

# Borel summability of WKB theoretic transformation to the Weber equation

By

Shinji SASAKI\*

## Abstract

We take a Schrödinger equation which has a pair of simple turning points connected by a Stokes line, and consider WKB theoretic transformation series to the Weber equation. Under suitable conditions, the transformation series is Borel summable, and analysis of so-called fixed singularities can be reduced to the Weber equation.

## § 1. Introduction

We consider a second order linear differential equation with a large parameter

$$(1.1) \quad \left( \frac{d^2}{dx^2} - \eta^2 Q(x) \right) \psi = 0.$$

The coefficient  $Q(x)$  is a holomorphic function (typically rational function or polynomial). The equation has formal solutions (WKB solutions) of the form

$$(1.2) \quad \psi(x, \eta) = \exp \left( \int^x S(x, \eta) dx \right),$$

where  $S(x, \eta) = \eta S_{-1} + S_0 + \eta^{-1} S_1 + \dots$  is a formal power series satisfying the Riccati equation

$$(1.3) \quad S^2 + \frac{dS}{dx} = \eta^2 Q(x).$$

The WKB solutions  $\psi(x, \eta)$  (or  $S(x, \eta)$ ) are divergent in general, and we apply Borel resummation method. (See e.g., [12], [9].) Under generic assumptions, the WKB solutions (of suitable normalization) is Borel summable (see [4], [8]), but in some cases not.

---

2010 Mathematics Subject Classification(s): Primary 34M60; Secondary 34M25.

*Key Words:* exact WKB analysis, fixed singularity, transformation series, Borel summability

Supported by JSPS Grants-in-Aid No.22-1398

\*RIMS, Kyoto University, Kyoto 606-8502, Japan. JSPS Research Fellow.

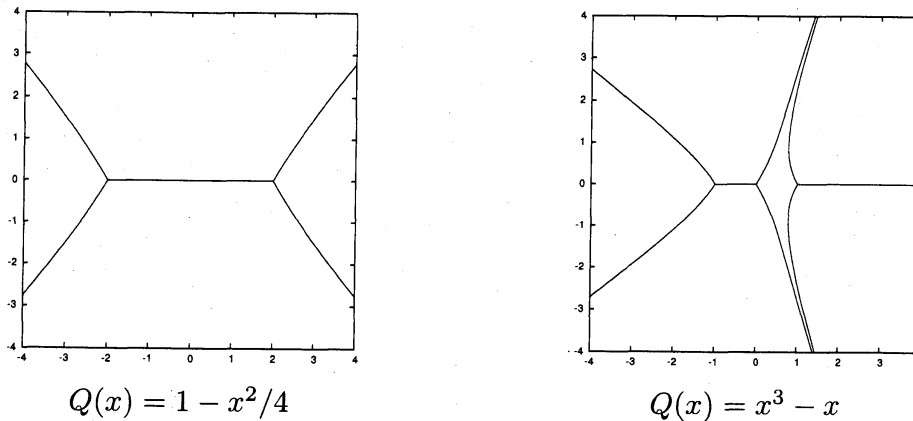


Figure 1. two examples in which Stokes lines connect simple turning points.

For example if  $Q(x) = E - x^2/4$  (namely the Weber equation) with a positive constant  $E > 0$ , WKB solutions are not Borel summable (depending on the normalization). See e.g., [11]. This is a general phenomenon if the equation has a pair of turning points (zeros of  $Q(x)$ ) connected by a Stokes line  $\Im \int^x \sqrt{Q(x)} dx = 0$ . See e.g., [2], [3]. In Figure 1, we give examples of Stokes lines connecting turning points.

In such cases, the Borel transform of a WKB solution has singularities on the real axis (in the Borel plane), which we call “fixed singularities”. They are fixed in the sense that the location is independent of  $x$ . To analyze such singularities, in [1], WKB theoretic transformation to the Weber equation is constructed, and Borel transformability is given. Here transformation series is a formal power series  $x(q, \eta) = x_0(q) + \eta^{-1}x_1(q) + \eta^{-2}x_2(q) + \dots$  which transforms the equation

$$(1.4) \quad \left( \frac{d^2}{dq^2} - \eta^2 Q(q) \right) \psi = 0$$

to the Weber equation (with an infinite power series  $E = E(\eta) = E_0 + \eta^{-1}E_1 + \eta^{-2}E_2 + \dots$ )

$$(1.5) \quad \left( \frac{d^2}{dx^2} - \eta^2 \left( E - \frac{x^2}{4} \right) \right) \phi = 0,$$

with a gauge transform  $\psi = x^{-1/2}\phi$ . This is equivalent to that  $x(q, \eta)$  satisfies the following:

$$(1.6) \quad Q(q) = \left( \frac{dx}{dq} \right)^2 \left( E - \frac{x^2}{4} \right) - \frac{1}{2}\eta^{-2}\{x; q\}.$$

Here  $\{x; q\}$  is the Schwarzian derivative. Though this is a transformation between equations, this also connects WKB solutions of certain normalization. (See [1] and the following section.)

