## Symmetry-breaking bifurcation of positive solutions to a one-dimensional Liouville type equation

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We consider the two-point boundary value problem for the one-dimensional Liouville type equation

(1) 
$$\begin{cases} u'' + \lambda |x|^l e^u = 0, & x \in (-1, 1), \\ u(-1) = u(1) = 0, \end{cases}$$

where  $\lambda > 0$  and l > 0.

Jacobsen and Schmitt [2] studied the exact multiplicity of radial solutions of the problem for the multi-dimensional Liouville type equation

(2) 
$$\begin{cases} \Delta u + \lambda |x|^l e^u = 0 & \text{in } B, \\ u = 0 & \text{on } \partial B, \end{cases}$$

where  $l \ge 0$  and  $B := \{x \in \mathbb{R}^n : |x| < 1\}$ . They proved the following (i)-(iii):

- (i) if  $1 \leq N \leq 2$ , then there exists  $\lambda_* > 0$  such that (2) has exactly two radial solutions for  $0 < \lambda < \lambda_*$ , a unique radial solution for  $\lambda = \lambda_*$  and no radial solution for  $\lambda > \lambda_*$ ;
- (ii) if  $3 \le N < 10 + 4l$ , then (2) has infinitely many radial solutions for  $\lambda = (l + 2)(N-2)$  and a finite but large number of radial solutions when  $|\lambda (l+2)(N-2)|$  is sufficiently small;
- (iii) if  $N \ge 10 + 4l$ , then (2) has a unique radial solution for  $0 < \lambda < (l+2)(N-2)$  and no radial solution for  $\lambda \ge (l+2)(N-2)$ .

We note here that every solution of (2) is positive in B, by the strong maximum principle. Result (i)–(iii) were established by Joseph and Lundgren [3] for the case l=0, that is, for the Liouville equation

(3) 
$$\begin{cases} \Delta u + \lambda e^{u} = 0 & \text{in } \Omega, \\ u = 0 & \text{on } \partial \Omega, \end{cases}$$

when  $\Omega = B$ . Gidas, Ni and Nirenberg's theorem ([1]) shows that every positive solution of (3) is radially symmetric when  $\Omega = B$ . However, when  $\Omega$  is an annulus  $A := \{x \in \mathbf{R}^N : a < |x| < b\}$ , a > 0, problem (3) may has non-radial solutions. Indeed, Lin [4] proved that (3) has infinitely many symmetry-breaking bifurcation points when N = 2 and  $\Omega = A$ . Nagasaki and Suzuki [6] found that large non-radial solutions of (3) when N = 2 and  $\Omega = A$ . More precisely, for each sufficiently large  $\mu > 0$ , there exist  $(\lambda, u)$  such that  $\lambda > 0$ , u is a non-radial solution of (3) and  $\int_A e^u dx = \mu$  when N = 2 and  $\Omega = A$ .

Recently, Miyamoto [5] considered the problem for the Liouville type equation (2) and proved the following result.

**Theorem A** ([5]). Let  $n_0$  be the largest integer that is smaller than  $1 + \frac{l}{2}$  and let  $\alpha_n := 2 \log \frac{2l+4}{l+2-2n}$ . All the radial solutions of (2) with N=2 can be written explicitly as

$$\lambda(\alpha) = 2(l+2)^2(e^{-\alpha/2} - e^{-\alpha}), \quad U(r;\alpha) = \log \frac{e^{\alpha}}{(1 + (e^{\alpha/2} - 1)r^{l+2})^2}.$$

The radial solutions can be parameterized by the  $L^{\infty}$ -norm, it has one turning point at  $\lambda = \lambda(\alpha_0) = (l+2)/2$ , and it blows up as  $\lambda \downarrow 0$ . For each  $n \in \{1, 2, \dots, n_0\}$ ,  $(\lambda(\alpha_n), U(r; \alpha_n))$  is a symmetry breaking bifurcation point from which an unbounded branch consisting of non-radial solutions of (2) with N=2 emanates, and  $U(r; \alpha)$  is nondegenerate if  $\alpha \neq \alpha_n$ ,  $n=0,1,\dots,n_0$ . Each non-radial branch is in  $(0,\lambda(\alpha_0))$   $\times \{u>0\} \subset \mathbf{R} \times H_0^2(B)$ .

