

MONOTONICITY ESTIMATE AND GLOBAL EXISTENCE
FOR THE P-HARMONIC FLOW

(p 調和写像流に対する単調性評価と大域存在)

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1 Introduction

Let \mathcal{N} be a n -dimensional smooth compact Riemannian manifold without boundary and isometrically embedded in \mathbb{R}^l ($l > n$). For a map u from $\mathbb{R}_\infty^m := (0, \infty) \times \mathbb{R}^m$ to \mathbb{R}^l we consider the p -harmonic flow

$$\begin{cases} \partial_t u - \operatorname{div}(|Du|^{p-2} Du) + |Du|^{p-2} A(u)(Du, Du) = 0 \\ u \in \mathcal{N} \end{cases}$$

where $p \geq 2$, $u(t, x) = (u^i(t, x))$, $i = 1, \dots, l$, is a vector-valued function, defined for $(t, x) \in \mathbb{R}_\infty^m$ with values into \mathbb{R}^l . $D_\alpha = \partial/\partial x_\alpha$, $\alpha = 1, \dots, m$, $Du = (D_\alpha u^i)$ is the spatial gradient of a map u , $|Du|^2 = \sum_{\alpha=1}^m \sum_{i=1}^l (D_\alpha u^i)^2$ and $\partial_t u$ is the derivative on time t . The second fundamental form $A(u)(Du, Du)$ of $\mathcal{N} \subset \mathbb{R}^l$ is on the orthogonal complement of the tangent space $\mathcal{T}_u \mathcal{N}$ (if necessary, the manifold \mathcal{N} is assumed to be orientable). Since $u = u(t, x)$, $(t, x) \in \mathbb{R}_\infty^m$, moves on the manifold \mathcal{N} , $\partial_t u \in \mathcal{T}_u \mathcal{N}$, and thus, $\partial_t u \cdot A(u)(Du, Du) = 0$ and, by multiplying the equation by $\partial_t u$ and the divergence theorem

$$\begin{aligned} |\partial_t u|^2 - \operatorname{div}(|Du|^{p-2} Du \cdot \partial_t u) + \partial_t \frac{1}{p} |Du|^p &= 0, \\ E(u) := \int_{\mathbb{R}^m} \frac{1}{p} |Du|^p dx, \quad \frac{d}{dt} E(u(t)) &= -\|\partial_t u(t)\|_2^2 \end{aligned}$$

and thus, $E(u(t)) \searrow 0$ and $u(t)$ may converge to a constant map as $t \nearrow \infty$.

Theorem 1 (A global existence and regularity for the p -harmonic flow) *Let $p > 2$ and let u_0 be a smooth map defined on \mathbb{R}^m with values to \mathcal{N} , satisfying $E(u_0) < \infty$. Then, there exists a global weak solution u of the Cauchy problem for the p -harmonic flow with initial data u_0 , satisfying the energy inequality*

$$\|\partial_t u\|_{L^2(\mathbb{R}_\infty^m)}^2 + \sup_{0 < t < \infty} E(u(t)) \leq E(u_0).$$

Moreover, the solution u is partial regular in the following sense : For any positive number γ_0 , $2 < \gamma_0 < p$, there exists a relatively closed set \mathcal{S} in \mathbb{R}_∞^m such that u and its gradient Du are locally in time-space continuous in the complement $\mathbb{R}_\infty^m \setminus \mathcal{S}$, and the size of \mathcal{S} is also estimated by the Hausdorff measure : The set \mathcal{S} is of at most locally zero m -dimensional Hausdorff measure with respect to the time-space metric $|t|^{1/\gamma_0} + |x|$, and, furthermore, for any positive time $\tau < \infty$, the $(m - \gamma_0)$ -dimensional Hausdorff measure of $\{\tau\} \times \mathcal{S}$ with respect to the usual Euclidean metric is locally zero.

Remark. The exponent γ_0 can be as close to p as possible.

In this note we report on the global existence of a partial regular weak solution of the Cauchy problem for p -harmonic flow. We use the so-called penalty approximating equation for the p -harmonic flow, and devise new monotonicity type formulas of a local scaled

(*) The work is partially supported by JSPS KAKENHI Grant number 15K04962.

energy and establish a uniform local regularity estimate for regular solutions of those equation. The regularity criterion obtained is almost optimal, comparing with that of the corresponding stationary case.

2 Penalty approximation

In this section we explain the approximation scheme for the p -harmonic flow. We will approximate the p -harmonic flow by the solutions of the gradient flow for the so-called penalized functional, introduced in [3] for the harmonic flow case $p = 2$.

Since the manifold \mathcal{N} is smooth and compact, there exists a tubular neighborhood $\mathcal{O}_{2\delta_{\mathcal{N}}}$ with width $2\delta_{\mathcal{N}}$ of \mathcal{N} in \mathbb{R}^l such that any point $u \in \mathcal{O}_{2\delta_{\mathcal{N}}}$ has a unique nearest point $\pi_{\mathcal{N}}(u) \in \mathcal{N}$ satisfying $\text{dist}(u, \mathcal{N}) = |u - \pi_{\mathcal{N}}(u)|$ for the Euclidean distance $|\cdot|$ in \mathbb{R}^l , where the projection $\pi_{\mathcal{N}} : \mathcal{O}_{2\delta_{\mathcal{N}}} \rightarrow \mathcal{N}$ is smooth, since the manifold \mathcal{N} is smooth. The distance function $\text{dist}(u, \mathcal{N})$ is Lipschitz continuous on $u \in \mathcal{O}_{2\delta_{\mathcal{N}}}$.

