

# Anticipating Quantum Stochastic Integrals for Basic Quantum Martingales

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## 1 Introduction

Since the quantum stochastic integrals of adapted quantum stochastic processes have been introduced by Hudson and Parthasarathy [10] as a quantum extension of the Itô (stochastic) integral, the quantum stochastic calculus has been studied extensively with wide applications (see [28, 32]).

The Hudson-Parthasarathy quantum stochastic integrals has been extended to the quantum stochastic integrals of nonadapted quantum stochastic processes by Belavkin [3], Lindsay [24] and Attal & Lindsay [2]. Since then the nonadapted quantum stochastic integral has been studied systematically in terms of quantum stochastic gradients by Ji & Obata [16, 18]. Based on the quantum white noise theory [12], the notion of quantum white noise derivatives has been introduced by Ji & Obata (see [14, 15, 16, 17, 19, 20]). The explicit formulas [16] of integrands for quantum stochastic integral representation of quantum martingales [11] has been derived in terms of the quantum white noise derivatives. Also, the notion of quantum stochastic gradients [18] has been introduced based on the notion of the quantum white noise derivatives. Recently, Ji & Sinha [21] studied the quantum stochastic integrals for quadratic quantum noises.

On the other hand, based on the white noise theory [8, 22, 30] introduced by Hida, Kuo & Russek [23] studied anticipating (classical) stochastic integrals by applying the quantum decomposition of a Brownian motion.

In this paper, we study some regularity properties of the quantum Hitsuda-Skorohod integrals as anticipating quantum stochastic integrals. Also, motivated by the results in [23], we discuss new types of anticipating quantum stochastic integrals in terms of pointwisely defined quantum white noise derivatives.

## 2 Admissible Generalized Operators

### 2.1 Admissible Rigging of Fock Space

We now review a construction of admissible rigging of Fock space which provides the basic structure of this paper. Let  $H = L^2(\mathbb{R}_+, dt)$  be the Hilbert space of complex valued

square integrable functions on  $\mathbb{R}_+ = [0, \infty)$  with respect to the Lebesgue measure  $dt$  and let  $\Gamma(H)$  be the Fock space over  $H$  defined by

$$\Gamma(H) = \left\{ \phi = (f_n)_{n=0}^\infty; f_n \in H^{\otimes n}, \sum_{n=0}^\infty n! |f_n|^2 < \infty \right\},$$

where  $H^{\otimes n}$  is the  $n$ -fold symmetric tensor product of  $H$  and  $|\cdot|$  is the Hilbertian norm on  $H$  and  $H^{\otimes n}$ . For  $p \geq 0$ , we set

$$\mathcal{G}_p = \left\{ \phi = (f_n)_{n=0}^\infty \in \Gamma(H); \|\phi\|_p^2 = \sum_{n=0}^\infty n! e^{2pn} |f_n|^2 < \infty \right\}$$

and  $\mathcal{G}_{-p}$  to be the completion of  $\Gamma(H)$  with respect to the norm  $\|\cdot\|_{-p}$  defined by

$$\|\phi\|_{-p}^2 = \sum_{n=0}^\infty n! e^{-2pn} |f_n|^2.$$

Then  $\{\mathcal{G}_p; p \in \mathbb{R}\}$  forms a chain of weighted Fock spaces and so we have

$$\mathcal{G} = \text{proj lim}_{p \rightarrow \infty} \mathcal{G}_p \subset \mathcal{G}_p \subset \mathcal{G}_0 = \Gamma(H) \subset \mathcal{G}_{-p} \subset \mathcal{G}^* \cong \text{ind lim}_{p \rightarrow \infty} \mathcal{G}_{-p}$$

for  $p \geq 0$ , where the strong dual space  $\Gamma(H)^*$  of  $\Gamma(H)$  is identified with  $\Gamma(H)$ , and the strong dual space  $\mathcal{G}^*$  of  $\mathcal{G}$  is topologically isomorphic to the inductive limit space  $\text{ind lim}_{p \rightarrow \infty} \mathcal{G}_{-p}$ . The canonical  $\mathbb{C}$ -bilinear form  $\langle\langle \cdot, \cdot \rangle\rangle$  on  $\mathcal{G}^* \times \mathcal{G}$  takes the form:

$$\langle\langle \Phi, \phi \rangle\rangle = \sum_{n=0}^\infty n! \langle f_n, g_n \rangle, \quad \Phi = (f_n) \in \mathcal{G}^*, \quad \phi = (g_n) \in \mathcal{G},$$

where  $\langle f_n, g_n \rangle$  is the canonical  $\mathbb{C}$ -bilinear form on  $H^{\otimes n} \times H^{\otimes n}$ . Note that  $\mathcal{G}$  is a countable Hilbert space but not necessarily a nuclear space. An element in  $\mathcal{G}$  is said to be *admissible* or *regular*.

**Remark 2.1** Let  $E_{\mathbb{R}} \subset H_{\mathbb{R}} \subset E_{\mathbb{R}}^*$  be a Gelfand triple, i.e.  $E_{\mathbb{R}}$  is a nuclear space, where  $H_{\mathbb{R}} = L_{\mathbb{R}}^2(\mathbb{R}_+, dt)$  is the Hilbert space of real valued square integrable functions on  $\mathbb{R}_+$  with respect to  $dt$ . Then for the standard Gaussian measure  $\mu$  on  $E_{\mathbb{R}}^*$  characterized by

$$\int_{E_{\mathbb{R}}^*} e^{i\langle x, \xi \rangle} d\mu(x) = e^{-\frac{1}{2}|\xi|^2}, \quad \xi \in E_{\mathbb{R}},$$

where  $\langle \cdot, \cdot \rangle$  is the canonical bilinear form on  $E_{\mathbb{R}}^* \times E_{\mathbb{R}}$  again, by the Wiener-Itô-Segal isomorphism,  $\Gamma(H)$  is unitarily equivalent with the Hilbert space  $L^2(E_{\mathbb{R}}^*, \mu)$  of complex valued square integrable functions on  $E_{\mathbb{R}}^*$  with respect to the Gaussian measure  $\mu$ . In this sense, the elements of  $\mathcal{G}$  are considered as admissible Gaussian functionals. The spaces  $\mathcal{G}$  and  $\mathcal{G}^*$  were introduced by Belavkin [3] and have appeared along with classical and quantum stochastic analysis, see e.g., [1, 4, 7, 11, 13, 14, 24, 25, 26, 33, 34].

## 2.2 Multiplications of Admissible Gaussian Functionals

Let  $\phi = (f_n), \psi = (g_n) \in \mathcal{G}$  be given. Suppose that  $f_n = 0$  and  $g_m = 0$  except for finite numbers of  $n$  and  $m$ . Then the *Wiener product* (or *pointwise multiplication*)  $\phi\psi \in \mathcal{G}$  of  $\phi$  and  $\psi$  is defined by

$$\phi\psi = (h_n), \quad h_n = \sum_{l+m=n} \sum_{k=0}^{\infty} k! \binom{l+k}{k} \binom{m+k}{k} f_{l+k} \widehat{\otimes}_k g_{m+k}, \quad (2.1)$$

where  $f_{l+k} \widehat{\otimes}_k g_{m+k}$  is the  $k$ -contraction of  $f_{l+k}$  and  $g_{m+k}$ ; see [30].

The following lemma is useful to study the continuities of Wiener product of admissible Gaussian functionals and similar estimates can be found in [25] (see also [30, 33]).

