

Multidimensional Central and Local Limit Theorems  
for the Phase Separation Line in the  
Two-Dimensional Ising Model

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§1. The Two-Dimensional Ising Model

Let  $\mathbb{Z}^2$  be the square lattice and  $L$  be its dual lattice; *i.e.*

$$\mathbb{Z}^2 \equiv \{(x_1, x_2); x_1 \text{ and } x_2 \text{ are integers}\},$$

$$L \equiv \{(x_1, x_2) + (\frac{1}{2}, \frac{1}{2}); (x_1, x_2) \in \mathbb{Z}^2\}.$$

We consider  $\mathbb{Z}^2$  and  $L$  also as graphs. Let  $\Omega \equiv \{-1, +1\}^L$  be the space of possible spin configurations on  $L$ . For each positive integer  $N$ , we define a square box  $V_N$  by

$$V_N \equiv \{(x_1, x_2) + (\frac{1}{2}, \frac{1}{2}); (x_1, x_2) \in \mathbb{Z}^2, 0 \leq x_1 \leq N-1, -[\frac{N}{2}] \leq x_2 \leq [\frac{N-1}{2}]\},$$

where  $[u]$  is the largest integer smaller than  $u$ . For given

$\omega \in \Omega$  and  $\beta > 0$ , the finite Gibbs state on  $\Omega_N \equiv \{-1, +1\}^{V_N}$  with boundary condition  $\omega$ , inverse temperature  $\beta$  (in the absence of the exterior magnetic field) is given by

$$P_N^\omega(\eta) = Z_N(\omega)^{-1} \exp\{-\beta U_N(\eta; \omega)\} \quad \eta \in \Omega_N$$

where  $Z_N(\omega) \equiv \sum_{\eta' \in \Omega_N} \exp\{-\beta U_N(\eta'; \omega)\}$ , and

$$U_N(\eta; \omega) \equiv - \sum_{x, y \in V_N}^* \eta(x)\eta(y) - \sum_{\substack{x \in V_N \\ y \in \partial V_N}}^* \eta(x)\omega(y),$$

where the summation  $\sum_{x \in A, y \in B}^*$  is taken over all pairs  $(x, y)$

such that (i)  $x$  and  $y$  are nearest neighbours in  $L$ , and (ii)  $x \in A$  and  $y \in B$ . From now on, we fix the boundary condition  $\omega$  as

$$\omega\left((x_1, x_2) + \left(\frac{1}{2}, \frac{1}{2}\right)\right) = \begin{cases} +1 & \text{if } x_2 \geq 0 \\ -1 & \text{if } x_2 < 0, \end{cases}$$

and we write simply  $P_N$  instead of  $P_N^\omega$ .

## §2. The Phase Separation Line

Let us fix  $N > 0$  arbitrarily. Define

$$\tilde{V}_N \equiv \{(x_1, x_2) \in \mathbb{Z}^2; 0 \leq x_1 \leq N, -[\frac{N+1}{2}] \leq x_2 \leq [\frac{N+1}{2}]\}.$$

We can regard  $\tilde{V}_N$  as a subgraph of  $\mathbb{Z}^2$ . We call a segment of length one connecting two points which are nearest neighbours in  $\mathbb{Z}^2$  (or in  $L$ ) by "a bond in  $\mathbb{Z}^2$  (or in  $L$ )". Then for each  $\eta \in \Omega_N$ , we can define a subgraph  $C_N(\eta)$  of  $\tilde{V}_N$  in the following way. A bond in  $\mathbb{Z}^2$  belongs to  $C_N(\eta)$  if and only if it crosses a bond in  $L$  connecting two points  $x, y \in L$  such that  $\eta(x)\eta(y) = -1$  (or  $\eta(x)\omega(y) = -1$ ).  $C_N(\eta)$  consists of some connected components, and each vertex of  $C_N(\eta)$  belongs to two of four bonds of  $C_N(\eta)$  unless the vertex is  $A = (0, 0)$  or  $B = (N, 0)$ . It is easy to see that the connected component of  $C_N(\eta)$  containing  $A$  also contains  $B$ . We denote this component by  $\lambda_N(\eta)$ , and call it by "the phase separation line". Let  $\Lambda_N \equiv \{\lambda_N(\eta) : \eta \in \Omega_N\}$ , and  $S \equiv \{\gamma; \text{subgraph of } \mathbb{Z}^2 \text{ such that (i) the length } |\gamma| \text{ is finite, and (ii) each vertex of } \gamma \text{ belongs to two or four bonds of } \gamma\}$ .

