Periodicizing Functions

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1 Introduction

We denote by \mathbb{R} and by \mathbb{C} the set of real numbers and complex numbers, respectively. Let us consider a linear differential equation of the form

$$\frac{dx}{dt} = Ax + f(t), \quad x(0) = w \in \mathbb{C}^d, \tag{1}$$

where $A \in M_d(\mathbb{C})$, the set of all $d \times d$ complex matrices, and $f : \mathbb{R} \to \mathbb{C}^d$ is a nontrivial continuous τ -periodic function.

The purpose of this paper is to find a periodicizing function for Equation (1). It is well known that the solution of the above equation is expressed as

$$x(t) := x(t; 0, w) = e^{At}w + \int_0^t e^{A(t-s)}f(s)ds.$$
 (2)

However, from this representation it is not easy to see asymptotic behaviors of solutions of Equation (1). In the paper [2], we gave a new representation of the solution of Equation (1), from which asymptotic behaviors of solutions are seen.

Its representation is essentially related to a periodicizing function for Equation (1), as stated below. We take a continuous function z(t) such that the function

$$h(t):=z(t)+\int_0^t e^{A(t-s)}f(s)ds, \ \ t\in \mathbb{R},$$

becomes a continuous τ -periodic function. Then the solution x(t) of Equation (1) is rewritten as

$$x(t) = (e^{At}w - z(t)) + h(t),$$

the first term of the right hand side in which is well known. We call such a function z(t) a "periodicizing" function (for Equation (1)). Therefore, to find a periodicizing function is very important in obtaining the representation of solutions and in

studying asymptotic behaviors of solutions for Equation (1). In this paper we will construct a periodicizing function z(t) as follows. Put

$$b_f = \int_0^{ au} e^{A(au-s)} f(s) ds.$$

In the first step, we prove that a periodicizing function z(t) satisfies

$$\Delta_{\tau}z(t) := z(t+\tau) - z(t) = -e^{tA}b_f, \quad t \in [0, \infty).$$

In the next step, we calculate an indefinite sum of $-e^{tA}b_f$; that is, $z(t) = \Delta_{\tau}^{-1}(-e^{At}b_f)$ by using a new representation of the solution obtained in [2] for the linear difference equation of the form $x_{n+1} = e^{\tau A}x_n + b_f$.

2 Discrete linear difference equations and the indefinite sum

2.1 Discrete linear difference equations

Let $\sigma(A)$ be the set of all eigenvalues of A and m the index of $\lambda \in \sigma(A)$. Let $M_{\lambda} = \mathcal{N}((A - \lambda E)^m)$ be the generalized eigenspace of $\lambda \in \sigma(A)$, where $E \in M_d(\mathbb{C})$ stands for the unite matrix. Then we have the direct sum decomposition

$$\mathbb{C}^d = \sum_{\lambda \in \sigma(A)} \oplus M_{\lambda}.$$

Let P_{λ} be the projection on \mathbb{C}^d to M_{λ} induced from this decomposition. Set $\mathbb{N} := \{1, 2, 3, \cdots\}$.

Now, we solve the discrete linear difference equation of the form

$$x_{n+1} = e^{\tau A} x_n + b, \quad x_0 = w,$$
 (3)

where $n \in \mathbb{N} \cup \{0\}$. For the simplicity of the description, we set

$$\varepsilon(z) = \frac{1}{e^z - 1}, \quad \varepsilon^{(i)}(z) = \frac{d^i}{dz^i} \frac{1}{e^z - 1}.$$

Moreover, we define $X_{\lambda}(A)$ and $Y_{\lambda}(A)$ for $\lambda \in \sigma(A)$ as

$$X_{\lambda}(A) = \sum_{i=0}^{m-1} \varepsilon^{(i)}(\tau \lambda) \frac{\tau^{i}}{i!} (A - \lambda E)^{i} \text{ if } e^{\tau \lambda} \neq 1$$

and

$$Y_{\lambda}(A) = \sum_{i=0}^{m-1} B_i \frac{\tau^i}{i!} (A - \lambda E)^i \text{ if } e^{\tau \lambda} = 1,$$

where $B_i, i \in \mathbb{N} \cup \{0\}$, stand for Bernoulli's numbers, refer to [3].