When N=2, radial solutions of problems (2) and (3) can be written explicitly, and hence, Lin [4] and Miyamoto [5] succeeded to show the existence of bifurcation points. That is difficult even if we know exact solutions, much more difficult if we do not know them usually. When  $N \neq 2$ , we do not find exact radial solutions of (2). However, the structure of eigenvalues and eigenfunctions of the linearized problem in the dimension 1 is well-known, and then, by the comparison function introduced in [7], we can find the Morse indices of even solutions of (1). Then we obtain the existence of a symmetry-breaking bifurcation point of (1).

Let m(U) be the Morse index of a solution U to (1), that is, the number of negative eigenvalues  $\mu$  of

(4) 
$$\begin{cases} \phi'' + \lambda |x|^l e^{U(x)} \phi + \mu \phi = 0, & x \in (-1, 1), \\ \phi(-1) = \phi(1) = 0, \end{cases}$$

A solution U of (1) is said to be degenerate if  $\mu = 0$  is an eigenvalue of (4). Otherwise, it is said to be nondegenerate.

The main result is as follows.

**Theorem 1.** For each  $\alpha > 0$ , there exists a unique  $(\lambda(\alpha), U(x; \alpha))$  such that (1) with  $\lambda = \lambda(\alpha)$  has a unique positive even solution  $U = U(x; \alpha)$  such that  $||U||_{\infty} = \alpha$ . Moreover, there exist  $\alpha_*$ ,  $\alpha_1$ ,  $\alpha_2$  and  $\alpha_3$  such that  $\alpha_* < \alpha_1 \le \alpha_2 \le \alpha_3$  and the following (i)-(vii) hold:

- (i) if  $0 < \alpha < \alpha_*$ , then m(U) = 0 and  $U(x; \alpha)$  is nondegenerate;
- (ii) if  $\alpha = \alpha_*$ , then m(U) = 0 and  $U(x; \alpha)$  is degenerate;
- (iii) if  $\alpha_* < \alpha < \alpha_1$ , then m(U) = 1 and  $U(x; \alpha)$  is nondegenerate;
- (iv) if  $\alpha = \alpha_1$ , then m(U) = 1 and  $U(x; \alpha)$  is degenerate;
- (v) if  $\alpha = \alpha_2$ , then m(U) = 1,  $U(x;\alpha)$  is degenerate and  $(U,\lambda)$  is a non-even bifurcation point, that is, for each  $\varepsilon > 0$ , there exists  $(\lambda, u)$  such that u is a non-even positive solution of (1) and  $|\lambda \lambda(\alpha_2)| + ||u U(\cdot, \alpha_2)||_{\infty} < \varepsilon$ ;
- (vi) if  $\alpha = \alpha_3$ , then m(U) = 2 and  $U(x; \alpha)$  is degenerate;
- (vii) if  $\alpha > \alpha_3$ , then m(U) = 2 and  $U(x; \alpha)$  is nondegenerate.

Here and Hereafter, we use the notation  $||U||_{\infty} = \sup_{x \in [-1,1]} U(x)$ .

For the proof of Theorem 1, see [8]. Here, we give a sufficient condition for the second eigenvalue of the linearized problem to be negative for the following problem

(5) 
$$\begin{cases} u'' + \lambda h(x) f(u) = 0, & x \in (-1, 1), \\ u(-1) = u(1) = 0. \end{cases}$$

where  $\lambda > 0$  and  $h \in C^1([-1,0) \cup (0,1]) \cap C[-1,1]$ , h(-x) = h(x), h(x) > 0 and  $h'(x) \geq 0$  for x > 0,  $f \in C^1[0,\infty)$ , f(s) > 0 and  $f'(s) \geq 0$  for s > 0. Namely we will show the following result, which plays a crucial role in the proof of Theorem 1.

