# Residual vanishing of concentration arising in the mean field equations

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#### **Abstract**

In this short report, we study the Sawada-Suzuki equation. In the positive case, we prove the property called *Residual vanishing* which means that a blow-up solution sequence (more precisely, its subsequence) converges to a finite sum of Dirac's measures in the sense of measure.

## 1 Introduction

In this report, we consider the Sawada-Suzuki equation ([6]):

$$\begin{cases} -\Delta v_n = \lambda_n \int_I \alpha \left( \frac{e^{\alpha v_n}}{\int_{\Omega} e^{\alpha v_n}} - \frac{1}{|\Omega|} \right) \mathcal{P}(d\alpha) & \text{in } \Omega \\ \int_{\Omega} v_n = 0, \end{cases}$$
 (1.1)

where  $(\lambda_n, v_n)$  is a solution sequence to (1.1),  $\lambda_n$  a non-negative number sequence tending to some non-negative number  $\lambda_0$ , I = [-1, 1],  $\Omega = (\Omega, g)$  a two dimensional orientable compact Riemannian manifold, and  $\mathcal{P}(d\alpha)$  a Borel probability measure on I. According to the result of [4], the following alternative holds:

(i) (Compactness)  $\limsup_{n\to\infty} ||v_n||_{\infty} < +\infty$ , namely, there exist  $v \in \mathcal{E}$  and a subsequence  $\{v_{n_k}\} \subset \{v_n\}$  such that  $v_{n_k} \to v$  in  $\mathcal{E}$  as  $k \to \infty$ , where

$$\mathcal{E} = \left\{ v \in H^1(\Omega) \mid \int_{\Omega} v = 0 \right\}.$$

(ii) (Concentration)  $\limsup_{n\to\infty} \|v_n\|_{\infty} = +\infty$ , namely, the set  $\mathcal{S} = \mathcal{S}_+ \cup \mathcal{S}_-$  is a non-empty and finite set, and there exists  $0 \le s_{\pm} \in L^1(\Omega)$  such that

$$\nu_{\pm,n} := \lambda_n \int_{I_{\pm}} \frac{\alpha e^{\alpha \nu_n}}{\int_{\Omega} e^{\alpha \nu_n}} \mathcal{P}(d\alpha) dx \stackrel{*}{\rightharpoonup} \nu_{\pm} = s_{\pm} dx + \sum_{x_0 \in \mathcal{S}_{\pm}} m(x_0) \delta_{x_0}(dx) \quad (1.2)$$

in  $\mathcal{M}(\Omega)$  with  $m(x_0) \geq 4\pi$  for all  $x_0 \in \mathcal{S}_{\pm}$ , where  $I_+ = (0,1], I_- = [-1,0),$   $\delta_x$  is the Dirac measure supported at x,  $\mathcal{M}(\Omega) = C(\Omega)^*$  and

$$S_{\pm} = \{x_0 \in \Omega \mid \text{there exists } \{x_n\} \subset \Omega \text{ such that } x_n \to x_0 \text{ and } v_n(x_n) \to \pm \infty \}.$$
(1.3)

It is natural to ask whether  $s_{\pm}$  is zero or not in (1.2). If this is the case, we call this property residual vanishing in this report. In the positive case, we obtain

**Proposition 1.** If (ii) above holds and  $I = I_+$ , then  $s = s_+ = 0$ .

**Remark 1.** We note that  $S = S_+$  in the case  $I = I_+$ , see [4] for details. The proof of this fact is based on the boundedness from below of the Green function associated to  $-\Delta$  on  $\Omega$ , i.e.,

$$egin{cases} -\Delta_x G(x,y) = \delta_y - rac{1}{|\Omega|} & in \ \Omega \ \int_\Omega G(x,y) dx = 0, & orall y \in \Omega, \end{cases}$$

see [1].

**Remark 2.** Residual vanishing also holds in the case  $I = I_{-}$ .

**Remark 3.** It is open whether residual vanishing is true or not in the general case. On the contrary, the problem is not solved even in the simple case  $\mathcal{P}(d\alpha) = \frac{1}{2}(\delta_{-1} + \delta_1)$  treated in [5].