The following result is found in [2].

Theorem 1 [2] Let $\lambda \in \sigma(A)$. The component $P_{\lambda}x_n$ of the solution $x_n, n \in \mathbb{N}$, of Equation (3) is given as follows:

1) If $e^{\tau\lambda} \neq 1$, then

$$P_{\lambda}x_{n} = e^{n\tau\lambda} \sum_{i=0}^{m-1} n^{i} \frac{\tau^{i}}{i!} (A - \lambda E)^{i} [P_{\lambda}w + X_{\lambda}(A)P_{\lambda}b] - X_{\lambda}(A)P_{\lambda}b$$
$$= e^{n\tau A} [P_{\lambda}w + X_{\lambda}(A)P_{\lambda}b] - X_{\lambda}(A)P_{\lambda}b$$

2) If $e^{\tau\lambda} = 1$, then

$$P_{\lambda}x_n = \sum_{i=0}^{m-1} \frac{n^{i+1}}{i+1} \frac{\tau^i}{i!} (A - \lambda E)^i [\tau(A - \lambda E)P_{\lambda}w + Y_{\lambda}(A)P_{\lambda}b] + P_{\lambda}w.$$

2.2 The indefinite sum

We prepare fundamental results on the indefinite sum. Let $\tau > 0$ and $h : [0, \infty) \to \mathbb{C}^d$ be a continuous function.

First, we consider the problem of finding a continuous solution of the following equation

$$\Delta_{\tau} z(t) := z(t+\tau) - z(t) = h(t), \quad t \in [0, \infty), \tag{4}$$

that is, the indefinite sum $z(t) = \Delta_{\tau}^{-1}h(t)$. If $z_0(t)$ is one of solutions of Equation (4), then any other solution z(t) is given by

$$z(t) = z_0(t) + c(t)$$

with an arbitrary continuous τ -periodic function c(t) (it is called the periodic constant). The following lemma is easily proved, refer to [3].

Lemma 2.1

1) Let $\varphi:[0,\tau]\to\mathbb{C}^d$ be a continuous function such that

$$\varphi(\tau) = \varphi(0) + h(0). \tag{5}$$

Then a continuous solution z(t) of Equation (4) satisfying the the initial condition $z(s) = \varphi(s), s \in [0, \tau]$, exists uniquely on $[0, \infty)$. Moreover, it is given by

$$z(s+n\tau) = \varphi(s) + \sum_{i=0}^{n-1} h(s+i\tau), \quad (s \in [0,\tau), \quad n = 1, 2, \cdots).$$
 (6)

2) Conversely, if a continuous function z(t) is a solution of Equation (4), then $\varphi(t) := z(t), \ t \in [0, \tau],$ satisfies the condition (5) and z(t) is given by (6).

Next, we consider a special case of Equation (4); that is,

$$z(t+\tau) - z(t) = -B(t)b, \quad t \in [0, \infty), \tag{7}$$

where $B(t), t \in [0, \infty)$, is a continuous matrix function such that

$$B(s+k\tau) = B(s)B^{k}(\tau), \ k \in \mathbb{N}. \tag{8}$$

In this case the continuous variable t in Equation (7) is reduced to the discrete variable.

Lemma 2.2

1) Let $\varphi:[0,\tau]\to\mathbb{C}^d$ be a continuous function such that

$$\varphi(\tau) = \varphi(0) - B(0)b. \tag{9}$$

Then a continuous solution z(t) of Equation (7) satisfying the initial condition $z(t) = \varphi(t), t \in [0, \tau]$, exists uniquely on $[0, \infty)$. Moreover, it is given by

$$z(s+n\tau) = \varphi(s) - B(s)x_n(0), \quad (s \in [0,\tau), \ n \in \mathbb{N}), \tag{10}$$

where $x_n(0)$ is the solution of the difference equation of the form

$$x_{m+1} = B(\tau)x_m + b, \quad x_0 = 0, \tag{11}$$

2) Conversely, if a continuous function z(t) is a solution of Equation (7), then $\varphi(t) := z(t), \ t \in [0, \tau],$ satisfies the condition (9) and z(t) is given by (10).