Let χ be a smooth, non-decreasing real-valued function defined on $[0, \infty)$ such that $\chi(s) = s$ for $s \leq (\delta_{\mathcal{N}})^2$ and $\chi(s) = 2(\delta_{\mathcal{N}})^2$ for $s \geq 4(\delta_{\mathcal{N}})^2$. Then, the function $\chi(\text{dist}^2(u, \mathcal{N}))$ is smooth on $u \in \mathbb{R}^l$ (for the proof we refer to the recent study of the squared distance function to manifold, due to Ambrosio et al. [1, Theorem 2.1]). Its gradient at $u \in \mathcal{O}_{2\delta_{\mathcal{N}}}$ is computed as

$$D_u \chi(\text{dist}^2(u, \mathcal{N})) = 2\chi'(\text{dist}^2(u, \mathcal{N})) \text{dist}(u, \mathcal{N}) D_u \text{dist}(u, \mathcal{N}) \quad ;$$

$$D_u \text{dist}(u, \mathcal{N}) = \frac{u - \pi_{\mathcal{N}}(u)}{|u - \pi_{\mathcal{N}}(u)|}$$

parallel to the vector field $u - \pi_{\mathcal{N}}(u)$ and orthogonal to $\mathcal{T}_{\pi_{\mathcal{N}}(u)}\mathcal{N}$. We also have that, for any $u \in \mathcal{N}$ and any tangent vector $\tau \in \mathcal{T}_u\mathcal{N}$,

$$|\tau^i \tau^j D_{u^i} D_{u^j} \text{dist}(u, \mathcal{N})| \leq C(\mathcal{N}) |\tau|^2$$

(See [1, Theorem 2.2]).

For positive parameters $1 \leq K \nearrow \infty$ and $1 > \epsilon \searrow 0$, we consider the Cauchy problem in \mathbb{R}_{∞}^m with initial data u_0 for the gradient flow, called the *penalized equation*,

$$(2.1) \quad \begin{cases} \partial_t u - \Delta_{p, \epsilon} u + C_0 K \chi'(\text{dist}^2(u, \mathcal{N})) \text{dist}(u, \mathcal{N}) D_u \text{dist}(u, \mathcal{N}) = 0 \\ u(0) = u_0 \end{cases}$$

associated with the *penalized functional*, defined by

$$(2.2) \quad F_{K, \epsilon}(u) := E_{\epsilon}(u) + C_0 \frac{K}{2} \int_{\mathbb{R}^m} \chi(\text{dist}^2(u, \mathcal{N})) \, dx,$$

where the positive constant C_0 will be stipulated later, depending only on p, m and \mathcal{N} (See Lemma 8). The partial differential operator $\Delta_{p, \epsilon}$ and its corresponding energy, called the regularized p -Laplace operator and the regularized p -energy, respectively, are defined as

$$(2.3) \quad \Delta_{p, \epsilon} u := \text{div} \left((\epsilon + |Du|^2)^{\frac{p-2}{2}} Du \right) \quad ; \quad E_{\epsilon}(u) := \int_{\mathbb{R}^m} \frac{1}{p} (\epsilon + |Du|^2)^{\frac{p}{2}} \, dx.$$

We have the global existence for (2.1), by the usual Galerkin method and monotonicity of the p -Laplace operator (refer to [2]). The regularity of solutions are obtained from Hölder regularity estimates for the evolutionary p -Laplace operator, with a boundedness of the derivative of the penalty term, the last term in (2.1).

Lemma 2 (Existence for the penalty approximation) *Let $p > 2$ and let u_0 be a smooth map defined on \mathbb{R}^m with values to \mathcal{N} , satisfying $E(u_0) < \infty$. For each positive numbers K and ϵ , there exists a weak solution $u = u_{K,\epsilon}$ of the Cauchy problem for the penalized equation (2.1) such that $u = u_{K,\epsilon}$ satisfies the energy inequality*

$$(2.4) \quad \|\partial_t u\|_{L^2(\mathbb{R}^m_\infty)}^2 + \sup_{0 < t < \infty} F_{K,\epsilon}(u) \leq E_\epsilon(u_0)$$

and, that $u, Du, \partial_t u$ and D^2u are locally (Hölder) continuous on time and space (with some Hölder exponent) in \mathbb{R}^m_∞ and u satisfies the penalized equation everywhere in \mathbb{R}^m_∞ .

We will call a solution having the regularity properties as in Lemma 2, a regular solution.

3 Uniform regularity estimate

In this section we show some regularity estimates for solutions $u = u_{K,\epsilon}$ of the penalized equations (2.1).

