in time.*

(b)  $p = 1$  :

(i) *If  $\|b_0\|_{L^1} < 8\pi/\chi$ , then the solution is globally bounded in time.*

(ii) *If  $\|b_0\|_{L^1} > 8\pi/\chi$  and  $M_2(0)$  is sufficiently small, then the solution blows up in finite time.*

(c) *If  $p > 1$  and  $M_2(0)$  is sufficiently small, then the solution blows up in finite time.*

(2) *If  $N \geq 3$  and  $M_{(N-2)p+2}(0)$  is sufficiently small, then the solution blows up in finite time.*

**Theorem 3** *Assume  $\phi(s) = \log s$ , (A1) and (A2).*

(1) *If  $N = 2$ , then the solution is globally bounded in time.*

(2)  $N \geq 3$  :

(a) *If  $\chi < 2/(N-2)$ , then the solution is globally bounded in time.*

(b) *If  $\chi > 2N/(N-2)$  and  $M_2(0)$  is sufficiently small, then the solution blows up in finite time.*

## 2. Time-global existence and boundedness

The purpose in this section is to sketch the proofs of Theorem 1 and (i) in Theorems 2 and 3.

Let  $G$  be the Green function of  $-\Delta + 1$  in  $\Omega$  with homogeneous Neumann boundary conditions. For  $N \geq 2$  we put

$$E(r) = (2\pi)^{-N/2} r^{(2-N)/2} \kappa_{(N-2)/2}(r) \quad \text{for } r > 0,$$

where  $\kappa_\nu$  is the modified Bessel function of the second kind of order  $\nu$  (see [13]).  $E$  is a fundamental solution of  $-\Delta + 1$ .

For the solution  $(b, s)$  to (1.1), (1.2) define the functions  $S$  and  $B$  by

$$S(r, t) = \int_{|x| \leq r} s(x, t) dx, \quad B(r, t) = \int_{|x| \leq r} b(x, t) dx \quad (2.1)$$

for  $0 \leq r \leq L$  and  $0 \leq t < T_{max}$ , respectively.  $B$  and  $S$  satisfy

$$\frac{\partial B}{\partial t} = r^{N-1} \frac{\partial}{\partial r} \left( r^{1-N} \frac{\partial B}{\partial r} \right) + \frac{\chi}{\omega_N} (B - S) \phi'(s) r^{1-N} \frac{\partial B}{\partial r}, \quad (2.2)$$

$$0 = r^{N-1} \frac{\partial}{\partial r} \left( r^{1-N} \frac{\partial S}{\partial r} \right) - S + B, \quad (2.3)$$

for  $0 < r < L$  and  $0 < t < T_{max}$ , and

$$B(0, t) = S(0, t) = 0, \quad B(L, t) = S(L, t) = \|b_0\|_{L^1},$$

where  $\omega_N$  is the surface area of the unit sphere  $S^{N-1}$  in  $\mathbf{R}^N$ .

In order to show the boundedness and time-global existence of solutions  $(b, s)$  to (1.1), (1.2), we begin with the following lemmas. These lemmas are shown by the arguments similar to those in [14] and [18], respectively, so we omit the proofs. In what follows,  $C$  denotes a generic positive constant depending on  $L$  and  $N$ .

**Lemma 2.1** *Let  $N \geq 2$ . Then*

$$s(x, t) \geq C \|b_0\|_{L^1} \quad \text{for } x \in \bar{\Omega} \text{ and } t \in (0, T_{max}).$$

**Lemma 2.2** *If the following condition*

$$\sup_{0 \leq t < T_{max}} \|s(\cdot, t)\|_{L^\infty} < \infty, \quad \sup_{0 \leq t < T_{max}} \|\nabla \phi(s(\cdot, t))\|_{L^\infty} < \infty,$$

*holds, then  $T_{max} = \infty$  and*

$$\sup_{t > 0} \|b(\cdot, t)\|_{L^\infty} < \infty.$$

For the following lemma, see [18].