**Lemma 2.2** *Let  $\phi = (f_n), \psi = (g_n) \in \mathcal{G}$  be given. Suppose that  $f_n = 0$  and  $g_m = 0$  except for a finite numbers of  $n$  and  $m$ . Then for any  $p, r, s \in \mathbb{R}$  with  $r + s > 0$  and*

$$(n+1) \left( \frac{e^{(s-3r)/2} + e^{(r-3s)/2}}{e^{-2p}(r+s)} \right)^n \leq c^n \quad (2.2)$$

for some  $0 < c < 1$ , it holds that

$$\|\phi\psi\|_p^2 \leq \frac{1}{1-c} \|\phi\|_r^2 \|\psi\|_s^2. \quad (2.3)$$

PROOF. For given  $h_n$  as in (2.1), we obtain that

$$\begin{aligned} n!|h_n|^2 &= n! \left( \sum_{l+m=n} \sum_{k=0}^{\infty} k! \binom{l+k}{k} \binom{m+k}{k} |f_{l+k}| |g_{m+k}| \right)^2 \\ &\leq n! \left( \sum_{l+m=n} \sum_{k=0}^{\infty} M_{l,m,k} \sqrt{(l+k)!} e^{r(l+k)} |f_{l+k}| \sqrt{(m+k)!} e^{s(m+k)} |g_{m+k}| \right)^2, \end{aligned} \quad (2.4)$$

where

$$M_{l,m,k} = \frac{e^{-rl-sm}}{l!m!} \frac{\sqrt{(l+k)!(m+k)!}}{k!} e^{-(r+s)k} \leq \frac{e^{-rl-sm}}{l!m!} \sqrt{C_{l,m;r+s}},$$

where

$$C_{l,m;q} = \sup_{n \geq 0} \left\{ \frac{(l+n)! (n+m)!}{n! n!} e^{-2qn} \right\} \leq e^{ql} m^m \left( \frac{e^{q/2}}{eq} \right)^{l+m} < \infty \quad (2.5)$$

for  $q > 0$  (see e.g., [30]: Section 4.1). Therefore, for any  $r \in \mathbb{R}$  and  $s \in \mathbb{R}$  with  $r + s > 0$ , from (2.4), by applying Cauchy-Schwarz inequality we obtain that

$$\begin{aligned} n!|h_n|^2 &\leq n! \left( \sum_{l+m=n} \frac{e^{-rl-sm}}{l!m!} \sqrt{C_{l,m;r+s}} \right)^2 \|\phi\|_r^2 \|\psi\|_s^2 \\ &\leq \left( \sum_{l+m=n} \frac{\sqrt{n!}}{l!m!} e^{-rl-sm} \sqrt{C_{l,m;r+s}} \right)^2 \|\phi\|_r^2 \|\psi\|_s^2. \end{aligned} \quad (2.6)$$

By applying a simple inequality  $n^n \leq e^n n!$ , from (2.5) we see that

$$\sqrt{C_{l,m;r+s}} \leq e^{(r+s)/2} \sqrt{l!m!} \left( \frac{e^{(r+s)/2}}{e(r+s)} \right)^{(l+m)/2} \leq e^{(r+s)/2} \sqrt{l!m!} \left( \frac{e^{(r+s)/2}}{r+s} \right)^{(l+m)/2}$$

Therefore, for any  $r \in \mathbb{R}$  and  $s \in \mathbb{R}$  with  $e^{r+s} \geq 2$ , from (2.6) we obtain that

$$\begin{aligned} n!|h_n|^2 &\leq \left( \sum_{l+m=n} \frac{\sqrt{n!}}{l!m!} e^{-rl-sm} \sqrt{C_{l,m;r+s}} \right)^2 \|\phi\|_r^2 \|\psi\|_s^2 \\ &\leq e^{r+s} \left( \sum_{l+m=n} \frac{\sqrt{n!}}{\sqrt{l!m!}} e^{-rl-sm} \left( \frac{e^{(r+s)/2}}{r+s} \right)^{(l+m)/2} \right)^2 \|\phi\|_r^2 \|\psi\|_s^2 \\ &\leq (n+1) \left( \sum_{l+m=n} \frac{n!}{l!m!} \left( \frac{e^{(s-3r)/2}}{r+s} \right)^l \left( \frac{e^{(r-3s)/2}}{r+s} \right)^m \right) \|\phi\|_r^2 \|\psi\|_s^2 \\ &= (n+1) \left( \frac{e^{(s-3r)/2} + e^{(r-3s)/2}}{r+s} \right)^n \|\phi\|_r^2 \|\psi\|_s^2. \end{aligned} \quad (2.7)$$

Therefore, from (2.1) and (2.7) we obtain that

$$\begin{aligned} \|\phi\psi\|_p^2 &= \sum_{n=0}^{\infty} n! e^{2pn} |h_n|^2 \leq \left[ \sum_{n=0}^{\infty} (n+1) \left( \frac{e^{(s-3r)/2} + e^{(r-3s)/2}}{e^{-2p}(r+s)} \right)^n \right] \|\phi\|_r^2 \|\psi\|_s^2 \\ &\leq \frac{1}{1-c} \|\phi\|_r^2 \|\psi\|_s^2 \end{aligned} \quad (2.8)$$

for some  $0 < c < 1$  satisfying (2.3), which gives the proof.  $\square$

The following two theorem are obvious consequences of Lemma 2.2.

**Theorem 2.3** ([33]) *The Wiener product of admissible Gaussian functionals is continuous from  $\mathcal{G} \times \mathcal{G}$  (equipped with the product topology) onto  $\mathcal{G}$ . In particular,  $\mathcal{G}$  is an algebra with respect to the Wiener product.*

**Theorem 2.4** *The Wiener product of admissible white noise functionals is continuous from  $\mathcal{G}^* \times \mathcal{G}$  (equipped with the product topology) onto  $\mathcal{G}^*$ .*

Let  $\phi = (f_n), \psi = (g_n) \in \mathcal{G}$  be given. Suppose that  $f_n = 0$  and  $g_m = 0$  except for finite numbers of  $n$  and  $m$ . Then the *Wick product* (or *normal-ordered product*)  $\phi \diamond \psi$  of  $\phi$  and  $\psi$  is defined by

$$\phi \diamond \psi = (k_n), \quad k_n = \sum_{l+m=n} f_l \widehat{\otimes} g_m, \quad (2.9)$$

see [6, 8, 22]

The following lemma is useful to study the continuities of Wick product of admissible Gaussian functionals and similar estimates can be found in [33].

**Lemma 2.5** Let  $\phi = (f_n), \psi = (g_n) \in \mathcal{G}$  be given. Suppose that  $f_n = 0$  and  $g_m = 0$  except for a finite numbers of  $n$  and  $m$ . Then for any  $p, r, s \in \mathbb{R}$  satisfying that

$$e^{2(p-r)} + e^{2(p-s)} < 1, \quad (2.10)$$

it holds that

$$\|\phi \diamond \psi\|_p^2 \leq \|\phi\|_r^2 \|\psi\|_s^2. \quad (2.11)$$

PROOF. For given  $k_n$  as in (2.9), we obtain that

$$\begin{aligned} n!|k_n|^2 &= n! \left( \sum_{l+m=n} |f_l| |g_m| \right)^2 \\ &\leq n! \left( \sum_{l+m=n} \frac{e^{-2rl-2sm}}{l!m!} \right) \left( \sum_{l+m=n} l!e^{2rl}|f_l|^2 m!e^{2sm}|g_m|^2 \right) \\ &\leq (e^{-2r} + e^{-2s})^n \left( \sum_{l+m=n} l!e^{2rl}|f_l|^2 m!e^{2sm}|g_m|^2 \right), \end{aligned} \quad (2.12)$$

Therefore, for any  $p, r, s \in \mathbb{R}$  satisfying (2.10), from (2.12) we obtain that

$$\begin{aligned} \|\phi \diamond \psi\|_p^2 &= \sum_{n=0}^{\infty} n!e^{2pn}|k_n|^2 \leq \sum_{n=0}^{\infty} (e^{2(p-r)} + e^{2(p-s)})^n \left( \sum_{l+m=n} l!e^{2rl}|f_l|^2 m!e^{2sm}|g_m|^2 \right) \\ &\leq \|\phi\|_r^2 \|\psi\|_s^2, \end{aligned}$$

which gives the proof.  $\square$

The following theorem is an obvious consequence of Lemma 2.5.

**Theorem 2.6 ([33])** The Wick product is continuous from  $\mathcal{G} \times \mathcal{G}$  (equipped with the product topology) onto  $\mathcal{G}$ , and from  $\mathcal{G}^* \times \mathcal{G}^*$  onto  $\mathcal{G}^*$ . In particular,  $\mathcal{G}$  and  $\mathcal{G}^*$  are algebras under the Wick product.

### 3 Admissible Generalized Operators

We denote by  $\mathcal{L}(\mathfrak{X}, \mathfrak{Y})$  the space of all continuous linear operators from a locally convex space  $\mathfrak{X}$  into another locally convex space  $\mathfrak{Y}$  equipped with the topology of bounded convergence. An operator in  $\mathcal{L}(\mathcal{G}, \mathcal{G}^*)$  is called an *admissible generalized operator* [14] or simply *admissible operator*.

