Periodic solutions of Boussinesq equations

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Let Ω be a bounded domain in \mathbb{R}^2 with the boundary $\partial\Omega$ such that

$$\partial\Omega=\Gamma_1\cup\Gamma_2,\ \Gamma_1\cap\Gamma_2=\phi.$$

We consider the following initial boundary value problem:

$$\begin{cases} \frac{\partial u}{\partial t} + (u \cdot \nabla)u = -\frac{1}{\rho} \nabla p + \nu \Delta u + \beta g \theta \\ \operatorname{div} u = 0, & x \in \Omega, t > 0, \\ \frac{\partial \theta}{\partial t} + (u \cdot \nabla)\theta = \chi \Delta \theta, \end{cases}$$
 (1)

$$\begin{cases} u(x,t) = 0, \ \theta(x,t) = \xi(x,t), & x \in \Gamma_1, t > 0, \\ u(x,t) = 0, \frac{\partial}{\partial n} \theta(x,t) = \eta(x,t), & x \in \Gamma_2, t > 0, \end{cases}$$
 (2)

$$\begin{cases}
 u(x,0) = a_0(x), \\
 \theta(x,0) = \tau_0(x),
\end{cases} \qquad x \in \Omega, \tag{3}$$

where $u=(u_1,u_2)$ is the fluid velocity, p is the pressure, θ is the temperature, $u\cdot \nabla = \sum_{j=1}^2 u_j \frac{\partial}{\partial x_j}$, $\frac{\partial \theta}{\partial n}$ denotes the outer normal derivative of θ at x to $\partial \Omega$, g(x,t) is the gravitational vector function, and ρ (density), ν (kinematic viscosity), β (coefficient of volume expansion), χ (thermal diffusivity) are positive constants. $\xi(x,t)$ (resp. $\eta(x,t)$) is a function defined on $\Gamma_1 \times (0,T)$ (resp. $\Gamma_2 \times (0,T)$) and $a_0(x)$ (resp. $\tau_0(x)$) is a vector (resp. scalar) function defined on Ω .

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In order to state our results, we introduce some **Function spaces**([1],[2],[3]).

 $L^p(\Omega)$ and the Sobolev space $W_p^\ell(\Omega)$ are defined as usual. We also denote $\mathbf{L}^p(\Omega) = L^p(\Omega) \times L^p(\Omega)$, $H^\ell(\Omega) = W_2^\ell(\Omega)$. Whether the elements of the space are scalar or vector functions is understood from the contexts unless stated explicitly.

 $D_{\sigma} = \{ \text{vector function } \varphi \in C^{\infty}(\Omega) \mid \text{supp} \varphi \subset \Omega, \text{div} \varphi = 0 \text{ in } \Omega \},$

 $H = \text{completion of } D_{\sigma} \text{ under the } L^{2}(\Omega)\text{-norm},$

 $V = \text{completion of } D_{\sigma} \text{ under the } H^1(\Omega)\text{-norm};$

 $D_0 = \{ \text{scalar function } \varphi \in C^{\infty}(\overline{\Omega}) \mid \varphi \equiv 0 \text{ in a neighborhood of } \Gamma_1 \},$

 $W = \text{completion of } D_0 \text{ under the } H^1(\Omega)\text{-norm},$

V', W' are dual space of V, W.

Definition 1

 $\{u,\theta\}$ is called a weak solution of evolutional problem (1),(2) if, for some function θ_0 such that

$$\theta_0 \in L^2(0,T:H^1(\Omega)), \quad \theta_0 = \xi \text{ on } \Gamma_1,$$

 $\{u, \theta\}$ satisfies following conditions:

$$u \in L^2(0,T:V), \quad \theta - \theta_0 \in L^2(0,T:W),$$

$$\begin{cases}
\frac{d}{dt}(u,v) + \nu(\nabla u, \nabla v) + ((u \cdot \nabla)u, v) - (\beta g\theta, v) = 0, & \forall v \in V, \\
\frac{d}{dt}(\theta,\tau) + \chi(\nabla \theta, \nabla \tau) + ((u \cdot \nabla)\theta, \tau) - \chi(\eta,\tau)_{\Gamma_2} = 0, & \forall \tau \in W,
\end{cases}$$
(4)

where

$$(\eta, au)_{\Gamma_2} = \int_{\Gamma_2} \eta(x') au(x') d\sigma.$$

As for the smoothness of $\partial\Omega$, we suppose Condition (H)

 $\partial\Omega$ is of class C^1 and devided as follows:

$$\partial\Omega=\Gamma_1\cup\Gamma_2,\quad \Gamma_1\cap\Gamma_2=\phi,\quad \text{measure of }\Gamma_1\neq0,$$

and the intersection $\overline{\Gamma}_1 \cap \overline{\Gamma}_2$ consists of finite points.