Proof 1) If $s \in [0, \tau)$ and $n \in \mathbb{N}$, then from (6) in Lemma 2.1 and (8) it follows that

$$z(s+n\tau) = z(s) - \sum_{i=0}^{n-1} B(s+i\tau)b$$

= $z(s) - B(s) \sum_{i=0}^{n-1} B^{i}(\tau)b$.

Clearly, we have that $\sum_{i=0}^{n-1} B^i(\tau)b = x_n(0)$. 2) is obvious.

We note that $B(t) = e^{tA}$, $A \in M_d(\mathbb{C})$, satisfies the condition (8).

3 A periodicizing function

In this section we construct a periodicizing function for Equation (1); that is,

$$\frac{dx}{dt} = Ax(t) + f(t), \ x(0) = w \in \mathbb{C}^d.$$

Let $\lambda \in \sigma(A)$. If an M_{λ} valued function y(t) satisfies the equation

$$\frac{dy}{dt} = Ay(t) + P_{\lambda}f(t),$$

we say that y(t) is a solution of Equation (1) in M_{λ} . Clearly, if x(t) is a solution of Equation (1), then $P_{\lambda}x(t)$ is a solution of Equation (1) in M_{λ} .

To apply our idea for Equation (1), we will translate the solution x(t) := x(t; 0, w) of Equation (1) as follows:

$$x(t) = e^{At}w - z(t) + h(t),$$

where

$$h(t) = z(t) + \int_0^t e^{A(t-s)} f(s) ds.$$
 (12)

The condition that h(t) is τ -periodic is equivalent to the condition that

$$z(t+ au)+\int_0^{t+ au}e^{A(t+ au-s)}f(s)ds=z(t)+\int_0^te^{A(t-s)}f(s)ds.$$

Since

$$\int_0^{t+\tau} e^{A(t+\tau-s)} f(s) ds = e^{At} b_f + \int_0^t e^{A(t-s)} f(s) ds,$$

we have

$$\Delta_{\tau}z(t) := z(t+\tau) - z(t) = -e^{At}b_f. \tag{13}$$

Therefore z(t) is an indefinite sum of $-e^{At}b_f$; that is, $z(t) = \Delta_{\tau}^{-1}(-e^{At}b_f)$. Summarizing these, we obtain the following result.

Lemma 3.1 A periodicizing function for Equation (1) is an indefinite sum of $-e^{At}b_f$. Moreover, the solution x(t) of Equation (1) is expressed as follows:

$$x(t) = e^{At}w - \Delta_{\tau}^{-1}(-e^{At}b_f) + h(t),$$

where

$$h(t) = \Delta_{\tau}^{-1}(-e^{At}b_f) + \int_0^t e^{A(t-s)}f(s)ds$$

is a τ -periodic function.

Since h(t) is a τ -periodic function and the second term of the right hand side in (12) is defined on \mathbb{R} , the periodicizing function z(t) is well defined on \mathbb{R} provided z(t) is defined on $[0,\infty)$.

Now, we are in a position to state the main theorem in this paper.

Theorem 2 Let $\lambda \in \sigma(A)$.

1) If $e^{\tau\lambda} \neq 1$, then

$$\Delta_{\tau}^{-1}(-e^{At}P_{\lambda}b) = -e^{tA}X_{\lambda}(A)P_{\lambda}b + c(t), \quad t \ge 0,$$

where c(t) is periodic constant.

2) If $e^{\tau\lambda} = 1$, then

$$\Delta_{ au}^{-1}(-e^{At}P_{\lambda}b) = -rac{e^{\lambda t}}{ au}\sum_{j=0}^{m-1}rac{t^{j+1}}{(j+1)!}(A-\lambda E)^{j}Y_{\lambda}(A)P_{\lambda}b + d(t), \quad t\geq 0,$$

where d(t) is periodic constant.