**Lemma 2.3** *Let  $N \geq 2$ . Then the following holds :*

$$B(|x|, t) E(|x|) \leq s(x, t) \leq C \|b_0\|_{L^1} E(|x|) \quad \text{in } \Omega \setminus \{0\} \times (0, T_{max}).$$

**Sketch of proofs of Theorem 1 and (i) in Theorems 2 and 3.** By Lemmas 2.1, 2.2 and Appendixes in [14] and [18], it suffices to show that

$$\sup_{0 \leq t < T_{max}} \|s(\cdot, t)\|_{L^\infty} < \infty, \quad \sup_{0 \leq t < T_{max}} \|\nabla s(\cdot, t)\|_{L^\infty} < \infty. \quad (2.4)$$

In the case of  $N = 1$ , (2.4) is shown by the arguments similar to those in [14]. Hence we will prove (2.4) in the case of  $N \geq 2$ .

We put

$$\Phi(u) = \begin{cases} p \|b_0\|_{L^1}^p u^{p-1} & \text{in the case of Theorem 2,} \\ u^{-1} & \text{in the case of Theorem 3} \end{cases}$$

for  $u > 0$ . It follows from Lemma 2.3, (2.2) and  $\partial B/\partial r \geq 0$  that  $B$  satisfies

$$\frac{\partial B}{\partial t} \leq r^{N-1} \frac{\partial}{\partial r} \left( r^{1-N} \frac{\partial B}{\partial r} \right) + \frac{\chi}{\omega_N} \Phi(E) r^{1-N} \frac{\partial B}{\partial r}.$$

We can construct the function  $W(r)$  such that

$$W(r) \sim r^N \text{ as } r \rightarrow 0,$$

$$\|b_0\|_{L^1} < W(L) \quad \text{and} \quad B(r, 0) \leq W(r) \quad \text{for } 0 \leq r \leq L,$$

and that

$$0 \geq r^{N-1} \frac{d}{dr} \left( r^{1-N} \frac{dW}{dr} \right) + \frac{\chi}{\omega_N} B \phi'(s) r^{1-N} \frac{dW}{dr} \text{ for } 0 < r < L.$$

Hence, the comparison theorem yields that

$$B(r, t) \leq W(r) \quad \text{for } 0 \leq r \leq L, 0 \leq t < T_{max},$$

which implies  $B(r, t) \leq Cr^N$ .

Since  $B(r, t) \leq Cr^N$  for  $0 \leq r \leq L$  and  $0 \leq t < T_{max}$ , it follows from (2.3) that

$$S(r, t) \leq Cr^N \quad \text{for } 0 \leq r \leq L, 0 \leq t < T_{max}.$$

Then we have that

$$|\nabla s(x, t)| = \frac{|S(|x|, t) - B(|x|, t)|}{\omega_N |x|^{N-1}} \leq C$$

for  $x \in \Omega$  and  $0 \leq t < T_{max}$ . The boundedness of  $\|s(\cdot, t)\|_{L^\infty}$  with respect to  $t \in [0, T_{max})$  follows from the estimate above of  $|\nabla s|$  and

$$\min_{x \in \Omega} s(x, t) \leq \frac{\|b_0\|_{L^1}}{|\Omega|} \quad \text{for } 0 \leq t < T_{max},$$

where  $|\Omega|$  is the volume of  $\Omega$ . Thus the proofs of (i) of Theorems 2 and 3 are complete.

### 3. Blow-up of solutions

The purpose in this section is to show the blow-up of solutions for the system (1.1), (1.2) in the case of  $N \geq 2$ .

In order to show the blow-up of solutions  $(b, s)$  to (1.1), (1.2) in [14] and [18], a differential inequality on a moment  $M_k(t)$  of  $b$  is constructed by use of some estimates of  $s$ , and under some conditions on  $b_0$  it is shown that the moment of  $b$  converges to 0 as  $t$  tends some  $T_0 \in (0, \infty)$  by use of the differential inequality.

