Yang-Mills equations and holomorphic structures of vector bundles

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- <u>l</u>. Main subject of Yang-Mills theory is to study variational problem associated to connections on a principal fibre bundle. Yang-Mills functional is defined over a set of connections on a given bundle in terms of square norm of curvature. Euler-Lagrange equation of this functional is called Yang-Mills equation. With respect to this equation we have two problems;
  - 1. when are there solutions ?
- what sort of geometric structure does the moduli space ( parameter space ) of solutions admit ?

In this lecture we discuss these over a four-dimensional manifold, especially over a complex 2-dim complex manifold with positive definite metric.

 $\underline{\underline{2}}$ . Let P be a principal fibre bundle with gauge group SU(n) over a compact oriented 4-dim manifold M with metric h.

<u>Definition</u>. A connection on P is a system  $A = \{A_{\alpha}, A_{\beta}, A_{\beta}, A_{\beta}, A_{\beta}, A_{\alpha}\}$  where each  $A_{\alpha}$  is an  $\mathcal{SU}(n)$ -valued 1-form defined over a trivializing neighborhood  $U_{\alpha}$  of P such that

these  $A_{\alpha}$ 's satisfy over  $U_{\alpha} \wedge U_{\beta} \neq \phi$ 

$$(*) A_{\beta} = g^{-1} \cdot dg + g^{-1} \cdot A_{\alpha} \cdot g.$$

Here  $g=g_{\alpha\beta}$  denotes the transition function of P:  $U_{\alpha} \cap U_{\beta} \xrightarrow{C^{\infty}} SU(n)$  and dot means multiplication of matrices.

We have another global definition of connection, equivalent to the above. But we adopt the above for a convenience.

For each Lie algebra valued function on M the covariant derivative  $\nabla_A$  is defined:  $\nabla_A \phi = d \phi + [A, \phi] = d \phi + A \cdot \phi - \phi \cdot A$  and we have also the covariant exterior derivative  $d_A \psi = d \psi + [A \wedge \psi] = d \psi + A \wedge \psi + (-1)^{p+1} \cdot \psi \wedge A$  ( $\psi$  is a p-form).

A 2-form  $F(A)=dA+A\Lambda A$  is said curvature form of A. The curvature form takes values in  $\mathcal{M}(n)$ . In local expression F(A) is represented by  $F(A)=\frac{1}{2}\sum_{\mu,\nu}F_{\mu\nu}\,dx^{\mu}\Delta dx^{\nu}$  ( $F_{\mu\nu}=-F_{\nu\mu}$ ),  $F_{\mu\nu}=\partial_{\mu}A_{\nu}-\partial_{\nu}A_{\mu}+[A_{\mu},A_{\nu}]$ .

We have an identity  $d_A^{\ } \nabla_A^{\ } \phi \ = \ [F(A), \phi] \,, \quad called$  Ricci formula which will be used later to define an elliptic complex.

Definition. Over the set of connections  $\mathcal{C}_{\mathtt{P}}$  Yang-Mills functional  $\mathcal{Y}^{\mathfrak{M}}$  is defined by

$$\mathcal{Y}\mathcal{M}(A) = \frac{1}{2} \int_{M} |F(A)|^{2} dv,$$

where  $|F(A)|^2 = \frac{1}{2} \sum_{i} h^{\mu \sigma} h^{\nu \tau} (-\text{Tr } F_{\mu \nu} \cdot F_{\sigma \tau})$  (  $(h^{\sigma \tau})$  is the inverse matrix of the metric  $(h_{\mu \nu})$ ).

Its Euler-Lagrange equation is written locally by

$$\sum_{\sigma,\tau} h^{\sigma\tau} \nabla_{\sigma} F_{\tau\mu} = \sum_{\sigma,\tau} h^{\sigma\tau} (D_{\sigma} F_{\tau\mu} + [A_{\sigma}, F_{\tau\mu}]) = 0$$

where  $D_{\sigma}$  denotes the covariant derivative with respect to Levi-Civita connection of M.

<u>Definition</u>. A connection is called Yang-Mills if it is a critical point of ym, that is, it is a solution to the Yang-Mills equation.