Proof Let us consider the equation

$$P_{\lambda}z(t+\tau) - P_{\lambda}z(t) = -P_{\lambda}e^{tA}b. \tag{14}$$

It follows from Lemma 2.2 that there exists a continuous solution $P_{\lambda}z(t)$ of Equation (14), which satisfies the relation

$$P_{\lambda}z(s+n\tau) = P_{\lambda}z(s) - P_{\lambda}e^{sA}x_n(0), \quad (s \in [0,\tau), n = 0, 1, 2, \cdots), \tag{15}$$

where $x_n(0)$ is the solution of Equation (3) with w = 0.

1) Assume that $e^{\lambda \tau} \neq 1$. Put $X = X_{\lambda}(A)P_{\lambda}b$. Using Theorem 1 we have

$$P_{\lambda}x_n(0) = e^{n\tau A}X - X,$$

from which yields that

$$P_{\lambda}e^{sA}x_n(0) = e^{sA} \left(e^{n\tau A}X - X\right)$$
$$= -e^{sA}X + e^{(s+n\tau)A}X.$$

Hence the relation (15) is reduced to

$$P_{\lambda}z(s+n\tau) = (P_{\lambda}z(s) + e^{sA}X) - e^{(s+n\tau)A}X.$$

Since

$$P_{\lambda}z(s+n\tau) + e^{(s+n\tau)A}X = P_{\lambda}z(s) + e^{sA}X,$$

 $c(t) := P_{\lambda}z(t) + e^{tA}X$ is τ -periodic. Therefore we obtain

$$P_{\lambda}z(t) = -e^{-tA}X + c(t).$$

2) Assume that $e^{\lambda \tau} = 1$. Put $Y = Y_{\lambda}(A)P_{\lambda}b$. Using Theorem 1 again, we have

$$\begin{split} &P_{\lambda}e^{sA}x_{n}(0)\\ &= e^{\lambda s}\sum_{k=0}^{m-1}\frac{s^{k}}{k!}(A-\lambda E)^{k}\sum_{j=0}^{m-1}\frac{n^{j+1}}{j+1}\frac{\tau^{j}}{j!}(A-\lambda E)^{j}Y\\ &= e^{\lambda s}\sum_{k=0}^{m-1}\sum_{j=0}^{m-1}\frac{s^{k}n^{j+1}\tau^{j}}{k!(j+1)!}(A-\lambda E)^{j+k}Y\\ &= \frac{e^{\lambda s}}{\tau}\sum_{i=0}^{m-1}\sum_{k+j=i}\frac{s^{k}(n\tau)^{j+1}}{k!(j+1)!}(A-\lambda E)^{i}Y\\ &= \frac{e^{\lambda s}}{\tau}\sum_{i=0}^{m-1}\sum_{k=0}^{i}\frac{s^{k}(n\tau)^{i-k+1}}{k!(i-k+1)!}(A-\lambda E)^{i}Y\\ &= \frac{e^{\lambda s}}{\tau}\sum_{i=0}^{m-1}\sum_{k=0}^{i+1}\frac{s^{k}(n\tau)^{i-k+1}}{k!(i+1-k)!}(A-\lambda E)^{i}Y - \frac{e^{\lambda s}}{\tau}\sum_{i=0}^{m-1}\frac{s^{i+1}}{(i+1)!}(A-\lambda E)^{i}Y\\ &= \frac{e^{\lambda(s+n\tau)}}{\tau}\sum_{i=0}^{m-1}\frac{(s+n\tau)^{i+1}}{(i+1)!}(A-\lambda E)^{i}Y - \frac{e^{\lambda s}}{\tau}\sum_{i=0}^{m-1}\frac{s^{i+1}}{(i+1)!}(A-\lambda E)^{i}Y. \end{split}$$

Thus the relation (15) becomes

$$P_{\lambda}z(s+n\tau) = P_{\lambda}z(s) + \frac{e^{\lambda s}}{\tau} \sum_{j=0}^{m-1} \frac{s^{j+1}}{(j+1)!} (A - \lambda E)^{j} Y$$
$$-\frac{e^{\lambda(s+n\tau)}}{\tau} \sum_{j=0}^{m-1} \frac{(s+n\tau)^{j+1}}{(j+1)!} (A - \lambda E)^{j} Y.$$