**Proposition 1.** Assume that, for each sufficiently large  $\alpha > 0$ , there exist  $\lambda(\alpha) > 0$  and  $U(x;\alpha)$  such that  $U(x;\alpha)$  is a positive even solution of (5) at  $\lambda = \lambda(\alpha)$  and  $\|U(\cdot;\alpha)\|_{\infty} = \alpha$ . Assume moreover that there exist  $s_0 > 0$  and  $\delta > 0$  such that

(6) 
$$\frac{l(x)(g(s)-1)-4}{g(s)+l(x)+3} \ge \delta, \quad x \in (0,1], \ s \ge s_0,$$

where l(x) = xh'(x)/h(x) and g(s) = sf'(s)/f(s). Let  $\mu_2(\alpha)$  be the second eigenvalue of

(7) 
$$\begin{cases} \phi'' + \lambda(\alpha)h(x)f'(U(x;\alpha))\phi + \mu\phi = 0, & x \in (-1,1), \\ \phi(-1) = \phi(1) = 0. \end{cases}$$

Then  $\mu_2(\alpha) < 0$  for all sufficiently large  $\alpha > 0$ .

In the case where  $h(x) = |x|^l$ , l > 0 and  $f(s) = e^s$ , it follows that l(x) = xh'(x)/h(x) = l for  $x \in (0,1]$  and g(s) = sf'(s)/f(s) = s, and hence (6) is satisfied.

We conclude that if U is a positive even solution of (1) and  $||U||_{\infty} \leq 1$ , then m(U) = 0. Indeed, let  $\mu_1$  be the first eigenvalue of (4) and let  $\phi_1$  be an eigenfunction corresponding to  $\mu_1$ . We may assume that  $\phi_1(x) > 0$  on (-1, 1). Integrating the equality

$$(\phi_1(x)U'(x) - \phi_1'(x)U(x))' = \mu_1\phi_1(x)U(x) + \lambda |x|^l e^{U(x)}\phi_1(x)(U(x) - 1)$$

on [-1,1], we have

$$\mu_1 \int_{-1}^1 \phi_1(x) U(x) dx = \lambda \int_{-1}^1 |x|^l e^{U(x)} \phi_1(x) (1 - U(x)) dx > 0.$$

Consequently, we have  $\mu_1 > 0$ , which means m(U) = 0. By applying Proposition 1, we can conclude that  $m(U(\cdot; \alpha)) = 0$  for  $0 < \alpha \le 1$  and  $m(U(\cdot; \alpha)) \ge 2$  for all sufficiently large  $\alpha > 1$ . Then, using the Leray-Schauder degree, we can find a bifurcation point.

To prove Proposition 1, we need the following two lemmas.

**Lemma 1.** Let  $\phi_2$  be an eigenfunction corresponding to the second eigenvalue  $\mu_2(\alpha)$  of (7). Then  $\phi_2$  is odd,  $\phi_2(0) = \phi_2(1) = 0$  and  $\phi_2(x) \neq 0$  for  $x \in (0,1)$ .

*Proof.* Let  $M_1$  be the first eigenvalue of

$$\begin{cases} \Phi'' + \lambda(\alpha)h(x)f'(U(x;\alpha))\Phi + M\Phi = 0, & x \in (0,1), \\ \Phi(0) = \Phi(1) = 0, & \end{cases}$$

and let  $\Phi_1$  be an eigenfunction corresponding to  $M_1$ . Then  $\Phi_1(0) = \Phi_1(1) = 0$  and  $\Phi_1(x) \neq 0$  on (0,1). Set

$$\Phi(x) = \begin{cases} \Phi_1(x), & x \in [0, 1], \\ -\Phi_1(-x), & x \in [-1, 0). \end{cases}$$