The following lemma is an immediate consequence of Hölder's inequality.

**Lemma 3.1** *Let  $f$  be an integrable function on  $\Omega$ , and  $p_1, p_2$  and  $p_3$  be numbers satisfying  $0 \leq p_1 < p_2 < p_3$ . Then*

$$\int_{\Omega} |f||x|^{p_2} dx \leq \left\{ \int_{\Omega} |f||x|^{p_1} dx \right\}^{(p_3-p_2)/(p_3-p_1)} \left\{ \int_{\Omega} |f||x|^{p_3} dx \right\}^{(p_2-p_1)/(p_3-p_1)}$$

Let  $S$  and  $B$  be the same functions as in (2.1). The following lemmas are stated in [15] and [18].

**Lemma 3.2** *The inequality holds :*

$$\begin{aligned} \frac{d}{dt} M_k(t) &\leq k(k+N-2) \int_{\Omega} b(x,t) |x|^{k-2} dx \\ &\quad + \frac{k\chi}{\omega_N} \int_{\Omega} \phi'(s(x,t)) b(x,t) \{S(|x|,t) - B(|x|,t)\} |x|^{k-N} dx \end{aligned}$$

on  $(0, T_{max})$ , where  $k \geq 2$ .

**Lemma 3.3** *Let  $N \geq 3$ . There exists a positive constant  $\delta$  such that*

$$\frac{\partial}{\partial r} (r^{N-1} s(x,t)) \geq 0 \quad \text{in } \{x \in R^N : |x| \leq \delta\} \times (0, T_{max}),$$

where  $r = |x|$ .

**Lemma 3.4** *Let  $N \geq 2$ . Then the following holds :*

$$s(x,t) \leq \frac{1}{\omega_N |x|^{N-1}} \int_{|y|=|x|} E(|x-y|) d\sigma \|b_0\|_{L^1} + \int_{\Omega} K(x,y) b(y,t) dy$$

in  $\Omega \setminus \{0\} \times (0, T_{max})$ .

**Sketch of proof of (ii) of Theorem 2.** Let  $k = (N-2)p + 2$ . In order to prove the theorem, it suffices to show the following inequality

$$\begin{aligned} \frac{d}{dt} M_k(t) &\leq k(k+N-2) \|b_0\|_{L^1}^{2/k} M_k(t)^{(k-2)/k} \\ &\quad + C \|b_0\|_{L^1}^{p+(k-2)/k} M_k(t)^{2/k} - C \|b_0\|_{L^1}^{p+1} \end{aligned} \quad (3.1)$$

for  $t \in (0, T_{max})$ . In fact, if  $M_k(0)$  is sufficiently small so that the right-hand side of (3.1) is negative at  $t = 0$ , there exists  $T_0 \in (0, \infty)$  such that

$$M_k(t) \rightarrow 0 \quad \text{as } t \rightarrow T_0.$$

Hence,  $T_{max}$  must be finite and  $T_{max} \leq T_0$ . By Appendixes in [14] and [18], we have

$$\limsup_{t \rightarrow T_{max}} \|b(\cdot, t)\|_{L^\infty} = \infty.$$

Let us first show (3.1) in the case of  $p \geq 1$ . Using Lemmas 2.3 and the properties of the fundamental solution, we obtain that

$$\int_{\Omega} s^{p-1}(x,t) b(x,t) B(|x|,t) |x|^{k-N} dx \geq C \|b_0\|_{L^1}^{p+1}. \quad (3.2)$$

and that

$$S(|x|,t) \leq C \|b_0\|_{L^1} |x|^2. \quad (3.3)$$

It follows from Lemma 2.3 and (3.3) and the properties of the fundamental solution that

$$\int_{\Omega} s^{p-1}(x, t) b(x, t) S(|x|, t) |x|^{k-N} dx \leq C \|b_0\|_{L^1}^p M_2(t). \quad (3.4)$$

Lemma 3.2 together with (3.2), (3.4) and Lemma 3.1 yields (3.1).