Lemma 3 (Energy inequality) *Let u_0 be a smooth map on \mathbb{R}^m with values to \mathcal{N} , satisfying $E(u_0) < \infty$, and $u = u_{K,\epsilon}$ be a regular solution of (2.1). Then, (2.4) holds.*

Proof. The energy inequality (2.4) is shown to be valid in the proof of Lemma 2. However, as a priori estimates for regular solutions of (2.1), we naturally multiply (2.1) by $\partial_t u$ and integrate by parts on space variable in \mathbb{R}^m_T for any $T > 0$. □

Lemma 4 (Boundedness) *Let $u = u_{K,\epsilon}$ be a regular solution of (2.1). Then it holds that $\sup_{\mathbb{R}^m_\infty} |u| \leq H$, where the positive number H is so large that $B(H) \supset \mathcal{O}_{2\delta_{\mathcal{N}}}(\mathcal{N})$ in \mathbb{R}^l , where $B(H) = B(H, 0)$ is a ball in \mathbb{R}^l of radius H with center of origin 0.*

Proof. We multiply (2.1) by $u(|u|^2 - H^2)_+$ and integrate in \mathbb{R}^m_∞ , where $(f)_+$ is the positive part of a function f . Since the support of χ' is in $\mathcal{O}_{2\delta_{\mathcal{N}}}(\mathcal{N}) \subset B(H)$, $\chi'(\text{dist}^2(u, \mathcal{N}))$ is zero in $\mathbb{R}^l \setminus B(H)$. Also $u_0 \in \mathcal{N} \subset B(H)$. Hence, we have

$$\begin{aligned} & \frac{1}{4} \int_{\mathbb{R}^m} (|u(t)|^2 - H^2)_+ dx \\ & \quad + \int_{\mathbb{R}^m} (\epsilon + |Du|^2)^{\frac{p-2}{2}} \left(\frac{1}{2} |D(|u|^2 - H^2)_+|_g^2 + |Du|^2 (|u|^2 - H^2)_+ \right) dz = 0 \quad ; \\ & \frac{1}{4} \int_{\mathbb{R}^m} (|u(t)|^2 - H^2)_+^2 dx \leq 0 \end{aligned}$$

and thus, $|u(t)| \leq H$ in \mathbb{R}^m and any $t \geq 0$. □

The partial regularity is based on the so-called *small energy regularity estimate* (refer to [9, Theorems 5.1, 5.3, 5.4 ; their proofs, pp. 491-494]). The small energy regularity estimate for the p -harmonic flow in the case $p > 2$ has been recently established in [7, 8]. Our main assertion here is that the small energy regularity estimate holds uniformly for solutions of the penalized equations.

Let us denote the penalized energy density for a map u by

$$(3.1) \quad e_{K,\epsilon}(u) := \frac{1}{p} (\epsilon + |Du|^2)^{\frac{p}{2}} + \frac{K}{2} \chi(\text{dist}^2(u, \mathcal{N})) .$$

Theorem 5 (Small energy regularity estimate) *Let $p > 2$. Let λ_0, B_0 and a_0 be positive numbers satisfying the conditions*

$$(3.2) \quad \frac{6p-4}{p+2} < \lambda_0 = B_0 < p \quad ; \quad \frac{\lambda_0-2}{p-2} < a_0 \leq 1.$$

Let $u = u_{K,\epsilon}$ be a regular solution of (2.1) on $\mathbb{R}_T^m = (0, T) \times \mathbb{R}^m$ for a positive $T < \infty$, satisfying the energy bound

$$(3.3) \quad \|\partial_t u\|_{L^2(\mathbb{R}_T^m)}^2 + \sup_{0 < t < T} F_{K,\epsilon}(u) \leq C$$

for a positive number C depending only on m, p and N . Then, there exists a small positive number $R_0 < 1$, depending only on m, N, p, B_0 and a_0 , and the following holds true : Let γ_0 be any positive number satisfying

$$2 < \gamma_0 < \frac{B_0(p+2) - 4p}{p-2}.$$

If, for some small positive $R < \min\{R_0, T^{1/\lambda_0}\}$,

$$(3.4) \quad \limsup_{r \searrow 0} r^{\gamma_0 - m} \int_{\{t=T-R\lambda_0\} \times B(r,0)} e_{K,\epsilon}(u(t, x)) dx \leq 1$$

then, there holds

$$(3.5) \quad \sup_{(T-(R/4)\lambda_0, T) \times B(R/4, 0)} e_{K,\epsilon}(u(t, x)) \leq C R^{-a_0 p},$$

where the positive constant C depends only on $\gamma_0, \lambda_0, B_0, a_0, p, m$ and N .

Remark. The positive number γ_0 can be as close to p as possible, if B_0 is close to p .

The novelty here is a new *monotonicity* type estimate of a *localized* scaled energy, which may be of its own interest. Let us define our localized scaled energy in the following way: Let $T \geq 0$ and $X \in \mathbb{R}^m$ be given, and (t_0, x_0) in the parabolic like envelope

$$\{(t, x) \in \mathbb{R}_\infty^m : t - T \geq |x - X|^{\lambda_0}\} \quad ; \quad \lambda_0 > 2.$$