It is not difficult to show residual vanishing in the case  $\mathcal{P}(d\alpha) = \delta_p$  for  $p \in I$  by a direct application of the result (Theorem 3) of [2]. Just to be safe, we show it here, assuming p = 1 for simplicity, i.e.,

$$-\Delta v_n = \lambda_n \left( \frac{e^{v_n}}{\int_{\Omega} e^{v_n}} - \frac{1}{|\Omega|} \right).$$

Fix  $x_0 \in \mathcal{S}$ . If it fails then it holds that

$$\liminf_{n\to\infty}\int_{\Omega}e^{v_n}<+\infty.$$

We introduce

$$z_n = v_n - \log \int_{\Omega} e^{v_n}$$

and obtain

$$-\Delta z_n = \lambda_n e^{z_n} - \frac{\lambda_n}{|\Omega|} \quad \text{in } \Omega.$$

It follows from the assumption of contradiction that  $z_n \to +\infty$  (for some subsequence still denoted by the same notation). Since  $\lambda_n$  is uniformly bounded and  $-\lambda_n/|\Omega|$  can be regarded as a simple perturbed term, we can safely apply the result of [2] to the equation of  $z_n$  to find that  $z_n \to -\infty$  in  $B(x_0, r_0) \setminus \{x_0\}$  for  $0 < r_0 \ll 1$ , where B(x, r) denotes a disk centerd at x with radius r for  $x \in \mathbf{R}^2$  and r > 0, in particular,  $B_r$  in the case x = 0. On the other hand,  $z_n$  is bounded below in  $B(x_0, r_0) \setminus \{x_0\}$  since  $S = S_+ \neq \emptyset$ , a contradiction.

Still, it seems to be difficult to directly apply the result of [2] to the general positive case. To overcome this difficulty, we introduce the key transformation, see (2.3) below, and then develop a blowup analysis.

This report consists of three sections. We prove Proposition 1 in Section 2, and several lemmas stated there are shown in Section 3.

# 2 Proof of Proposition 1

In this section, we write I and S by  $I_+$  and  $S_+$ , respectively, in order to stress that we treat the positive case.

To prove the proposition, we have only to show

$$\mathcal{P}(\{\alpha \in I_{+} \mid \liminf_{n \to \infty} \int_{\Omega} e^{\alpha v_{n}} = +\infty\}) = \mathcal{P}(I_{+}). \tag{2.1}$$

To confirm this, we fix  $\omega \subset\subset \Omega \setminus \mathcal{S}_+$ . Then, it holds that

$$0 \leq \int_{\omega} s_{+} dx = \lim_{n \to \infty} \int_{\omega} \nu_{+,n} = \lim_{n \to \infty} \lambda_{n} \int_{\omega} \int_{I_{+}} \left( \frac{\alpha e^{\alpha v_{n}}}{\int_{\Omega} e^{\alpha v_{n}}} - \frac{1}{|\Omega|} \mathcal{P}(d\alpha) \right)$$
  
$$\leq (\lambda_{0} + 1) C(\omega) \lim_{n \to \infty} \int_{I_{+}} \frac{\mathcal{P}(d\alpha)}{\int_{\Omega} e^{\alpha v_{n}}} = 0$$

because  $\lambda_n \to \lambda_0$  and  $v_n$  is uniformly bounded in  $\omega$ . Hence, we obtain s = 0 in  $\omega$  by  $0 \le s_{+,n} \in L^1(\Omega)$ . Since  $\omega \subset\subset \Omega \setminus \mathcal{S}_+$  is arbitrary, the proposition holds if (2.1) is true.