Since

$$d(t) := P_{\lambda} z(t) + \frac{e^{\lambda t}}{\tau} \sum_{j=0}^{m-1} \frac{t^{j+1}}{(j+1)!} (A - \lambda E)^{j} Y$$

is τ -periodic, we obtain

$$P_{\lambda}z(t) = -\frac{e^{\lambda t}}{\tau} \sum_{i=0}^{m-1} \frac{t^{j+1}}{(j+1)!} (A - \lambda E)^{j} Y + d(t).$$

Combining Lemma 3.1 and Theorem 2, we can obtain the following result, which slightly modifies the one given in [2].

Theorem 3 Let $\lambda \in \sigma(A)$ and x(t) := x(t; 0, w) be the solution of Equation (1). 1) If $e^{\tau \lambda} \neq 1$, then

$$P_{\lambda}x(t) = e^{At}[P_{\lambda}w + X_{\lambda}(A)P_{\lambda}b_{f}] + u_{\lambda}(t,b_{f})$$

$$= e^{\lambda t}\sum_{j=0}^{m-1} \frac{t^{j}}{j!}(A - \lambda E)^{j}[P_{\lambda}w + X_{\lambda}(A)P_{\lambda}b_{f}] + u_{\lambda}(t,b_{f}),$$

where

$$u_{\lambda}(t,b_f) = -e^{At}X_{\lambda}(A)P_{\lambda}b_f + \int_0^t e^{(t-s)A}P_{\lambda}f(s)ds$$

is a τ -periodic solution of Equation (1) in M_{λ} .

2) If $e^{\tau\lambda} = 1$, then

$$P_{\lambda}x(t) = \frac{e^{\lambda t}}{\tau} \sum_{j=0}^{m-1} \frac{t^{j+1}}{(j+1)!} (A - \lambda E)^{j} [\tau(A - \lambda E)P_{\lambda}w + Y_{\lambda}(A)P_{\lambda}b_{f}]$$
$$+ e^{\lambda t}P_{\lambda}w + v_{\lambda}(t, b_{f}),$$

where $e^{\lambda t}P_{\lambda}w$ and

$$v_{\lambda}(t,b_f):=-\frac{e^{\lambda t}}{\tau}\sum_{j=0}^{m-1}\frac{t^{j+1}}{(j+1)!}(A-\lambda E)^jY_{\lambda}(A)P_{\lambda}b_f+\int_0^te^{(t-s)A}P_{\lambda}f(s)ds$$

are τ -periodic functions, which are not necessarily a solution of Equation (1) in M_{λ} .

Proof 1) Assume that $e^{\lambda \tau} \neq 1$. Combining Lemma 3.1 and Theorem 2, we have

$$P_{\lambda}x(t) = e^{At}P_{\lambda}w - \Delta_{\tau}^{-1}(-e^{At}P_{\lambda}b_f) + u_{\lambda}(t,b_f)$$

= $e^{tA}[P_{\lambda}w + X_{\lambda}(A)P_{\lambda}b_f] + u_{\lambda}(t,b_f),$

where

$$u_{\lambda}(t,b_f) = -e^{tA}X_{\lambda}(A)P_{\lambda}b_f + \int_0^t e^{(t-s)A}P_{\lambda}f(s)ds.$$

Notice that the periodic constant c(t) is canceled. It is easy to see that $u_{\lambda}(t, b_f)$ is a τ -periodic solution of Equation (1) in M_{λ} .