Let us consider the case  $0 < p < 1$ . By Lemmas 2.3 and the properties of the fundamental solution, we have

$$\begin{aligned} & \int_{\Omega} s^{p-1}(x, t) b(x, t) B(|x|, t) |x|^{k-N} dx \\ & \geq C \|b_0\|_{L^1}^{p-1} \int_{\Omega} b(x, t) B(|x|, t) dx = \frac{C}{2} \|b_0\|_{L^1}^{p+1}. \end{aligned} \quad (3.5)$$

It follows from Lemmas 2.3 and 3.3 and the properties of the fundamental solution that

$$\int_{|x| \leq \delta} s^{p-1}(x, t) b(x, t) S(|x|, t) |x|^{k-N} dx \leq C \|b_0\|_{L^1}^p \int_{|x| \leq \delta} b(x, t) |x|^2 dx. \quad (3.6)$$

By Lemma 2.1 and (3.3), we have

$$\begin{aligned} & \int_{\delta \leq |x| \leq L} s^{p-1}(x, t) b(x, t) S(|x|, t) |x|^{k-N} dx \\ & \leq C \|b_0\|_{L^1}^p \int_{\delta \leq |x| \leq L} b(x, t) |x|^{k-N+2} dx \\ & \leq C \|b_0\|_{L^1}^p \delta^{k-N} \int_{\delta \leq |x| \leq L} b(x, t) |x|^2 dx. \end{aligned} \quad (3.7)$$

Combining (3.6) with (3.7) yields that

$$\int_{\Omega} s^{p-1}(x, t) b(x, t) S(|x|, t) |x|^{k-N} dx \leq C \|b_0\|_{L^1}^p M_2(t). \quad (3.8)$$

By (3.5) and (3.8), the similar argument to that in the case of  $p \geq 1$  gives us (3.1). Thus the proof is complete.

**Sketch of proof of (ii) in Theorem 3.** Observe that it follows from Lemma 3.4 and the properties of the fundamental solution that for  $0 \leq t < T_{max}$  and  $0 < |x| \leq L/2$ ,

$$s(x, t) \leq \left\{ \frac{1}{\omega_N (N-2) |x|^{N-2}} + C(|x|^{3-N} + 1) \right\} \|b_0\|_{L^1}.$$

For  $0 \leq t < T_{max}$  and  $0 < \delta \leq L/2$ , we then have that

$$\begin{aligned} & \int_{\Omega} b(x, t) B(|x|, t) \frac{1}{|x|^{N-2} s(x, t)} dx \\ & \geq \frac{1}{\|b_0\|_{L^1}} \left\{ \frac{1}{(N-2)\omega_N} + C\delta \right\}^{-1} \int_{|x| \leq \delta} b(x, t) B(|x|, t) dx \end{aligned}$$

$$\begin{aligned}
&\geq \frac{(N-2)\omega_N}{2(1+C\delta)\|b_0\|_{L^1}} B(\delta, t)^2 & (3.9) \\
&\geq \frac{(N-2)\omega_N}{2(1+C\delta)\|b_0\|_{L^1}} \left( \|b_0\|_{L^1} - \frac{1}{\delta^2} M_2(t) \right)_+^2 \\
&\geq \frac{(N-2)\omega_N}{2(1+C\delta)} \left( \|b_0\|_{L^1} - \frac{2}{\delta^2} M_2(t) \right),
\end{aligned}$$

where  $(\cdot)_+ = \max\{\cdot, 0\}$ . It follows from Lemma 3.3 that

$$\begin{aligned}
&\int_{\Omega} b(x, t) S(|x|, t) \frac{1}{|x|^{N-2} s(x, t)} dx \\
&= \omega_N \int_{|x| \leq \delta} b(x, t) |x|^2 dx + \int_{\delta \leq |x| \leq L} b(x, t) S(|x|, t) \frac{1}{|x|^{N-2} s(x, t)} dx & (3.10) \\
&\leq CM_2(t)
\end{aligned}$$