Now, we suppose that (2.1) is false. Then, there exists a number  $\alpha_*$  such that

$$0 < \alpha_* := \sup\{\alpha \in I_+ \mid \liminf_{n \to \infty} \int_{\Omega} e^{\alpha v_n} < +\infty\} \quad \text{and} \quad \mathcal{P}((0, \alpha_*]) > 0.$$
(2.2)

Fix  $x_0 \in \mathcal{S}_+$  and take  $r_0 > 0$  satisfying  $\overline{B(x_0, r_0)} \cap \mathcal{S}_+ = \{x_0\}$ . It is possible to take such an  $r_0$  because  $\mathcal{S}$  is a finite set. We may assume  $x_0 = 0$  by a translation. Then, there exist  $x_n \in B_{r_0}$  and  $\alpha_n \in \mathbf{R}$  such that

$$x_n \to 0 \quad v_n(x_n) = \max_{\overline{B_{3r_0}}} v_n \to +\infty,$$

$$e^{\alpha_n v_n(x_n)} = \int_{I_+} \frac{\alpha e^{\alpha v_n(x_n)}}{\int_{\Omega} e^{\alpha v_n}} \mathcal{P}(d\alpha). \tag{2.3}$$

For this  $\alpha_n$ , we obtain the following lemmas shown in next section.

**Lemma 1.** There exists  $C_1 > 0$ , independent of n, such that

$$\int_{I_{\perp}} \frac{\alpha e^{(\alpha - \alpha_n)v_n(x)}}{\int_{\Omega} e^{\alpha v_n}} \mathcal{P}(d\alpha) \le C_1$$

for all  $x \in \overline{B_{2r_0}}$ .

### Lemma 2. We have

$$\alpha_n \to \alpha_0 \in [\alpha_*, 1],$$

passing to a subsequence.

Here, we develop a blow-up argument. Set

$$\begin{cases} w_n(x) = \alpha_n v_n(x_n) - L, \\ \tilde{w}_n(x) = w_n(\sigma_n x + x_n) + 2\log\sigma_n, \\ \sigma_n = e^{-w_n(x_n)/2} \ (\to 0 \text{ by Lemma 2}), \end{cases}$$

where  $L \gg 1$  will be determined later on. The function  $\tilde{w}_n = \tilde{w}_n(x)$  is a solution to

$$\begin{cases}
-\Delta \tilde{w}_{n} = \alpha_{n} \tilde{V}_{n}(x) e^{\tilde{w}_{n}} - \sigma_{n}^{2} \frac{\alpha_{n} \lambda_{n}}{|\Omega|} \int_{I_{+}} \alpha \mathcal{P}(d\alpha) & \text{in } B_{r_{0}/\sigma_{n}} \\
\tilde{w}_{n} \leq \tilde{w}_{n}(0) = 0 & \text{in } B_{r_{0}/\sigma_{n}} \\
\int_{B_{r_{0}/\sigma_{n}}} \tilde{V}_{n} e^{\tilde{w}_{n}} \leq m(0),
\end{cases}$$
(2.4)

where

$$\tilde{V}_n(x) = e^L \cdot \lambda_n \int_{I_+} \frac{\alpha e^{(\alpha - \alpha_n)v_n(\sigma_n x + x_n)}}{\int_{\Omega} e^{\alpha v_n}} \mathcal{P}(d\alpha).$$

**Lemma 3.** There exist  $\tilde{w} \in C^2(\mathbf{R}^2)$  and  $0 < \tilde{V} \in C^2(\mathbf{R}^2) \cap L^{\infty}(\mathbf{R}^2)$  such that

$$\tilde{w}_n \to \tilde{w}, \quad \tilde{V}_n \to \tilde{V} \qquad in \mathbf{R}^2$$

and

$$\begin{cases}
-\Delta \tilde{w} = \alpha_0 \tilde{V}(x) e^{\tilde{w}} & \text{in } \mathbf{R}^2 \\
\tilde{w} \leq \tilde{w}(0) = 0 & \text{in } \mathbf{R}^2 \\
\int_{\mathbf{R}^2} \tilde{V} e^{\tilde{w}} \leq m(0).
\end{cases}$$
(2.5)

Lemma 3 is also shown in next section.