Noting the fact that  $\lim_{x\to -0} \Phi''(x) = \lim_{x\to -0} (-\Phi_1''(-x)) = -\Phi_1''(0) = 0$ , we easily check that  $\Phi$  is a solution of

$$\begin{cases} \Phi'' + \lambda(\alpha)h(x)f'(U(x;\alpha))\Phi + M_1\Phi = 0, & x \in (-1,1), \\ \Phi(-1) = \Phi(1) = 0, & \end{cases}$$

and  $\Phi$  is odd,  $\Phi(x) \neq 0$  on (0,1) and  $\Phi(0) = 0$ . Therefore,  $M_1$  is an eigenvalue of (7) and  $\Phi$  is an eigenfunction corresponding to  $M_1$ . Since  $\Phi$  has exactly one zero in (-1,1),  $M_1$  must be  $\mu_2$  and hence  $\phi_2(x)$  must be  $c\Phi(x)$  for some  $c \neq 0$ .

**Lemma 2.** Assume that  $w \in C[a,b]$  is positive and concave on (a,b). Let  $\rho \in (0,1/2)$ . Then  $w(x) \ge \rho \max_{\xi \in [a,b]} w(\xi)$  for  $x \in [(1-\rho)a + \rho b, \rho a + (1-\rho)b]$ .

*Proof.* We take  $c \in [a, b]$  for which  $w(c) = \max_{\xi \in [a, b]} w(\xi)$ . Then w(c) > 0. Since w is positive and concave on (a, b), we have

$$w(x) \ge \frac{w(c)(x-a)}{c-a} \ge \frac{w(c)(x-a)}{b-a} =: l_1(x), \quad x \in [a,c],$$

and

$$w(x) \ge \frac{w(c)(b-x)}{b-c} \ge \frac{w(c)(b-x)}{b-a} =: l_2(x), \quad x \in [c,b].$$

Hence  $w(x) \ge \min\{l_1(x), l_2(x)\}$  on [a, b]. We conclude that if  $x \in [(1-\rho)a+\rho b, (a+b)/2]$ , then

$$\min\{l_1(x), l_2(x)\} = l_1(x) \ge l_1((1-\rho)a + \rho b) = \rho w(c),$$

and if  $x \in [(a+b)/2, \rho a + (1-\rho)b]$ , then

$$\min\{l_1(x), l_2(x)\} = l_2(x) \ge l_2(\rho a + (1 - \rho)b) = \rho w(c).$$

The proof is complete.

Now we are ready to show Proposition 1.

Proof of Proposition 1. Let  $\alpha > 0$  be sufficiently large. We use the following comparison function y(x) introduced in [7]:

$$y(x) = xU(x;\alpha) - (x-1)^2 U'(x;\alpha).$$

This function y(x) satisfies y(0) = y(1) = 0, y(x) > 0 on (0,1), and

$$y'' + \lambda(\alpha)h(x)f'(U(x;\alpha))y = \lambda(\alpha)x^{-1}h(x)H(x;\alpha)f(U(x;\alpha))$$

for  $x \in (0,1]$ , where

$$H(x;\alpha) = (1-x)^2 l(x) + x(3x-4) + x^2 g(U(x;\alpha)).$$

Let  $\phi_2(x;\alpha)$  be an eigenfunction corresponding to  $\mu_2(\alpha)$ . From Lemma 1, it follows that  $\phi_2(0;\alpha) = \phi_2(1;\alpha) = 0$  and  $\phi_2(x;\alpha) \neq 0$  for  $x \in (0,1)$ . Without loss of generality, we may assume that  $\phi_2(x;\alpha) > 0$  for  $x \in (0,1)$  and  $\max_{\xi \in [0,1]} \phi_2(\xi;\alpha) = 1$ . We observe that

$$(y'\phi_2 - y\phi_2')' = \mu_2(\alpha)\phi_2 y + \lambda(\alpha)x^{-1}h(x)H(x;\alpha)f(U(x;\alpha))\phi_2, \quad x \in (0,1].$$

Integrating this equality on (0,1), we obtain

$$(8) \qquad \mu_2(\alpha) \int_0^1 \phi_2(x;\alpha) y(x) dx + \lambda(\alpha) \int_0^1 x^{-1} h(x) H(x;\alpha) f(U(x;\alpha)) \phi_2(x;\alpha) dx = 0.$$