#### 3.1 Integral Kernel Operators

Let  $l, m$  be non-negative integers. Let  $K_{l,m} \in \mathcal{L}(H^{\otimes m}, H^{\otimes l})$  and  $\Phi = (f_n)_{n=0}^{\infty} \in \mathcal{G}^*$ . For each  $n \geq 0$ , we put

$$g_{l+n} = \frac{(n+m)!}{n!} (K_{l,m} \otimes I^{\otimes n} f_{n+m})_{\text{sym}}. \quad (3.1)$$

Then from Lemma 4.1 in [11], for any  $p \in \mathbb{R}$  and  $q > 0$ , we obtain that

$$\begin{aligned} \sum_{n=0}^{\infty} (l+n)! e^{2p(l+n)} |g_{l+n}|^2 &\leq \|K_{l,m}\|^2 \sum_{n=0}^{\infty} (n+m)! \frac{(l+n)!}{n!} \frac{(n+m)!}{n!} e^{2p(l+n)} |f_{n+m}|^2 \\ &\leq \|K_{l,m}\|^2 e^{2(pl-(p+q)m)} C_{l,m;q} \|\phi\|_{p+q}^2, \end{aligned} \quad (3.2)$$

where  $\|K_{l,m}\|$  is the operator norm and  $C_{l,m;q}$  is given as in (2.5). Therefore, we define an linear operator  $\Xi_{l,m}(K_{l,m})$  on  $\mathcal{G}^*$  by

$$\Xi_{l,m}(K_{l,m})\Phi = (g_{l+n})_{n=0}^{\infty}, \quad \Phi = (f_n)_{n=0}^{\infty} \in \mathcal{G}^*, \quad (3.3)$$

where  $g_{l+n}$  is given as in (3.1). Then for any  $p \in \mathbb{R}$  and  $q > 0$  it holds that

$$\|\Xi_{l,m}(K_{l,m})\Phi\|_p \leq \|K_{l,m}\| e^{(pl-(p+q)m)} \sqrt{C_{l,m;q}} \|\Phi\|_{p+q}, \quad \Phi \in \mathcal{G}^*, \quad (3.4)$$

which implies that  $\Xi_{l,m}(K_{l,m}) \in \mathcal{L}(\mathcal{G}_{p+q}, \mathcal{G}_p)$ . The operator  $\Xi_{l,m}(K_{l,m})$  is called the *integral kernel operator* with kernel  $K_{l,m}$  (see [13, 9, 22, 30]).

Now the following theorem is obvious.

**Theorem 3.1** ([11]) *Let  $l, m$  be non-negative integers and let  $K_{l,m} \in \mathcal{L}(H^{\otimes m}, H^{\otimes l})$ . Then it holds that*

$$\Xi_{l,m}(K_{l,m}) \in \mathcal{L}(\mathcal{G}, \mathcal{G}) \cap \mathcal{L}(\mathcal{G}^*, \mathcal{G}^*).$$

Let  $\eta \in H$  and let  $K_\eta \in \mathcal{L}(H, \mathbb{C})$  be defined by  $K_\eta(f) = \langle \eta, f \rangle$  for any  $f \in H$ . For simple notation, we identify  $\eta = K_\eta = K_\eta^*$ , where  $K_\eta^*$  is the adjoint operator of  $K_\eta$  with respect to the canonical bilinear form  $\langle \cdot, \cdot \rangle$ , i.e.,  $K_\eta^*(a) = a\eta$  for all  $a \in \mathbb{C}$ . Then the *annihilation operator*  $a(\eta)$  and the *creation operator*  $a^*(\eta)$  associated with  $\eta$  are defined by  $a(\eta) = \Xi_{0,1}(\eta)$  and  $a^*(\eta) = \Xi_{1,0}(\eta)$ , respectively, and then from Theorem 3.1, it holds that

$$a(\eta), a^*(\eta) \in \mathcal{L}(\mathcal{G}, \mathcal{G}) \cap \mathcal{L}(\mathcal{G}^*, \mathcal{G}^*).$$

It is straightforward to verify the canonical commutation relation:

$$[a(\xi), a(\eta)] = 0, \quad [a^*(\xi), a^*(\eta)] = 0, \quad [a(\xi), a^*(\eta)] = \int_{\mathbb{R}_+} \xi(t)\eta(t)dt = \langle \xi, \eta \rangle \quad (3.5)$$

for  $\xi, \eta \in H$ .

The exponential vector  $\phi_\xi$  associated with  $\xi \in H$  is defined by  $\phi_\xi = (\xi^{\otimes n}/n!)_{n=0}^{\infty}$ . Then  $\{\phi_\xi; \xi \in H\}$  spans a dense subspace of  $\mathcal{G}$ .

**Proposition 3.2** ([14, 5]) *Let  $\zeta \in H$  be given. Then it holds that*

$$a(\zeta)(\Phi\psi) = (a(\zeta)\Phi)\psi + \Phi(a(\zeta)\psi), \quad \Phi \in \mathcal{G}^*, \quad \psi \in \mathcal{G}, \quad (3.6)$$

$$a(\zeta)(\Phi \diamond \Psi) = (a(\zeta)\Phi) \diamond \Psi + \Phi \diamond (a(\zeta)\Psi), \quad \Phi, \Psi \in \mathcal{G}^*. \quad (3.7)$$

PROOF. (i) For any  $\xi, \eta \in H$ , we obtain that

$$\begin{aligned} a(\zeta)(\phi_\xi\phi_\eta) &= a(\zeta)(\phi_{\xi+\eta})e^{\langle \zeta, \eta \rangle} = \langle \zeta, \xi + \eta \rangle \phi_{\xi+\eta}e^{\langle \zeta, \eta \rangle} = \langle \zeta, \xi + \eta \rangle \phi_\xi\phi_\eta \\ &= (a(\zeta)\phi_\xi)\phi_\eta + \phi_\xi(a(\zeta)\phi_\eta). \end{aligned}$$

Therefore, by the continuity property  $a(\zeta) \in \mathcal{L}(\mathcal{G}, \mathcal{G}) \cap \mathcal{L}(\mathcal{G}^*, \mathcal{G}^*)$  and the fact that exponential vectors span a dense subspace of  $\mathcal{G}$  and  $\mathcal{G}^*$ , we complete the proof.

(ii) The proof is similar to the proof of (i). In fact, we obtain that

$$\begin{aligned} a(\zeta)(\phi_\xi \diamond \phi_\eta) &= a(\zeta)(\phi_{\xi+\eta}) = \langle \zeta, \xi + \eta \rangle \phi_{\xi+\eta} = \langle \zeta, \xi + \eta \rangle \phi_\xi \diamond \phi_\eta \\ &= (a(\zeta)\phi_\xi) \diamond \phi_\eta + \phi_\xi \diamond (a(\zeta)\phi_\eta). \end{aligned}$$

□

Let  $K \in \mathcal{L}(H, H)$ . Then from Theorem 3.1, it holds that

$$\Lambda(K) := \Xi_{1,1}(K) \in \mathcal{L}(\mathcal{G}, \mathcal{G}) \cap \mathcal{L}(\mathcal{G}^*, \mathcal{G}^*).$$

The operator  $\Lambda(K)$  is called the *conservation operator*. for any  $p \in \mathbb{R}$  and  $q > 0$ , from (3.4) we obtain that

$$\|\Lambda(K)\Phi\|_p \leq e^{-q} \sqrt{C_{1,1;q}} \|K\| \|\Phi\|_{p+q}, \quad \Phi \in \mathcal{G}^*. \quad (3.8)$$

## 3.2 Multiplication Operators

**Theorem 3.3** For any  $\Phi \in \mathcal{G}^*$  and  $\phi, \psi \in \mathcal{G}$ , it holds that

$$\langle\langle \Phi\phi, \psi \rangle\rangle = \langle\langle \Phi, \phi\psi \rangle\rangle. \quad (3.9)$$

PROOF. For given  $\Phi = (F_n) \in \mathcal{G}^*$  and any  $\xi, \eta \in H$ , from (2.1) we obtain that

$$\Phi\phi_\xi = \left( \sum_{l+m=n} \sum_{k=0}^{\infty} \binom{l+k}{k} (F_{l+k} \widehat{\otimes}_k \xi^{\otimes k}) \otimes \frac{\xi^{\otimes m}}{m!} \right)_{n=0}^{\infty}$$

and

$$\begin{aligned} \langle\langle \Phi\phi_\xi, \phi_\eta \rangle\rangle &= \sum_{n=0}^{\infty} \left\langle \sum_{m=0}^n \sum_{k=0}^{\infty} \binom{n-m+k}{k} (F_{n-m+k} \widehat{\otimes}_k \xi^{\otimes k}) \otimes \frac{\xi^{\otimes m}}{m!}, \eta^{\otimes n} \right\rangle \\ &= \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \left\langle \sum_{k=0}^{\infty} \binom{n+k}{k} (F_{n+k} \widehat{\otimes}_k \xi^{\otimes k}) \otimes \frac{\xi^{\otimes m}}{m!}, \eta^{\otimes(n+m)} \right\rangle \\ &= e^{\langle \xi, \eta \rangle} \sum_{l=0}^{\infty} \left\langle F_l, \sum_{n+k=l} \binom{n+k}{k} \eta^{\otimes n} \widehat{\otimes}_k \xi^{\otimes k} \right\rangle \\ &= e^{\langle \xi, \eta \rangle} \langle\langle \Phi, \phi_{\xi+\eta} \rangle\rangle \\ &= \langle\langle \Phi, \phi_\eta \phi_\xi \rangle\rangle. \end{aligned}$$