Hereafter the notation of double sign correspondence is used. The localized scaled energy is defined by

$$(3.6) \quad E_\pm(r) = \frac{1}{\Lambda^p} \int_{\{t=t_0 \pm \Lambda^{2-p} r^2\} \times \mathbb{R}^m} \bar{e}_{K,\epsilon}(u(t, x)) \mathcal{B}_\pm(t_0, x_0; t, x) \mathcal{C}^q(t, x) dx \quad ;$$

$$\bar{e}_{K,\epsilon}(u) := \frac{1}{p} (\epsilon + |Du|^2)^{\frac{p}{2}} + C_0 \frac{K}{2} \chi(\text{dist}^2(u, N))$$

and $\Lambda = \Lambda(r)$ is a function of a scale radius r , defined as

$$(3.7) \quad \Lambda = \Lambda(r) = r^{\frac{B_0-2}{2-p}} \quad ; \quad B_0 > \frac{6p-4}{p+2}$$

for any $r > 0$. The *forward* or *backward* in time Barenblatt like function, denoted by \mathcal{B}_+ and \mathcal{B}_- , respectively, are defined by

$$(3.8) \quad \mathcal{B}_\pm(t_0, x_0; t, x) = \frac{1}{(\mp t_0 \pm t)^{\frac{m}{B_0}}} \left(1 - \left(\frac{|x - x_0|}{2(\mp t_0 \pm t)^{\frac{1}{B_0}}} \right)^{\frac{p-1}{p-2}} \right)^{\frac{p-1}{p-2}}, \quad \mp t < \mp t_0.$$

The localized function \mathcal{C} is defined and used as

$$(3.9) \quad \mathcal{C}(t, x) := \left((t - T)^{1/\lambda_0} - |x - X| \right)_+ \quad ; \quad q > 2.$$

We call $E_+(r)$ and $E_-(r)$ the forward and backward localized scaled p -energy, respectively.

Our main ingredient is the following monotonicity type estimate of a scaled energy.

Lemma 6 (Monotonicity estimate for the backward localized scaled p -energy) *Let $p > 2$ and $q > 2$. Suppose that $t_0 - T \leq 1$. For any regular solution to (2.1) the following estimate holds for all positive numbers r, ρ , $r^{B_0} = \Lambda(r)^{2-p}r^2 < \rho^{B_0} = \Lambda(\rho)^{2-p}\rho^2 \leq \min\{1, (t_0 - T)/2\}$,*

$$(3.10) \quad E_-(r) \leq E_-(\rho) + C(\rho^\mu - r^\mu) + C \int_{t_0 - \rho^{B_0}}^{t_0 - r^{B_0}} \|\mathcal{C}^{\bar{q}}(t) \bar{e}_{K, \epsilon}(u(t))\|_{L^\infty(B((t_0 - t)^{1/B_0}, x_0))} dt,$$

where $\bar{q} = \min\{q - 2, q(p - 1)/p\}$, B_0 as in (3.7), and the positive exponent μ depends only on \mathcal{N} , m , p and B_0 , and the positive constant C depends only on the same ones as μ and q .

Lemma 7 (Monotonicity estimate for the forward localized scaled p -energy) *Let $p > 2$ and $q > 2$. Suppose that $t_0 - T \leq 1$. For any regular solution to (2.1) the following estimate holds for all positive numbers r, ρ , $r^{B_0} = \Lambda(r)^{2-p}r^2 < \rho^{B_0} = \Lambda(\rho)^{2-p}\rho^2 \leq 1$*

$$(3.11) \quad E_+(\rho) \leq (1 + r^{-c_0 B_0}) E_+(r) + C(\rho^\mu - r^\mu) + C \int_{t_0 + r^{B_0}}^{t_0 + \rho^{B_0}} \|\mathcal{C}^{\bar{q}}(t) \bar{e}_{K, \epsilon}(u(t))\|_{L^\infty(B((t - t_0)^{1/B_0}, x_0))} dt,$$

where c_0 is a positive number satisfying $c_0 > 2(p - B_0)/B_0(p - 2)$, which can be as close to $2(p - B_0)/B_0(p - 2)$ as possible, $\bar{q} = \min\{q - 2, q(p - 1)/p\}$, B_0 as in (3.7), and the positive constants μ and C have the same dependence as those in Lemma 6.

Remark. In Lemma 7, the positive number c_0 can be as close to 0 as possible, if B_0 is close to p .

We need the so-called Bochner type estimate for the penalized energy density. Here the positive constant C_0 in (2.1) is appropriately chosen.

Lemma 8 (Bochner type estimate) *Let $p > 2$ and $u = u_{K, \epsilon}$ be a regular solution to (2.1). For brevity, put $e(u) = e_{K, \epsilon}(u)$. Then, it holds in \mathbb{R}_∞^m that*

$$(3.12) \quad \begin{aligned} & \partial_t e(u) - \sum_{\alpha, \beta=1}^m D_\alpha \left((\epsilon + |Du|^2)^{\frac{p-2}{2}} \mathcal{A}^{\alpha\beta} D_\beta e(u) \right) \\ & + C_1 (\epsilon + |Du|^2)^{\frac{p-2}{2}} |D^2 u|^2 + C_2 \left| 2^{-1} K D_u \chi(\text{dist}^2(u, \mathcal{N})) \right|^2 \\ & \leq C_3 \left(1 + e(u)^{\frac{2}{p}} \right) e(u)^{2(1 - \frac{1}{p})}, \end{aligned}$$

where

$$\mathcal{A}^{\alpha\beta} := \delta^{\alpha\beta} + (p - 2) \frac{D_\alpha u \cdot D_\beta u}{\epsilon + |Du|^2},$$

the positive constants C_i ($i = 1, 2, 3$) depend on m , p and \mathcal{N} .

4 Passing to the limit

In this section we present the proof of Theorem 1, based on Theorem 5.