Since

$$\begin{split} H(x) &= \left[ g(U(x;\alpha)) + l(x) + 3 \right] \left( x - \frac{l(x) + 2}{g(U(x;\alpha)) + l(x) + 3} \right)^2 \\ &+ \frac{l(x) \left[ g(U(x;\alpha)) - 1 \right] - 4}{g(U(x;\alpha)) + l(x) + 3} \\ &\geq \frac{l(x) \left[ g(U(x;\alpha)) - 1 \right] - 4}{g(U(x;\alpha)) + l(x) + 3}, \end{split}$$

we have

(9) 
$$\int_{0}^{1} x^{-1}h(x)H(x;\alpha)f(U(x;\alpha))\phi_{2}(x;\alpha)dx$$

$$\geq \int_{0}^{1} x^{-1}h(x)\frac{l(x)[g(U(x;\alpha))-1]-4}{g(U(x;\alpha))+l(x)+3}f(U(x;\alpha))\phi_{2}(x;\alpha)dx.$$

Since  $U''(x;\alpha) = -\lambda(\alpha)h(x)f(U(x;\alpha)) < 0$  on (0,1], we find that  $U'(x;\alpha)$  is decreasing in  $x \in (0,1]$ . From  $U'(0;\alpha) = 0$  it follows that  $U'(x;\alpha) < 0$  for  $x \in (0,1]$ , which implies that  $U(x;\alpha)$  is also decreasing in  $x \in (0,1]$ . Then there exists  $x(\alpha) \in (0,1)$  such that  $U(x;\alpha) \geq s_0$  for  $x \in [0,x(\alpha)]$  and  $U(x;\alpha) < s_0$  for  $x \in (x(\alpha),1]$ . Since  $U(x;\alpha)$  is concave on (0,1), we conclude that

$$U(x;\alpha) \ge \alpha(1-x), \quad x \in [0,1],$$

which shows that if  $x \in [0, (\alpha - s_0)/\alpha]$ , then  $U(x; \alpha) \ge s_0$ . Therefore,  $x(\alpha) \ge (\alpha - s_0)/\alpha$ , which implies

$$\lim_{\alpha \to \infty} x(\alpha) = 1.$$

We take  $s_1 \geq s_0$  for which  $x(\alpha) \geq 3/4$  for  $\alpha \geq s_1$ . If  $\alpha \geq s_1$ , then (6) implies

(11) 
$$\int_{0}^{x(\alpha)} x^{-1}h(x) \frac{l(x)[g(U(x;\alpha)) - 1] - 4}{g(U(x;\alpha)) + l(x) + 3} f(U(x;\alpha))\phi_{2}(x;\alpha)dx$$

$$\geq \delta f(s_{0}) \int_{0}^{x(\alpha)} x^{-1}h(x)\phi_{2}(x;\alpha)dx$$

$$\geq \delta f(s_{0}) \int_{1/4}^{3/4} x^{-1}h(x)\phi_{2}(x;\alpha)dx.$$

Recalling  $\max_{\xi \in [0,1]} \phi_2(\xi) = 1$ , we have

(12) 
$$\int_{x(\alpha)}^{1} x^{-1}h(x) \frac{l(x)[g(U(x;\alpha)) - 1] - 4}{g(U(x;\alpha)) + l(x) + 3} f(U(x;\alpha))\phi_{2}(x;\alpha)dx$$

$$\geq -\int_{x(\alpha)}^{1} x^{-1}h(x) \frac{(l(x) + 4)f(U(x;\alpha))\phi_{2}(x;\alpha)}{g(U(x;\alpha)) + l(x) + 3} dx$$

$$\geq -f(s_{0}) \int_{x(\alpha)}^{1} x^{-1}h(x) \frac{l(x) + 4}{l(x) + 3} dx.$$

Now we will show that there exists  $s_2 \geq s_1$  such that  $\mu_2(\alpha) < 0$  for  $\alpha \geq s_2$ . Assume to the contrary that there exists  $\{\alpha_n\}_{n=1}^{\infty}$  such that  $\mu_2(\alpha_n) \geq 0$  and  $\alpha_n \geq s_1$  for  $n \in \mathbb{N}$  and  $\lim_{n\to\infty} \alpha_n = \infty$ .