Since the exponential vectors span a dense subspace of  $\mathcal{G}$ , by the continuity of the Wiener product (see Theorems 2.3 and 2.4), the proof is immediate. □

Let  $\Phi \in \mathcal{G}^*$  be given. Then we consider the Wiener multiplication operator  $M_\Phi : \mathcal{G} \rightarrow \mathcal{G}^*$  and then from (3.9),  $\phi, \psi \in \mathcal{G}$ , we obtain that it holds that

$$\langle\langle M_\Phi\phi, \psi \rangle\rangle = \langle\langle \Phi\phi, \psi \rangle\rangle = \langle\langle \Phi, \phi\psi \rangle\rangle.$$

**Theorem 3.4 ([30])** For each  $\zeta \in H$ ,  $X_\zeta = (0, \zeta, 0, \dots) \in \mathcal{G}$  as a Wiener multiplication operator is represented as the sum of  $a(\zeta)$  and  $a^*(\zeta)$ , i.e.,

$$X_\zeta = a(\zeta) + a^*(\zeta), \quad (3.10)$$

which is called the quantum decomposition of  $X_\zeta$ .

PROOF. Since  $X_\zeta = (0, \zeta, 0, \dots) \in \mathcal{G}$ , from (3.9) we obtain that

$$\begin{aligned} \langle X_\zeta \phi_\xi, \phi_\eta \rangle &= \langle X_\zeta, \phi_\xi \phi_\eta \rangle = \langle X_\zeta, \phi_{\xi+\eta} \rangle e^{\langle \xi, \eta \rangle} = \langle \zeta, \xi + \eta \rangle e^{\langle \xi, \eta \rangle} \\ &= \langle (a(\zeta) + a^*(\zeta)) \phi_\xi, \phi_\eta \rangle, \end{aligned}$$

which gives the proof.  $\square$

For each  $t \geq 0$ , put  $B_t = X_{\mathbf{1}_{[0,t]}}$ . Then  $\{B_t\}_{t \geq 0}$  becomes a Brownian motion which is called a *realization of Brownian motion* and so from Theorem 3.4 we have the following quantum decomposition of Brownian motion:

$$B_t = a(\mathbf{1}_{[0,t]}) + a^*(\mathbf{1}_{[0,t]}), \quad t \geq 0. \quad (3.11)$$

**Remark 3.5** The operators  $\mathcal{L}(\mathcal{G}, \mathcal{G}^*)$  on admissible Gaussian functionals play an essential role in the study of quantum martingales and integral representations [11, 14, 16, 17].

## 4 Quantum White Noise Derivatives

In this section, we briefly review some basic properties of quantum white noise derivatives [14, 15, 16, 17, 18, 20].

### 4.1 Annihilation and Creation Derivatives

For any admissible operator  $\Xi \in \mathcal{L}(\mathcal{G}, \mathcal{G}^*)$  and  $\zeta \in H$ , from Theorem 3.1 the commutators

$$[a(\zeta), \Xi] = a(\zeta)\Xi - \Xi a(\zeta), \quad -[a^*(\zeta), \Xi] = \Xi a^*(\zeta) - a^*(\zeta)\Xi$$

are well defined as compositions of admissible operators, i.e., belong to  $\mathcal{L}(\mathcal{G}, \mathcal{G}^*)$ . We define

$$D_\zeta^+ \Xi = [a(\zeta), \Xi], \quad D_\zeta^- \Xi = -[a^*(\zeta), \Xi].$$

These are called the *creation derivative* and *annihilation derivative* of  $\Xi$ , respectively. Both together are referred to as the *quantum white noise derivatives* (*qwn-derivatives* for brevity) of  $\Xi$ . By the definitions, it is obvious that

$$\begin{aligned} (D_\zeta^+ \Xi)^* &= ([a(\zeta), \Xi]^*)^* = (a(\zeta)\Xi^* - \Xi^* a(\zeta))^* = \Xi a^*(\zeta) - a^*(\zeta)\Xi \\ &= D_\zeta^- \Xi. \end{aligned} \quad (4.1)$$

For each admissible operator  $\Xi \in \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q)$ , we operator norm of  $\Xi$  is denoted by  $\|\Xi\|_{p,q}$ .

**Theorem 4.1 ([14])** Let  $\zeta \in H$  be given. Then  $D_\zeta^\pm$  are continuous linear operators from  $\mathcal{L}(\mathcal{G}, \mathcal{G}^*)$  itself.



PROOF. Suppose that  $\Xi \in \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q)$ . Then for any  $r > 0$ , by applying (3.4), we obtain that

$$\begin{aligned} \|D_\zeta^+ \Xi\|_{p-r; q+r} &= \|[a(\zeta), \Xi]\|_{p-r; q+r} = \|a(\zeta)\Xi - \Xi a(\zeta)\|_{p-r; q+r} \\ &\leq \|a(\zeta)\|_{q; q+r} \|\Xi\|_{p; q} + \|\Xi\|_{p; q} \|a(\zeta)\|_{p-r; p}, \end{aligned}$$

which implies that  $D_\zeta^+$  is a continuous linear operator on  $\mathcal{L}(\mathcal{G}, \mathcal{G}^*)$ . Similarly, we see that  $D_\zeta^-$  is a continuous linear operator on  $\mathcal{L}(\mathcal{G}, \mathcal{G}^*)$ .  $\square$

**Proposition 4.2** *For each  $\zeta \in H$  and  $\Phi \in \mathcal{G}^*$ , it holds that*

$$(D_\zeta^+ M_\Phi) \phi_0 = (D_\zeta^- M_\Phi) \phi_0 = a(\zeta)\Phi.$$

PROOF. We obtain that

$$\begin{aligned} (D_\zeta^+ M_\Phi) \phi_0 &= (a(\zeta)M_\Phi - M_\Phi a(\zeta)) \phi_0 = a(\zeta)\Phi, \\ (D_\zeta^- M_\Phi) \phi_0 &= (M_\Phi a^*(\zeta) - a^*(\zeta)M_\Phi) \phi_0 = \Phi X_\zeta - a^*(\zeta)\Phi = a(\zeta)\Phi, \end{aligned}$$

where we used the quantum decomposition as  $\Phi X_\zeta = X_\zeta \Phi = (a(\zeta) + a^*(\zeta))\Phi$ .  $\square$

## 4.2 Pointwise QWN-Derivatives

Let  $\phi = (f_n) \in \mathcal{G}$  and  $t \in \mathbb{R}_+$  be given. Suppose that  $f_n = 0$  except for a finite number of  $n$ . We define

$$D_t \phi := (n f_n(t, \cdot))_{n=1}^\infty,$$

where  $f_n(t, \cdot) \in H^{\widehat{\otimes}(n-1)}$ , and then  $D$  is called the *classical stochastic gradient*. The classical stochastic gradient is denoted by  $\nabla$  in some literatures see [8, 16, 18, 22, 29]. We now extend the domain of  $D$  to the space  $\mathcal{G}^*$ .

**Lemma 4.3** ([16]) *For any  $p \in \mathbb{R}$  and  $r > 0$  we have*

$$\|D\phi\|_{L^2(\mathbb{R}_+, \mathcal{G}_{-p-r})}^2 = \int_{\mathbb{R}_+} \|D\phi(t)\|_{-p-r}^2 dt \leq K(p, r) \|\phi\|_{-p}^2, \quad \phi \in \mathcal{G}, \quad (4.2)$$

where  $K(p, r) = \sup_n (n+1)e^{2p-2rn} < \infty$ . In particular, the classical stochastic gradient

$$D : \mathcal{G}_{-p} \rightarrow L^2(\mathbb{R}_+, \mathcal{G}_{-p-r}) \cong L^2(\mathbb{R}_+) \otimes \mathcal{G}_{-p-r} \quad (4.3)$$

is a continuous linear map.