2) Assume that $e^{\lambda \tau} = 1$. In view of Lemma 3.1 we have

$$P_{\lambda}x(t) = e^{At}P_{\lambda}w - \Delta_{\tau}^{-1}(-e^{At}P_{\lambda}b_f) + v_{\lambda}(t,b_f),$$

where

$$v_{\lambda}(t,b_f) = \Delta_{ au}^{-1}(-e^{At}P_{\lambda}b_f) + \int_0^t e^{(t-s)A}P_{\lambda}f(s)ds.$$

Furthermore, from Theorem 2 we have

$$e^{tA}P_{\lambda}w - \Delta_{\tau}^{-1}(-e^{At}P_{\lambda}b_{f})$$

$$= e^{t\lambda}P_{\lambda}w + e^{t\lambda}\sum_{j=1}^{m-1}\frac{t^{j}}{j!}(A - \lambda E)^{j}P_{\lambda}w$$

$$+ \frac{e^{\lambda t}}{\tau}\sum_{j=0}^{m-1}\frac{t^{j+1}}{(j+1)!}(A - \lambda E)^{j}Y_{\lambda}(A)P_{\lambda}b_{f} - d(t)$$

$$= e^{t\lambda}P_{\lambda}w + \frac{e^{t\lambda}}{\tau}\sum_{j=0}^{m-1}\frac{t^{j+1}}{(j+1)!}(A - \lambda E)^{j}(\tau(A - \lambda E)P_{\lambda}w + Y_{\lambda}(A)P_{\lambda}b_{f}) - d(t).$$

Therefore

$$P_{\lambda}x(t) = \frac{e^{t\lambda}}{\tau} \sum_{j=0}^{m-1} \frac{t^{j+1}}{(j+1)!} (A - \lambda E)^{j} (\tau(A - \lambda E)P_{\lambda}w + Y_{\lambda}(A)P_{\lambda}b_{f})$$
$$+e^{t\lambda}P_{\lambda}w + v_{\lambda}(t,b_{f}).$$

We note that the periodic constant d(t) is canceled. \square

Notice that from this result we can easily obtain asymptotic behaviors of solutions of Equation (1), for details, refer to [2].

Example We will explain Theorem 2 and Theorem 3 through a simple one dimensional linear differential equation

$$\frac{dx}{dt} = ax(t) + f(t), \quad x(0) = w \in \mathbb{C}, \tag{16}$$

where $a \in \mathbb{C}$ and f is a continuous τ -periodic scalar function. Then (13) is reduced to

$$\Delta_{\tau} z(t) := z(t+\tau) - z(t) = -e^{at} b_f.$$

Using Theorem 2 with $B_0 = 1$, we have

$$z(t) := \Delta_{\tau}^{-1}(-e^{at}b_f) = \begin{cases} \frac{e^{at}}{1 - e^{a\tau}}b_f, & (e^{a\tau} \neq 1) \\ -\frac{e^{at}}{\tau}tb_f, & (e^{a\tau} = 1). \end{cases}$$
(17)

Therefore, by Theorem 3 the solution x(t) of Equation (16) is expressed as follows.

1) If $e^{a\tau} \neq 1$, then

$$x(t;0,w) = e^{at} \left(w - \frac{1}{1 - e^{a\tau}} b_f \right) + u(t,b_f),$$

where

$$u(t,b_f)=e^{at}rac{1}{1-e^{a au}}b_f+\int_0^t e^{a(t-s)}f(s)ds$$

is a τ -periodic solution of Equation (16).

2) If $e^{a\tau} = 1$, then

$$x(t;0,w) = \frac{e^{at}}{\tau}tb_f + e^{at}w + v(t,b_f),$$

where

$$v(t,b_f) = -e^{at}\frac{t}{\tau}b_f + \int_0^t e^{a(t-s)}f(s)ds$$

is a τ -periodic function, however, which is not necessary a solution of Equation (16).

References

- [1] J. Kato, T. Naito and J.S. Shin, Bounded solutions and periodic solutions to linear differential equations in Banach spaces, Vietnam J. of Math. (Proceeding in DEAA), 30 (2002) 561-575.
- [2] J. Kato, T. Naito and J.S. Shin, A characterization of solutions in linear differential equations with periodic forcing functions, J. Difference Equations and Applications, 11 (2005) 1-19.
- [3] K. S. Miller, An Introduction to the Calculus of Finite Differences and Difference Equations, Dover Publications, New York, 1960.