For a solution  $\tilde{w}$  to (2.5), we set

$$\tilde{\phi}(x) = \frac{\alpha_0}{2\pi} \int_{\mathbf{R}^2} \tilde{V}(y) e^{\tilde{w}(y)} \log \frac{|x-y|}{1+|y|} dy, \tag{2.6}$$

complying [3]. Noting that

$$\tilde{V}e^{\tilde{w}} \in L^1 \cap L^{\infty}(\mathbf{R}^2), \tag{2.7}$$

we find that the function  $\tilde{\phi}$  set by (2.6) is well-defined in  $\mathbb{R}^2$ , and can show the following lemma because the proof of Lemma 1.1 of [3] is applicable to our case, see also Remark below.

**Lemma 4.** There exists  $C_2 > 0$ , independent of L, such that

$$\tilde{w}(x) \ge -\beta \log(1+|x|) - C_2 \tag{2.8}$$

for  $x \in \mathbf{R}^2$ , where

$$\beta = \frac{\alpha_0}{2\pi} \int_{\mathbf{R}^2} \tilde{V} e^{\tilde{w}}.$$
 (2.9)

Remark 4. In Lemma 1.1 of [3], the integrability condition  $\int_{\mathbf{R}^2} e^{\tilde{w}} dx < +\infty$  is assumed to show the estimates from above and below for solutions and the estimate from below for  $\beta$ . However, it is not required if one only needs the estimate from below (2.8).

*Proof of Proposition 1*: Fix  $R \gg 1$ . It follows from Lemmas 3-4 that

$$v_n(x) \ge v_n(x_n) - \frac{\beta}{\alpha_n} \log \left(1 + \left| \frac{x - x_n}{\sigma_n} \right| \right) - \frac{C_2}{\alpha_n} + \varepsilon_n$$

for all  $x \in B(x_n, \sigma_n R)$ , where  $\varepsilon_n$  is a quantity converging to 0 as  $n \to \infty$ . This  $\varepsilon_n$  may be changed in the following but keeps the property that  $\varepsilon_n \to 0$ . We obtain

$$\int_{B(x_n,\sigma_n)} e^{\alpha v_n} \ge e^{\alpha v_n(x_n) - \alpha C_2/\alpha_n - 1} \int_{B(x_n,\sigma_n R)} \left( 1 + \left| \frac{x - x_n}{\sigma_n} \right| \right)^{-\alpha \beta/\alpha_n} dx$$

$$= e^{(\alpha - \alpha_n)v_n(x_n)} \cdot e^{L - \alpha C_2/\alpha_n - 1} \int_{B_R} (1 + |x|)^{-\alpha \beta/\alpha_n} dx \quad (2.10)$$

for all  $\alpha \in I_+$ . Thus, (2.3) and (2.10) yield

$$1 = \int_{I_{+}} \frac{\alpha e^{(\alpha - \alpha_{n})v_{n}(x_{n})}}{\int_{\Omega} e^{\alpha v_{n}}} \mathcal{P}(d\alpha)$$

$$\leq \varepsilon_{n} + \int_{[\alpha_{n}, 1]} \frac{\int_{B(x_{n}, \sigma_{n})} e^{\alpha v_{n}}}{\int_{\Omega} e^{\alpha v_{n}}} \cdot \frac{\alpha}{e^{L - \alpha C_{2}/\alpha_{n} - 1} \int_{B_{R}} (1 + |x|)^{-\alpha \beta/\alpha_{n}} dx} \mathcal{P}(d\alpha)$$

$$\leq \varepsilon_{n} + \frac{1}{e^{L - C_{2}/\alpha_{n} - 1} \int_{B_{R}} (1 + |x|)^{-\beta/\alpha_{n}} dx}.$$
(2.11)

Since  $\beta/\alpha_n \leq (\alpha_0/\alpha_n) \cdot (m(0)/2\pi)$  by (2.9) and Lemma 3, inequality (2.11) implies

$$1 \le \varepsilon_n + \frac{1}{e^{L - C_2/\alpha_n - 1} \int_{B_R} (1 + |x|)^{-\frac{\alpha_0}{\alpha_n} \cdot \frac{m(0)}{2\pi}} dx},$$

or

$$1 \le \frac{e^{1 + C_2/\alpha_0} - L}{\int_{B_R} (1 + |x|)^{-\frac{m(0)}{2\pi}} dx},$$

which is a contradiction if L is sufficiently large. The proof is complete.  $\square$ 

# 3 Proof of Lemmas 1-3

As having announced in the previous sections, we show Lemmas 1-3 in this section. We again consider the positive case (i.e.,  $S = S_+$  and  $I = I_+$ ) in what follows.

