

Representations of the Cuntz algebra \mathcal{O}_3 arising from real cubic transformations

Katsunori Kawamura *
Research Institute for Mathematical Sciences
Kyoto University, Kyoto 606-8502, Japan

Abstract

We construct representations of the Cuntz algebra \mathcal{O}_3 from real cubic transformations on closed intervals. By intertwining relations of transformations, we have those of operators of representations of \mathcal{O}_2 and \mathcal{O}_3 . By these relations, we show that such representations are unitarily equivalent to barycentric representations.

1 Introduction

We study representations of the Cuntz algebra \mathcal{O}_N naturally arising from dynamical systems with non injective transformations. In [6], we show that a representation of \mathcal{O}_2 arising from a real quadratic transformation $x^2 - 2$ on a closed interval $[-2, 2]$ is equivalent to one of GP representation in [5] by using intertwining relations of operators corresponding to semiconjugacies of dynamical systems. This means that representations of \mathcal{O}_N arising from non linear transformations are not always hard to treat.

In this paper, we consider a real cubic transformation T_3 on a closed interval $[-1, 1]$ defined by

$$T_3(x) \equiv 4x^3 - 3x \quad (x \in [-1, 1]). \quad (1.1)$$

Put a representation $(L_2[-1, 1], \pi_t)$ of \mathcal{O}_3 by

$$(\pi_t(s_i)\phi)(x) \equiv m_i(x)\phi(T_3(x)) \quad (i = 1, 2, 3) \quad (1.2)$$

for $\phi \in L_2[-1, 1]$ where $m_i(x) \equiv \chi_{D_i}(x)\sqrt{|12x^2 - 3|}$, χ_{D_i} the characteristic function on D_i , $i = 1, 2, 3$, $D_1 = [-1, -1/2]$, $D_2 = [-1/2, 1/2]$, $D_3 = [1/2, 1]$ and s_1, s_2, s_3 are generators of \mathcal{O}_3 . Put $\mathbf{N} = \{1, 2, 3, \dots\}$.

*e-mail : kawamura@kurims.kyoto-u.ac.jp.

Theorem 1.1 *The representation $(L_2[-1, 1], \pi_t)$ in (1.2) is unitarily equivalent to the following representation $(l_2(\mathbf{N}), \pi')$ of \mathcal{O}_3 :*

$$\frac{1}{\sqrt{3}} \sum_{i=1}^3 e^{2\pi\sqrt{-1}\tau_{ij}} \pi'(s_j) e_n \equiv e_{3(n-1)+i} \quad (i = 1, 2, 3, n \in \mathbf{N})$$

where

$$\tau_{ij} \equiv \frac{(i-1)(j-1)}{3} \quad (i, j = 1, 2, 3). \quad (1.3)$$

Specially, $(L_2[-1, 1], \pi_t)$ is irreducible.

$$\Omega(x) = \frac{1}{\sqrt{\pi}} \frac{1}{(1-x^2)^{1/4}}$$

is unique positive normalized eigen function of $\pi_t(s_1) + \pi_t(s_2) + \pi_t(s_3)$ with eigen value $\sqrt{3}$.

In § 2, we review the representation theory of the Cuntz algebra. In § 3, we show intertwining relations between cubic transformations and piecewise linear transformations by the cosine function. In § 4, we show Theorem 1.1 by using of the intertwining relation derived in § 3. In § 5, we compare these results with those in [6] and discuss about remaining problems.

2 Representations of \mathcal{O}_N

For $N \geq 2$, let \mathcal{O}_N be the Cuntz algebra([3]), that is, it is a C^* -algebra which is universally generated by generators s_1, \dots, s_N satisfying

$$s_i^* s_j = \delta_{ij} I \quad (i, j = 1, \dots, N), \quad \sum_{i=1}^N s_i s_i^* = I. \quad (2.1)$$

In this paper, any representation means a unital $*$ -representation. By simplicity and uniqueness of \mathcal{O}_N , it is sufficient to define operators S_1, \dots, S_N on an infinite dimensional Hilbert space which satisfy (2.1) in order to construct a representation of \mathcal{O}_N .

2.1 The barycentric representation of \mathcal{O}_N

In order to characterize the representation in (1.2), we introduce a representation of \mathcal{O}_N which is defined by an eigen equation.

Definition 2.1 $(\mathcal{H}, \pi, \Omega)$ is a barycentric representation of \mathcal{O}_N if (\mathcal{H}, π) is a cyclic representation of \mathcal{O}_N with a cyclic vector Ω which satisfies the following condition:

$$\frac{1}{\sqrt{N}}\pi(s_1 + \cdots + s_N)\Omega = \Omega.$$

Proposition 2.2 Let $(\mathcal{H}, \pi, \Omega)$ be a barycentric representation of \mathcal{O}_N .

- (i) The representation (\mathcal{H}, π) exists uniquely up to unitary equivalences.
- (ii) (\mathcal{H}, π) is irreducible.

Proof. See Lemma A.3 (i),(ii),(iii). ■

The barycentric representation of \mathcal{O}_N is the GP representation by the parameter $(\frac{1}{\sqrt{N}}, \dots, \frac{1}{\sqrt{N}})$ ([5]). The naming of “barycentric” is that a representation of \mathcal{O}_N on a simplex Δ_{N-1} which is equivalent to the barycentric representation, is naturally defined by using the barycenter of Δ_{N-1} ([8]).

Proposition 2.3 The representation $(l_2(\mathbf{N}), \pi', e_1)$ in Theorem 1.1 is the barycentric representation of \mathcal{O}_3 .

Proof. See Corollary A.4. ■

2.2 Isometries arising from transformations on measure spaces

We give a method of construction of representation of \mathcal{O}_N from a branching function system over a measure space ([2, 6, 7]). We introduce an easy method to construct partial isometries from maps on a measure space.