We have assumed that M is oriented and 4-dimensional. Then we have the operator \* which is an involutive endomorphism of the bundle of 2-forms  $\Lambda^2$ . Hence  $\Lambda^2 \quad \text{splits into} \quad \Lambda_+^2 + \Lambda_-^2 \quad \text{corresponding to eigenvalues.}$  A 2-form  $\alpha$  is said (anti-)self-dual if  $\alpha \in \Lambda_+^2 \text{ (or } \Lambda_-^2 \text{).}$  Curvature form splits also into  $F(A) = F_+ + F_-.$ 

<u>Definition</u>. A connection is said self-dual ( or anti-self-dual ) if it is a solution of self-dual equation  $F_{-}=0$  ( or anti-self-dual equation  $F_{+}=0$  ).

From Chern-Weil homomorphism theorem  $\ensuremath{\mathcal{Y}}\ensuremath{\mathcal{M}}$  has the lowest bounds;

$$\mathcal{Y}\mathcal{M}(A) = \frac{1}{2} \int_{M} \{|F_{+}|^{2} + |F_{-}|^{2}\} dv \ge \frac{1}{2} |\int_{M} \{|F_{+}|^{2} - |F_{-}|^{2}\} dv |$$

for all A  $\varepsilon$   $C_p$  and the integral of the right hand side represents  $(-c) c_2(P)[M]$  for a universal constant c>0 and the second Chern number  $c_2(P)[M]$ . Then the equality "=" holds if and only if A is self-dual or anti-self-dual according to  $c_2(P)[M] \leq 0$  or  $\geq 0$ . Thus we have

<u>Proposition 2.1</u>. Every self-dual connection ( or antiself-dual connection ) minimizes ym. Hence it is a Yang-Mills connection.

Note. The Yang-Mills equation is anon-linear second order equation with respect to unknown connection. But the (anti-) self-duality equation is of first order.

For Problem 1  $c_2(P)[M] \ge 0$  is a necessary condition for P to admit an anti-self-dual connection. Conversely we have

Theorem (Taubes [11]). If  $H_+^2(M) = \{\text{ harmonic 2-form which is anti-self-dual }\}$  vanishes, then P with  $c_2(P)[M] \ge 0$  admits an irreducible anti-self-dual connection provided that a principal fibre bundle over  $S^4$  with identical Chern number carries an irreducible anti-self-dual connection.

The irreducibility of connection will be defined at § 4. It is seen that a connection is irreducible if and only if the gauge group can not be compressed into  $S(U(n_1)\times U(n_2))$ ,  $(n_1 + n_2 = n)$ .

With respect to Problem 2 we have by fundamental ideas of Atiyah-Hitchin-Singer [2] and also of Donaldson [3] the moduli space of anti-self-dual connections carries a structure of real analytic set. Donaldson applied the structure to 4-dimensional topology.

3. Assume that the base space M is a complex surface with complex coordinates  $z^1$ ,  $z^2$  and of a Kähler metric  $h = \sum_{\mu,\nu} h_{\mu\nu}(z^1,z^2) dz^{\mu} dz^{\nu}$ . Then the space M carries an orientation induced naturally from the complex structure.

Fact(Atiyah[1],Itoh[6]) A connection over a complex surface is anti-self-dual if and only if F(A) is of type (1,1) and primitive, that is, F(A) is written by  $F(A) = \sum_{\mu,\nu} F_{\mu\nu} dz^{\mu} \Lambda dz^{\bar{\nu}} \quad \text{and satisfies}$   $\sum_{\sigma,\tau} h^{\sigma\bar{\tau}} F_{\sigma\bar{\tau}} = 0.$ 

Koiso generalizes this anti-self-duality to higher dimensional complex manifolds [9].

Theorem (Kobayashi[8]). If the bundle P admits an antiself-dual connection, then the associated vector bundle  $E = P \times_{\rho} \mathbb{C}^n \quad \text{must be a (semi-)stable holomorphic vector}$  bundle in the sense of Mumford and Takemoto.

Recently Donaldson proved in his preprint [4] that over an algebraic surface the stability is indeed a sufficient condition.

Over a Kähler surface the following is an answer to Problem 2.

Theorem [7]. The moduli space of irreducible anti-self-dual connections admits over a compact Kähler surface a structure of complex analytic set, that is, it is a zero point set of local several holomorphic functions.

- Remarks. 1. An anti-self-dual connection is just an Einstein Hermitian structure of the associated vector bundle in the sense of Kobayashi.
- 2. Over an algebraic surface the moduli of stable holomorphic structure of a smooth vector bundle is a quasi-projective variety (Maruyama[10]).