Since  $\phi_2(x; \alpha_n) > 0$  and

$$\phi_2''(x;\alpha_n) = -h(x)f'(U(x;\alpha_n))\phi_2(x;\alpha_n) - \mu_2(\alpha_n)\phi_2(x;\alpha_n) \le 0, \quad x \in (0,1),$$

we find that  $\phi_2(x; \alpha_n)$  is concave on (0,1). From Lemma 2 with  $\rho = 1/4$ , a = 0 and b = 1, it follows that

$$\phi_2(x; \alpha_n) \ge \frac{1}{4} \max_{\xi \in [0,1]} \phi_2(\xi; \alpha_n) = \frac{1}{4}, \quad x \in \left[\frac{1}{4}, \frac{3}{4}\right].$$

By (11), we have

(13) 
$$\int_{0}^{x(\alpha)} x^{-1}h(x) \frac{l(x)[g(U(x;\alpha)) - 1] - 4}{g(U(x;\alpha)) + l(x) + 3} f(U(x;\alpha))\phi_{2}(x;\alpha) dx$$

$$\geq \frac{\delta f(s_{0})}{4} \int_{1/4}^{3/4} x^{-1}h(x) dx.$$

Combining (8) with (9), (12) and (13), we have

$$0 \ge -\mu_2(\alpha_n) \int_0^1 \phi_2(x; \alpha_n) y(x) dx$$

$$\ge \lambda(\alpha_n) f(s_0) \left[ \frac{\delta}{4} \int_{1/4}^{3/4} x^{-1} h(x) dx - \int_{x(\alpha_n)}^1 x^{-1} h(x) \frac{l(x) + 4}{l(x) + 3} dx \right],$$

which implies

$$\int_{x(\alpha_n)}^1 x^{-1}h(x)\frac{l(x)+4}{l(x)+3}dx \ge \frac{\delta}{4} \int_{1/4}^{3/4} x^{-1}h(x)dx > 0, \quad n \in \mathbf{N}.$$

This contradicts the fact (10). Consequently, there exists  $s_2 \geq s_1$  such that  $\mu_2(\alpha) < 0$  for  $\alpha \geq s_2$ . This completes the proof.

## REFERENCES

- [1] B. Gidas, W.-M. Ni and L. Nirenberg, Symmetry and related properties via the maximum principle. Comm. Math. Phys. 68 (1979) 209–243.
- [2] J. Jacobsen and K. Schmitt, The Liouville-Bratu-Gelfand Problem for Radial Operators, J. Differential Equations 184 (2002) 283–298.

- [3] D.D. Joseph and T.S. Lundgren, Quasilinear Dirichlet problems driven by positive sources, Arch. Rational Mech. Anal. 49 (1972/73) 241-269.
- [4] S.-S. Lin, On non-radially symmetric bifurcation in the annulus, J. Differential Equations 80 (1989) 251–279.
- [5] Y. Miyamoto, Nonradial maximizers for a Hénon type problem and symmetry breaking bifurcations for a Liouville-Gel'fand problem with a vanishing coefficient, *Math. Ann.* **361** (2015) 787–809.
- [6] K. Nagasaki and T. Suzuki, Radial and nonradial solutions for the nonlinear eigenvalue problem  $\Delta u + \lambda e^u = 0$  on annuli in  $\mathbb{R}^2$ , J. Differential Equations 87 (1990) 144–168.
- [7] S. Tanaka, Morse index and symmetry-breaking for positive solutions of one-dimensional Hénon type equations, J. Differential Equations 255 (2013) 1709–1733.
- [8] S. Tanaka, Symmetry-breaking bifurcation for the one-dimensional Liouville type equation, in preparation.

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