Let (X, μ) be a measure space and $Y \subset X$ a measurable subset of X .

Definition 2.4 (i) $RN(Y, X)$ is the set of measurable maps on Y defined by

$$RN(Y, X) \equiv \left\{ f : Y \rightarrow X \mid \begin{array}{l} f \text{ is injective and} \\ \text{there exists } \Phi_f \text{ and } \Phi_f > 0 \text{ a.e. } Y \end{array} \right\}$$

where Φ_f is the Radon-Nikodým derivative of $\mu \circ f$ with respect to μ on Y .

- (ii) $RN_{loc}(X) \equiv \bigcup_{Y \subset X} RN(Y, X)$ where Y is taken from all measurable subsets of X .

For $f, g \in RN_{loc}(X)$, we denote the domain and the range of f by $D(f)$ and $R(f)$, respectively and $(f \circ g)(x) \equiv f(g(x))$ when $D(f) \supset R(g)$.

Lemma 2.5 (i) *If $f \in RN_{loc}(X)$, then $f^{-1} \in RN_{loc}(X)$.*

(ii) *For $f, g \in RN_{loc}(X)$, $f \circ g \in RN_{loc}(X)$ when $D(f) \supset R(g)$.*

(iii) *For $f, g \in RN_{loc}(X)$, $\Phi_{f \circ g} = \Phi_g \cdot ((\Phi_f) \circ g)$.*

Proof. These are easily checked by direct computation and property of Radon-Nikodým derivative. ■

Note that $RN_{loc}(X)$ is a groupoid by Lemma 2.5 (ii).

Definition 2.6 *For $f \in RN_{loc}(X)$, define an operator $S(f) : L_2(X, \mu) \rightarrow L_2(X, \mu)$ by*

$$(S(f)\phi)(x) \equiv \begin{cases} \{\Phi_f(f^{-1}(x))\}^{-1/2} \phi(f^{-1}(x)) & (\text{when } x \in R(f)), \\ 0 & (\text{otherwise}) \end{cases} \quad (2.2)$$

for $\phi \in L_2(X, \mu)$ and $x \in X$.

We simply denote

$$S(f) = \mathcal{L}_f M_{\Phi_f^{-1/2}} \quad (2.3)$$

where M_g is the multiplication operator of $g \in L_\infty(X, \mu)$ and \mathcal{L}_f is defined by

$$(\mathcal{L}_f \phi)(x) \equiv \chi_{R(f)}(x) \phi(f^{-1}(x)) \quad (x \in X)$$

and χ_Y is the characteristic function on $Y \subset X$.

Lemma 2.7 (i) *For $f \in RN_{loc}(X)$, $S(f)$ is a partial isometry on $L_2(X, \mu)$ with the initial projection $M_{\chi_{D(f)}}$ and the range projection $M_{\chi_{R(f)}}$.*

(ii) *For $f \in RN_{loc}(X)$, $S(f)^* = S(f^{-1})$.*

(iii) *$S(id_Y) = M_{\chi_Y}$.*

(iv) *For $f, g \in RN_{loc}(X)$, $\mathcal{L}_f \mathcal{L}_g = \mathcal{L}_{f \circ g}$ when $D(f) \supset R(g)$.*

(v) *For $f \in RN_{loc}(X)$ and $g \in L_\infty(X)$, $\mathcal{L}_f M_g = M_{\chi_{R(f)}} M_{g \circ f^{-1}} \mathcal{L}_f$.*

Proof. (i), (iv) and (v) follow by simple computation. (ii) Since $R(f^{-1}) = D(f)$, $D(f^{-1}) = R(f)$ and the property of Radon-Nikodým derivative, it follows. (iii) Since $\Phi_{id_Y} = \chi_Y$, it follows. \blacksquare

Let $\text{PIso}(L_2(X, \mu))$ be the groupoid of partial isometries on $L_2(X, \mu)$ by usual product of operators.

Lemma 2.8 *A map $S : RN_{loc}(X) \rightarrow \text{PIso}(L_2(X, \mu))$ defined in (2.2) is a groupoid homomorphism, that is,*

$$S(f)S(g) = S(f \circ g) \quad (2.4)$$

when $f, g \in RN_{loc}(X)$ and $D(f) \supset D(g)$.

Proof. By Lemma 2.5 (iii) and Lemma 2.7 (iv),(v),

$$S(f)S(g) = \mathcal{L}_f M_{\Phi_f^{-1/2}} \mathcal{L}_g M_{\Phi_g^{-1/2}} = \mathcal{L}_{f \circ g} M_{(\Phi_f \circ \Phi_g)^{-1/2}} M_{\Phi_g^{-1/2}} = S(f \circ g).$$

Remark that $f \circ g$ in rhs of (2.4) is only the composition of two transformations f and g but not special product of them. By Lemma 2.8, we see that a map S realizes the iteration of transformations on a measure space as the product of operators on a Hilbert space naturally.

The notion of branching function system was introduced in [2] in order to construct a representation of \mathcal{O}_N from a family of maps.

Let $N \geq 2$.

Definition 2.9 (i) $f = \{f_i\}_{i=1}^N$ is a branching function system over (X, μ) if $f_i \in RN_{loc}(X)$ and $D(f_i) = X$, put $R_i \equiv R(f_i)$, then $\mu(R_i \cap R_j) = 0$, $i \neq j$ and $\mu\left(X \setminus \bigcup_{i=1}^N R_i\right) = 0$.

(ii) F is the coding map of a branching function system $f = \{f_i\}_{i=1}^N$ on (X, μ) if $(F \circ f_i)(x) = x$ for a.e. $x \in X$ and $i = 1, \dots, N$.