PROOF. For each  $\phi = (f_n)_{n=0}^\infty \in \mathcal{G}$  consisting of continuous functions  $f_n$  on  $\mathbb{R}_+^n$ , we have  $D\phi(t) = ((n+1)f_{n+1}(t, \cdot))_{n=0}^\infty$ , where the right-hand side has a pointwise meaning. Then we obtain that

$$\begin{aligned} \int_{\mathbb{R}_+} \|D\phi(t)\|_{-p-r}^2 dt &= \sum_{n=0}^\infty n! e^{-2(p+r)n} \int_{\mathbb{R}_+} |(n+1)f_{n+1}(t, \cdot)|_0^2 dt \\ &= \sum_{n=0}^\infty (n+1)e^{2p-2rn} \times (n+1)! e^{-2p(n+1)} |f_{n+1}|_0^2 \\ &\leq K(p, r) \|\phi\|_{-p}^2, \end{aligned}$$

which implies the proof of (4.2). □

Put

$$L^2(\mathbb{R}, \mathcal{G}) := \text{proj} \lim_{p \rightarrow \infty} L^2(\mathbb{R}, \mathcal{G}_p) \cong \text{proj} \lim_{p \rightarrow \infty} L^2(\mathbb{R}) \otimes \mathcal{G}_p,$$

$$L^2(\mathbb{R}_+, \mathcal{G}^*) := \text{ind} \lim_{p \rightarrow \infty} L^2(\mathbb{R}_+, \mathcal{G}_{-p}) \cong \text{ind} \lim_{p \rightarrow \infty} L^2(\mathbb{R}) \otimes \mathcal{G}_{-p}.$$

Then by Lemma 4.3, the classical stochastic gradient  $D$  is a continuous linear map from  $\mathcal{G}$  into  $L^2(\mathbb{R}_+, \mathcal{G})$  and from  $\mathcal{G}^*$  into  $L^2(\mathbb{R}_+, \mathcal{G}^*)$ .

We see from (4.3) that  $D\Phi(t)$  has a meaning as  $\mathcal{G}_{-p-r}$ -valued  $L^2$ -function in  $t \in \mathbb{R}_+$ . Given  $\zeta \in L^2(\mathbb{R}_+)$ , the linear map  $\mathcal{G}_{p+r} \ni \psi \mapsto \langle\langle D\Phi, \zeta \otimes \psi \rangle\rangle$  is continuous. Therefore there exists a unique  $\Psi \in \mathcal{G}_{-p-r}$  such that

$$\langle\langle D\Phi, \zeta \otimes \psi \rangle\rangle = \langle\langle \Psi, \psi \rangle\rangle, \quad \psi \in \mathcal{G}_{p+r}.$$

It is reasonable to write

$$\Psi = \int_{\mathbb{R}_+} \zeta(t) D\Phi(t) dt.$$

As is easily seen, the Schwarz inequality holds:

$$\left\| \int_{\mathbb{R}_+} \zeta(t) D\Phi(t) dt \right\|_{-p-r} \leq |\zeta|_0 \| D\Phi \|_{L^2(\mathbb{R}_+, \mathcal{G}_{-p-r})}, \tag{4.4}$$

which implies that the map

$$\mathcal{G}_{-p} \ni \Phi \mapsto \int_{\mathbb{R}_+} \zeta(t) D\Phi(t) dt \in \mathcal{G}_{-(p+r)}$$

is continuous. On the other hand, for any  $\xi \in H$ , we obtain that

$$\begin{aligned} \int_{\mathbb{R}_+} \zeta(t) D\phi_\xi(t) dt &= \int_{\mathbb{R}_+} \zeta(t) a_t \phi_\xi dt = \int_{\mathbb{R}_+} \zeta(t) \xi(t) \phi_\xi dt = \langle \zeta, \xi \rangle \phi_\xi \\ &= a(\zeta) \phi_\xi. \end{aligned}$$

Therefore, we obtain that

$$\int_{\mathbb{R}_+} \zeta(t) D\Phi(t) dt = a(\zeta) \Phi, \quad \Phi \in \mathcal{G}^*, \tag{4.5}$$

see [16].

**Remark 4.4** The space  $\mathcal{G}^*$  as a domain of the classical gradient  $D$  appeared in Aase–Øksendal–Privault–Ubøe [1]. For a standard domain see e.g., Kuo [22], Malliavin [27], Nualart [29].

Let  $\Xi \in \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q)$  for some  $p, q \in \mathbb{R}$ . Then for any  $r > 0$  and  $\phi \in \mathcal{G}$ , from (4.2) we obtain that

$$\begin{aligned} \int_{\mathbb{R}_+} \| \Xi D_t \phi \|_q^2 dt &\leq \int_{\mathbb{R}_+} \| \Xi \|_{p;q}^2 \| D_t \phi \|_p^2 dt \\ &\leq K(-p, r) \| \Xi \|_{p;q}^2 \| \phi \|_{p+r}^2, \end{aligned}$$

which implies that

$$\int_{\mathbb{R}_+} \|\Xi D_t\|_{p+r;q}^2 dt \leq K(-p, r) \|\Xi\|_{p;q}^2,$$

and so the map

$$\mathcal{L}(\mathcal{G}_p, \mathcal{G}_q) \ni \Xi \longmapsto \Xi D \in L^2(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_{p+r}, \mathcal{G}_q))$$

is continuous. Similarly, we obtain that

$$\int_{\mathbb{R}_+} \|D_t \Xi \phi\|_{q-r}^2 dt \leq K(-q, r) \|\Xi \phi\|_q^2 \leq K(-q, r) \|\Xi\|_{p;q}^2 \|\phi\|_p^2,$$

which implies that

$$\int_{\mathbb{R}_+} \|\Xi D_t\|_{p+r;q}^2 dt \leq K(-p, r) \|\Xi\|_{p;q}^2,$$

and so the map

$$\mathcal{L}(\mathcal{G}_p, \mathcal{G}_q) \ni \Xi \longmapsto D \Xi \in L^2(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_{q-r}))$$

is continuous. Therefore, the *pointwise creation derivative*  $D_t^+$  is defined by

$$D_t^+ \Xi = D_t \Xi - \Xi D_t, \quad \Xi \in \mathcal{L}(\mathcal{G}, \mathcal{G}^*)$$

and  $D_t^+ \Xi$  is an  $\mathcal{L}(\mathcal{G}, \mathcal{G}^*)$ -valued  $L^2$ -function in  $t \in \mathbb{R}_+$ . Motivated by (4.1), the *pointwise annihilation derivative*  $D_t^-$  is defined by

$$D_t^- \Xi = (D_t^+ \Xi^*)^*, \quad \Xi \in \mathcal{L}(\mathcal{G}, \mathcal{G}^*),$$

see [16, 18]. In fact, for given  $\Xi \in \mathcal{L}(\mathcal{G}, \mathcal{G}^*)$  and  $\zeta \in H$ , from (4.5) we obtain that

$$\int_{\mathbb{R}_+} \zeta(t) D_t^+ \Xi dt = D_\zeta^+ \Xi$$

and

$$D_\zeta^- \Xi = (D_\zeta^+ \Xi^*)^* = \int_{\mathbb{R}_+} \zeta(t) (D_t^+ \Xi^*)^* dt.$$

**Proposition 4.5** *For each  $t \geq 0$  and  $\Phi \in \mathcal{G}^*$ , it holds that*

$$(D_t^+ M_\Phi) \phi_0 = (D_t^- M_\Phi) \phi_0 = D_t \Phi.$$

**PROOF.** We obtain that

$$\begin{aligned} (D_t^+ M_\Phi) \phi_0 &= (D_t M_\Phi - M_\Phi D_t) \phi_0 = D_t \Phi, \\ (D_t^- M_\Phi) \phi_0 &= (D_t^+ M_\Phi^*)^* \phi_0 = (D_t^+ M_\Phi)^* \phi_0 = (M_{D_t \Phi})^* \phi_0 = D_t \Phi, \end{aligned}$$

which gives the proof. □

## 5 Anticipating Quantum Stochastic Integrals

For each  $t \geq 0$ , let  $\mathcal{F}_t$  be the  $\sigma$ -field generated by  $\{B_s; 0 \leq s \leq t\}$ . A one-parameter family  $\Phi = \{\Phi_t\}_{t \geq 0} \subset \mathcal{G}^*$  is called a *generalized stochastic process* [4, 11, 31] if there exists a  $p \geq 0$  (independent of  $t \geq 0$ ) such that  $\Phi_t \in \mathcal{G}_{-p}$  for all  $t \geq 0$  and the map  $t \mapsto \Phi_t \in \mathcal{G}_{-p}$  is Borel measurable on  $\mathbb{R}_+$ . A generalized stochastic process  $\{\Phi_t = (F_{t;n})\}_{t \geq 0}$  is said to be *adapted* (w.r.t.  $\mathcal{F}_t$ ) if for all  $t \geq 0$  and  $n \geq 0$ ,  $\text{supp} F_{t;n} \subset [0, t]^n$ .