Let $\{\epsilon_k\}$ and $\{K_k\}$ be sequences such that $\epsilon_k \searrow 0$ and $K_k \nearrow \infty$ as $k \rightarrow \infty$. Let u_{K_k, ϵ_k} , $k = 1, 2, \dots$, be a sequence of solutions of the Cauchy problem with initial data u_0 for the penalized equations (2.1) with approximating numbers $\epsilon = \epsilon_k$ and $K = K_k$, obtained in Lemma 2. Hereafter we put $u_k = u_{K_k, \epsilon_k}$, $e_k(u_k) = e_{K_k, \epsilon_k}(u_{K_k, \epsilon_k})$, for brevity.

By the energy inequality (2.4), there exist a subsequence of $\{u_k\}$, denoted by the same notation, and the limit map u such that, as $k \rightarrow \infty$,

$$(4.1) \quad u_k \rightarrow u \quad \text{weakly } * \text{ in } L^\infty(0, \infty; W^{1,p}(\mathbb{R}^m, \mathbb{R}^l)),$$

$$(4.2) \quad \partial_t u_k \rightarrow \partial_t u \quad \text{weakly in } L^2(\mathbb{R}_\infty^m, \mathbb{R}^l),$$

$$(4.3) \quad Du_k \rightarrow Du \quad \text{weakly in } L^p_{\text{loc}}(\mathbb{R}_\infty^m, \mathbb{R}^{ml}),$$

$$(4.4) \quad \chi(\text{dist}^2(u_k, \mathcal{N})) \rightarrow 0 \quad \text{strongly in } L^2_{\text{loc}}(\mathbb{R}_\infty^m, \mathbb{R}^l),$$

$$(4.5) \quad u_k \rightarrow u \quad \text{strongly in } L^q_{\text{loc}}(\mathbb{R}_\infty^m, \mathbb{R}^l) \text{ for any } q, 1 \leq q < \frac{mp}{(m-p)_+},$$

where the strong convergence in (4.5) follows from (4.1) and (4.2) (see [2, Lemma 1.4, p. 28]). Thus, furthermore, for a subsequence $\{u_k\}$ denoted by the same notation,

$$(4.6) \quad u_k \rightarrow u, \quad \text{dist}(u_k, \mathcal{N}) \rightarrow 0 \quad \text{almost everywhere in } \mathbb{R}_\infty^m.$$

We demonstrate that the limit map u is a *partial regular* weak solution of the p -harmonic flow, as in the statement of Theorem 1. The proof is divided to several steps and proceeded.

Size estimate of the singular set Let R_0 be a sufficient small positive number, determined in Theorem 5. For τ , $0 < \tau < \infty$, and R , $0 < R < \min\{R_0, \tau^{1/\lambda_0}\}$, we put two subsets in \mathbb{R}^m as

$$(4.7) \quad \begin{aligned} \mathcal{S}(\tau, R) &:= \left\{ x_0 \in \mathbb{R}^m : \limsup_{k \rightarrow \infty} \left(\limsup_{r \searrow 0} r^{\gamma_0 - m} \int_{\{t = \tau - R^{\lambda_0}\} \times B(r, x_0)} e_k(u_k(t, x)) dx \right) \geq 1 \right\} \quad ; \\ \mathcal{T}(\tau, R) &:= \bigcap_{l=1}^\infty \bigcup_{k=l}^\infty \left\{ x_0 \in \mathbb{R}^m : \limsup_{r \searrow 0} r^{\gamma_0 - m} \int_{\{t = \tau - R^{\lambda_0}\} \times B(r, x_0)} e_k(u_k(t, x)) dx > 1/2 \right\}. \end{aligned}$$

From the definition of limit supremum on k and (4.7), we see that, for every τ , $0 < \tau < \infty$, and R , $0 < R < \min\{R_0, \tau^{1/\lambda_0}\}$,

$$(4.8) \quad \mathcal{S}(\tau, R) \subset \mathcal{T}(\tau, R).$$

Here we have the estimation of size (see [5, Theorem 2.2 ; its proof, pp. 101-103] for the proof) : It holds that, for every τ , $0 < \tau < \infty$, and R , $0 < R < \min\{R_0, \tau^{1/\lambda_0}\}$,

$$\mathcal{H}^{m-\gamma_0}(\mathcal{T}(\tau, R)) = 0$$

and so, by (4.8),

$$\mathcal{H}^{m-\gamma_0}(\mathcal{S}(\tau, R)) = 0 \quad ; \quad \mathcal{H}^{m-\gamma_0} \left(\bigcap_{0 < R < \min\{R_0, \tau^{1/\lambda_0}\}} \mathcal{S}(\tau, R) \right) = 0.$$

Let us define the *singular set* as

$$(4.9) \quad \mathcal{S} = \bigotimes_{0 < \tau < \infty} \bigcap_{0 < R < \min\{R_0, \tau^{1/\lambda_0}\}} \mathcal{S}(\tau, R),$$

where $\bigotimes_{0 < \tau < \infty}$ means the direct product of sets on positive time $\tau < \infty$. Then, for any positive $T < \infty$ and any open set K compactly contained in \mathbb{R}^m , letting $K_T = (0, T) \times K$, with respect to the time-space metric $|t|^{1/\gamma_0} + |x|$,

$$\mathcal{H}^m(\mathcal{S} \cap K_T) = \int_0^T \mathcal{H}^{m-\gamma_0} \left(\bigcap_{0 < R < R_0} \mathcal{S}(\tau, R) \cap K \right) d\tau = 0.$$