Proposition 2.10 *For a branching function system $f = \{f_i\}_{i=1}^N$ on (X, μ) ,*

$$\pi_f(s_i) \equiv S(f_i) \quad (i = 1, \dots, N),$$

defines a representation $(L_2(X, \mu), \pi_f)$ of \mathcal{O}_N .

Proof. It is straightforward to show that $S(f_1), \dots, S(f_N)$ satisfy (2.1) by Lemma 2.7, Lemma 2.8 and Definition 2.9. \blacksquare

We show several examples of branching function system associated with 1-dimensionanl dynamical systems in [8].

Lemma 2.11 *Let (Y, ν) be another measure space. Assume that φ from X to Y is a measurable bijection a.e. X and Y and its Radon-Nikodým derivative Φ_φ is positive a.e. X .*

- (i) *An operator $U_\varphi : L_2(X, \mu) \rightarrow L_2(Y, \nu)$ defined by $U_\varphi \equiv \mathcal{L}_\varphi M_{\Phi_\varphi^{1/2}}$ is a unitary.*
- (ii) *For $f \in RN_{loc}(X)$, $\varphi \circ f \circ \varphi^{-1} \in RN_{loc}(Y)$ and $U_\varphi S(f) U_\varphi^* = S(\varphi \circ f \circ \varphi^{-1})$.*
- (iii) *If $f = \{f_i\}_{i=1}^N$ is a branching function system over (X, μ) , then $\{\varphi \circ f_i \circ \varphi^{-1}\}_{i=1}^N$ is a branching function system over (Y, ν) , too.*

Proof. These are easily proved by direct computation. ■

Proposition 2.12 *Let $f = \{f_i\}_{i=1}^N$ and $g = \{g_i\}_{i=1}^N$ be branching function systems over measure spaces (X, μ) and (Y, ν) , respectively. Assume that there is a map φ from X to Y which satisfies the assumption in Lemma 2.11 and map identities $g_i = \varphi \circ f_i \circ \varphi^{-1}$ for $i = 1, \dots, N$ hold. Then $(L_2(X, \mu), \pi_f)$ and $(L_2(Y, \nu), \pi_g)$ are unitarily equivalent.*

Proof. By Lemma 2.11 (ii), we can show $S(g_i) = U_\varphi S(f_i) U_\varphi^*$ for $i = 1, \dots, N$. These relations induce a unitary equivalence between π_f and π_g in Proposition 2.10 immediately. ■

3 The cosine function as intertwining map for cubic transformations

In order to show an intertwining relation of T_3 in (1.1), we prepare other transformations.

On a closed interval $[-1, 1]$, we consider transformations $V_3, C : [-1, 1] \rightarrow [-1, 1]$ by

$$V_3(x) \equiv \begin{cases} 3x + 2 & \left(-1 \leq x \leq -\frac{1}{3}\right), \\ -3x & \left(-\frac{1}{3} \leq x \leq \frac{1}{3}\right), \\ 3x - 2 & \left(\frac{1}{3} \leq x \leq 1\right), \end{cases} \quad C(x) \equiv \cos \pi x.$$

We denote $T \equiv T_3$ and $V \equiv V_3$ here. Then the following intertwining relation holds:

$$C \circ V = T \circ C. \quad (3.1)$$

Remark that C is not injective. We can not say that T and V are conjugate. C is a semiconjugacy between V and T ([4]).

We divide the interval $[-1, 1]$ by the following families of subintervals of $[-1, 1]$:

$$\left(D'_i\right)_{i=1}^3 \equiv \left(\left[-1, -\frac{1}{3}\right], \left[-\frac{1}{3}, \frac{1}{3}\right], \left[\frac{1}{3}, 1\right] \right),$$

$$(D_i)_{i=1}^3 \equiv \left(\left[-1, -\frac{1}{2}\right], \left[-\frac{1}{2}, \frac{1}{2}\right], \left[\frac{1}{2}, 1\right] \right),$$

$$\left(R'_i\right)_{i=1}^6 \equiv \left(\left[-1, -\frac{2}{3}\right], \left[-\frac{2}{3}, -\frac{1}{3}\right], \left[-\frac{1}{3}, 0\right], \left[0, \frac{1}{3}\right], \left[\frac{1}{3}, \frac{2}{3}\right], \left[\frac{2}{3}, 1\right] \right),$$

$$(R_i)_{i=1}^6 \equiv \left(\left[-1, -\frac{\sqrt{3}}{2}\right], \left[-\frac{\sqrt{3}}{2}, -\frac{1}{2}\right], \left[-\frac{1}{2}, 0\right], \left[0, \frac{1}{2}\right], \left[\frac{1}{2}, \frac{\sqrt{3}}{2}\right], \left[\frac{\sqrt{3}}{2}, 1\right] \right),$$

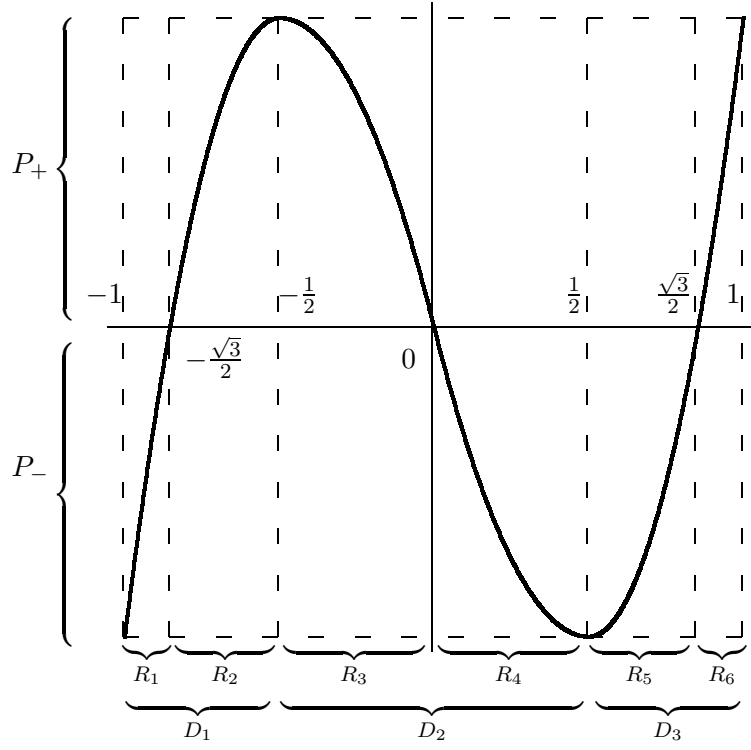
$$(P_-, P_+) \equiv ([-1, 0], [0, 1]).$$