Proof of Lemma 1: Since  $S = S_+$ , there exists  $C_3 > 0$ , independent of n, such that  $v_n > -C_3$  in  $\Omega$ . We use (2.3) and Jensen's inequality to calculate

$$\int_{I_{+}} \frac{\alpha e^{(\alpha - \alpha_{n})v_{n}(x)}}{\int_{\Omega} e^{\alpha v_{n}}} \mathcal{P}(d\alpha)$$

$$\leq \int_{I'_{+,n}} \frac{\alpha e^{-(\alpha_{n} - \alpha)v_{n}(x)}}{\int_{\Omega} e^{\alpha v_{n}}} + \int_{I_{+}} \frac{\alpha e^{(\alpha - \alpha_{n})v_{n}(x_{n})}}{\int_{\Omega} e^{\alpha v_{n}}} \mathcal{P}(d\alpha)$$

$$\leq \frac{\alpha_{n} \mathcal{P}(I'_{+,n}) e^{\alpha_{n} C_{3}}}{|\Omega|} + 1 \leq \frac{e^{C_{3}}}{|\Omega|} + 1$$

for all  $x \in \overline{B_{2r_0}}$  and n, where

$$I'_{+,n} = \begin{cases} (0, \alpha_n) & \text{if } \alpha_n > 0\\ \emptyset & \text{if } \alpha_n \le 0. \end{cases}$$

The lemma is completely shown.

*Proof of Lemma 2*: Put  $\alpha_0 = \lim_{n \to \infty} \alpha_n$ .

Assume that  $\alpha_0 > 1$ . Then, there exists  $\delta > 0$  such that

$$e^{(1+\delta)v_n(x_n)} \le e^{\alpha_n v_n(x_n)},$$

that is, by Jensen's inequality,

$$e^{\frac{\delta}{2}v_n(x_n)} \le \int_{I_+} \frac{\alpha e^{(\alpha - 1 - \delta/2)v_n(x_n)}}{\int_{\Omega} e^{\alpha v_n}} \mathcal{P}(d\alpha) \le e^{-\frac{\delta}{2}v_n(x_n)} |\Omega|^{-1}$$

for  $n \gg 1$ , which is a contradiction because  $v_n(x_n) \to +\infty$ .

Next, assume that  $\alpha_0 \leq 0$ . In the case that  $\mathcal{P}((0,\alpha_*)) > 0$ , there exists  $0 < \varepsilon \ll 1$  such that  $\mathcal{P}([\varepsilon, \alpha_* - \varepsilon]) > 0$ , and therefore

$$1 = \int_{I_{+}} \frac{\alpha e^{(\alpha - \alpha_{n})v_{n}(x_{n})}}{\int_{\Omega} e^{\alpha v_{n}}} \mathcal{P}(d\alpha)$$

$$\geq \int_{[\varepsilon, \alpha_{*} - \varepsilon]} \frac{\alpha e^{(\alpha - \varepsilon/2)v_{n}(x_{n})}}{\int_{\Omega} e^{\alpha v_{n}}} \mathcal{P}(d\alpha)$$

$$\geq c(\varepsilon) e^{\frac{\varepsilon}{2}v_{n}(x_{n})} \mathcal{P}([\varepsilon, \alpha_{*} - \varepsilon]) \rightarrow +\infty$$

as  $n \to \infty$ , a contradiction. In the case that  $\mathcal{P}(\{\alpha_*\}) = \mathcal{P}((0, \alpha_*]) > 0$ , it holds that  $\liminf_{n \to \infty} \int_{\Omega} e^{\alpha_* v_n} < +\infty$ , and hence

$$1 = \int_{I_{+}} \frac{\alpha e^{(\alpha - \alpha_{n})v_{n}(x_{n})}}{\int_{\Omega} e^{\alpha v_{n}}} \mathcal{P}(d\alpha)$$

$$\geq \alpha_{*} e^{(\alpha_{*} - \alpha_{n})v_{n}(x_{n})} \left( \int_{\Omega} e^{\alpha_{*} v_{n}} \right)^{-1} \mathcal{P}(\{\alpha_{*}\}) \to +\infty$$

as  $n \to \infty$ , a contradiction.