A one-parameter family  $\{\Xi_t\}_{t \in \mathbb{R}_+} \subset \mathcal{L}(\mathcal{G}, \mathcal{G}^*)$  is called a *quantum stochastic process*. Our approach covers a wide class of classical and quantum stochastic processes in the sense that  $\mathcal{G}^*$  and  $\mathcal{L}(\mathcal{G}, \mathcal{G}^*)$  involve distributions. As examples, for each  $t \geq 0$ , we put

$$A_t = a(\mathbf{1}_{[0,t]}), \quad A_t^* = a^*(\mathbf{1}_{[0,t]}), \quad \Lambda_t = \Xi_{1,1}(\mathbf{1}_{[0,t]}).$$

For the definition of  $\Lambda_t$ , the indicator function  $\mathbf{1}_{[0,t]}$  is considered as a multiplication operator on  $H$ , i.e.,  $\mathbf{1}_{[0,t]}(\xi) = \mathbf{1}_{[0,t]}\xi =: \xi_{[0,t]}$  for any  $\xi \in H$ . Then for each  $t \geq 0$ ,  $A_t, A_t^*, \Lambda_t \in \mathcal{L}(\mathcal{G}, \mathcal{G}) \cap \mathcal{L}(\mathcal{G}^*, \mathcal{G}^*)$ . The processes  $\{A_t\}_{t \geq 0}$ ,  $\{A_t^*\}_{t \geq 0}$  and  $\{\Lambda_t\}_{t \geq 0}$  are called the *annihilation*, *creation* and *conservation* (or *gauge*) *processes*, respectively.

### 5.1 Quantum Hitsuda–Skorohod Integrals

In this section, we study the Hitsuda–Skorohod type quantum stochastic integrals with their regular properties.

**Theorem 5.1** *Let  $p, q \in \mathbb{R}$  be given and  $\Xi \in L^2(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q))$  be a quantum stochastic process. Then there exists an admissible operator, denoted by  $\delta^-(\Xi)$ , in  $\mathcal{L}(\mathcal{G}_{p+r}, \mathcal{G}_q)$  for any  $r > 0$  such that*

$$\delta^-(\Xi)\phi = \int_{\mathbb{R}_+} \Xi(t)(D_t\phi) dt \tag{5.1}$$

for any  $\phi \in \mathcal{G}$ .

PROOF. For any  $\phi \in \mathcal{G}$  and  $r > 0$ , by applying (4.2), we obtain that

$$\begin{aligned} \left\| \int_{\mathbb{R}_+} \Xi(t)(D_t\phi) dt \right\|_q &\leq \int_{\mathbb{R}_+} \|\Xi(t)\|_{p;q} \|D_t\phi\|_p dt \\ &\leq \left( \int_{\mathbb{R}_+} \|\Xi(t)\|_{p;q}^2 dt \right)^{1/2} \left( \int_{\mathbb{R}_+} \|D_t\phi\|_p^2 dt \right)^{1/2} \\ &\leq \sqrt{K(-p, r)} \left( \int_{\mathbb{R}_+} \|\Xi(t)\|_{p;q}^2 dt \right)^{1/2} \|\phi\|_{p+r}^2, \end{aligned}$$

which implies that the linear operator

$$\mathcal{G}_{p+r} \ni \phi \longmapsto \int_{\mathbb{R}_+} \Xi(t)(D_t\phi) dt \in \mathcal{G}_q$$

is continuous. □

For given  $\Xi \in L^2(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q))$ , the admissible operator  $\delta^-(\Xi)$  satisfying (5.1) is called the *annihilation integral* of  $\Xi$ , see [3, 24, 16, 18].

**Remark 5.2** Let  $p, q \in \mathbb{R}$  be given and  $\Xi \in L^2(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q))$  be a quantum stochastic process. Then for any  $\xi \in H$ , we obtain that

$$\delta^-(\Xi)\phi_\xi = \int_{\mathbb{R}_+} \Xi(t)(D_t\phi_\xi) dt = \int_{\mathbb{R}_+} \xi(t)\Xi(t)\phi_\xi dt = \left( \int_{\mathbb{R}_+} \Xi(t) dA_t \right) \phi_\xi,$$

which implies that

$$\delta^-(\Xi) = \int_{\mathbb{R}_+} \Xi(t) dA_t$$

on a certain domain. Furthermore, if  $\Xi$  is adapted, then  $\delta^-(\Xi)$  coincides with the annihilation integral of Hudson-Parthasarathy. For the definition of the adaptedness of quantum stochastic processes, we refer to [11]. Also, for more study on quantum Hitsuda-Skorohod integrals, we refer to [3, 24, 18].

As for a criterion for  $\delta^-(\Xi)$  being a bounded operator on  $\Gamma(H)$ , we have the following corollary. A similar result can be found in [18].

**Corollary 5.3** For any  $r > 0$  and  $\Xi \in L^2(\mathbb{R}, \mathcal{L}(\mathcal{G}_{-r}, \Gamma(H)))$ , the annihilation integral  $\delta^-(\Xi)$  is a bounded operator on  $\Gamma(H)$ .

PROOF. The proof is immediate from Theorem 5.1.  $\square$

**Theorem 5.4** Let  $p, q \in \mathbb{R}$  be given and  $\Xi \in L^2(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q))$  be a quantum stochastic process. Then there exists an admissible operator, denoted by  $\delta^+(\Xi)$ , in  $\mathcal{L}(\mathcal{G}_p, \mathcal{G}_{q-r})$  for any  $r > 0$  such that

$$\langle\langle \delta^+(\Xi)\phi, \psi \rangle\rangle = \int_{\mathbb{R}_+} \langle\langle \Xi(t)\phi, D_t\psi \rangle\rangle dt \quad (5.2)$$

for  $\phi, \psi \in \mathcal{G}$ .

PROOF. For any  $\phi, \psi \in \mathcal{G}$  and  $r > 0$ , by applying (4.2), we obtain that

$$\begin{aligned} \left| \int_{\mathbb{R}_+} \langle\langle \Xi(t)\phi, D_t\psi \rangle\rangle dt \right| &\leq \int_{\mathbb{R}_+} \|\Xi(t)\phi\|_q \|D_t\psi\|_{-q} dt \\ &\leq \left( \int_{\mathbb{R}_+} \|\Xi(t)\|_{p;q}^2 dt \right)^{1/2} \left( \int_{\mathbb{R}_+} \|D_t\psi\|_{-q}^2 dt \right)^{1/2} \|\phi\|_p \\ &\leq \sqrt{K(q, r)} \left( \int_{\mathbb{R}_+} \|\Xi(t)\|_{p;q}^2 dt \right)^{1/2} \|\phi\|_p \|\psi\|_{-q+r}, \end{aligned}$$

which implies that the bilinear form

$$\mathcal{G}_p \times \mathcal{G}_{-q+r} \ni (\phi, \psi) \longmapsto \int_{\mathbb{R}_+} \langle\langle \Xi(t)\phi, D_t\psi \rangle\rangle dt \in \mathbb{C}$$

is continuous. Therefore, there exists a unique admissible operator  $\delta^+(\Xi) \in \mathcal{L}(\mathcal{G}_p, \mathcal{G}_{p-r})$  such that (5.2) holds.  $\square$

For given  $\Xi \in L^2(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q))$ , the admissible operator  $\delta^+(\Xi)$  satisfying (5.2) is called the *creation integral* of  $\Xi$ , see [3, 24, 16, 18].

As for a criterion for  $\delta^+(\Xi)$  being a bounded operator on  $\Gamma(H)$ , we have the following corollary. A similar result can be found in [18].

**Corollary 5.5** For any  $r > 0$  and  $\Xi \in L^2(\mathbb{R}, \mathcal{L}(\Gamma(H), \mathcal{G}_r))$ , the creation integral  $\delta^+(\Xi)$  is a bounded operator on  $\Gamma(H)$ .