The relations between these divisions are followings:

$$\left(D'_i\right)_{i=1}^3 = \left(R'_1 \cup R'_2, R'_3 \cup R'_4, R'_5 \cup R'_6\right),$$

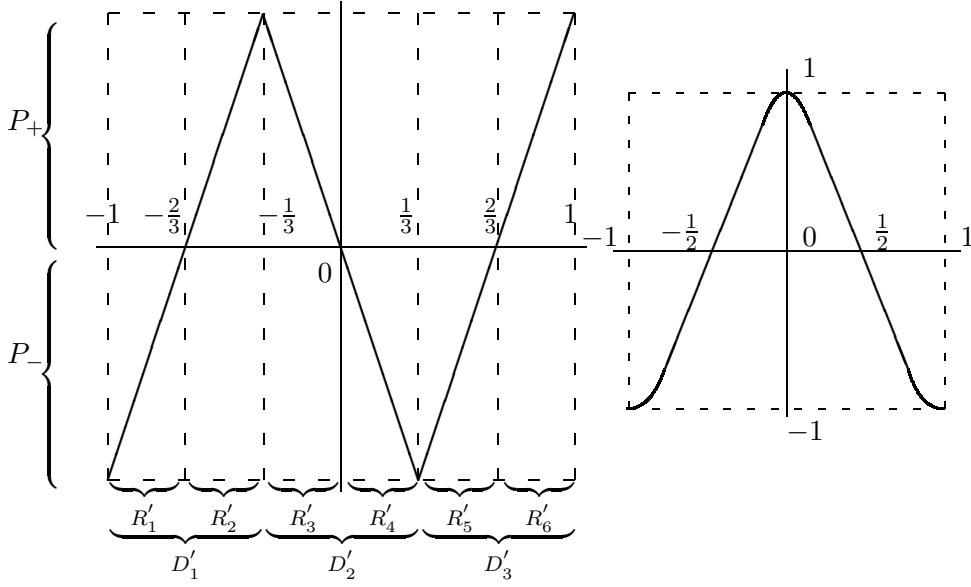
$$(P_-, P_+) = \left(R'_1 \cup R'_2 \cup R'_3, R'_4 \cup R'_5 \cup R'_6\right).$$

$T_3(x)$



$V_3(x)$

$C(x)$



The images of transformations T, V, C on these divisions are followings:

$$\begin{aligned} \left(T(R_i)\right)_{i=1}^6 &= (P_-, P_+, P_+, P_-, P_-, P_+), \\ \left(C(R'_i)\right)_{i=1}^6 &= (D_1, D_2, D_3, D_3, D_2, D_1), \\ \left(V(R'_i)\right)_{i=1}^6 &= (P_+, P_-, P_-, P_+, P_+, P_-). \end{aligned}$$

For these divisions, we have the following six relations by (3.1):

$$\begin{aligned} C|_{P_+} \circ V|_{D'_1} &= T|_{D_1} \circ C|_{P_-}, & C|_{P_+} \circ V|_{D'_2} &= T|_{D_3} \circ C|_{P_+}, \\ C|_{P_-} \circ V|_{D'_1} &= T|_{D_2} \circ C|_{P_-}, & C|_{P_+} \circ V|_{D'_3} &= T|_{D_2} \circ C|_{P_+}, \\ C|_{P_-} \circ V|_{D'_2} &= T|_{D_3} \circ C|_{P_-}, & C|_{P_-} \circ V|_{D'_3} &= T|_{D_1} \circ C|_{P_+}. \end{aligned}$$

Note that restrictions $C|_{P_\pm}$, $V|_{D'_i}$ and $T|_{D_i}$, $i = 1, 2, 3$ are injective. Let

$$c_\pm \equiv (C|_{P_\pm})^{-1}, \quad v_i \equiv (V|_{D'_i})^{-1}, \quad t_i \equiv (T|_{D_i})^{-1} \quad (3.2)$$

for $i = 1, 2, 3$. Then

$$\begin{aligned} c_- \circ t_1 &= v_1 \circ c_+, & c_+ \circ t_3 &= v_2 \circ c_+, \\ c_- \circ t_2 &= v_1 \circ c_-, & c_+ \circ t_2 &= v_3 \circ c_+, \\ c_- \circ t_3 &= v_2 \circ c_-, & c_+ \circ t_1 &= v_3 \circ c_-. \end{aligned}$$

From these, we have the following relations:

$$\begin{aligned} v_1 &= \begin{cases} c_- \circ t_1 \circ c_+^{-1} & (\text{on } P_+), \\ c_- \circ t_2 \circ c_-^{-1} & (\text{on } P_-), \end{cases} & v_2 &= \begin{cases} c_+ \circ t_3 \circ c_+^{-1} & (\text{on } P_+), \\ c_- \circ t_3 \circ c_-^{-1} & (\text{on } P_-), \end{cases} \\ v_3 &= \begin{cases} c_+ \circ t_2 \circ c_+^{-1} & (\text{on } P_+), \\ c_+ \circ t_1 \circ c_-^{-1} & (\text{on } P_-). \end{cases} \end{aligned} \quad (3.3)$$

4 Representations of \mathcal{O}_3 arising from dynamical systems

We show intertwining relations of operators in this section.