*Theorem (Gallavotti)*

There exists a function  $\varphi : \mathcal{N} = \bigoplus_{n=0}^{\infty} S^n \rightarrow \mathbb{R}$  such that

(i)  $\varphi$  is symmetric, (ii)  $\sum_{\substack{\Gamma \in \mathcal{N} \\ \Gamma \ni 0}} |\varphi(\Gamma)| < +\infty$  and (iii) for each

$\lambda \in \Lambda_N$ ,

$$P_N(\{\eta \in \Omega_N; \lambda_N(\eta) = \lambda\}) = \exp\{-2\beta|\lambda| - \sum_{\substack{\Gamma \in \mathcal{N} \\ \Gamma \cap \lambda \neq \emptyset \\ \Gamma \subset V_N}} \varphi(\Gamma)\} / \sum_{\lambda' \in \Lambda_N} \exp\{-2\beta|\lambda'| - \sum_{\substack{\Gamma \in \mathcal{N} \\ \Gamma \cap \lambda' \neq \emptyset \\ \Gamma \subset \tilde{V}_N}} \varphi(\Gamma)\}.$$

We denote the left hand side of the above equality simply by  $P_N(\lambda)$ .

Thus we are given a sequence of probability spaces  $(\Lambda_N, P_N)$ .

To state our result, we need some notations. For each  $\lambda \in \Lambda_N$ ,

$0 \leq \ell \leq N$ , let  $\bar{X}_N(\ell)$  and  $\underline{X}_N(\ell)$  be  $\bar{X}_N(\ell) \equiv \max\{k; (k, \ell) \in \lambda\}$ ,

$\underline{X}_N(\ell) \equiv \min\{k; (k, \ell) \in \lambda\}$ .

*Theorem 1.*

Let an integer  $k \geq 1$ , and  $0 < t_1 < t_2 < \dots < t_k < 1$  and the numbers  $-\infty < T_j < T'_j < \infty$ ,  $j = 1, 2, \dots, k$  be given arbitrarily.

Let  $a_j^{(N)} \equiv [t_j \cdot N]$   $j = 1, 2, \dots, k$ , where  $[u]$  is the largest integer smaller than  $u$ . Then,

$$\begin{aligned} & \lim_{N \rightarrow \infty} P_N \left( \bigcap_{j=1}^k \left\{ T_j \leq \frac{\bar{X}_N(a_j^{(N)})}{\sigma\sqrt{N}} \leq T'_j \right\} \right) \\ &= \lim_{N \rightarrow \infty} P_N \left( \bigcap_{j=1}^k \left\{ T_j \leq \frac{\underline{X}_N(a_j^{(N)})}{\sigma\sqrt{N}} \leq \frac{\bar{X}_N(a_j^{(N)})}{\sigma\sqrt{N}} \leq T'_j \right\} \right) \\ &= P_0^1; 0^0 (X_{t_j} \in [T_j, T'_j], j = 1, 2, \dots, k) \end{aligned}$$

where  $\sigma = \sigma(\beta) > 0$  is a constant depending only on  $\beta > 0$  and  $(\{X_t\}_{0 \leq t \leq 1}, P_0^1; 0^0)$  is a one-dimensional Brownian Bridge such that  $P_0^1; 0^0(X_0 = X_1 = 0) = 1$ .