Regularity of the limit map We now show the regularity of limit map u in the complement of \mathcal{S} . Let (t_0, x_0) be in the complement of \mathcal{S} . Thus, there exist a positive $R < \min\{R_0, (t_0)^{1/\lambda_0}\}$ and an infinite family $\{u_k\}$ of regular solutions such that

$$\limsup_{r \searrow 0} r^{\gamma_0 - m} \int_{\{t=t_0 - R^{\lambda_0}\} \times B(r, x_0)} e_k(u_k(t, x)) dx < 1.$$

Then we can apply Theorem 5 for each u_k above to obtain

$$(4.10) \quad \sup_{(t_0 - (R/4)^{\lambda_0}, t_0) \times B(R/4, x_0)} e(u_k) \leq C R^{-\alpha_0 p},$$

where the positive constant C depends only on λ_0, B_0, m, p and \mathcal{N} .

Now we will show the uniform continuity of $\{u_k\}$ in $Q := (t_0 - (R/8)^{\lambda_0}, t_0) \times B(R/8, x_0)$. For this purpose we will have a local L^2 estimate of derivative of the penalty term. For any smooth function ϕ of compact support in Q , we multiply the Bochner type estimate (3.12) by ϕ^2 and integrate by parts in Q to have, letting $K = K_k, u = u_k$ and $e(u) = e_k(u_k)$,

$$(4.11) \quad \int_Q \phi^2 \left(\frac{C_1}{2} (\epsilon + |Du|^2)^{\frac{p-2}{2}} |D^2u|^2 + \frac{C_2}{2} \left| \frac{K}{2} Du \chi(\text{dist}^2(u, \mathcal{N})) \right|^2 \right) dz \\ \leq \int_Q \left(\phi |\partial_t \phi| e(u) + |D\phi|^2 \left(\frac{2p}{C_1} e(u) + \frac{2}{C_2} \epsilon(u)^{\frac{2}{p}} \right) + C_3 \phi^2 \left(1 + \epsilon(u)^{\frac{2}{p}} \right) \epsilon(u)^{2(1-\frac{1}{p})} \right) dz,$$

where we use the Cauchy inequality in the first inequality.

Let $(t_0, x_0) \subset Q$ be any point and $r \leq R/8$ be any positive number, and $Q(r) = (t_0 - r^q, t_0) \times B(r, x_0)$ with $q > 1$. In (4.11) we choose a smooth function ϕ such that $0 \leq \phi \leq 1, \phi = 1$ in $Q(r), \phi = 0$ outside $Q(2r)$, and $|D\phi| \leq C/r$ and $|\partial_t \phi| \leq C/r^q$. Thus we have, by (4.10),

$$(4.12) \quad \int_{Q(r)} \left(\frac{C_1}{2} (\epsilon + |Du|^2)^{\frac{p-2}{2}} |D^2u|^2 + \frac{C_2}{2} \left| \frac{K}{2} Du \chi(\text{dist}^2(u, \mathcal{N})) \right|^2 \right) dz \\ \leq C (r^m + r^{m+q-2} + r^{m+q}) \leq C r^m.$$

We also need the Poincaré inequality of parabolic type (refer to [6]) : Let $u = u_k$. There exists a positive constant C , depending only on m and p , such that, for any $Q(r) \subset Q$,

$$(4.13) \quad \|u - \bar{u}_{Q(r)}\|_{L^2(Q(r))}^2 \leq C \left(r^2 \|Du\|_{L^2(Q(r))}^2 + r^{-m+q-2} \|(\epsilon + |Du|^2)^{1/2}\|_{L^{p-1}(Q(r))}^{2(p-1)} \right. \\ \left. + r^{2q} \|2^{-1} K Du \chi^2(u, \mathcal{N})\|_{L^2(Q(r))}^2 \right),$$

where $\bar{u}_{Q(r)}$ is the integral mean of u in $Q(r)$.

Substituting (4.10) and (4.12) into (4.13), we have, for any $(t_0, x_0) \in Q$, any positive $r \leq R/8$, and $Q(r) = (t_0 - r^q, t_0) \times B(r, x_0)$,

$$(4.14) \quad \|u - \bar{u}_{Q(r)}\|_{L^2(Q(r))}^2 \leq C (r^{m+q+2} + r^{m+3q-2} + r^{m+2q})$$

and thus, choosing $q > 1$ in (4.14), we obtain from Campanato's isomorphism theorem (refer to [5]) that $\{u_k\}$ is uniformly Hölder continuous in Q with exponent $\min\{1, q-1, \frac{q}{2}\}$, uniformly on u_k . Thus, we see that $\{u_k\}$ is equicontinuous, and uniformly bounded in Q by Lemma 4. Therefore, by Arzela-Ascoli theorem we find for a subsequence denoted by the same notation $\{u_k\}$ and the limit map u that, as $k \rightarrow \infty$,

$$(4.15) \quad u_k \rightarrow u \quad \text{uniformly in } Q$$

and that the limit map u is uniformly continuous in Q . From (4.10) and (4.15), we obtain that, as $k \rightarrow \infty$,

$$(4.16) \quad \chi(\text{dist}^2(u_k, \mathcal{N})) \leq C/K_k \rightarrow 0 \quad \text{uniformly in } Q \implies u \in \mathcal{N} \quad \text{in } Q$$