Example. In the case of  $M = P_2(\mathbb{C})$ , SU(2) and  $c_2(P)$  = k (> 1) the moduli space is a complex manifold without singularity of  $\dim_{\mathbb{C}}$  4 k - 3.

The theorem is first conjectured by Atiyah.

A sketch of proof of Theorem will be given in \ \ \ 5.

To define precisely the moduli space over a general 4-manifold we need notion of gauge transformations. f of P (that is, f;  $P \longrightarrow P$ A bundle automorphism is a diffeomorphism and satisfies f(ua) = f(u)awhich induces on M the identity transformation is called a guage transformation. A gauge transformation can be identified with a global section of the automorphism bundle  $G_{p} = P \times_{conj} SU(n)$ . A gauge transformation operates on a connection A to induce a new connection f\*A =  $f^{-1}.df + f^{-1}.A.f$  with curvature form  $F(f*A) = f^{-1}.F(A).f.$ Therefore ym is gauge-invariant, because  $|F(f*A)|^2 =$  $|F(A)|^2$  and an anti-self-dual connection is transformed into new one which is also anti-self-dual. But in physical and geometrical meaning they represent the same thing. So we have a quotient space of the set of anti-self-dual connections modulo gauge transformations and call it moduli of anti-self-dual connections: space  $\mathcal{M}_{a}$ 

{ anti-self-dual connections } 
$$\xrightarrow{\pi} \mathcal{M}_{a}$$
.

The adjoint representation of the gauge group induces a vector bundle associated to P which we call the adjoint bundle and denote it by  $\mathcal{G}_{\text{D}}.$ 

Almost everything of Yang-Mills theory must be discussed over this bundle because curvature forms are  $\mathcal{G}_p$ -valued 2-forms and an infinitesimal gauge transformation is just a global section of  $\mathcal{G}_p$  etc.

Now we assume that a connection A is anti-self-dual over a general compact oriented Riemannian 4-manifold M. By the aid of the Ricci formula we have the sequence:

$$(*) \quad 0 \longrightarrow \Omega^{0}(\mathcal{O}_{P}) \xrightarrow{\nabla_{A}} \Omega^{1}(\mathcal{O}_{P}) \xrightarrow{d_{A}^{+} = P_{+} \circ d_{A}} \Omega_{+}^{2}(\mathcal{O}_{P}) \longrightarrow 0$$

is an elliptic complex, that is,  $d_A^+ \circ V_A^- = 0$  and the symbol sequence of this is exact at any covector  $\xi \neq 0$ . Here  $\Omega^k(\mathcal{O}_P)$  denotes the space of smooth  $\mathcal{O}_P^-$ -valued k-forms ( k = 1, 2 ) and  $\Omega_+^2(\mathcal{O}_P)$  the space of  $\mathcal{O}_P^-$ -valued self-dual 2-forms. Further  $P_+$  is the orthogonal projection of  $\Lambda^2$  onto  $\Lambda_+^2$ . Cohomologies  $H^0$ ,  $H^1$  and  $H_+^2$  associated to (\*) are all finite dimensional. The Atiyah-Singer index theorem shows that the index of (\*), that is,  $\dim H^0 - \dim H^1$  +  $\dim H_+^2$  is given by  $\int_M$  { characteristic classes of M and P }

According to either  $H^0=0$  or  $H^0\neq 0$  and either  $H^2_+=0$  or  $H^2_+\neq 0$  we have four cases:

(i)  $H^0 = 0$  and  $H_+^2 = 0$ ; the moduli space  $\mathcal{M}_a$  is smooth at  $[A] = \pi(A)$  and has dimension equal to - index (\*)

- (ii)  $H^0 = 0$  and  $H^2_+ \neq 0$ ;  $\mathcal{M}_a$  is at [A] a zero point set of  $a \ C^\omega map \ \Phi \quad from \ a \ neighborhood \ of \ H^1$  to  $H^2_+$
- (iii)  $H^0 \neq 0$  and  $H^2_+ = 0$ ;  $\mathcal{M}_a$  is written as a  $\Gamma_A$ -quotient of a 0-neighborhood of  $H^1$
- (iv)  $H^0 \neq 0$  and  $H_+^2 \neq 0$ ;  $\mathcal{M}_a$  has at [A] a structure of  $\Gamma_A$ -quotient of zero point set of a  $C^\omega$ -map  $\Phi$  from a 0-neighborhood of  $H^1$  to  $H_+^2$

Here  $\Gamma_A$  is the isotropy group of A, that is,  $\Gamma_A$  = { gauge transformations f; f\*A = A }.