in  $(0, T_{max})$ . Hence, combining Lemma 3.2 with (3.9) and (3.10) concludes that

$$\frac{d}{dt} M_2(t) \leq \left\{ 2N - \frac{(N-2)\chi}{1+C\delta} \right\} \|b_0\|_{L^1} + C\chi(1+\delta^{-2})M_2(t)$$

on  $(0, T_{max})$ . Suppose that  $\delta$  is sufficiently small so that  $2N(1+C\delta) - (N-2)\chi < 0$ . Using the argument similar to that in the sketch of proof of Theorem 2, then we have the proof.

## References

- [1] P. Biler and T. Nadzieja, Existence and nonexistence of solutions for a model of gravitational interactions of particles, I, *Colloq. Math.*, **66** (1994), 319–334.
- [2] P. Biler, Existence and nonexistence of solutions for a model of gravitational interactions of particles, III, *Colloq. Math.*, **68** (1995), 229–239.
- [3] S. Childress, Chemotactic collapse in two dimensions, *Lecture Notes in Biomath.*, vol. 55, Springer, Berlin-Heidelberg-New York, 1984, 61–66.
- [4] S. Childress and J. K. Percus, Nonlinear aspects of chemotaxis, *Math. Biosci.*, **56** (1981), 217–237.
- [5] J. I. Diaz and T. Nagai, Symmetrization in a parabolic-elliptic system related to chemotaxis, *Adv. Math. Sci. Appl.*, **5** (1995), 659–680.
- [6] G. B. Folland, *Introduction to Partial Differential Equations*, Princeton University Press, Princeton, New Jersey, 1976.
- [7] P. R. Garabedian, *Partial Differential Equations*, John Wiley & Sons, New York, 1964.
- [8] D. Henry, *Geometric Theory of Semilinear Parabolic Equations*, *Lecture Notes in Math.*, vol. 840, Springer, Berlin-Heidelberg-New York, 1981.

- [9] M. A. Herrero and J. J. L. Velázquez, Singularity patterns in a chemotaxis model, *Math. Ann.*, to appear.
- [10] W. Jäger and S. Luckhaus, On explosions of solutions to a system of partial differential equations modelling chemotaxis, *Trans. Amer. Math. Soc.*, **329** (1992), 819–824.
- [11] E. F. Keller and L. A. Segel, Initiation of slime mold aggregation viewed as an instability, *J. Theor. Biol.*, **26** (1970), 399–415.
- [12] I. R. Lapidus and M. Levandowsky, Modeling chemosensory responses of swimming eukaryotes, *Biological Growth and Spread, Proceedings, Heidelberg 1979, Lecture Notes in Biomathematics*, vol. 38, Springer-Verlag, 1980.
- [13] T. Myint-U and L. Debnath, *Partial Differential Equations for Scientists and Engineers*, North-Holland, New York, 1987.
- [14] T. Nagai, Blow-up of radially symmetric solutions to a chemotaxis system, *Adv. Math. Sci. Appl.*, **5** (1995), 581–601.
- [15] T. Nagai and T. Senba, Global existence and blow-up of radially symmetric solutions to a parabolic-elliptic system to chemotaxis, preprint.
- [16] V. Nanjundiah, Chemotaxis, signal relaying, and aggregation morphology, *J. Theor. Biol.*, **42** (1973), 63–105.
- [17] R. Schaaf, Stationary solutions of chemotaxis systems, *Trans. Amer. Math. Soc.*, **292** (1985), 531–556.
- [18] T. Senba, Blow-up of radially symmetric solutions to some systems of partial differential equations modelling chemotaxis, *Adv. Math. Sci. Appl.*, to appear.
- [19] A. Yagi, Norm behavior of solutions to the parabolic system of chemotaxis, preprint.