Now we will show that the limit map u satisfies the p -harmonic flow equation in Q . From (4.10) and (4.11) we also see that $\left\{ (K_k/2) D_u \chi(\text{dist}^2(u, \mathcal{N})|_{u=u_k}) \right\}$ is bounded in $L^2(Q, \mathbb{R}^l)$ and thus, there exists a vector-valued function $\nu \in L^2(Q, \mathbb{R}^l)$ such that, as $k \rightarrow \infty$,

$$(4.17) \quad (K_k/2) D_u \chi(\text{dist}^2(u, \mathcal{N})|_{u=u_k}) \rightarrow \nu \quad \text{weakly in } L^2(Q).$$

By the continuity of u in Q the image $u(Q)$ of Q is an open subset of \mathcal{N} . Let $\mathcal{P}_{\mathcal{N}}(u(Q))$ be a neighborhood of $u(Q)$ in \mathcal{N} . Let $\tau(v)$ be any smooth tangent vector field of \mathcal{N} on $\mathcal{P}_{\mathcal{N}}(u(Q))$, $\tau(v) \in \mathcal{T}_v \mathcal{N}$ for any $v \in \mathcal{P}_{\mathcal{N}}(u(Q))$. By (4.15), we can choose a sufficiently large k_0 such that, for any $k \geq k_0$, $u_k \in \mathcal{O}_{\delta_{\mathcal{N}}}$ in Q , and $\pi_{\mathcal{N}}(u_k) \in \mathcal{P}_{\mathcal{N}}(u(Q)) \subset \mathcal{N}$ and $\tau(\pi_{\mathcal{N}}(u_k)) \in \mathcal{T}_{\pi_{\mathcal{N}}(u_k)} \mathcal{N}$ in Q , where $\mathcal{O}_{\delta_{\mathcal{N}}}$ is a tubular neighborhood of \mathcal{N} with width $\delta_{\mathcal{N}}$, and $\pi_{\mathcal{N}}$ is the nearest point projection to \mathcal{N} from the tubular neighborhood of \mathcal{N} . Thus, we have that

$$\begin{aligned} D_u \chi(\text{dist}^2(u, \mathcal{N})|_{u=u_k}) \cdot \tau(\pi_{\mathcal{N}}(u_k)) &= 2\chi' \text{dist}(u_k, \mathcal{N}) D_u \text{dist}(u, \mathcal{N})|_{u=u_k} \cdot \tau(\pi_{\mathcal{N}}(u_k)) \\ &= 0 \quad \text{in } Q, \end{aligned}$$

because $D_u \text{dist}(u, \mathcal{N})|_{u=u_k}$ is orthogonal to $\mathcal{T}_{\pi_{\mathcal{N}}(u_k)} \mathcal{N}$ for any $z \in Q$, and thus,

$$(4.18) \quad \int_Q \frac{K_k}{2} D_u \chi(\text{dist}^2(u, \mathcal{N})|_{u=u_k}) \cdot \tau(\pi_{\mathcal{N}}(u_k)) dz = 0.$$

By (4.15) and (4.17), we can take the limit as $k \rightarrow \infty$ in (4.18) to have, for any smooth tangent vector field $\tau(v)$ of \mathcal{N} on $\mathcal{P}_{\mathcal{N}}(u(Q)) \subset \mathcal{N}$, as $k \rightarrow \infty$,

$$\begin{aligned} 0 &= \int_Q \frac{K_k}{2} D_u \chi(\text{dist}^2(u, \mathcal{N})|_{u=u_k}) \cdot \tau(\pi_{\mathcal{N}}(u_k)) dz \rightarrow \int_Q \nu \cdot \tau(u) dz \\ &\implies \int_Q \nu \cdot \tau(u) dz = 0 \\ (4.19) \quad &\iff \nu(z) \perp \mathcal{T}_{u(z)} \mathcal{N} \quad \text{for any } z \in Q. \end{aligned}$$

and, thus, $\nu(z)$ is a normal vector field along $u(z)$ for any $z \in Q$. In the weak form of (2.1), for any smooth map ϕ with compact support in Q ,

$$\int_Q \left(\partial_t u_k \cdot \phi + (\epsilon_k + |Du_k|^2)^{\frac{p-2}{2}} Du_k \cdot D\phi + \frac{K_k}{2} D_u \chi(\text{dist}^2(u, \mathcal{N})|_{u=u_k}) \cdot \phi \right) dz = 0,$$

we pass to the limit as $k \rightarrow \infty$ to find that the limit map u satisfies

$$(4.20) \quad \int_Q (\partial_t u \cdot \phi + |Du|^{p-2} Du \cdot D\phi + \nu \cdot \phi) dz = 0,$$

where we use the convergence in the 1st line of (4.19) and, the strong convergence of gradients $\{Du_k\}$, obtained from (2.1) with the convergence (4.1), (4.2) and (4.17) (see [2, Theorem 2.1, pp. 31-33]). Therefore, we have that

$$(4.21) \quad \partial_t u - \Delta_p u + \nu = 0 \quad \text{almost everywhere in } Q \text{ as } L^2(Q)\text{-map.}$$