*Remark* Gallavotti has proved the above theorem in the case when  $k = 1$  in [2], and it is announced that Cammarota has proved it for general  $k \geq 1$ , but we have got nothing in print yet.

### §3. An Auxiliary Ensemble and the Central Limit Theorem

In order to prove Theorem 1, we need another sequence of probability spaces  $(\hat{\Lambda}_N, \hat{P}_N)$ . Let  $I_N \equiv \{(x_1, x_2) \in \mathbb{Z}^2; 0 \leq x_1 \leq N\}$ , and  $\hat{\Lambda}_N \equiv \{\text{connected subgraph } \lambda \text{ of } I_N \text{ such that (i) } |\lambda| < +\infty, \text{ (ii) } \lambda \ni A, \text{ (iii) there exists a point } B' \text{ in } \{(N, \ell): \ell \in \mathbb{Z}\}\}$

such that each vertex of  $\lambda$  belongs to two or four bonds of  $\lambda$  unless it is  $A$  or  $B'$ , and  $A$  and  $B'$  belong to one or three bonds of  $\lambda$ . For each  $\lambda \in \hat{\Lambda}_N$ , we define

$$\hat{P}_N(\lambda) \equiv \exp\{-2\beta|\lambda| - \sum_{\substack{\Gamma \in \mathcal{N} \\ \Gamma \subset I_N \\ \Gamma \cap \lambda \neq \emptyset}} \varphi(\Gamma)\} / \sum_{\tilde{\lambda} \in \hat{\Lambda}_N} \exp\{-2\beta|\tilde{\lambda}| - \sum_{\substack{\Gamma \in \mathcal{N} \\ \Gamma \subset I_N \\ \Gamma \cap \tilde{\lambda} \neq \emptyset}} \varphi(\Gamma)\}.$$

Then, Gallavotti has shown that for every  $A_N \subset \Lambda_N$ ,

$$P_N(A_N) = \hat{P}_N(A_N | \tilde{\Lambda}_N) + o\left(\frac{1}{N^\ell}\right) \quad \text{as } N \rightarrow \infty \quad \forall \ell \in \mathbb{N},$$

where  $\tilde{\Lambda}_N \equiv \{\lambda \in \hat{\Lambda}_N; B' = B\}$ .

Hence we only have to investigate the probability space  $(\hat{\Lambda}_N, \hat{P}_N)$ . we obtain the following theorem first, which is much easier to prove than Theorem 1.

*Theorem 2.*

For each integer  $k \geq 1$ ,  $0 = t_0 < t_1 < \dots < t_k < t_{k+1} = 1$  and  $-\infty < T_j < T'_j < +\infty$ ,  $j = 1, 2, \dots, k+1$ , we have

$$\begin{aligned} & \lim_{N \rightarrow \infty} \hat{P}_N \left( \bigcap_{j=1}^{k+1} \left\{ T_j \leq \frac{\bar{X}_N(a_j^{(N)})}{\sigma\sqrt{N}} \leq T'_j \right\} \right) \\ &= \lim_{N \rightarrow \infty} \hat{P}_N \left( \bigcap_{j=1}^{k+1} \left\{ T_j \leq \frac{X_N(a_j^{(N)})}{\sigma\sqrt{N}} \leq \frac{X_N(a_j^{(N)})}{\sigma\sqrt{N}} \leq T'_j \right\} \right) \\ &= P_0(X_{t_j} \in [T_j, T'_j], j = 1, 2, \dots, k+1), \end{aligned}$$

where  $(\{X_t\}, P_0)$  is a one-dimensional Brownian Motion starting at 0.

References

- [1] Del Grosso, G., On the Local Central Limit Theorem for Gibbs Processes. *Comm. Math. Phys.*, 37, 1974.
- [2] Gallavotti, G., The Phase Separation Line in the Two-Dimensional Ising Model. *Comm. Math. Phys.*, 27, 1972.
- [3] Gallavotti, G. and Martin-Löf, A., Surface Tension in the Ising Model. *Comm. Math. Phys.*, 25, 1972.