PROOF. The proof is immediate from Theorem 5.4.  $\square$

**Remark 5.6** The classical Hitsuda–Skorohod integral  $\delta$  is defined as the adjoint map of the classical stochastic gradient  $D$  (see [8, 18, 22, 29]), i.e., for given  $\Psi \in L^2(\mathbb{R}_+, \mathcal{G}^*)$ , the classical Hitsuda–Skorohod integral  $\delta(\Psi) \in \mathcal{G}^*$  of  $\Psi$  is defined by

$$\langle\langle \delta(\Psi), \phi \rangle\rangle = \int_{\mathbb{R}_+} \langle\langle \Psi(t), D_t \phi \rangle\rangle dt, \quad \phi \in \mathcal{G}. \quad (5.3)$$

Therefore, by denoting  $(\Xi\phi)(t) = \Xi(t)\phi$ , from (5.2) we have

$$\delta^+(\Xi)\phi = \delta(\Xi\phi), \quad \phi \in \mathcal{G}, \quad (5.4)$$

see [2, 24, 18].

The creation and annihilation integrals are related directly. The following corollary gives a relation between creation and annihilation integrals.

**Corollary 5.7** ([18]) Let  $p, q \in \mathbb{R}$  be given and  $\Xi \in L^2(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q))$  be a quantum stochastic process. Then it holds that

$$(\delta^-(\Xi))^* = \delta^+(\Xi^*). \quad (5.5)$$

PROOF. For any  $\phi, \psi \in \mathcal{G}$ , we obtain that

$$\begin{aligned} \langle\langle \delta^-(\Xi)\phi, \psi \rangle\rangle &= \int_{\mathbb{R}_+} \langle\langle \Xi(t)(D_t \phi), \psi \rangle\rangle dt = \int_{\mathbb{R}_+} \langle\langle \Xi^*(t)\psi, (D_t \phi) \rangle\rangle dt \\ &= \langle\langle \delta^+(\Xi^*)\psi, \phi \rangle\rangle, \end{aligned}$$

which proves (5.5).  $\square$

**Theorem 5.8** Let  $p, q \in \mathbb{R}$  be given and  $\Xi \in L^\infty(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q))$  be a quantum stochastic process. Then there exists an admissible operator, denoted by  $\delta^0(\Xi)$ , in  $\mathcal{L}(\mathcal{G}_{p+r}, \mathcal{G}_{q-r})$  for any  $r > 0$  such that

$$\langle\langle \delta^0(\Xi)\phi, \psi \rangle\rangle = \int_{\mathbb{R}_+} \langle\langle \Xi(t)D_t \phi, D_t \psi \rangle\rangle dt \quad (5.6)$$

for  $\phi, \psi \in \mathcal{G}$ .

PROOF. The proof is a simple modification of the proofs of Theorems 5.1 and 5.4.  $\square$

For given  $\Xi \in L^2(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q))$ , the admissible operator  $\delta^0(\Xi)$  satisfying (5.6) is called the *conservation integral* of  $\Xi$ , see [3, 24, 16, 18].

As for a criterion for  $\delta^0(\Xi)$  being a bounded operator on  $\Gamma(H)$ , we have the following corollary. A similar result can be found in [18].

**Corollary 5.9** For any  $r > 0$  and  $\Xi \in L^2(\mathbb{R}, \mathcal{L}(\mathcal{G}_{-r}, \mathcal{G}_r))$ , the conservation integral  $\delta^0(\Xi)$  is a bounded operator on  $\Gamma(H)$ .

PROOF. The proof is immediate from Theorem 5.8.  $\square$

## 5.2 Extensions of Anticipating Quantum Stochastic Integrals

In this section, motivated by the results in [23], we discuss extensions of the quantum Hitsuda-Skorohod integrals studied in Section 5.1. Based on the quantum white noise calculus [12], we have the following integral representations:

$$A_t = \int_0^t a_s ds, \quad A_t^* = \int_0^t a_s^* ds, \quad \Lambda_t = \int_0^t a_s^* a_s ds,$$

where  $a_t$  and  $a_t^*$  are the pointwisely defined annihilation and creation operators. On the other hand, the pointwisely defined annihilation operator  $a_t$  and the stochastic gradient  $D_t$  coincide on a certain domain. Hence, the following informal computations gives motivations for extensions of the quantum Hitsuda-Skorohod integrals: for a given quantum stochastic process  $\{\Xi_t\}_{t \geq 0} \subset \mathcal{L}(\mathcal{G}, \mathcal{G}^*)$  of enough regular operators  $\Xi_t$ , we may write as

$$\begin{aligned} \int_0^t \Xi_s dA_s &= \int_0^t \Xi_s D_s ds = \delta^-(\mathbf{1}_{[0,t]}\Xi), \\ \int_0^t (dA_s) \Xi_s &= \int_0^t D_s \Xi_s ds = \int_0^t \Xi_s D_s ds + \int_0^t D_s^+ \Xi_s ds = \delta^-(\mathbf{1}_{[0,t]}\Xi) + \int_0^t D_s^+ \Xi_s ds, \\ \int_0^t (dA_s^*) \Xi_s &= \int_0^t D_s^* \Xi_s ds = \delta^+(\mathbf{1}_{[0,t]}\Xi), \\ \int_0^t \Xi_s dA_s^* &= \int_0^t \Xi_s D_s^* ds = \delta^+(\mathbf{1}_{[0,t]}\Xi) + \int_0^t D_s^- \Xi_s ds, \\ \int_0^t \Xi_s d\Lambda_s &= \int_0^t \Xi_s D_s^* D_s ds = \delta^0(\mathbf{1}_{[0,t]}\Xi) + \delta^-(\mathbf{1}_{[0,t]}D^- \Xi), \\ \int_0^t (d\Lambda_s) \Xi_s &= \int_0^t D_s^* D_s \Xi_s ds = \delta^0(\mathbf{1}_{[0,t]}\Xi) + \delta^+(\mathbf{1}_{[0,t]}D^+ \Xi). \end{aligned} \quad (5.7)$$

However,  $D_t^\pm \Xi_t$  has no meaning directly. For example, we consider the annihilation process  $A_t = a(\mathbf{1}_{[0,t]})$  and then

$$D_t^- A_s = \mathbf{1}_{[0,s]}(t).$$

But the annihilation process  $A_t$  can be defined as  $a(\mathbf{1}_{[0,t]})$  and then we would have  $D_t^- A_s = \mathbf{1}_{[0,s]}(t)$ . Therefore,  $D_t^- A_t$  cannot be defined in a unique way [23]. From the above example, if we deal with quantum stochastic processes, then it is natural to consider two kinds of pointwisely defined annihilation derivative,  $D_{t+}^\pm$  and  $D_{t-}^\pm$ . Let  $\{\Xi_t\}_{s \geq 0} \subset \mathcal{L}(\mathcal{G}, \mathcal{G}^*)$  be a quantum stochastic process. We define

$$D_{t+}^\pm \Xi_t = \lim_{s \downarrow t} D_s^\pm \Xi_t, \quad D_{t-}^\pm \Xi_t = \lim_{s \uparrow t} D_s^\pm \Xi_t,$$

if the limits exist.

**Definition 5.10** Let  $\{\Xi_t\}_{t \geq 0} \subset \mathcal{L}(\mathcal{G}, \mathcal{G}^*)$  be a quantum stochastic process.

- (1) Suppose that  $\delta^+(\Xi)$  exists, and  $D_{t+}^- \Xi_t$  exists and it is integrable on  $\mathbb{R}_+$ . Then we define

$$\int_{\mathbb{R}_+} \Xi_t dA_{t+}^* = \delta^+(\Xi) + \int_{\mathbb{R}_+} D_{t+}^- \Xi_t dt. \quad (5.8)$$

- (2) Suppose that  $\delta^+(\Xi)$  exists, and  $D_{t-}^-\Xi_t$  exists and it is integrable on  $\mathbb{R}_+$ . Then we define

$$\int_{\mathbb{R}_+} \Xi_t dA_{t-}^* = \delta^+(\Xi) + \int_{\mathbb{R}_+} D_{t-}^-\Xi_t dt. \tag{5.9}$$

- (3) Suppose that  $\delta^+(\Xi)$  exists, and  $D_{t+}^-\Xi_t, D_{t-}^-\Xi_t$  exist as integrable functions on  $\mathbb{R}_+$ . Then we define

$$\langle \alpha \rangle \int_{\mathbb{R}_+} \Xi_t \circ dA_t^* = \delta^+(\Xi) + \alpha_1 \int_{\mathbb{R}_+} D_{t+}^-\Xi_t dt + \alpha_2 \int_{\mathbb{R}_+} D_{t-}^-\Xi_t dt \tag{5.10}$$

for  $\alpha = (\alpha_1, \alpha_2) \in \mathbb{R}^2$ , which is called the  $\langle \alpha \rangle$ -creation integral.