The purpose of this paper is to present Borel summability of transformation series in a simple case, and add a brief explanation of the consequence (the following section).

In ending of this introduction, we refer to the work of Kamimoto and Koike ([5]), which shows Borel summability of transformation series to the Airy equation ( $Q(x) = x$ ). The Airy equation has only one simple turning point, and is the simplest equation whose (Borel transformed) WKB solution has so-called “movable singularities”. The basic idea of the proof of the Weber case follows the Airy case [5], while one additional problem arises which we should overcome. In this paper, we do not give a proof of Borel summability of transformation series to the Weber equation. A detailed proof will be given elsewhere.

## § 2. Borel summability of transformation series

In this section, for simplicity we assume that the coefficient  $Q(q)$  in (1.4) is polynomial. Let  $q_{\pm}$  be simple turning points of the equation (1.4). Assume  $q_{\pm}$  are connected by a Stokes line and the other Stokes lines emanating from the two points tend to infinity. For example, if  $Q(q) = q(q^2 - 1)$  and we take  $q_+ = 0$  and  $q_- = -1$ , these conditions are satisfied (See Figure 1). Take a neighborhood

$$(2.1) \quad D = \left\{ \left| \int_{q_+}^q \sqrt{Q} dq \right| < d \right\} \cup \left\{ \left| \int_{q_-}^q \sqrt{Q} dq \right| < d \right\}$$

of  $\{q_{\pm}\}$  and set

$$(2.2) \quad \hat{D} = \bigcup_{q' \in D} \left\{ \Im \int_{q'}^q \sqrt{Q} dq = 0 \right\}.$$

(cf. Figure 2.) We take  $d$  small enough so that  $\hat{D}$  does not contain any turning points except for  $q_{\pm}$ . Then there exist formal power series  $x(q, \eta) = x_0(q) + \eta^{-1}x_1(q) + \eta^{-2}x_2(q) + \cdots$  and  $E(\eta) = E_0 + \eta^{-1}E_1 + \eta^{-2}E_2 + \cdots$  with  $x_j(q)$  being holomorphic on  $\hat{D}$  ( $j = 0, 1, 2, \dots$ ) which satisfy the equation (1.6) and  $dx_0/dq \neq 0$ .  $x(q, \eta)$  and  $E(\eta)$  are uniquely determined up to the choice of  $x_0(q)$ . See [1], [9].

*Remark.*  $x_0(q)$  is a map which maps a turning point to a turning point, a level curve (Stokes line)  $\Im \int^q \sqrt{Q} dq = 0$  to a level curve (Stokes line)  $\Im \int^x \sqrt{(E_0 - x^2/4)} dx = 0$ . There are two turning points  $q_{\pm}$ , and we have two choices of  $x_0(q)$ .

The Borel summability of  $E(\eta)$  is known. See [8]. In addition we have the following theorem.

**Theorem 2.1.** *Under the assumptions above, the transformation series  $x(q, \eta)$  is Borel summable uniformly on  $\hat{D}$ .*

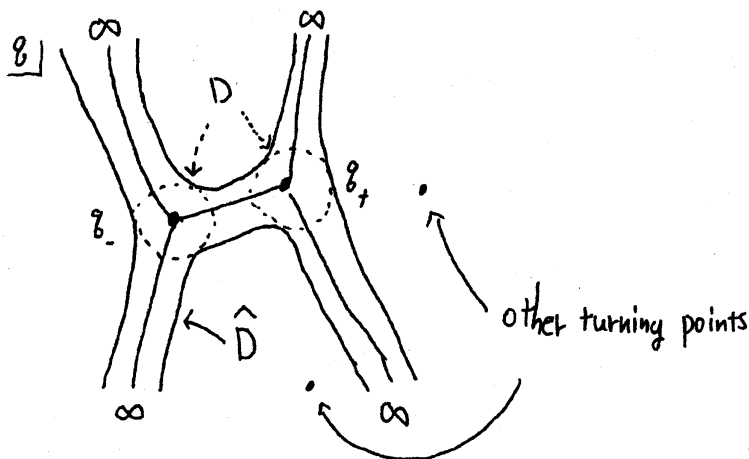


Figure 2. Domains  $D$  and  $\hat{D}$ .