We have shown that  $\alpha_0 \in (0,1]$ . It is left to show that  $\alpha_0 \geq \alpha_*$ . To prove this, we finally assume that  $\alpha_0 \in (0,\alpha_*)$ . Consider

$$\varphi_n = \alpha_n v_n - \log \int_{\Omega} e^{\alpha_n v_n}.$$

Passing to a subsequence, we have

$$\varphi_n(x_n) \to +\infty.$$
 (3.1)

The function  $\varphi_n = \varphi_n(x)$  satisfies

$$\begin{cases}
-\Delta \varphi_n = K_n(x)e^{\varphi_n} - \frac{\alpha_n \lambda_n}{|\Omega|} \int_{I_+} \alpha \mathcal{P}(d\alpha) & \text{in } B_{2r_0} \\
\int_{\Omega} e^{\varphi_n} = 1,
\end{cases}$$
(3.2)

where

$$K_n(x) = \alpha_n \lambda_n \left( \int_{\Omega} e^{\alpha_n v_n} \right) \int_{I_+} \frac{\alpha e^{(\alpha - \alpha_n) v_n(x)}}{\int_{\Omega} e^{\alpha v_n}} \mathcal{P}(d\alpha).$$

Lemma 1 and the boundedness  $\liminf_{n\to\infty} \int_{\Omega} e^{\alpha_n v_n} < +\infty$  show that there exists  $C_4 > 0$ , independent of n, such that

$$0 \le K_n \le C_4 \quad \text{in } B_{2r_0}. \tag{3.3}$$

Consequently, (3.1)-(3.3) assure that

$$\varphi_n \to -\infty$$
 locally uniformly in  $B_{2r_0 \setminus \{0\}}$  (3.4)

by virtue of the result of [2]. However, (3.4) is false since  $S = S_+$  and  $\lim \inf_{n \to \infty} \int_{\Omega} e^{\alpha_n v_n} < +\infty$ .

Proof of Lemma 3: It follows from Lemma 2 that

$$0 \le \tilde{V}_n \le e^{L(\lambda_0 + 1)} C_1 \quad \text{in } B_{r_0/\sigma_n}$$

for  $n \gg 1$ . We also have

$$0 \le e^{\tilde{w}_n} \le 1$$
 in  $B_{r_0/\sigma_n}$ 

for all n, and

$$\sigma_n^2 \frac{\alpha_n \lambda_n}{|\Omega|} \int_{I_+} \alpha \mathcal{P}(d\alpha) \to 0$$

as  $n \to \infty$ . Combining these properties with  $\tilde{w}_n(0) = 0$ , we can safely apply the result of [2] to find that, for every R > 0, there exists  $C_5(R) > 0$  such that

$$\tilde{w}_n \ge -C_5(R) \quad \text{in } B_R \tag{3.5}$$

for  $n \gg 1$ . Thus, the elliptic regularity and a diagonal arugument show that there exists  $\tilde{w} \in C^{1+\alpha}(\mathbf{R}^2)$ ,  $\alpha \in (0,1)$ , such that

$$\tilde{w}_n \to \tilde{w} \quad \text{in } C_{loc}^{1+\alpha}(\mathbf{R}^2).$$
 (3.6)

Noting the definitions of  $\tilde{V}_n$  and  $\tilde{w}_n$ , we see that there exists  $\tilde{V} \in C^{1+\alpha}(\mathbf{R}^2)$ ,  $\alpha \in (0,1)$ , such that

$$\tilde{V}_n \to \tilde{V} \quad \text{in } C_{loc}^{1+\alpha}(\mathbf{R}^2).$$
 (3.7)

We again use the elliptic regularity, together with (3.6)-(3.7), and conclude the relation (2.5) and  $\tilde{w}, \tilde{V} \in C^2(\mathbf{R}^2)$ .