4.1 Intertwining relations of operators

By (3.2), $c \equiv \{c_+, c_-\}$ is a branching function system over $[-1, 1]$ for $N = 2$ and both $t \equiv \{t_1, t_2, t_3\}$ and $v \equiv \{v_1, v_2, v_3\}$ are those over $[-1, 1]$ for $N = 3$. Hence operators $\{S(c_\pm)\}$ gives a representation of \mathcal{O}_2 on $L_2[-1, 1]$ and $\{S(t_i)\}_{i=1}^3, \{S(v_i)\}_{i=1}^3$ give those of \mathcal{O}_3 on $L_2[-1, 1]$, too.

Then $\pi_t(s_i) = S(t_i)$ for $i = 1, 2, 3$ where π_t is the representation of \mathcal{O}_3 on $L_2[-1, 1]$ in (1.2) and

$$\begin{cases} (S(v_1)\phi)(x) = 3^{1/2}\chi_{[-1, -1/3]}(x)\phi(3x + 2), \\ (S(v_2)\phi)(x) = 3^{1/2}\chi_{[-1/3, 1/3]}(x)\phi(-3x), \\ (S(v_3)\phi)(x) = 3^{1/2}\chi_{[1/3, 1]}(x)\phi(3x - 2), \end{cases} \quad (4.1)$$

$$\begin{cases} (S(c_+)\phi)(x) = \sqrt{\pi \sin \pi|x|}\chi_{[0, 1]}(x)\phi(\cos \pi x), \\ (S(c_-)\phi)(x) = \sqrt{\pi \sin \pi|x|}\chi_{[-1, 0]}(x)\phi(\cos \pi x), \end{cases}$$

for $\phi \in L_2[-1, 1]$ and $x \in [-1, 1]$.

By Lemma 2.8 and (3.3), we have the following operator relations:

$$\begin{cases} S(v_1) = S(c_-)S(t_1)S(c_+)^* + S(c_-)S(t_2)S(c_-)^*, \\ S(v_2) = S(c_+)S(t_3)S(c_+)^* + S(c_-)S(t_3)S(c_-)^*, \\ S(v_3) = S(c_+)S(t_2)S(c_+)^* + S(c_+)S(t_1)S(c_-)^*. \end{cases} \quad (4.2)$$

Note that (4.2) is a relation of representations π_c, π_t, π_v and it is derived from (3.1) essentially. We explain these relations in § 5. Let

$$A_1 \equiv \sum_{i=1}^3 S(v_i), \quad A_2 \equiv \sum_{i=1}^3 S(t_i), \quad B \equiv S(c_+) + S(c_-).$$

Then we have the following intertwining relations of operators:

$$B^* A_1 = A_2 B^*. \quad (4.3)$$

Let $\mathbf{1}$ be the constant function of $[-1, 1]$ with value 1. Note $A_1\mathbf{1} = \sqrt{3}\mathbf{1}$. Let $\Omega_0 \equiv B^*\mathbf{1}$. By (4.3),

$$A_2\Omega_0 = A_2B^*\mathbf{1} = B^*A_1\mathbf{1} = \sqrt{3}B^*\mathbf{1} = \sqrt{3}\Omega_0.$$

Hence Ω_0 is an eigen function of A_2 with eigen value $\sqrt{3}$.

Lemma 4.1 *The operator $S(t_1)+S(t_2)+S(t_3)$ has an eigen vector $(S(c_+)^*+S(c_-)^*)\mathbf{1}$ with eigen value $\sqrt{3}$.*

Let Ω be the normalization of Ω_0 . Then

$$\Omega(x) = \frac{1}{\sqrt{\pi}} \frac{1}{(1-x^2)^{1/4}} \quad (x \in [-1, 1]).$$

Theorem 4.2 *$(L_2[-1, 1], \pi_t, \Omega)$ is the barycentric representation.*

Proof. By Lemma 4.1, it is sufficient to show the cyclicity of $(L_2[-1, 1], \pi_t)$. By (1.2), $\pi_t(s_I)\Omega = M_{\chi_{D_I}} 3^{k/2}\Omega$. Hence $K \equiv \{\chi_{D_I} \cdot \Omega : I \in \{1, 2, 3\}^k, k \geq 1\} \subset \pi_t(\mathcal{O}_3)\Omega$. Since $\Omega(x) > 0$ for any $x \in (0, 1)$, $L_2[-1, 1] = \overline{\text{Lin} \langle K \rangle} \subset \pi_t(\mathcal{O}_3)\Omega$. Therefore $(L_2[-1, 1], \pi_t)$ is cyclic. Hence $(L_2[-1, 1], \pi_t, \Omega)$ is the barycentric representation. \blacksquare

Proof of Theorem 1.1. By Proposition 2.3, $(l_2(\mathbf{N}), \pi')$ in Theorem 1.1 is equivalent to the barycentric representation. By Proposition 2.2 (i) and Theorem 4.2, the statement in Theorem 1.1 follows. \blacksquare

4.2 Generalization

We generalize our result slightly. Assume $a, b \in \mathbf{R}$, $a < b$. Let

$$\varphi_{(a,b)} : [-1, 1] \rightarrow [a, b]; \quad \varphi_{(a,b)}(x) \equiv \frac{b-a}{2}x + \frac{a+b}{2}.$$