Thus the equation (1.4) on  $\hat{D}$  is transformed to the canonical equation (1.5) by two Borel summable series  $x(q, \eta)$  and  $E(\eta)$ . Then as is explained in [1] and [9] (though mainly Airy case, not Weber case), a WKB solution of (1.4) is also transformed into a WKB solution of (1.5); Let  $\psi(q, \eta)$  be a WKB solution of (1.4) normalized at  $q_+$   $\phi(x, E, \eta)$  be a WKB solution of (1.5) normalized at  $2\sqrt{E}$ . Here we assume  $x_0(q_+) = 2\sqrt{E_0}$ . (For normalization, see e.g., [9].) Then the following relation holds:

$$(2.3) \quad \psi(q, \eta) = \left( \frac{dx}{dq}(q, \eta) \right)^{-1/2} \phi(x(q, \eta), E(\eta), \eta).$$

Though this is a formal relation, if Borel transformed, this becomes an analytic relation. Set  $x(q, \eta) = x_0(q) + X(q, \eta)$  and  $E(\eta) = E_0 + F(\eta)$ . By Taylor expansion, we have

$$(2.4) \quad \begin{aligned} \psi(q, \eta) &= \left( \frac{dx}{dq}(q, \eta) \right)^{-1/2} \sum_{n=0}^{\infty} \frac{X^n(q, \eta)}{n!} \frac{\partial^n \phi}{\partial x^n}(x_0(q), E(\eta), \eta) \\ &= \left( \frac{dx}{dq}(q, \eta) \right)^{-1/2} \sum_{n=0}^{\infty} \frac{X^n(q, \eta)}{n!} \left( \sum_{m=0}^{\infty} \frac{F^m(\eta)}{m!} \frac{\partial^{n+m} \phi}{\partial E^m \partial x^n}(x_0(q), E_0, \eta) \right). \end{aligned}$$

Then by Borel transform, we have

$$(2.5) \quad \psi_B(q, y) = \left( \left( \frac{dx}{dq} \right)^{-1/2} \right)_B(q, y) * \sum_{n=0}^{\infty} \frac{X_B^{*n}(q, y)}{n!} * \left( \sum_{m=0}^{\infty} \frac{F_B^{*m}(y)}{m!} * \frac{\partial^{n+m} \phi_B}{\partial E^m \partial x^n}(x_0(q), E_0, y) \right),$$

where the subscript B means Borel transform and  $*$  is convolution. Now let us take one term

$$\left( \left( \frac{dx}{dq} \right)^{-1/2} \right)_B(q, y) * \frac{X_B^{*n}(q, y)}{n!} * \frac{F_B^{*m}(y)}{m!} * \frac{\partial^{n+m} \phi_B}{\partial E^m \partial x^n}(x_0(q), E_0, y).$$

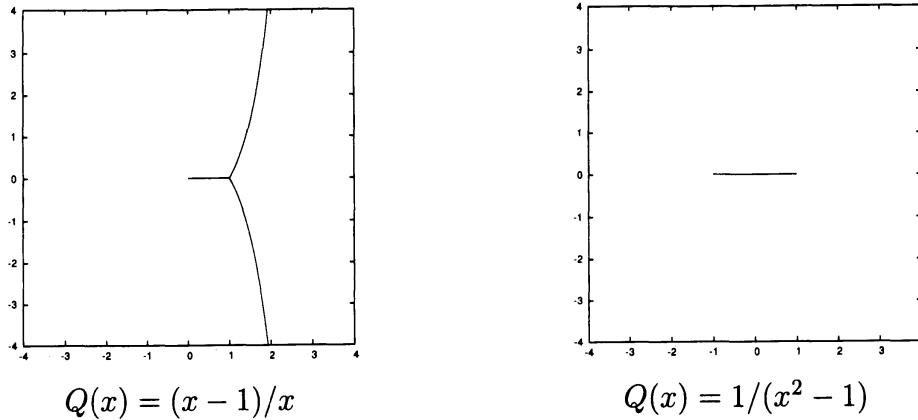


Figure 3. Stokes lines connecting a pair of a simple turning point and a simple pole(left), a pair of simple poles (right).