**Theorem 5.11** *Let  $p, q \in \mathbb{R}$  be given and let  $\Xi \in L^2(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q))$  be a quantum stochastic process.*

- (1) *Suppose that  $D_{t+}^-\Xi_t$  exists and  $D_{\cdot+}^-\Xi \in L^1(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_{q-r}))$  for some  $r > 0$ . Then the integral  $\int_{\mathbb{R}_+} \Xi_t dA_{t+}^*$  exists as an operator in  $\mathcal{L}(\mathcal{G}_p, \mathcal{G}_{q-r})$ .*
- (2) *Suppose that  $D_{t-}^-\Xi_t$  exists and  $D_{\cdot-}^-\Xi \in L^1(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_{q-r}))$  for some  $r > 0$ . Then the integral  $\int_{\mathbb{R}_+} \Xi_t dA_{t-}^*$  exists as an operator in  $\mathcal{L}(\mathcal{G}_p, \mathcal{G}_{q-r})$ .*

PROOF. (1) Since  $\Xi \in L^2(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_q))$ , by Theorem 5.4, the quantum Hitsuda-Skorohod creation integral  $\delta^+(\Xi)$  of  $\Xi$  exists as an admissible operator in  $\mathcal{L}(\mathcal{G}_p, \mathcal{G}_{q-s})$  for any  $s > 0$ . Also, since, by assumption,  $D_{t+}^-\Xi_t$  exists and  $D_{\cdot+}^-\Xi \in L^1(\mathbb{R}_+, \mathcal{L}(\mathcal{G}_p, \mathcal{G}_{q-r}))$  for some  $r > 0$ , for any  $\phi \in \mathcal{G}$  we obtain that

$$\left\| \int_{\mathbb{R}_+} D_{t+}^-\Xi_t \phi dt \right\|_{q-r} \leq \left( \int_{\mathbb{R}_+} \|D_{t+}^-\Xi_t\|_{p; q-r} dt \right) \|\phi\|_p,$$

which implies that the integral  $\int_{\mathbb{R}_+} D_{t+}^-\Xi_t dt$  exists as an admissible operator in  $\mathcal{L}(\mathcal{G}_p, \mathcal{G}_{q-r})$ . Finally, the integral  $\int_{\mathbb{R}_+} \Xi_t dA_{t+}^*$  exists as an operator in  $\mathcal{L}(\mathcal{G}_p, \mathcal{G}_{q-r})$ .

- (2) The proof is similar to the proof of (1). □

By similar arguments used in Definition 5.10, we can define the quantum stochastic integrals of types given as in (5.7) of which the study will be appear in some other papers.

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## References

[1] K. Aase, B. Øksendal, N. Privault and J. Ubøe: *White noise generalizations of the Clark–Haussmann–Ocone theorem with application to mathematical finance*, Finance Stochast. **4** (2000), 465–496.

[2] S. Attal and J. M. Lindsay: *Quantum stochastic calculus with maximal operator domains*, Ann. Probab. **32** (2004), 488–529.



- [3] V. P. Belavkin: *A quantum nonadapted Ito formula and stochastic analysis in Fock scale*, J. Funct. Anal. **102** (1991), 414–447.
- [4] F. E. Benth and J. Potthoff: *On the martingale property for generalized stochastic processes*, Stoch. Stoch. Rep. **58** (1996), 349–367.
- [5] D. M. Chung and T. S. Chung: *Wick derivations on white noise functionals*, J. Korean Math. Soc. **33** (1996), 993–1008.
- [6] D. M. Chung, U. C. Ji and N. Obata: *Quantum stochastic analysis via white noise operators in weighted Fock space*, Rev. Math. Phys. **14** (2002), 241–272.
- [7] M. Grothaus, Yu. G. Kondratiev and L. Streit: *Complex Gaussian analysis and the Bargmann–Segal space*, Methods of Funct. Anal. and Topology **3** (1997), 46–64.
- [8] T. Hida, H.-H. Kuo, J. Potthoff and L. Streit: “White Noise: An Infinite Dimensional Calculus,” Kluwer Academic Publishers, 1993.
- [9] T. Hida, N. Obata and K. Saitô: *Infinite dimensional rotations and Laplacians in terms of white noise calculus*, Nagoya Math. J. **128** (1992), 65–93.
- [10] R. L. Hudson and K. R. Parthasarathy: *Quantum Ito’s formula and stochastic evolutions*, Commun. Math. Phys. **93** (1984), 301–323.
- [11] U. C. Ji: *Stochastic integral representation theorem for quantum semimartingales*, J. Funct. Anal. **201** (2003), 1–29.
- [12] U. C. Ji and N. Obata: *Quantum white noise calculus*, in “Non-Commutativity, Infinite-Dimensionality and Probability at the Crossroads (N. Obata, T. Matsui and A. Hora, Eds.),” pp. 143–191, World Scientific, 2002.
- [13] U. C. Ji and N. Obata: *A role of Bargmann–Segal spaces in characterization and expansion of operators on Fock space*, J. Math. Soc. Japan **56** (2004), 311–338.
- [14] U. C. Ji and N. Obata: *Admissible white noise operators and their quantum white noise derivatives*, in “Infinite Dimensional Harmonic Analysis III (H. Heyer, T. Hiraï, T. Kawazoe, K. Saito, Eds.),” pp. 213–232, World Scientific, 2005.
- [15] U. C. Ji and N. Obata: *Generalized white noise operators fields and quantum white noise derivatives*, Seminaires et Congrès **16** (2007), 17–33.
- [16] U. C. Ji and N. Obata: *Annihilation-derivative, creation-derivative and representation of quantum martingales*, Commun. Math. Phys. **286** (2009), 751–775.
- [17] U. C. Ji and N. Obata: *Quantum stochastic integral representations of Fock space operators*, Stochastics **81** (2009), 367–384.
- [18] U. C. Ji and N. Obata: *Quantum stochastic gradients*, Interdiscip. Inform. Sci. **14** (2009), 345–359.
- [19] U. C. Ji and N. Obata: *Implementation problem for the canonical commutation relation in terms of quantum white noise derivatives*, J. Math. Phys. **51** (2010), 123507.

- [20] U. C. Ji and N. Obata: *An implementation problem for boson fields and quantum Girsanov transform*, J. Math. Phys. **57** (2016), 083502.
- [21] U. C. Ji and K. B. Sinha: *Quantum stochastic calculus associated with quadratic quantum noises*, J. Math. Phys. **57** (2016), 022702.
- [22] H.-H. Kuo: “White Noise Distribution Theory,” CRC Press, 1996.
- [23] H.-H. Kuo and A. Russek: *White noise approach to stochastic integration*, J. Multivariate Anal. **24** (1988), 218–236.
- [24] J. M. Lindsay: *Quantum and non-causal stochastic integral*, Probab. Theory Related Fields **97** (1993), 65–80.
- [25] J. M. Lindsay and H. Maassen: *An integral kernel approach to noise*, in “Quantum Probability and Applications III”, Eds. Accardi, L., von Waldenfels, W., Lecture Notes in Math. **1303**, pp. 192–208, Springer-Verlag, 1988.
- [26] J. M. Lindsay and K. R. Parthasarathy: *Cohomology of power sets with applications in quantum probability*, Commun. Math. Phys. **124** (1989), 337–364.
- [27] P. Malliavin: “Stochastic Analysis,” Springer-Verlag, 1997.
- [28] P.-A. Meyer: “Quantum Probability for Probabilists,” Lecture Notes in Math. Vol. **1538**, Springer-Verlag, 1993.
- [29] D. Nualart: “The Malliavin Calculus and Related Topics,” Springer-Verlag, New York, 1995.
- [30] N. Obata: “White Noise Calculus and Fock Space,” Lecture Notes in Math. Vol. **1577**, Springer, 1994.
- [31] N. Obata: *Generalized quantum stochastic processes on Fock space*, Publ. RIMS, Kyoto Univ. **31** (1995), 667–702.
- [32] K. R. Parthasarathy: “An Introduction to Quantum Stochastic Calculus,” Birkhäuser, 1992.
- [33] J. Potthoff and M. Timpel: *On a dual pair of spaces of smooth and generalized random variables*, Potential Anal. **4** (1995), 637–654.
- [34] A. S. Üstünel and M. Zakai: *The composition of Wiener functionals with non-absolutely continuous shifts*, Probab. Theory Related Fields **98** (1994), 163–184.

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