We now observe that

$$(4.22) \quad |\nu(z)| = -|Du(z)|^{p-2} Du(z) \cdot (Du(z) \cdot D_u \gamma(u)|_{u=u(z)}) \quad \text{almost every } z \in Q.$$

Let $\bar{z} = (\bar{t}, \bar{x}) \in Q$ be arbitrarily taken and fixed. Let $\gamma(v)$ be a smooth unit normal vector field of \mathcal{N} in $u(Q) \subset \mathcal{N}$ such that $\gamma(v) \in (\mathcal{T}_v \mathcal{N})^\perp$, $|\gamma(v)| = 1$ for any $v \in u(Q)$ and $\gamma(u(\bar{z})) = \nu(\bar{z})/|\nu(\bar{z})|$. We take the composite map $\gamma(u)$ of $\gamma(\cdot)$ and the limit map u , and use a test function $\gamma(u)\eta$ for any smooth real-valued function η with compact support in Q to have

$$\begin{aligned} & \int_Q (\partial_t u \cdot \gamma(u)\eta + |Du|^{p-2} Du \cdot (D\gamma(u)\eta + \gamma(u)D\eta) + \nu \cdot \gamma(u)\eta) dz = 0 \quad ; \\ & \int_Q (|Du|^{p-2} Du \cdot D\gamma(u) + \nu \cdot \gamma(u)) \eta dz = 0, \\ & \implies \nu \cdot \gamma(u) = -|Du|^{p-2} Du \cdot D\gamma(u) \quad \text{almost everywhere in } Q, \end{aligned}$$

where, in the 2nd line, we use that $\partial_t u, D_\alpha u \in \mathcal{T}_u \mathcal{N}$, $\alpha = 1, \dots, m$, and $\gamma(u) \in (\mathcal{T}_u \mathcal{N})^\perp$ in Q . The last line yields, at $z = \bar{z}$,

$$|\nu(\bar{z})| = -|Du(\bar{z})|^{p-2} Du(\bar{z}) \cdot (Du(\bar{z}) \cdot D_u \gamma(u)|_{u=u(\bar{z})}).$$

Furthermore, we find that, for a positive constant C depending only on bounds of curvature of \mathcal{N} ,

$$(4.23) \quad |\nu| \leq C |Du|^p \quad \text{almost everywhere in } Q.$$

In fact, from (4.22) we obtain

$$|\nu(z)| \leq C \max_{v \in u(Q)} |D_v \gamma(v)| |Du(z)|^p \quad \text{for almost every } z \in Q.$$

Finally, we have by (4.23) and (4.10) that

$$(4.24) \quad \begin{aligned} & \partial_t u - \Delta_p u = -\nu \in L^\infty(Q) \quad \text{almost everywhere in } Q \\ & \implies Du \text{ is locally H\"older continuous in } Q, \end{aligned}$$

where, for the last statement of gradient continuity, we refer to [4, Theorem 1.1, p. 245 ; Sect 4, p. 291 ; Sect. 1 -(ii), pp. 217-218]. The use of convergence (4.3) and (4.2) in the energy boundedness (2.4) for u_k also yields

$$(4.25) \quad \|\partial_t u\|_{L^2(\mathbb{R}^m_\geq)}^2 + \sup_{0 < t < \infty} E(u(t)) \leq E(u_0).$$

Closedness of \mathcal{S} \mathcal{S} is actually closed set in \mathcal{M}_∞ . For any $z_0 = (t_0, x_0)$ in the complement of \mathcal{S} , we can take a positive $R \leq R_0$ and an neighborhood of z_0 , $Q' := (t_0 - (R/4)^{\lambda_0}, t_0) \times B(R/4)(x_0)$, and an infinite family $\{u_k\}$ of regular solutions of (2.1), and have the uniform boundedness in Q' of gradients as in (4.10). Thus, we have that, for any solution u_k , and any $z' = (t', x')$ in $Q := (t_0 - (R/8)^{\lambda_0}, t_0) \times B(R/8)(x_0)$ and all small positive $r < R/8$,

$$(4.26) \quad r^{\gamma_0 - m} \int_{\{t=t' - (R/8)^{\lambda_0}\} \times B(r, x')} e(u_k(t, x)) dx \leq C R^{-p\alpha_0} r^{\gamma_0}$$

and thus, for any $z' = (t', x')$ in Q ,

$$\limsup_{k \rightarrow \infty} \left(\limsup_{r \searrow 0} r^{\gamma_0 - m} \int_{\{t=t' - R^{\lambda_0}\} \times B(r, x')} e(u_k(t, x)) dx \right) = 0,$$

which implies that Q is a subset of the complement of \mathcal{S} . Therefore, we see that the complement of \mathcal{S} is open and thus, \mathcal{S} is closed.

Weak solution of the p -harmonic flow The proof is based on the size estimate of singular set \mathcal{S} above. A covering argument is applied for the singular set \mathcal{S} , by use of a family of parabolic cylinders under an intrinsic scaling, depending on a size of gradient of solution. For the details see [8]. □

Acknowledgments : The authour would like to express his sincere gratitude to Professor Katsuo Matsuoka for supporting and giving him an opportunity to talk at RIMS.

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