Put $T_3^{(a,b)} \equiv \varphi_{(a,b)} \circ T_3 \circ (\varphi_{(a,b)})^{-1}$, $V_3^{(a,b)} \equiv \varphi_{(a,b)} \circ V_3 \circ (\varphi_{(a,b)})^{-1}$. Then

$$V_3^{(a,b)}(x) = 2\alpha \left| \frac{1}{\alpha}x - \beta \right| + \alpha(\beta - 1),$$

$$T_3^{(a,b)}(x) = \frac{4}{\alpha^2}x^3 - \frac{12\beta}{\alpha}x^2 + 3(4\beta^2 - 1)x - 4\alpha\beta(\beta^2 - 1) \quad (4.4)$$

where $\alpha = \frac{b-a}{2}$, $\beta = \frac{b+a}{b-a}$. For example, $T_3^{(-1,1)}(x) = T_3(x)$, $T_3^{(-2,2)}(x) = x^3 - 3x$ and $T_3^{(0,1)}(x) = x(4x - 3)^2$.

By Proposition 2.12, for any $a, b \in \mathbf{R}$, $a < b$, the representation $(L_2[a, b], \pi_t^{(a,b)})$ of \mathcal{O}_3 arising from $T_3^{(a,b)}$ is given by

$$(\pi_t^{(a,b)}(s_i)\phi)(x) = m_i^{(a,b)}(x)\phi(T^{(a,b)}(x)) \quad (\phi \in L_2[a, b])$$

where

$$m_i^{(a,b)}(x) \equiv \chi_{D_i^{(a,b)}}(x) \cdot \sqrt{\left| \frac{12}{\alpha^2}x^2 - \frac{24\beta}{\alpha}x + 3(4\beta^2 - 1) \right|}$$

and $D_1^{(a,b)} = [a, \frac{3a+b}{4}]$, $D_2^{(a,b)} = [\frac{3a+b}{4}, \frac{a+3b}{4}]$, $D_3^{(a,b)} = [\frac{a+3b}{4}, b]$. Put

$$\Omega^{(a,b)}(x) \equiv \frac{1}{\sqrt{\pi}} \frac{1}{\{(3b - a - 2x)(b - 3a + 2x)\}^{1/4}}.$$

Then $(L_2[a, b], \pi_t^{(a,b)}, \Omega^{(a,b)})$ is the barycentric representation of \mathcal{O}_3 .

5 Discussion

Our results in this paper is similar to those in [6]. Although, the intertwining relations (4.2) are more complicated than the former. Their abstractions are followings: Put $\text{Rep}(\mathcal{O}_N, \mathcal{H})$ the set of all unital $*$ -representations of \mathcal{O}_N on a Hilbert space \mathcal{H} for $N \geq 2$. By [6], there is a map

$$\zeta^{(2)} : \text{Rep}(\mathcal{O}_2, \mathcal{H}) \times \text{Rep}(\mathcal{O}_2, \mathcal{H}) \rightarrow \text{Rep}(\mathcal{O}_2, \mathcal{H}); \quad (\pi', \pi) \mapsto \zeta_{\pi'}^{(2)}(\pi)$$

defined by

$$\begin{cases} \zeta_{\pi'}^{(2)}(\pi)(s_1) = \pi'(s_1)\pi(s_1)\pi'(s_1)^* + \pi'(s_1)\pi(s_2)\pi'(s_2)^*, \\ \zeta_{\pi'}^{(2)}(\pi)(s_2) = \pi'(s_2)\pi(s_1)\pi'(s_1)^* + \pi'(s_2)\pi(s_2)\pi'(s_2)^*. \end{cases}$$

On the other hand, by (4.2), there is a map

$$\zeta^{(3)} : \text{Rep}(\mathcal{O}_3, \mathcal{H}) \times \text{Rep}(\mathcal{O}_2, \mathcal{H}) \rightarrow \text{Rep}(\mathcal{O}_3, \mathcal{H}); \quad (\pi', \pi) \mapsto \zeta_{\pi'}^{(3)}(\pi)$$

defined by

$$\begin{cases} \zeta_{\pi'}^{(3)}(\pi)(t_1) = \pi'(s_1)\pi(t_1)\pi'(s_2)^* + \pi'(s_1)\pi(t_2)\pi'(s_1)^*, \\ \zeta_{\pi'}^{(3)}(\pi)(t_2) = \pi'(s_2)\pi(t_3)\pi'(s_2)^* + \pi'(s_1)\pi(t_3)\pi'(s_1)^*, \\ \zeta_{\pi'}^{(3)}(\pi)(t_3) = \pi'(s_2)\pi(t_2)\pi'(s_2)^* + \pi'(s_2)\pi(t_1)\pi'(s_1)^* \end{cases}$$

where $\{s_1, s_2\}$ and $\{t_1, t_2, t_3\}$ are generators of \mathcal{O}_2 and \mathcal{O}_3 , respectively. Apparently, $\zeta^{(3)}$ is not a simple generalization of $\zeta^{(2)}$. In fact, we can not guess the case of $N = 4, 5, 6, \dots$ for \mathcal{O}_N . Note that $\zeta^{(2)}$ and $\zeta^{(3)}$ are derived by intertwining relations of a quadratic transformation and a cubic transformation on a closed interval, respectively.

We can not yet make clear mechanism when “nice” intertwining relations occur. For example, we can not make representation theory from general cubic transformations $a_1x^3 + a_2x^2 + a_3x + a_4$ except $T_3^{(a,b)}$ in (4.4).

Acknowledgement: We would like to thank Cho Chien-Hong for the idea of the map $T_3(x) = 4x^3 - 3x$ by improving our results in [6].

Appendix A Proof of Proposition 2.2

Results in Proposition 2.2 are included in [5]. The barycentric representation is a kind of GP representation in [5]. Here we show the proof of Proposition 2.2 for convenience.