Those properties of  $\, \mathcal{M}_{\rm a} \,$  stated above are based on the following facts:

Fact 1. Since  $F(A + \alpha) = F(A) + d_A^{\alpha} + \alpha \Lambda \alpha$  ( $\alpha \in \Omega^1(\mathcal{O}_P)$ ), for a fixed anti-self-dual connection A a connection  $A + \alpha$  is anti-self-dual if and only if  $d_A^+ \alpha + P_+(\alpha \Lambda \alpha) = 0$ .

- Fact 2. Because  $\operatorname{Ker} \operatorname{d}_A^* = \{ \alpha \in \Omega^1(\mathcal{O}_P); \operatorname{d}_A^*\alpha = 0 \}$  is transversal in  $\mathcal{C}_P$  to the orbit through A of gauge transformations, where  $\operatorname{d}_A^*$  is the formal adjoint of  $\operatorname{d}_A$ ,  $\mathcal{M}_A$  has a neighborhood homeomorphic to a slice  $\mathcal{C}_A = \{ \alpha \in \Omega^1(\mathcal{O}_P); \operatorname{d}_A^*\alpha = 0, \operatorname{d}_A^+\alpha + \operatorname{P}_+(\alpha \Lambda \alpha) = 0 \}$  or a  $\Gamma_A$  quotient  $\mathcal{L}_A^{\prime}\Gamma_A$  according to the irreducibility of A. Fact 3. We combine these facts to derive so-called Kuranishi's map from the slice to  $\operatorname{H}^1$ . Define a map  $\Phi$ ;  $\Omega^1(\mathcal{O}_P) \longrightarrow \Omega^1(\mathcal{O}_P)$ ,  $\Phi(\alpha) = \alpha + (\operatorname{d}_A^+)^*\circ \operatorname{G}_A(\operatorname{P}_+\alpha \Lambda \alpha)$  where  $\operatorname{G}_A$  is the Green operator of the Laplacian  $\operatorname{d}_A^{\dagger}\circ (\operatorname{d}_A^{\dagger})^*$ . Since  $\operatorname{d}_A^{\dagger}\circ (\operatorname{d}_A^{\dagger}) = \operatorname{d}_A^{\dagger}\circ (\operatorname{d}_A^{\dagger})^*$ . Since  $\operatorname{d}_A^{\dagger}\circ (\operatorname{d}_A^{\dagger}) = \operatorname{d}_A^{\dagger}\circ (\operatorname{d}_A^{\dagger})^*$  has an inverse over a neighborhood of 0. We give a map  $\psi$ ;  $\{\beta \in \operatorname{Ker} \operatorname{d}_A^*; \|\beta\|_{L^2} < \varepsilon \} \longrightarrow \operatorname{H}^2$  by  $\psi(\beta) = -$  harmonic part of  $\operatorname{L}_A^{\dagger}\circ (\operatorname{d}_A^{\dagger})^*$  for  $\alpha = \Phi^{-1}(\beta)$  to assert that  $\Phi$ ;  $\mathcal{L}_A^{\dagger} = \operatorname{L}_A^{\dagger}\circ (\operatorname{d}_A^{\dagger})^*$ .
- Now we assume again in this section that the base space is Kähler. In order to give the moduli space  $\mathcal{M}_a$  a complex manifold structure we need finer considerations than the discussion of §4. For this purpose we complexify every real structures, SU(n) to SL(n; $\mathbb{C}$ ), the principal fibre bundle, the adjoint bundle, gauge transformations etc, except but the transition functions of the bundle. A connection A then splits into the (1,0)-part A' and the (0,1)-part A" so that A'' satisfies

$$A''_{\beta} = g^{-1}.\overline{\partial}g + g^{-1}.A''_{\alpha}.$$

<u>Definition</u>. A system  $\{\widetilde{A}_{\alpha}, \widetilde{A}_{\beta}, \ldots\}$  of locally defined  $\mathfrak{sl}(n;\mathfrak{C})$ -valued (0,1)-forms compatible with the transition functions of P,

$$\widetilde{A}_{\beta} = g^{-1} \cdot \overline{\partial} g + g^{-1} \cdot \widetilde{A}_{\alpha} \cdot g$$
 on  $U_{\alpha} \cap U_{\beta}$ ,

is called a (0,1)-connection.