It is clear that  $\tilde{V} \in L^{\infty}(\mathbf{R}^2)$  by Lemma 1, and therefore, we must show that  $\int_{\mathbf{R}^2} \tilde{V}e^{\tilde{w}} \leq m(0)$  and that  $\tilde{V} > 0$  in  $\mathbf{R}^2$ .

For every R > 0 and  $0 < r \ll 1$ ,

$$\int_{B_R} \tilde{V}e^{\tilde{w}} \le \liminf_{n \to \infty} \int_{B_R} \tilde{V}_n e^{\tilde{w}_n} \le \liminf_{n \to \infty} \int_{B_{r/\sigma_n}} \tilde{V}_n e^{\tilde{w}_n}$$

$$= \liminf_{n \to \infty} \int_{B(x_n, r)} \nu_{+,n} \le m(0) + \int_{B_{2r}} \nu_{+}$$

by the Fatou lemma, the definitions of  $w_n$ ,  $\tilde{w}_n$ ,  $\sigma_n$  and  $\tilde{V}_n$ , and (1.2). Letting  $R \uparrow +\infty$  and  $r \downarrow 0$ , we obtain  $\int_{\mathbb{R}^2} \tilde{V} e^{\tilde{w}} \leq m(0)$ .

Finally, we use the definitions of  $w_n$ ,  $\tilde{w}_n$ ,  $\sigma_n$  and  $\tilde{V}_n$ , (3.5),  $\tilde{w}_n \leq 0$  and (1.2) to obtain  $C_6(R) > 0$ , independent of  $n \gg 1$ , such that

$$\begin{split} \tilde{V}_n(x) &= e^L \lambda_n \int_{I_+} \frac{\alpha e^{\frac{\alpha - \alpha_n}{\alpha_n} (\tilde{w}_n(x) + \alpha_n v_n(x_n))}}{\int_{\Omega} e^{\alpha v_n}} \mathcal{P}(d\alpha) \\ &\geq e^{L - C_6(R)} \lambda_n \int_{I_+} \frac{\alpha e^{(\alpha - \alpha_n) v_n(x_n)}}{\int_{\Omega} e^{\alpha v_n}} \mathcal{P}(d\alpha) = e^{L - C_6(R)} \lambda_n \end{split}$$

for all  $x \in B_R$  and  $n \gg 1$ , and for every R > 0, which means  $\tilde{V} > 0$  in  $\mathbb{R}^2$  because  $\lambda_n \to \lambda_0 > 0$  by  $S = S_+ \neq \emptyset$ .

# References

[1] T. Aubin, Some Nonlinear Problems in Riemannian Geometry, Springer, Berlin, 1998.

- [2] H. Brezis and F. Merle, Uniform estimates and blowup behavior for solutions of  $-\Delta u = V(x)e^u$  in two dimensions, Commun. Partial Differential Equations 16 (1991) 1223-1253.
- [3] W. Chen and C. Li, Qualitative properties of solutions to some nonlinear elliptic equations in  $\mathbb{R}^2$ , Duke Math. J. **71** (1993) 427-439.
- [4] H. Ohtsuka, T. Ricciardi, and T. Suzuki, Blow-up analysis for an elliptic equation describing stationary vortex flows with variable intensities in 2D-turbulence, J. Differential Equations 249 (2010) 1436-1465.
- [5] H. Ohtsuka and T. Suzuki, Mean field equation for the equilibrium turbulence and a related functional inequality, Adv. Differential Equations 11 (2006) 281-304.
- [6] K. Sawada and T. Suzuki, Derivation of the equilibrium mean field equations of point vortex and vortex filament system, Theoret. Appl. Mech. Japan **56** (2008) 285-290.
- [7] T. Suzuki, Mean Field Theories and Dual Variation, Atlantis Press, Amsterdam, 2008.