Statistical closure using bulk properties for turbulent shear flows

Akira Yoshizawa

吉 澤 徴

Of a variety of turbulence models, the $k-\epsilon$ model is widely used for engineering purposes. Its importance is diminishing, however, from the viewpoint of scientific turbulence research. This fact is attributed to the deficiency that the $k-\epsilon$ model cannot predict anisotropy of turbulent intensities accurately.

Previously, the author studied the Reynolds stress with the aid of a two-scale direct-interaction approximation (TSDIA). This formalism is founded on an asymptotic expansion based on a scale parameter distinguishing the slow variation of mean flows from the fast variation of turbulent fluctuations. Consequently, TSDIA may be regarded as a derivative expansion formalism incorporated with DIA 2 , by Kraichnan (for details about the structure of scale expansion in TSDIA, see Ref.4). For the Reynolds stress $R^{\alpha\beta}$, TSDIA gives

$$R^{\alpha\beta} \equiv -\langle u^{\dagger \alpha} u^{\dagger \beta} \rangle$$

$$= -\frac{2}{3} k \delta^{\alpha\beta} + \nu_{e} (\frac{\partial U^{\alpha}}{\partial x^{\beta}} + \frac{\partial U^{\beta}}{\partial x^{\alpha}}) + \tau^{\alpha\beta}$$
(1)

(< > denotes the ensemble mean). Here, \overrightarrow{u} and \overrightarrow{u} are the

ensemble mean and fluctuation of velocity, k is the turbulent kinetic energy ($k \equiv \langle u'^a u'^a \rangle / 2$; the summation convention is applied to repeated superscripts), and $\delta^{\alpha\beta}$ is the Kronecker delta symbol. The first two terms on the right-hand side give the familiar eddy-viscosity representation. The appearance of isotropic eddy viscosity ν_{α} is due to an assumption about isotropy of the basic field or the lowestorder term in the two-scale expansion of the fluctuating field. On the other hand, $\tau^{\alpha\beta}$ expresses the deviation from the isotropic eddy-viscosity representation (the details are given in Ref.1). Equation (1) was incorporated into the $k-\epsilon$ model to be applied successfully to the study of anisotropic turbulent intensities in channel flows. 5 This improvement of the $k-\epsilon$ closure for turbulent shear flows indicates a possibility that models of $k\!-\!\epsilon$ type as well as the secondorder models can remain a useful method for simulating complicated flows such as recirculating flows.

At this time, it is a major concern about turbulence modeling (including the second-order models) that no statistical foundation has yet been given to a model equation for ϵ (the dissipation rate of turbulent kinetic energy), unlike a model equation for k. In reality, the structure of the transport equation for k is simple, and its modeling is not difficult. Some theoretical support can be given to the modeling by using turbulence theories such as DIA. 6

On the other hand, the structure of the transport equation for ϵ is extremely complicated. As a result, the

 ϵ equation in the k- ϵ model has not been constructed statistically so far on the basis of the exact transport equation for ϵ . This situation will become a big obstacle when we desire to incorporate anisotropic effects such as and buoyancy force into the ϵ equation. This point is a major motivation of present work.

In the present paper, we can make use of an analytical expression for k from TSDIA to show that a model ϵ equation is derived with the aid of the concept of transferability of model. The model equation is very similar to the counterpart in the k- ϵ model.

The details of this work are given in Ref.7.

- 1. A. Yoshizawa, Phys. Fluids 27, 1377 (1984).
- 2. R.H. Kraichnan, J. Fluid Mech. 5, 497 (1959).
- 3. R.H. Kraichnan, Phys. Fluids 7, 1048 (1964).
- 4. A. Yoshizawa, Phys. Fluids 28, 3226 (1985).
- 5. S. Nisijima and A. Yoshizawa, AIAA J. 25, 414 (1987).
- 6. A. Yoshizawa, Phys. Fluids 28, 59 (1985).
- 7. A. Yoshizawa, Phys. Fluids 30, 628 (1987).