A.1 The standard representation of \mathcal{O}_N

We introduce the standard representation of \mathcal{O}_N in order to prove Proposition 2.2. Let $l_2(\mathbf{N})$ be the Hilbert space with the canonical basis $\{e_n : n \in \mathbf{N}\}$, $\mathbf{N} = \{1, 2, 3, \dots\}$, and make the following representation $(l_2(\mathbf{N}), \pi_S)$ of the Cuntz algebra \mathcal{O}_N which is called *the standard representation of \mathcal{O}_N* ([1]):

$$\pi_S(s_i)e_n \equiv e_{N(n-1)+i} \quad (i = 1, \dots, N, n \in \mathbf{N}). \quad (\text{A.1})$$

From this, we have $\pi_S(s_i)^*e_{N(n-1)+j} = \delta_{ij}e_n$ for $i, j = 1, \dots, N$ and $n \in \mathbf{N}$. Note that this is a permutative representation of \mathcal{O}_N by [2]. By (A.1),

$$\pi_S(s_1)e_1 = e_1. \quad (\text{A.2})$$

Lemma A.1 $(l_2(\mathbf{N}), \pi_S)$ is irreducible.

Proof. By Theorem 2.7 in [2], this representation is irreducible. ■

Proposition A.2 For $N \geq 2$, put

$$\Lambda_N \equiv \{1, \dots, N\} \cup \bigcup_{k \geq 1} \left(\{1, \dots, N\}^k \times \{2, \dots, N\} \right).$$

Let (\mathcal{H}, π) be a representation of \mathcal{O}_N . Assume that

$$\text{there is a unit cyclic vector } \Omega \in \mathcal{H} \text{ such that } \pi(s_1)\Omega = \Omega. \quad (\text{A.3})$$

Then the followings hold:

- (i) $\{\pi(s_I)\Omega : I \in \Lambda_N\}$ is an orthonormal family in \mathcal{H} .
- (ii) The representation of \mathcal{O}_N which satisfies the condition (A.3) is unique up to unitary equivalences.

Proof. (i) We denote $|I|$ the length of $I \in \Lambda_N$. Put $I, J \in \Lambda_N$.

If $|I| = |J|$, then $\langle \pi(s_I)\Omega | \pi(s_J)\Omega \rangle = \delta_{I,J}$.

Assume that $|I| > |J|$. Put $I = (i_1, \dots, i_{k+l})$ and $J = (j_1, \dots, j_k)$, $k, l \geq 1$. Then

$$\langle \pi(s_I)\Omega | \pi(s_J)\Omega \rangle = \delta_{I_1, J} \langle \pi(s_{I_2})\Omega | \Omega \rangle = \delta_{I_1, J} \langle \pi(s_{I_2})\Omega | \pi(s_1^l)\Omega \rangle$$

where we use $\pi(s_1^l)\Omega = \Omega$ and $I_1 = (i_1, \dots, i_k)$ and $I_2 = (i_{k+1}, \dots, i_{k+l})$. By choice of I , $I_2 \neq (1, \dots, 1)$. Hence $\langle \pi(s_{I_2})\Omega | \pi(s_1^l)\Omega \rangle = \delta_{I_2, J_0} = 0$ where $J_0 \equiv \underbrace{(1, \dots, 1)}_l$. Therefore $\langle \pi(s_I)\Omega | \pi(s_J)\Omega \rangle = 0$. From this, we obtain

$\langle \pi(s_I)\Omega | \pi(s_J)\Omega \rangle = \delta_{I,J}$ for each $I, J \in \Lambda_N$.

(ii) Put $\{1, \dots, N\}^* \equiv \bigcup_{k \geq 0} \{1, \dots, N\}^k$. Let $(\mathcal{H}, \pi, \Omega)$ be a representation which satisfies (A.3). Because Ω is a cyclic vector, the linear span of $\{\pi(s_I s_J^*)\Omega : I, J \in \{1, \dots, N\}^*\}$ is dense in \mathcal{H} . Because $\pi(s_1)\Omega = \Omega$, $\{\pi(s_I s_J^*)\Omega : I, J \in \{1, \dots, N\}^*\} = \{\pi(s_I)\Omega : I \in \Lambda_N\}$. Hence $\{\pi(s_I)\Omega : I \in \Lambda_N\}$ is a complete orthonormal basis of \mathcal{H} by (i). In this way, we know that any representation which satisfies the condition (A.3) always has such basis. Therefore we have a natural unitary between any two such representations. The uniqueness is proved. \blacksquare

By this, we call the standard representation the representation which satisfies the condition (A.3), too.

A.2 The barycentric representation of \mathcal{O}_N

Recall Definition 2.1.

Lemma A.3 *Let $N \geq 2$.*

- (i) *The barycentric representation of \mathcal{O}_N exists.*

- (ii) *The barycentric representation of \mathcal{O}_N is unique up to unitary equivalences.*
- (iii) *The barycentric representation of \mathcal{O}_N is irreducible.*
- (iv) *The barycentric representation of \mathcal{O}_N is not equivalent to the standard representation.*
- (v) *The GNS representation of \mathcal{O}_N of the following state ρ*

$$\rho(s_J s_{J'}^*) = \frac{1}{N^{(|J|+|J'|)/2}} \quad (J, J' \in \{1, \dots, N\}^k, k \geq 0) \quad (\text{A.4})$$

is equivalent to the barycentric representation where we define $s_J = I$ when $J = \emptyset$. Specially, ρ is pure.