Since  $X$  and  $F$  are Borel summable, the front part

$$\left( \left( \frac{dx}{dq} \right)^{-1/2} \right)_B (q, y) * \frac{X_B^{*n}(q, y)}{n!} * \frac{F_B^{*m}(y)}{m!}$$

is holomorphic in a strip region containing the positive real axis. Since we know well about  $\phi_B$  (see e.g., [11], [10]), for this single term, we can see continuability avoiding singularities, discontinuity at a singularity, etc. Then by summing up with respect to  $m$  and  $n$  (with care on convergence), we see continuability etc. also for  $\psi_B(q, y)$ .

*Remark.*  $\phi_B(x_0(q), E_0, y)$  has infinitely many singularities in the  $y$ -plane with real period  $2\pi E_0$ , and with Borel summability we can analyze all singularities through transformation. Thus Borel summability of transformation is important in the analysis of fixed singularities.

*Remark.* In this paper, we considered only two simple turning points problem. On the other hand, simple poles (of  $Q$ ) are known to play a role similar to simple turning points ([6], [7]), and a pair of a simple turning point and a simple pole, or a pair of simple poles causes fixed singularities as well (cf. Figure 3). The former one can be treated in the same manner as a pair of simple turning points. However the latter one is difficult to treat with. Also, a sole simple turning point makes a pair in some sense, generating a loop of Stokes line (cf. Figure 4), and this has the same difficulty.

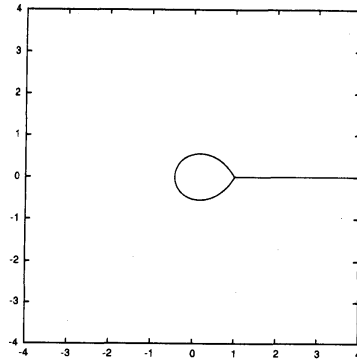


Figure 4. A loop of Stokes line ending a sole simple turning point.

### References

- [1] Aoki, T., Kawai, T. and Takei, Y., The Bender-Wu analysis and the Voros theory, *Special Functions: ICM-90 Satellite Conference Proceedings* (M. Kashiwara and T. Miwa, eds), Springer-Verlag, 1991, pp. 1–29.
- [2] Delabaere, E., Dillinger, H. and Pham, F., Résurgence de Voros et périodes des courbes hyperelliptiques, *Ann. Inst. Fourier (Grenoble)* **43** (1993), 163–199.
- [3] Delabaere, E. and Pham, F., Resurgent methods in semi-classical asymptotics, *Ann. Inst. Henri Poincaré* **71** (1999), 1–94.
- [4] Dunster, T. M., Lutz, D. A. and Schäfke, R., Convergent Liouville-Green expansions for second-order linear differential equations, with an application to Bessel functions, *Proc. R. Soc. Lond. Ser. A Math. Phys. Eng. Sci.* **440** (1993), 37–54.
- [5] Kamimoto, S. and Koike, T., On the Borel summability of WKB-theoretic transformation series, *RIMS preprint* **1726** (2011).
- [6] Koike, T., On a regular singular point in the exact WKB analysis, *Toward the Exact WKB Analysis of Differential Equations, Linear or Non-Linear* (C. J. Howls, T. Kawai and Y. Takei eds.), Kyoto Univ. Press, 2000, pp. 39–54.
- [7] Koike, T., On the exact WKB analysis of second order linear ordinary differential equations with simple poles, *Publ. Res. Inst. Math. Sci.* **36** (2000), 297–319.
- [8] Koike, T. and Schäfke, R., in preparation.
- [9] Kawai, T. and Takei, Y., *Algebraic Analysis of Singular Perturbation Theory*, Transl. Math. Monogr. **227**, Amer. Math. Soc., 2005.
- [10] Sasaki, S., Resurgence of WKB solutions of the Weber equation through integral representation, in preparation.
- [11] Takei, Y., Sato's conjecture for the Weber equation and transformation theory for Schrödinger equations with a merging pair of turning points, *Differential Equations and Exact WKB Analysis* (Y. Takei, ed), *RIMS Kôkyûroku Bessatsu* **B10**, 2008, pp. 205–224.
- [12] Voros, A., The return of the quartic oscillator — The complex WKB method, *Ann. Inst. Henri Poincaré* **39** (1983), 211–338.