A (0,1)-connection  $\widetilde{A}$  induces the partial covariant derivative  $\overline{\partial}_{\widetilde{A}} \colon \Omega^0(\mathfrak{G}_P^{\mathbb{C}}) \longrightarrow \Omega^{0,1}(\mathfrak{G}_P^{\mathbb{C}}), \qquad \phi \longmapsto \partial \phi + [A, \phi]$  and also the covariant exterior derivative  $\overline{\partial}_A \colon \Omega^{p,q}(\mathfrak{G}_P^{\mathbb{C}}) \longrightarrow \Omega^{p,q+1}(\mathfrak{G}_P^{\mathbb{C}})$  where  $\mathfrak{G}_P^{\mathbb{C}}$  denotes the complexification of  $\mathfrak{G}_P$ .

<u>Definition</u>. A (0,1)-connection  $\widetilde{A}$  is said holomorphic if its curvature  $F(\widetilde{A}) = \overline{\partial} \widetilde{A} + \widetilde{A} \wedge \widetilde{A}$  vanishes.

Remark. Each holomorphic (0,1)-connection induces by integrability condition a holomorphic structure on  $\mathcal{G}_{\mathbb{P}}^{\mathbb{C}}$  such that  $\bar{\partial}_{A}$  is just  $\bar{\partial}$ -operator(Atiyah-Hitchin-Singer [2] and Griffiths [5]).

Note. Holomorphic (0,1)-connection can be also defined over a higher dimensional complex manifold.

In a way similar to the case of anti-self-dual connections we can define moduli space  $\mathcal{M}_h$  of holomorphic (0,1)-connections with respect to complex gauge transformations.

Since each anti-self-dual connection A has curvature form of type (1,1), its (0,1)-component A" gives automatically a holomorphic (0,1)-connection because  $F(A") \quad \text{is just the } (0,2)\text{-component of full curvature form } F(A).$  Then we have a canonical map  $\phi$  from  $\mathcal{M}_a$  to  $\mathcal{M}_h$  by assigning [A"] to [A].

When  $\widetilde{A}$  is holomorphic, the sequence (\*\*) is an elliptic complex

$$(**) \quad 0 \longrightarrow \Omega^{0}(\mathcal{I}_{P}^{\mathbb{C}}) \xrightarrow{\overline{\partial}_{\widehat{A}}} \Omega^{0,1}(\mathcal{I}_{P}^{\mathbb{C}}) \xrightarrow{\overline{\partial}_{A}} \Omega^{0,2}(\mathcal{I}_{P}^{\mathbb{C}})$$

$$\longrightarrow 0.$$

Proposition 5.1. The moduli space of irreducible holomorphic
(0,1)-connections is a complex analytic set.

Here a (0,1)-connection is called irreducible when  $\operatorname{Ker} \{ \ \overline{\partial}_{\widetilde{A}} \colon \ \Omega^0(\ {\mathfrak G}_{\operatorname{P}}^{\ \mathbb{C}}) \longrightarrow \Omega^{0,1}(\ {\mathfrak G}_{\operatorname{P}}^{\ \mathbb{C}}) \} = 0$ . This proposition is shown by a way similar to the anti-self-dual case, while the complex analyticity is assured by the fact that the Kuranishi's map is in this case holomorphic.

Proposition 5.2. i) The spaces  $\mathcal{M}_a$  and  $\mathcal{M}_h$  of generic connections ( that is, connections with  $H^0 = 0$  and  $H^2 = 0$ , respectively ) are smooth manifolds with same real dimension  $2 \{ c_2(\mathcal{O}_P^{\mathbb{C}}) - \dim SU(n) \cdot p(M) \}$  where p(M) is the arithmetic genus of M and ii) the canonical map  $\phi$  is smooth and of maximal rank over  $\mathcal{M}_{a,gen} = \{ [A] \in \mathcal{M}_a; A \text{ is generic } \}$ 

Note. We use the moment map due to Donaldson [4] to conclude that  $\phi$  is one-to-one over  $\mathcal{M}_{a,ir} = \{ [A] \in \mathcal{M}_a; A \text{ is irreducible } \}$ 

As a consequence of these results the moduli space  $\mathcal{M}_{a,ir}$  has a structure of complex analytic set which around a generic one is smooth and is sigular at one with  $H_+^2 = 0$ . This is just the statement of Theorem.

We remark that around a reducible anti-self-dual connection the moduli space  $\, \mathcal{M}_{\rm a} \,$  is a  $\, \Gamma_{\rm A} \!$ -quotient of a real analytic set.

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