Proof. (i) Put $g = (g_{ij}) \in U(N)$ by

$$g_{ij} \equiv \frac{1}{\sqrt{N}} e^{2\pi\sqrt{-1}\tau_{ij}} \quad (i, j = 1, \dots, N). \quad (\text{A.5})$$

where $\tau_{ij} \equiv \frac{1}{N}(i-1)(j-1)$ for $i, j = 1, \dots, N$. Then specially, $g_{1j} = \frac{1}{\sqrt{N}}$ for $i = 1, \dots, N$. Let $(l_2(\mathbf{N}), \pi_S)$ be the standard representation of \mathcal{O}_N in (A.1). Put $\pi' \equiv \pi_S \circ \alpha_g$ where α is an action of $U(N)$ on \mathcal{O}_N defined by $\alpha_g(s_i) \equiv \sum_{j=1}^N g_{ji} s_j$ for $i = 1, \dots, N$.

Since $(l_2(\mathbf{N}), \pi_S)$ is an irreducible representation of \mathcal{O}_N , $(l_2(\mathbf{N}), \pi')$ is too. Specially, $(l_2(\mathbf{N}), \pi')$ is cyclic. Because $\alpha_{g^*}(s_1) = \frac{1}{\sqrt{N}}(s_1 + \dots + s_N)$,

$$\frac{1}{\sqrt{N}} \pi'(s_1 + \dots + s_N) e_1 = (\pi \circ \alpha_g)(\alpha_{g^*}(s_1)) e_1 = \pi_S(s_1) e_1 = e_1.$$

Hence $(l_2(\mathbf{N}), \pi', e_1)$ is a barycentric representation of \mathcal{O}_N .

(ii) Assume that $(\mathcal{H}_i, \pi_i, \Omega_i)$ is a barycentric representation of \mathcal{O}_N for $i = 1, 2$. By the similar argument in (i), $(\mathcal{H}_i, \pi_i \circ \alpha_{g^*})$ is the standard representation of \mathcal{O}_N with respect to $g \in U(N)$ in (i) for $i = 1, 2$. By uniqueness of the standard representation, $(\mathcal{H}_1, \pi_1 \circ \alpha_{g^*}) \sim (\mathcal{H}_2, \pi_2 \circ \alpha_{g^*})$. Hence $(\mathcal{H}_1, \pi_1) \sim (\mathcal{H}_2, \pi_2)$. Therefore the barycentric representation is unique up to unitary equivalences.

(iii) By (ii), the barycentric representation is equivalent to the action of the standard representation by suitable $g \in U(N)$. Since the standard representation is irreducible, the barycentric representation is too.

(iv) Let $(\mathcal{H}, \pi, \Omega)$ and $(\mathcal{H}', \pi', \Omega')$ be the standard representation and the barycentric representation of \mathcal{O}_N , respectively. Assume that (\mathcal{H}, π) and

(\mathcal{H}', π') are unitarily equivalent. Then we can identify $(\mathcal{H}', \pi') = (\mathcal{H}, \pi)$. By assumption,

$$\langle \Omega | \Omega' \rangle = \langle \pi(s_1)\Omega | \pi(N^{-1/2}(s_1 + \dots + s_N))\Omega' \rangle = N^{-1/2} \langle \Omega | \Omega' \rangle.$$

Hence $\langle \Omega | \Omega' \rangle = 0$. By cyclicity, there is $J = (j_1, \dots, j_k) \in \{1, \dots, N\}^k$, $k \geq 1$, such that $\langle \Omega | \pi(s_J)\Omega' \rangle \neq 0$. However,

$$\langle \Omega | \pi(s_J)\Omega' \rangle = \langle \pi(s_1)^k \Omega | \pi(s_J)\Omega' \rangle = \left(\prod_{l=1}^k \delta_{1, j_l} \right) \langle \Omega | \Omega' \rangle = 0.$$

This is contradiction. Therefore (\mathcal{H}, π) and (\mathcal{H}', π') are not unitarily equivalent.

(v) Let $(\mathcal{H}, \pi, \Omega)$ be the barycentric representation. Put $\rho(x) \equiv \langle \Omega | \pi(x)\Omega \rangle$ for $x \in \mathcal{O}_N$. Then ρ satisfies the condition (A.4). Since the uniqueness of the GNS-representation and (\mathcal{H}, π) is irreducible, the statement holds. ■

Corollary A.4 *The representation $(l_2(\mathbf{N}), \pi', e_1)$ of \mathcal{O}_3 in Theorem 1.1 is the barycentric representation.*

Proof. By the proof of Lemma A.3 (i), it follows when $N = 3$. ■

References

- [1] M.Abe and K.Kawamura, *Recursive Fermion System in Cuntz Algebra. I —Embeddings of Fermion Algebra into Cuntz Algebra —*, Comm. Math. Phys. **228** (2002) 85-101.
- [2] O. Bratteli and P. E. T. Jørgensen, *Iterated Function Systems and Permutation Representations of the Cuntz Algebra*, Memoirs Amer. Math. Soc. **39** No. 663 (1999).
- [3] J. Cuntz, *Simple C^* -algebras generated by isometries*, Comm. Math. Phys. **57** (1977), 173-185.
- [4] Robert L. Devaney, *A first course in chaotic dynamical systems: Theory and experiment*, Addison-Wesley Publishing Company (1997).

- [5] K.Kawamura, *Generalized permutative representation of Cuntz algebra. I —Generalization of cycle type—*, preprint RIMS-1380 (2002).
- [6] K.Kawamura, *Representations of the Cuntz algebra \mathcal{O}_2 arising from real quadratic transformations*, preprint RIMS-1396 (2003).
- [7] K.Kawamura and O.Suzuki, *Construction of orthonormal basis on self-similar sets by generalized permutative representations of the Cuntz algebras*, preprint RIMS-1408 (2003).
- [8] K.Kawamura, *Construction of representations of the Cuntz-Krieger algebra*, in preparation.