In [3], we showed the existence and the uniqueness of weak solution of evolutional problem for $2 \le n \le 4$. For n = 2, we have the following result: **Theorem A**

Let Ω be a bounded domain in R^2 with C^1 boundary satisfying Condition(H). If the function g is in $L^{\infty}(\Omega \times (0,T))$, $\xi \in C^1(\overline{\Gamma}_1 \times [0,T])$, $\eta \in L^2(\Gamma_2 \times (0,T))$, $a_0 \in H$, $\tau_0 \in L^2(\Omega)$, then there exists one and only one weak solution $\{u,\theta\}$ of (1), (2) satisfying the initial condition (3). Furthermore

$$u \in C([0,T]:H), \ \theta \in C([0,T]:L^2(\Omega)).$$

Definition 2

 $\{u,\theta\}$ is called a periodic weak solution of (1), (2) with period T_0 , if $\{u,\theta\}$ is a weak solution of (1), (2) for $T=T_0$ satisfying

$$u(x, T_0) = u(x, 0), \quad \theta(x, T_0) = \theta(x, 0).$$
 (5)

We also obtained the existence of periodic weak solutions([3]).

Theorem B

Let Ω be a bounded domain in R^2 with C^1 boundary satisfying Condition (H). Let g(x,t), $\xi(x,t)$, $\eta(x,t)$ be periodic with respect to t with period T_0 , satisfying $g \in L^{\infty}(\Omega \times (0,T_0))$, $\xi \in C^1(\overline{\Gamma}_1 \times [0,T_0])$ and $\eta \in L^2(\Gamma_2 \times (0,T_0))$. Set $g_{\infty} = ||g||_{L^{\infty}(\Omega \times (0,T_0))}$. If $\frac{\beta g_{\infty}}{\sqrt{\nu \chi}}$ is sufficiently small, then there exists a periodic weak solution of (1), (2) with period T_0 . Furthermore

$$u \in C([0,\infty): H), \ \theta \in C([0,\infty): L^2(\Omega)).$$

Now we can state our results. As for the uniqueness of periodic weak solutions, we obtained:

Theorem 1

Let $\{u_{\pi}, \theta_{\pi}\}$ be a weak periodic solution of (1), (2) with period T_0 such that for some p > 2,

$$\operatorname{ess.sup}_{t}\{c||u_{\pi}(t)||_{p} + \frac{1}{4\chi}(c||\theta_{\pi}(t)||_{p} + c'\beta g_{\infty})^{2}\} < \nu, \tag{6}$$

where c and c' are constants depending on Ω . If $\{u_{\pi} + u, \theta_{\pi} + \theta\}$ is a weak periodic solution of (1), (2) with period T_0 , then $u = 0, \theta = 0$.

Let $g \in L^{\infty}(\Omega \times (0,\infty))$, $\xi \in C^1(\overline{\Gamma}_1 \times [0,\infty))$ $\eta \in L^2(\Gamma_2 \times (0,\infty))$, $a_0 \in H, \tau_0 \in L^2(\Omega)$. Let T be any positive number. Then there exists one and only one weak solution $\{u_T, \theta_T\}$ of (1), (2) satisfying (3). Therefore, for T < T',

$$u_T(t) = u_{T'}(t), \quad \theta_T(t) = \theta_{T'}(t) \quad \text{ for } \forall t \in (0, T)$$

hold, and we can omit T. This solution is called a global weak solution. We obtained the asymptotic property of solutions of Boussinesq equations as follows.

Theorem 2

Let g, ξ, η satisfy the condition of Theorem B, $a_0 \in H$, $\tau_0 \in L^2(\Omega)$. Let $\{u, \theta\}$ be a global weak solution of (1), (2) satisfying (3), $\{u_{\pi}, \theta_{\pi}\}$ a periodic weak solution satisfying (6). Then

$$\lim_{t\to\infty} \{||u(t) - u_{\pi}(t)||^2 + ||\theta(t) - \theta_{\pi}(t)||^2\} = 0.$$

Remark

- (i) Since $u_{\pi} \in L^{2}(0,T:V) \cap C([0,T]:H)$, u_{π} belongs to the space $L^{2p/(p-2)}(0,T:\mathbf{L}^{p}(\Omega))$ for $\forall p > 2$. Similarly θ_{π} is in $L^{2p/(p-2)}(0,T:L^{p}(\Omega))$. The condition (6) is stronger than this one.
- (ii) When (6) holds, such periodic solution is unique (Theorem 1).

References

[1] Morimoto, H., On the existence of weak solutions of equations of natural convection, J. Fac. Sci. Univ. Tokyo, Sect.IA, 36 (1989), p.87-102.

- [2] Morimoto, H., On the existence of weak solutions of stationary Boussinesq equations, to appear in Sûrikaiseki Kenkyûsho Kôkyûroku.
- [3] Morimoto, H., Nonstationary Boussinesq equations and its reproductive property, Preprint.
- [4] Temam, R., Navier-Stokes Equations, North-Holland, Amsterdam, 1979.